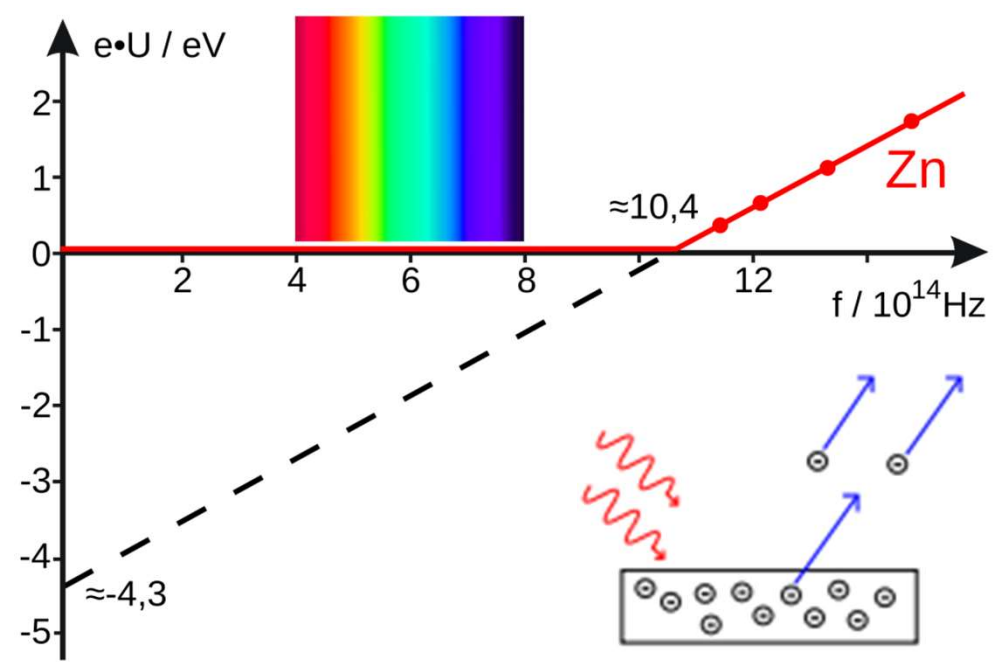
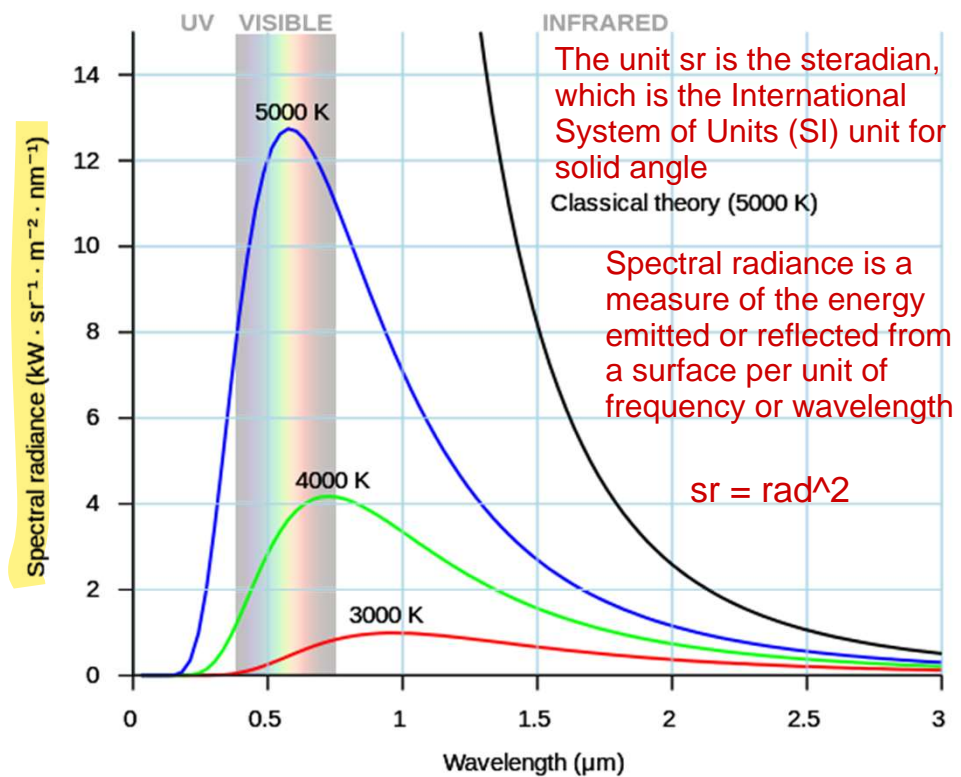


# “Old Quantum Theory”: historical perspective



Quantum  
Stat. Mech.

