CH111

Section D4

Lecture 8

31.08.2023

2p Orbitals

$$\psi_{2,1,0} = \psi_{2p_z} = \frac{1}{4\sqrt{2\pi}} \left(\frac{1}{a_o}\right)^{\frac{3}{2}} \left(\frac{r}{a_0}\right) e^{-\frac{r}{2a_o}} \cos\theta$$

$$\psi_{2,1,+1} = \psi_{2p_{+1}} = \frac{1}{8\sqrt{\pi}} \left(\frac{1}{a_o}\right)^{\frac{3}{2}} \left(\frac{r}{a_0}\right) e^{-\frac{r}{2a_o}} \sin\theta \ e^{i\phi}$$
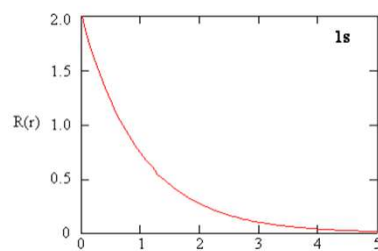
$$\psi_{2,1,-1} = \psi_{2p_{-1}} = \frac{1}{8\sqrt{\pi}} \left(\frac{1}{a_o}\right)^{\frac{3}{2}} \left(\frac{r}{a_0}\right) e^{-\frac{r}{2a_o}} \sin\theta \ e^{-i\phi}$$
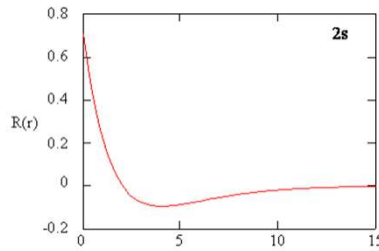
Linear combination

$$\psi_{2p_x} = \frac{1}{\sqrt{32\pi}} \left(\frac{1}{a_o}\right)^{\frac{3}{2}} \left(\frac{r}{a_0}\right) e^{-\frac{r}{2a_o}} \sin\theta \cos\phi = \frac{1}{\sqrt{2}} \left(\psi_{2,1,+1} + \psi_{2,1,-1}\right)$$

$$\psi_{2p_y} = \frac{1}{\sqrt{32\pi}} \left(\frac{1}{a_o}\right)^{\frac{3}{2}} \left(\frac{r}{a_0}\right) e^{-\frac{r}{2a_o}} \sin\theta \sin\phi = \frac{1}{\sqrt{2}i} \left(\psi_{2,1,+1} - \psi_{2,1,-1}\right)$$
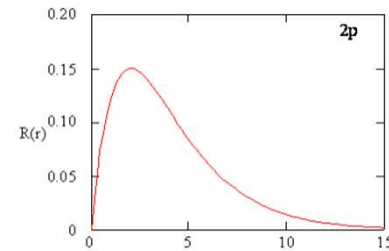
Radial functions

$$\rho = \frac{r}{a_0}$$

$$\psi_{1s}^{100} = N'e^{-\rho}$$
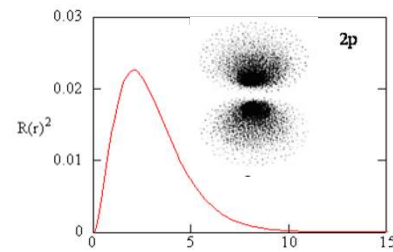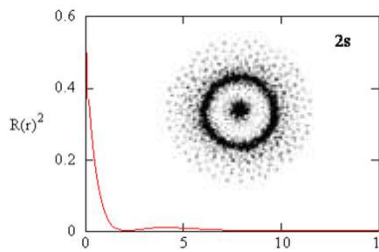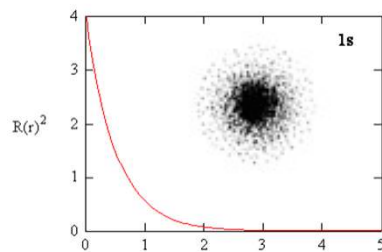
$$\psi_{2s}^{200} = N''(2 - \rho)e^{-\frac{\rho}{2}}$$

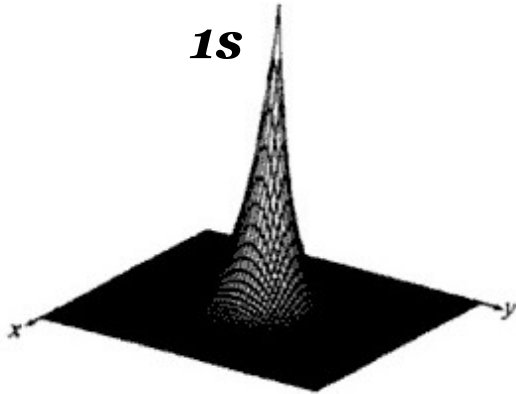$$\psi_{2p_z}^{210} = N'''\rho e^{-\frac{\rho}{2}}\cos\theta$$

For s-Orbitals the maximum probability denisty of finding the electron is on the nucleus

For s-Orbitals the probability of finding the electron on the nucleus zero
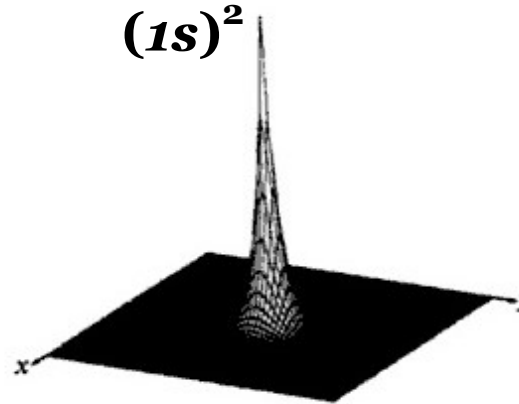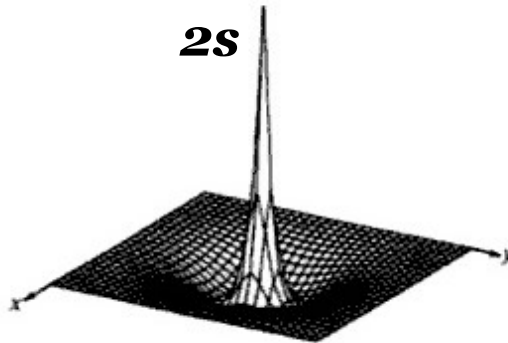
Surface plots

1s

Surface plot of the 1s wavefunction (orbital) of the hydrogen atom. The height of any point on the surface above the *xy* plane (the nuclear plane) represents the magnitude of the Ψ_{1s} function at the at point (*x,y*) in the nuclear plane. The nucleus is located in the *xy* place immediately below the 'peak'
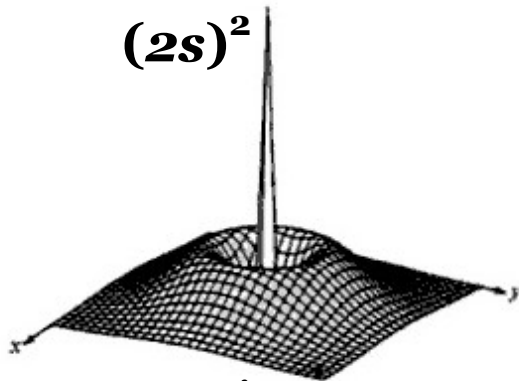
$(1s)^2$

Surface plot of the $|\Psi_{1s}|^2$; the probability density associated with the 1s wavefunction of the hydrogen atom.

2s

Surface plot of the Ψ_{2s} ; 2s wavefunction (orbital) of the hydrogen atom. The height of any point on the surface above the *xy* plane (the nuclear plane) represents the magnitude of the Ψ_{2s} function at the at point (*x,y*) in the nuclear plane. Note that there is a negative region (depression) about the nucleus; the negative region begins at $r=2a_0$ an goes asymptotically to zero at $r=\infty$.

$(2s)^2$

Surface plot of the $|\Psi_{2s}|^2$; the probability density associated with the 1s wavefunction of the hydrogen atom. Note that the negative region of the 2s plot on the left now appears as positive region.
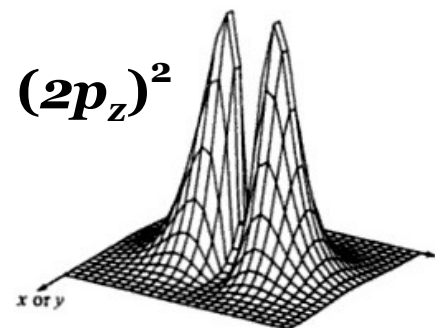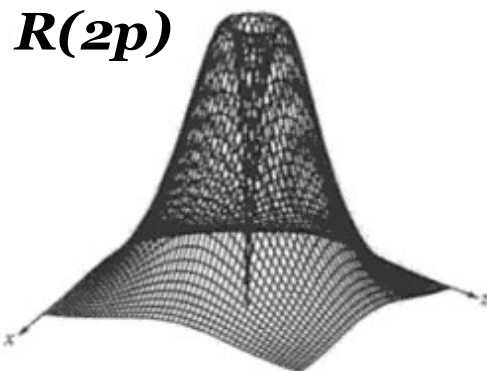
Surface plots

$2p_z$

$2p_z$

$(2p_z)^2$

-5 -4 -3 -2 -1 0 +1 +2 +3 +4 +5

z coordinate
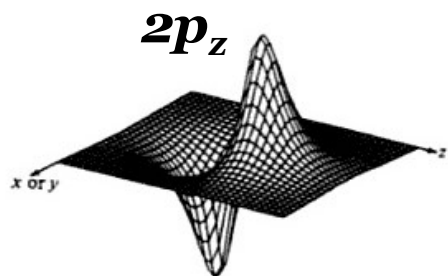
Surface plot of the *2p$_z$* wavefunction (orbital) in the *xz* (or *yz*) plane for the hydrogen atom. The 'pit' represents the negative lobe and the 'hill' the positive lobe of a 2p orbital.
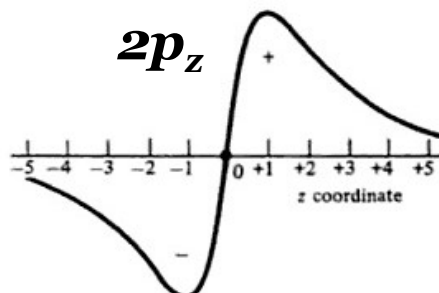
Profile of the 2p$_z$ orbital along the z-azis.

Surface plot of the $(2p_z)^2$; the probability density associated with the *2p$_z$* wavefunction of the hydrogen atom. Each of the hills represents and area in the *xz* (or *yz*) plane where the probability density is the highest, The probability density along the *x* (or *y*) axis passing through the nucleus (0,0) is everywhere zero.

$R(2p)$

$R(2p)$

Magnitude

0.50
0.45
0.40
0.35
0.30
0.25
0.20
0.15
0.10
0.05
0.00

0   1   2   3   4   5

Distance, *z*

Surface plot of radial portion of a 2p wavefunction of the hydrogen atom. The gird lines have been left transparent so that the inner 'hollow' portion is visible.
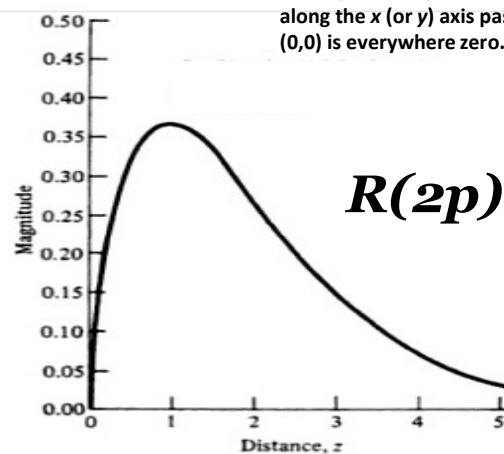
Profile of the radial portion of a 2p wavefunction of the hydrogen atom.

Surface plots

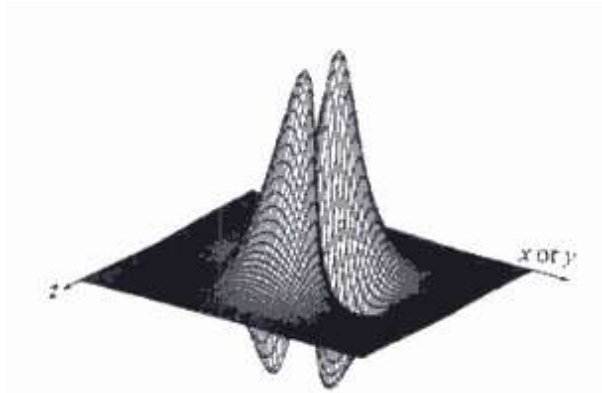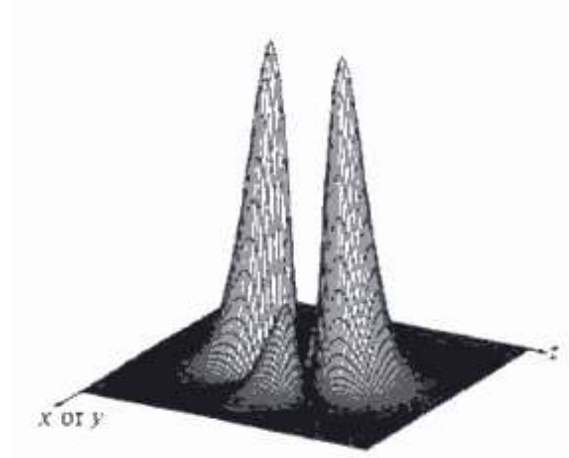Surface plot of the *3d*$_{z2}$ wavefunction (orbital) in the *xz* (or *yz*) plane for the hydrogen atom. The large hills correspond to the positive lobes and the small pits correspond to the negative lobes.

Surface plot of the (*3d*$_{z2}$)² the probability density associated with the *3d*$_{z2}$ orbital of the hydrogen atom. This figure is rotated with respect to the figure on the left so that the small hill will be clearly visible. Another smaller hill is hidden behind the large hill.
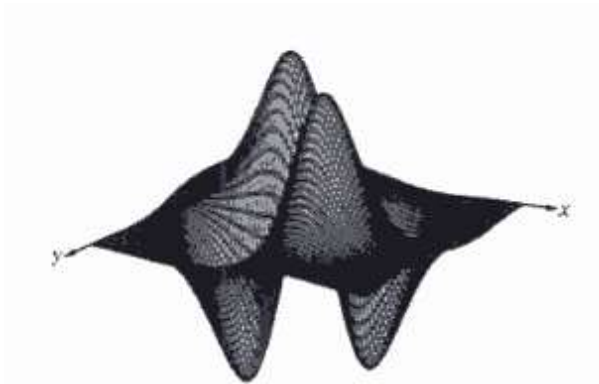
Surface plot of the *3d*$_{xy}$ wavefunction (orbital) in the *xz* plane for the hydrogen atom. The hills and the pits have same amplitude.
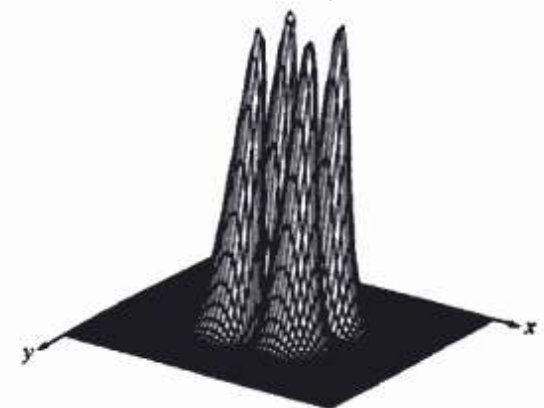
Surface plot of the (*3d*$_{xy}$)² the probability density associated with the *3d*$_{xy}$ orbital of the hydrogen atom. Pits in the figure to the left appear has hills.