**Black-body radiation**  
(Planck, 1900)

**Specific heats of solids**  
(Einstein, 1907)

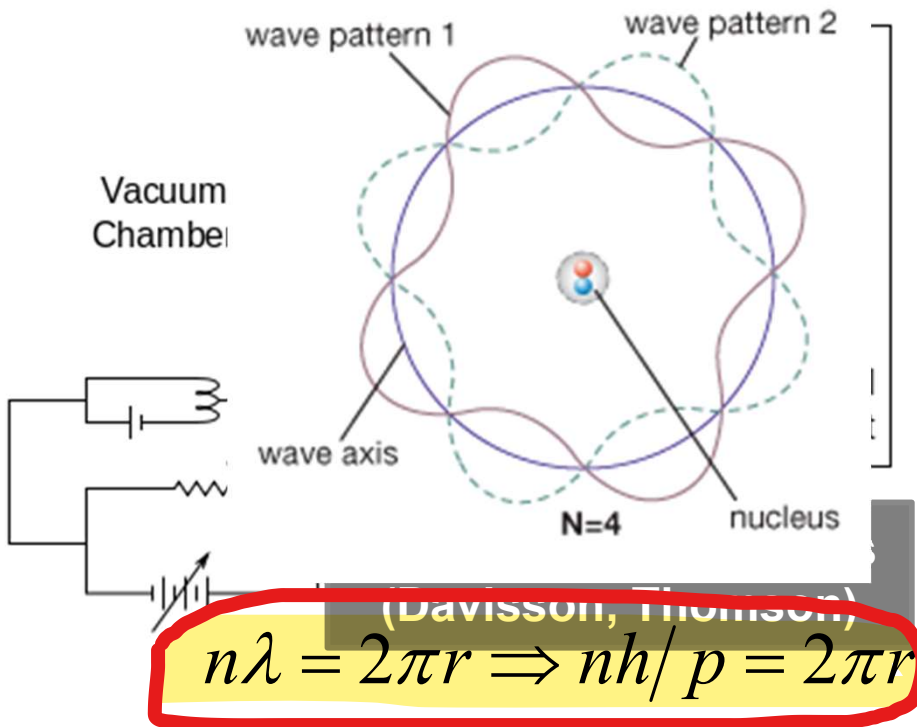
Light emitted and absorbed by  
'oscillators' in quanta:  $E_n - E_{n-1} = h\nu$

**Photoelectric Effect**  
(Einstein, 1905)

Special Relativity (Einstein, 1905):  $E^2 = (m_0 c^2)^2 + (pc)^2$

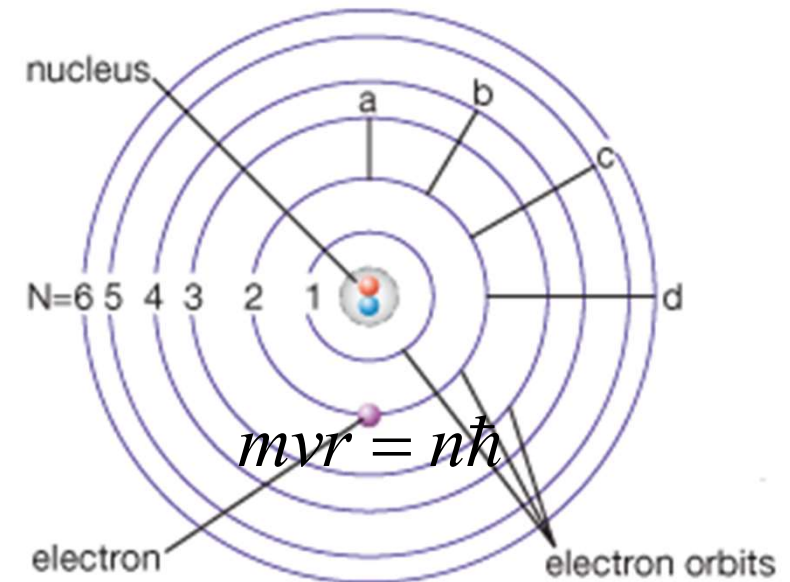
Light is a particle:  $E = h\nu$   
 $\rightarrow p = h/\lambda$

# “Old Quantum Theory”: historical perspective



**Matter waves** (de Broglie, 1924):  $\lambda = h/p$

**Electrons are waves!**



Hydrogen emission spectra (Lyman, Balmer, Paschen, Brackett, Pfund)

**Hydrogen atom**  
(Bohr, 1913)

$$E_n = -\frac{m_0 q^4}{2(4\pi\epsilon_0 n\hbar)^2} = \frac{-13.6\text{eV}}{n^2}$$