Radial and Radial Distribution Functions

Probability of finding the electron anywhere in a shell of thickness dr at radius r is $4\pi r^2 R_{nl}^2(r)dr$ (for s)

$r^2 \rightarrow$ increasing function

$4\pi r^2 R_{nl}^2(r)dr \rightarrow 0$ as $4\pi r^2 dr \rightarrow 0$

$n=1, \ell=0$

$n=2, \ell=0$

$n=2, \ell=1$

1 nanometer

Product of increasing function and decreasing function: MAXIMUM

- Radial wavefunction
- Radial distribution function
- Volume element
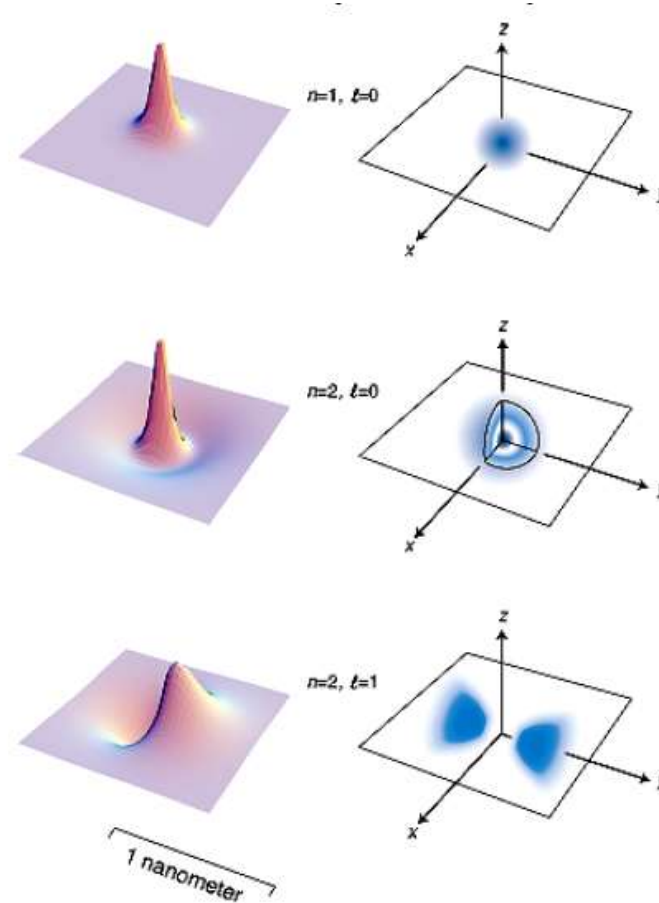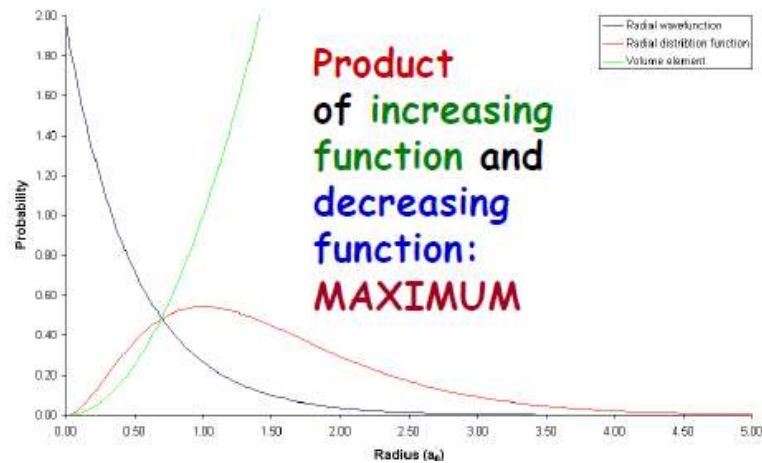
Radial Distribution Functions

$4\pi r^2 R_{nl}^2(r)$ (FOR ALL ORBITALS)

Number of radial nodes = n-l-1

3s: n=3, l=0
Nodes=2

3p: n=3, l=1
Nodes=1

3d: n=3, l=2
Nodes=2
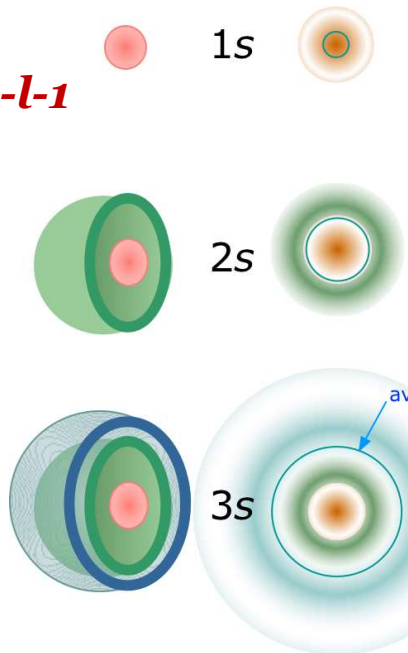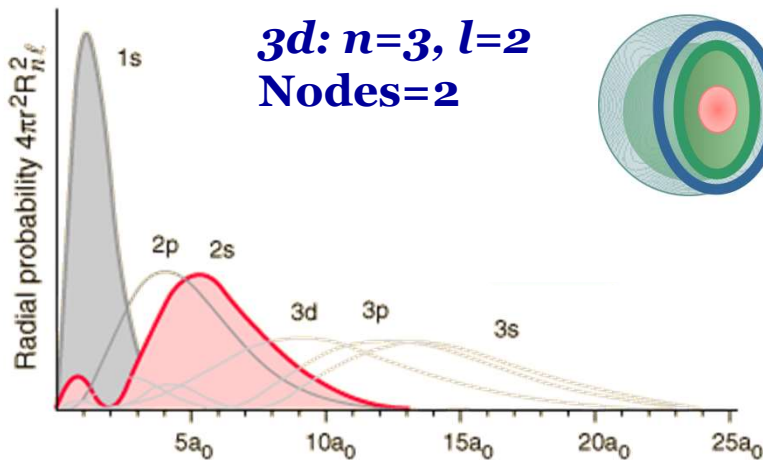
$\langle r \rangle = \langle \Psi_{ns} | r | \Psi_{ns} \rangle$
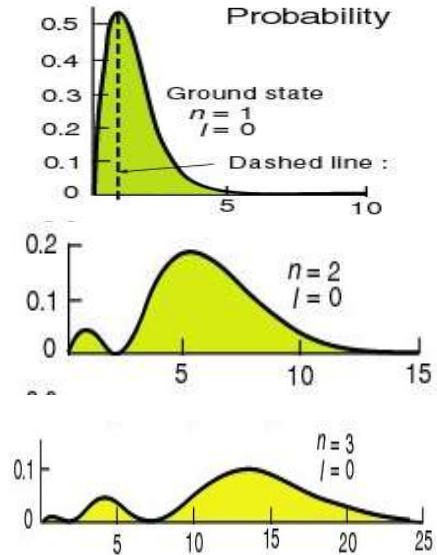
Shapes and Symmetries of the Orbitals

s-Orbitals

$$\psi_{1s} = \frac{1}{\sqrt{\pi}}\left(\frac{1}{a_o}\right)^{\frac{3}{2}} e^{-\frac{r}{a_o}} \qquad\qquad \psi_{2s} = \frac{1}{4\sqrt{2\pi}}\left(\frac{1}{a_o}\right)^{\frac{3}{2}}\left(2 - \frac{r}{a_0}\right) e^{-\frac{r}{2a_o}}$$

Function of only *r*; No angular dependence
⇒Spherical symmetric

$n\text{-}l\text{-}1=0 \longleftarrow$ radial nodes $\longrightarrow n\text{-}l\text{-}1=1$

$l=0 \longleftarrow$ angular nodes $\longrightarrow l=0$

$n\text{-}l=0 \longleftarrow$ Total nodes $\longrightarrow n\text{-}l=1$

Shapes and Symmetries of the Orbitals

p-Orbitals

$$\psi_{210} = \psi_{2p_z} = \frac{1}{4\sqrt{2\pi}} \left(\frac{1}{a_o}\right)^{\frac{3}{2}} \left(\frac{r}{a_0}\right) e^{-\frac{r}{2a_o}} \cos\theta$$

Function of only r , θ (and ϕ)
  $\Rightarrow$Not Spherical symmetric
$2p_z$ Orbital: No ϕ dependence
  $\Rightarrow$Symmetric around z-axis

radial nodes $n-l-1=0$
angular nodes $l=1$ xy nodal plane
Total nodes $n-l=1$ Zero amplitude at nucleus

Angular Distribution Functions

p-Orbitals

$$\psi_{210} = \psi_{2p_z} = \frac{1}{4\sqrt{2\pi}}\left(\frac{1}{a_o}\right)^{\frac{3}{2}}\left(\frac{r}{a_0}\right)e^{-\frac{r}{2a_o}}\cos\theta \quad m = 0 \text{ case}$$

Angular part: Polar plot of $2p_z$ --- $\cos\theta$
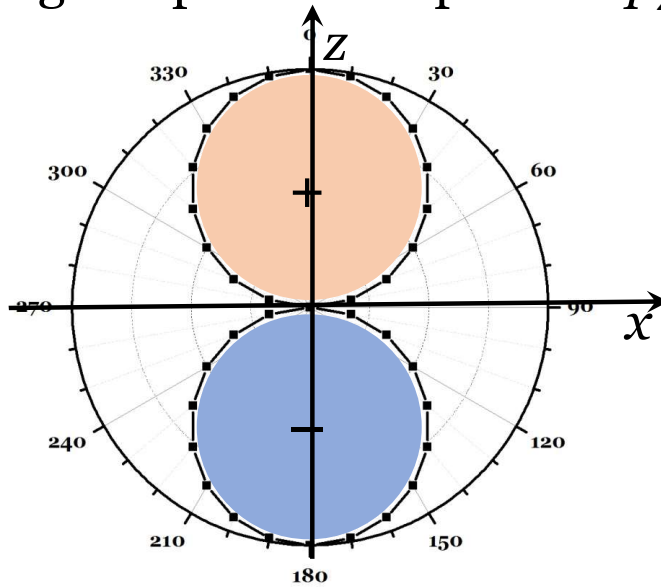
θ	cosθ
0	1.000
30	0.866
60	0.500
90	0.000
120	-0.500
150	-0.866
180	-1.000
210	-0.866
240	-0.500
270	0.000
300	0.500
330	0.866
360	1.000

Remember surface plot!
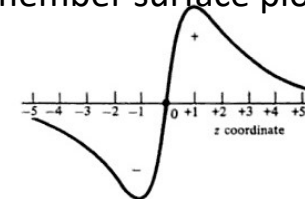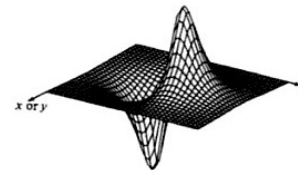
$$\psi_{210} = \psi_{2p_z} = N\rho e^{-\frac{\rho}{2}}\cos\theta$$

p-Orbitals

$$\psi_{210} = \psi_{2p_z} = N\rho e^{-\frac{\rho}{2}} \cos\theta$$

$$\psi_{2p_z} = N\rho e^{-\frac{\rho}{2}} \cos\theta$$

$$\psi_{2p_x} = N\rho e^{-\frac{\rho}{2}} \sin\theta \cos\phi$$

$$\psi_{2p_x} = N\rho e^{-\frac{\rho}{2}} \sin\theta \sin\phi$$

Color/shading are related to sign of the wavefunction

d-Orbitals

$\psi_{3d_z^2} = N_1\rho^2(3\cos^2\theta - 1)e^{-\frac{\rho}{3}}$

$\psi_{3d_{xz}} = N_2\rho^2(\sin\theta\cos\theta\cos\phi)e^{-\frac{\rho}{3}}$

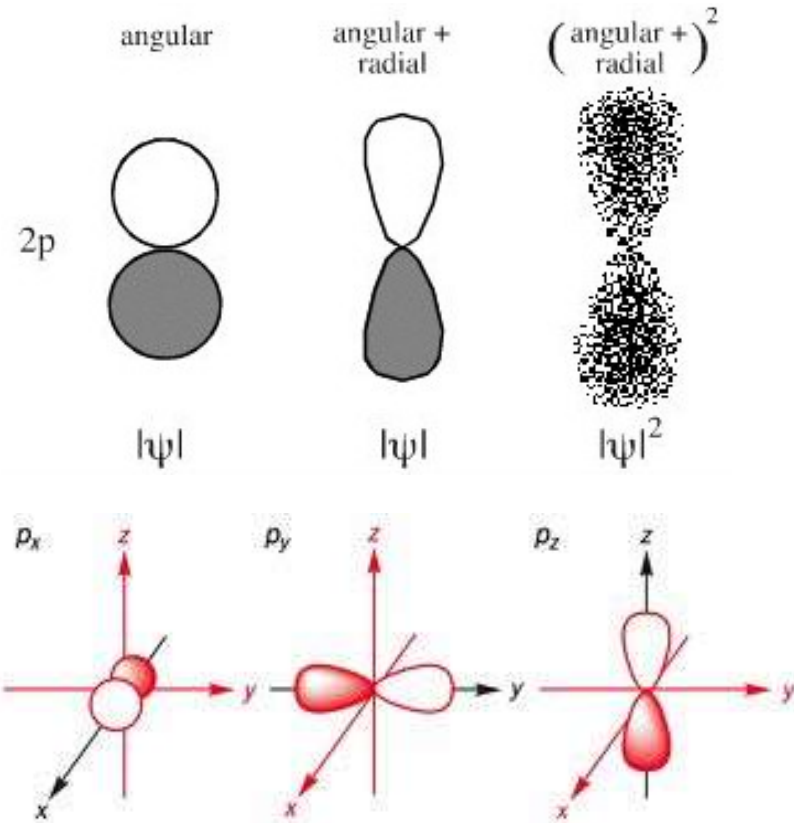$\psi_{3d_{yz}} = N_3\rho^2(\sin\theta\cos\theta\sin\phi)e^{-\frac{\rho}{3}}$

$\psi_{3d_{x^2-y^2}} = N_4\rho^2(\sin^2\theta\cos 2\phi)e^{-\frac{\rho}{3}}$

$\psi_{3d_{xy}} = N_5\rho^2(\sin^2\theta\sin 2\phi)e^{-\frac{\rho}{3}}$

Angular + Radial

Question:
Are the angular nodes for
3d x2-y2 and 3dz2 the same?