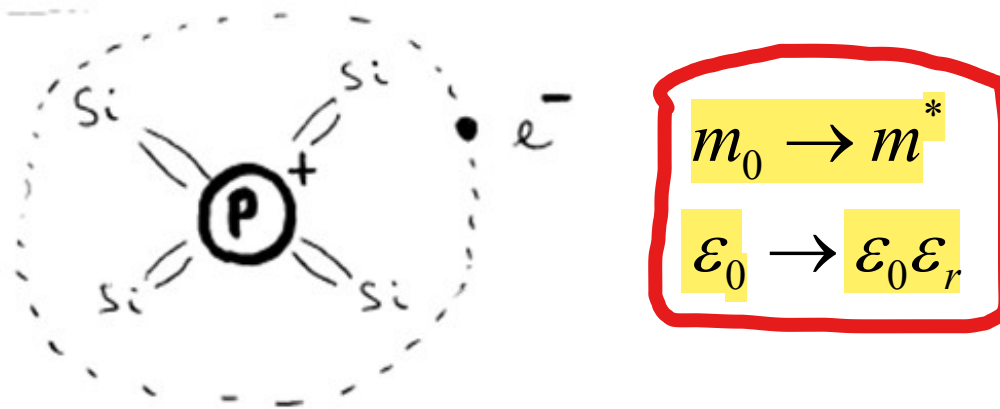
$$r_n = \frac{4\pi\epsilon_0 (n\hbar)^2}{m_0 q^2} = 0.53\text{\AA} \cdot n^2$$

## Question 0

### Bohr Model for hydrogen atom

$$r_n = \frac{4\pi\epsilon_0 (n\hbar)^2}{m_0 q^2} = 0.53\text{\AA} \cdot n^2 \quad E_n = -\frac{m_0 q^4}{2(4\pi\epsilon_0 n\hbar)^2} = \frac{-13.6\text{eV}}{n^2}$$

### "HYDROGENIC" MODEL FOR DOPANTS



Dopants give rise to carriers because:

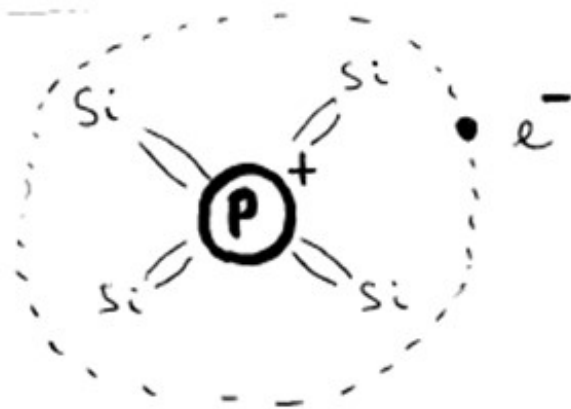
- a. That is their job, and they understand their job
- b. It is more difficult to ionize inside a semiconductor
- ☒ c. It is easier to ionize inside a semiconductor than in vacuum
- d. It is always thermodynamically favorable to give up the carrier

## Question 0

Typically, in a semiconductor:  $m^* \sim 0.1m_0, \epsilon_r \sim 10$   
 $\Rightarrow r_d \sim 5nm, E_d \sim 10meV$

Will dopant ionization happen at all temperatures?