Angular part

Blue: -ve
Yellow: +ve

d_{z^2} $d_{x^2-y^2}$ d_{yz} d_{zx} d_{xy}

d_{xz}

$d_{x^2-y^2}$

d_{z^2}

f-Orbitals

$n=4; \; l=3; \; m=0, \pm 1, \; \pm 2, \; \pm 3$

$f_{z^3 - \frac{3}{5} zr^2}$ $f_{y^3 - \frac{3}{5} yr^2}$ $f_{x^3 - \frac{3}{5} xr^2}$

f_{xyz} $f_{x(z^2 - y^2)}$ $f_{y(x^2 - z^2)}$ $f_{z(x^2 - y^2)}$
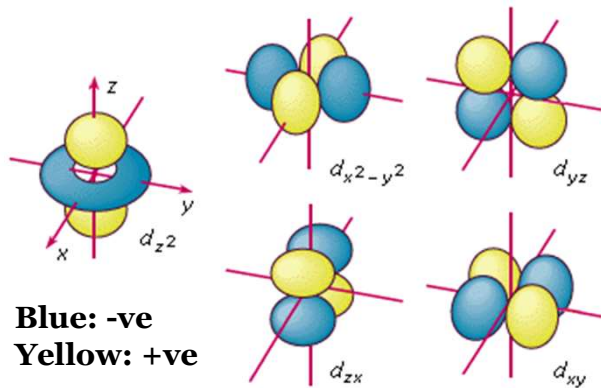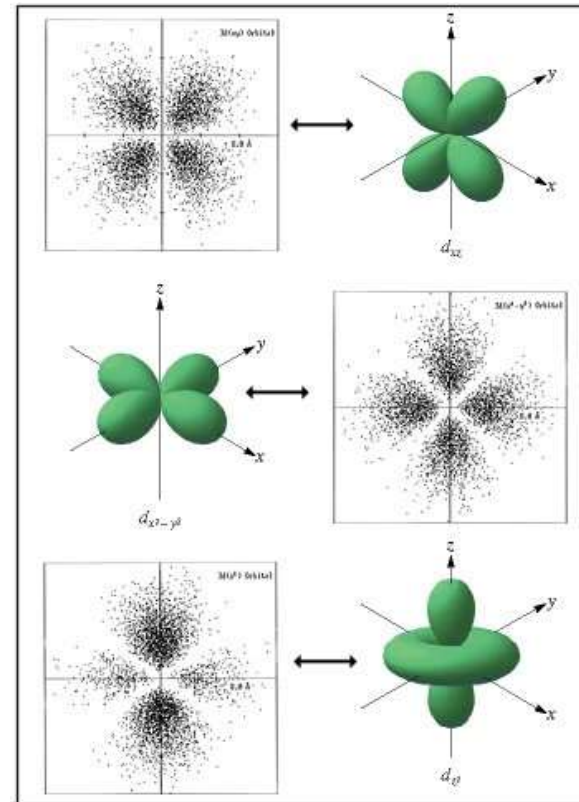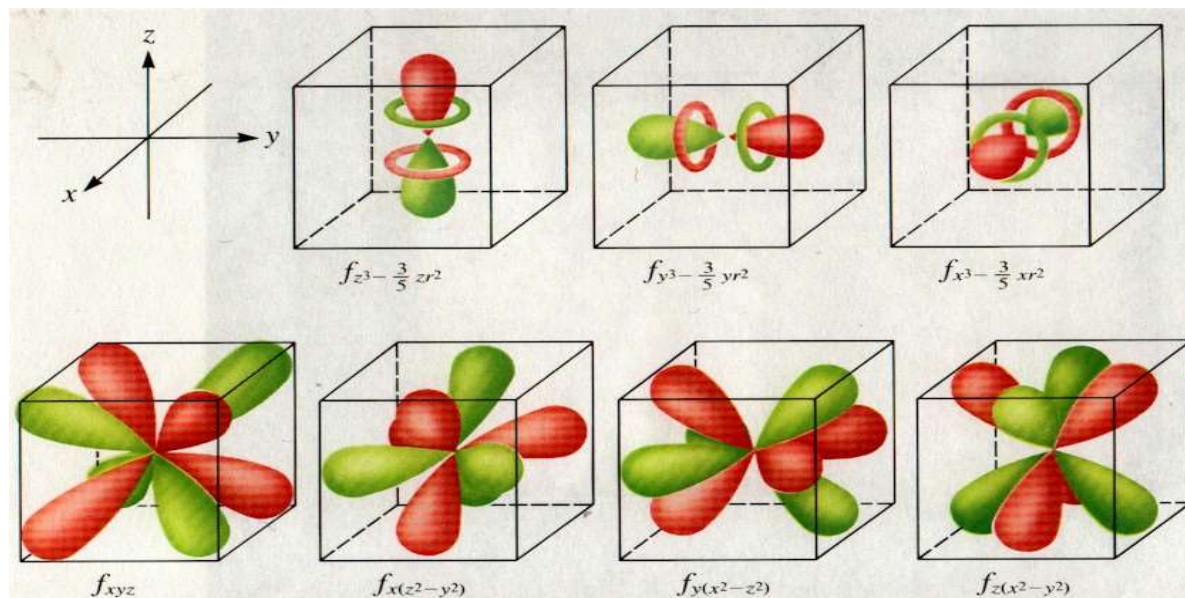
Green: -ve
Red: +ve

Cross-sections of Orbitals
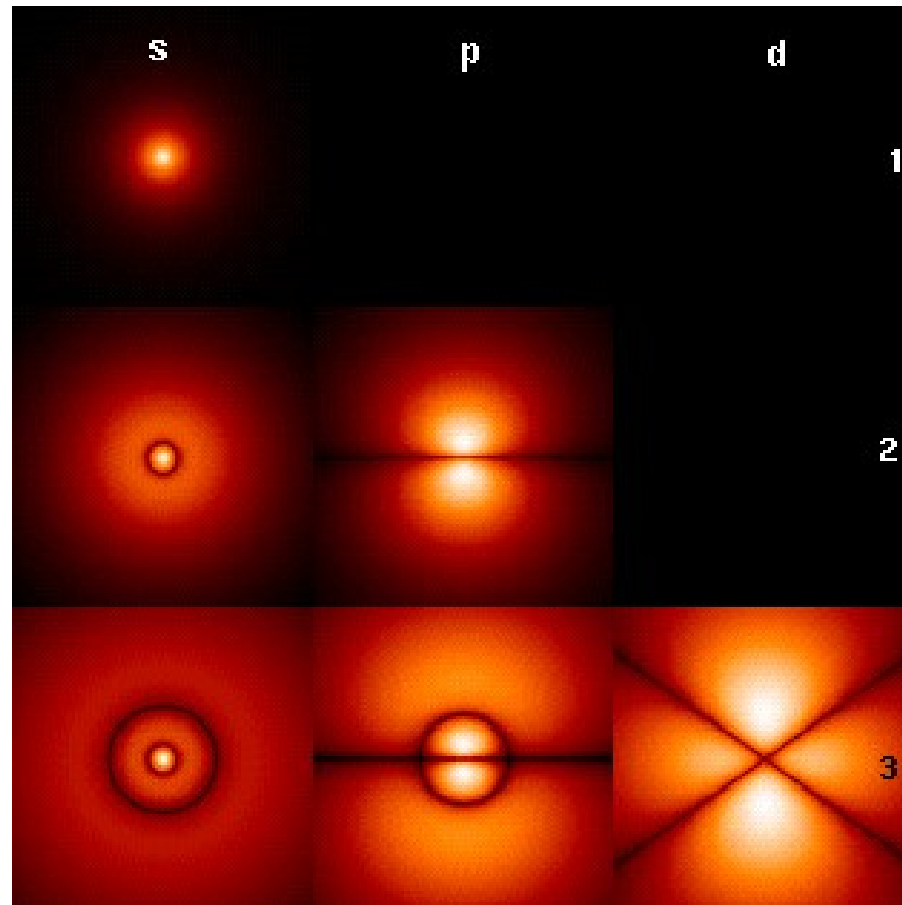
Hydrogen Wavefunctions
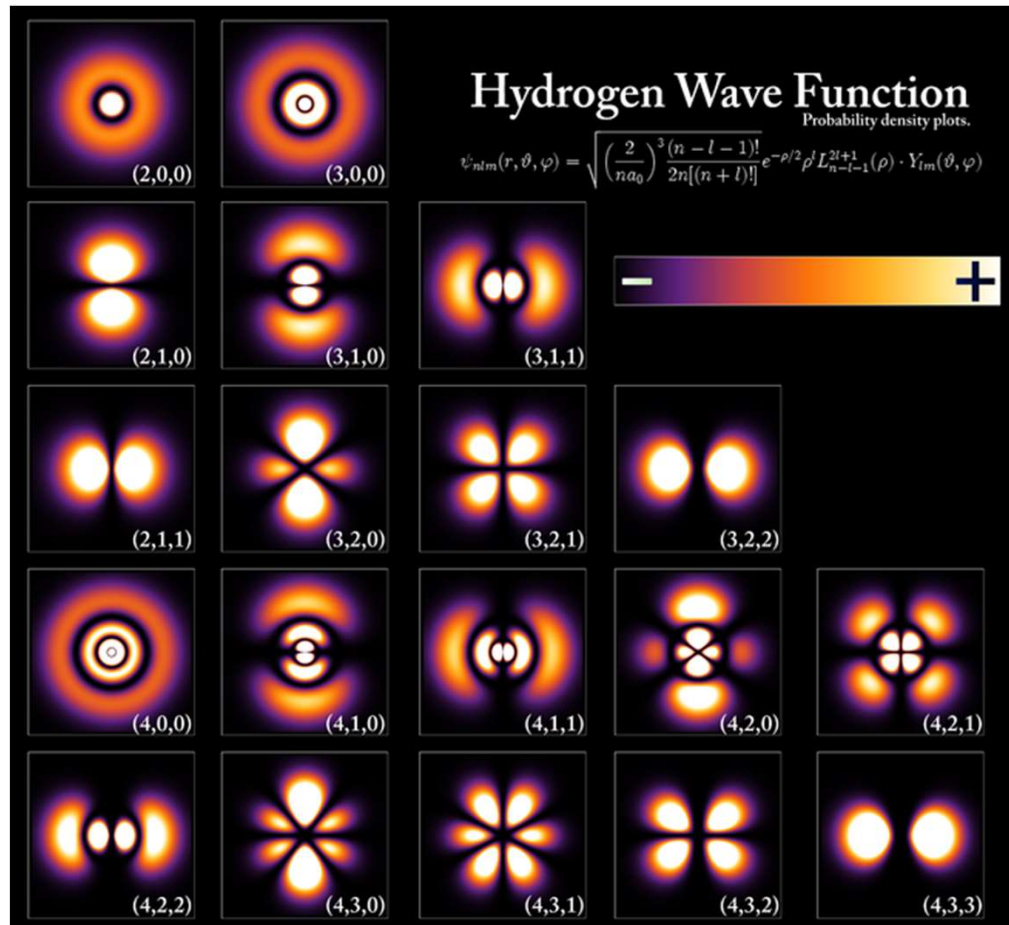
Orbitals: External Structure
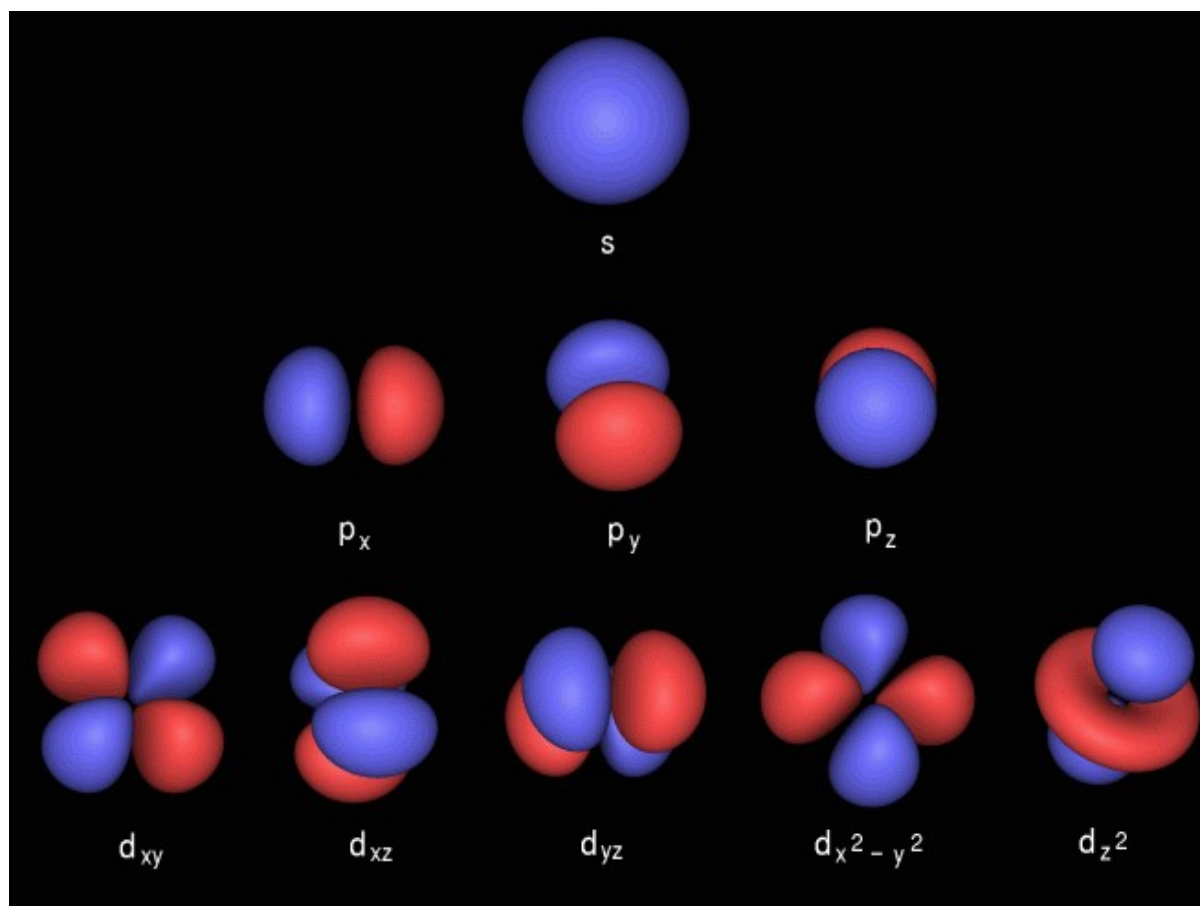
Orbitals: Internal Stucture
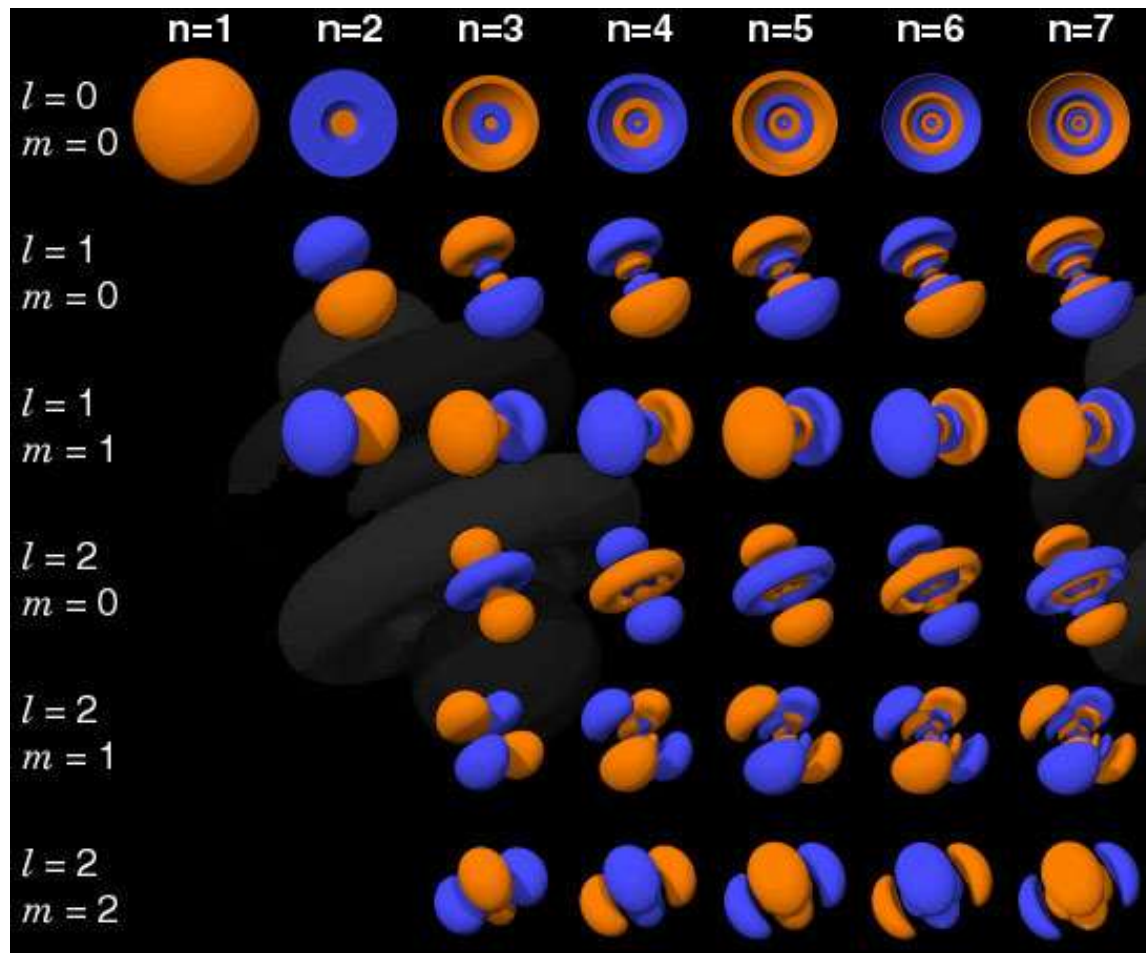
Hydrogen atom & Orbitals

Hydrogen atom has only one electron, so why bother about all these orbitals?

1. Excited states
2. Spectra
3. Many electron atoms

Many Electron Atoms

Helium is the simplest many electron atom
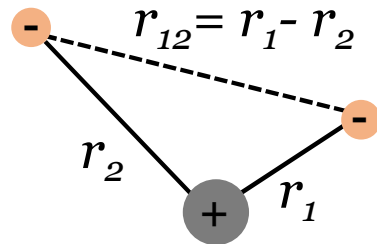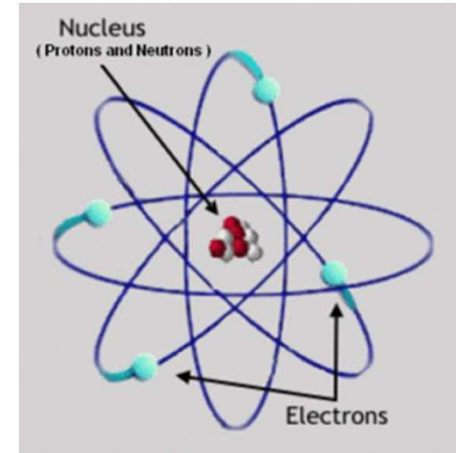
$r_{12} = r_1 - r_2$

KE of
Nucleus

KE of
Electron2

Attraction between
nucleus and Electron1

$$\hat{H} = -\frac{\hbar^2}{2m_N}\nabla_N^2 - \frac{\hbar^2}{2m_e}\nabla_1^2 - \frac{\hbar^2}{2m_e}\nabla_2^2 - \frac{1}{4\pi\varepsilon_0}\left(\frac{Z_N e^2}{r_1} + \frac{Z_N e^2}{r_2} - \frac{e^2}{r_{12}}\right)$$

KE of
Electron1

Attraction between
nucleus and Electron1

Repulsion between
Electron1 and Electron2

Helium Atom

$$\hat{H} = -\frac{\hbar^2}{2m_N}\nabla_N^2 - \frac{\hbar^2}{2m_e}\nabla_1^2 - \frac{\hbar^2}{2m_e}\nabla_2^2 - \frac{1}{4\pi\varepsilon_0}\left(\frac{Z_Ne^2}{r_1} + \frac{Z_Ne^2}{r_2} - \frac{e^2}{r_{12}}\right)$$

$$\hat{H} = -\frac{\hbar^2}{2m_N}\nabla_N^2 - \frac{\hbar^2}{2m_e}\nabla_1^2 - \frac{QZ_Ne^2}{r_1} - \frac{\hbar^2}{2m_e}\nabla_2^2 - \frac{QZ_Ne^2}{r_2} + \frac{Qe^2}{r_{12}}; Q = \frac{1}{4\pi\varepsilon_0}$$

$$\hat{H}_N = -\frac{\hbar^2}{2m_N}\nabla_N^2 \quad \hat{H}_e = -\frac{\hbar^2}{2m_e}\nabla_1^2 - \frac{QZ_Ne^2}{r_1} - \frac{\hbar^2}{2m_e}\nabla_2^2 - \frac{QZ_Ne^2}{r_2} + \frac{Qe^2}{r_{12}}$$

$$\hat{H}_N\chi_N = E_n\chi_N \qquad \hat{H}_e\psi_e = E_e\psi_e$$

Helium Atom

$$\widehat{H}_e = -\frac{\hbar^2}{2m_e}\nabla_1^2 - \frac{QZ_Ne^2}{r_1} - \frac{\hbar^2}{2m_e}\nabla_2^2 - \frac{QZ_Ne^2}{r_2} + \frac{Qe^2}{r_{12}}$$

$$\widehat{H}_e = \widehat{H}_1 + \widehat{H}_2 + \frac{Qe^2}{r_{12}}$$

$$\widehat{H}_1 = -\frac{\hbar^2}{2m_e}\nabla_1^2 - \frac{QZ_Ne^2}{r_1} \quad \text{and} \quad \widehat{H}_2 = -\frac{\hbar^2}{2m_e}\nabla_2^2 - \frac{QZ_Ne^2}{r_2}$$

The Hamiltonians $\widehat{H}_1$ and $\widehat{H}_1$ are one electron Hamiltonians similar to that of hydrogen atom

$$\widehat{H}_e\psi_e(r_1,\theta_1,\phi_1,r_2,\theta_2,\phi_2)$$
$$= \widehat{H}_1\psi_e(r_1,\theta_1,\phi_1,r_2,\theta_2,\phi_2) + \widehat{H}_2\psi_e(r_1,\theta_1,\phi_1,r_2,\theta_2,\phi_2)$$
$$+ \frac{Qe^2}{r_{12}}\psi_e(r_1,\theta_1,\phi_1,r_2,\theta_2,\phi_2)$$

Orbital Approximation

$$\psi_e(r_1, \theta_1, \phi_1, r_2, \theta_2, \phi_2) = \psi_{1e}(r_1, \theta_1, \phi_1)\psi_{2e}(r_2, \theta_2, \phi_2)$$

$$\psi_e(1,2,3,\ldots n) \approx \phi(1) \cdot \phi(2) \cdot \phi(3) \cdots \phi(n)$$

Orbital is a one electron wavefunction

The total electronic wavefunction of n number of electrons can be written as a product of n one electron wavefunctions

Helium Atom: Orbital Approximation

$$\hat{H}_e \psi_e(r_1, \theta_1, \phi_1, r_2, \theta_2, \phi_2)$$
$$= \hat{H}_1 \psi_e(r_1, \theta_1, \phi_1, r_2, \theta_2, \phi_2) + \hat{H}_2 \psi_e(r_1, \theta_1, \phi_1, r_2, \theta_2, \phi_2)$$
$$+ \frac{Qe^2}{r_{12}} \psi_e(r_1, \theta_1, \phi_1, r_2, \theta_2, \phi_2)$$

$$\psi_e(r_1, \theta_1, \phi_1, r_2, \theta_2, \phi_2) = \psi_{1e}(r_1, \theta_1, \phi_1)\psi_{2e}(r_2, \theta_2, \phi_2)$$

$$\hat{H}_e \psi_e = \hat{H}_1 \psi_{1e}(r_1, \theta_1, \phi_1)\psi_{2e}(r_2, \theta_2, \phi_2) + \hat{H}_2 \psi_{1e}(r_1, \theta_1, \phi_1)\psi_{2e}(r_2, \theta_2, \phi_2)$$
$$+ \frac{Qe^2}{r_{12}} \psi_{1e}(r_1, \theta_1, \phi_1)\psi_{2e}(r_2, \theta_2, \phi_2)$$

Helium Atom: Orbital Approximation

$$\widehat{H}_e\psi_e = \widehat{H}_1\psi_{1e}(r_1,\theta_1,\phi_1)\psi_{2e}(r_2,\theta_2,\phi_2) + \widehat{H}_2\psi_{1e}(r_1,\theta_1,\phi_1)\psi_{2e}(r_2,\theta_2,\phi_2)$$
$$+ \frac{Qe^2}{r_{12}}\psi_{1e}(r_1,\theta_1,\phi_1)\psi_{2e}(r_2,\theta_2,\phi_2)$$

$$\widehat{H}_e\psi_e$$
$$= \varepsilon_1\psi_{1e}(r_1,\theta_1,\phi_1)\psi_{2e}(r_2,\theta_2,\phi_2) + \varepsilon_2\psi_{1e}(r_1,\theta_1,\phi_1)\psi_{2e}(r_2,\theta_2,\phi_2)$$
$$+ \frac{Qe^2}{r_{12}}\psi_{1e}(r_1,\theta_1,\phi_1)\psi_{2e}(r_2,\theta_2,\phi_2)$$

$$\widehat{H}_e\psi_e = \left(\varepsilon_1 + \varepsilon_2 + \frac{Qe^2}{r_{12}}\right)[\psi_{1e}(r_1,\theta_1,\phi_1)\psi_{2e}(r_2,\theta_2,\phi_2)]$$

$$\varepsilon_1 = \varepsilon_2 = -\frac{Z^2\mu e^4}{8\varepsilon_0^2 h^2 n^2} = \frac{-13.6Z^2}{n^2}eV$$

Helium Atom: Orbital Approximation

$$\hat{H}_e \psi_e = \left(\varepsilon_1 + \varepsilon_2 + \frac{Qe^2}{r_{12}} \right) [\psi_{1e}(r_1, \theta_1, \phi_1)\psi_{2e}(r_2, \theta_2, \phi_2)]$$

If we ignore the term $\dfrac{Qe^2}{r_{12}}$

$$\hat{H}_e \psi_e = (\varepsilon_1 + \varepsilon_2)[\psi_{1e}(r_1, \theta_1, \phi_1)\psi_{2e}(r_2, \theta_2, \phi_2)]$$

$$E_{He} = \varepsilon_1 + \varepsilon_2 = -108.8 eV$$

$$\psi_e = \left(\frac{1}{\sqrt{\pi}} \left(\frac{Z}{a_o} \right)^{\frac{3}{2}} e^{-\frac{Zr}{a_o}} \right) \cdot \left(\frac{1}{\sqrt{\pi}} \left(\frac{Z}{a_o} \right)^{\frac{3}{2}} e^{-\frac{Zr}{a_o}} \right) = \psi_{1s}(1) \cdot \psi_{1s}(2)$$

Helium Atom: Orbital Approximation

$$E_{He} = \varepsilon_1 + \varepsilon_2 = -(54.4 + 54.4)eV = -108.8eV$$

$$E_{He} = -(24.59 + 54.4)eV = -78.99eV \text{ (Experimental)}$$

Ignoring $\dfrac{Qe^2}{r_{12}}$ is not justified! Need better approximation

Many Electron Atoms

$$\widehat{H} = -\frac{\hbar^2}{2m_\alpha}\nabla_N^2 - \frac{\hbar^2}{2m_e}\sum_{i=1}^{n}\nabla_i^2 - QZ_Ne^2\sum_{i=1}^{n}\frac{1}{r_i} + Qe^2\sum_{i=1}^{n}\sum_{j>i}^{n}\frac{1}{r_{ij}}$$

Nuclei are static

$$\widehat{H}_e = -\frac{\hbar^2}{2m_e}\sum_{i=1}^{n}\nabla_i^2 - QZ_Ne^2\sum_{i=1}^{n}\frac{1}{r_i} + Qe^2\sum_{i=1}^{n}\sum_{j>i}^{n}\frac{1}{r_{ij}}$$

Inter-electron repulsion term leads to deviation from the hydrogen atom. Unfortunately **CANNOT** be ignored

$$\widehat{H}_e = \sum_{i=1}^{n}H_i + Qe^2\sum_{i=1}^{n}\sum_{j>i}^{n}\frac{1}{r_{ij}}$$

$$f\left(\frac{1}{r_{ij}}\right) \neq g\left(\frac{1}{r_i}\right) + h\left(\frac{1}{r_j}\right)$$

$\frac{1}{r_{ij}}$ term in the Hamiltonian is not separable

Many Electron Atoms

$$\widehat{H}_e = -\frac{\hbar^2}{2m_e} \sum_{i=1}^{n} \nabla_i^2 - QZ_Ne^2 \sum_{i=1}^{n} \frac{1}{r_i} + Qe^2 \sum_{i=1}^{n} \sum_{j>i}^{n} \frac{1}{r_{ij}}$$

$$\widehat{H}_e = \sum_{i=1}^{n} H_i + Qe^2 \sum_{i=1}^{n} \sum_{j>i}^{n} \frac{1}{r_{ij}}$$

Hamiltonian is no longer spherically symmetric and the Time-Independent Schrodinger Equation (**TISE**) cannot be solved using analytical techniques

Numerical methods must be used solve the **TISE**