### "HYDROGENIC" MODEL FOR DOPANTS



$$m_0 \rightarrow m^*$$

$$\epsilon_0 \rightarrow \epsilon_0 \epsilon_r$$

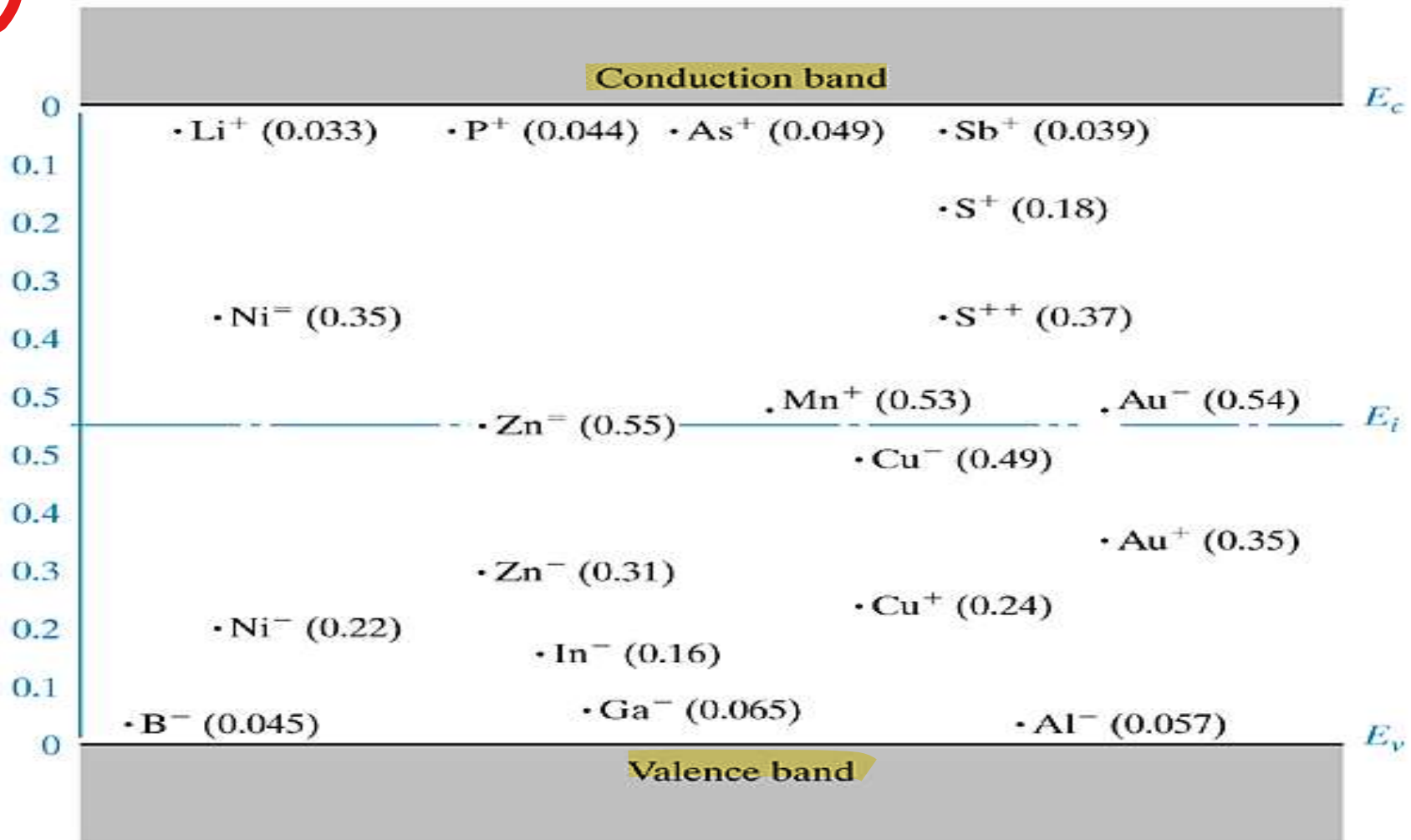
$$r_d = 0.53A \frac{\epsilon_r}{(m^*/m_0)}$$

$$E_d = -13.6eV \frac{(m^*/m_0)}{\epsilon_r^2}$$

Dopants give rise to carriers because:

- That is their job, and they understand their job
- It is more difficult to ionize inside a semiconductor
- It is easier to ionize inside a semiconductor than in vacuum
- It is always thermodynamically favorable to give up the carrier

# Question 0



Real impurities: shallow vs. deep, donor/acceptor vs. trap

# Schrödinger Wave Equation

Plane wave amplitude

$$\Psi \sim \exp[i(kx - \omega t)]$$

Momentum

$$p = \frac{h}{\lambda} = \frac{2\pi\hbar}{2\pi/k} = \hbar k$$

Energy

$$E = h\nu = \hbar\omega$$

$$-i\hbar \frac{\partial \Psi}{\partial x} = (\hbar k) \Psi$$

$$i\hbar \frac{\partial \Psi}{\partial t} = (\hbar\omega) \Psi$$

$$H = \frac{p^2}{2m} + V$$

$$p \rightarrow -i\hbar \frac{\partial}{\partial x}$$

$$E \rightarrow i\hbar \frac{\partial}{\partial t} \cong H$$

$$i\hbar \frac{\partial \Psi}{\partial t} = H\Psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V\Psi$$

$$H\psi = E\psi$$

Time-dependent  
Schrodinger Equation

Time-independent  
Schrodinger Equation

# Solutions to the Schrödinger Equation

$$H_n \psi_n = E_n \psi_n \quad \Psi = \psi_n(x) e^{-iE_n t/\hbar} \rightarrow \sum_n c_n \psi_n(x) e^{-iE_n t/\hbar}$$

Probability of finding particle in the neighborhood of  $x$ :

$$d^3x |\psi_n|^2 = d^3x \psi_n^* \psi_n$$

Orthonormal basis

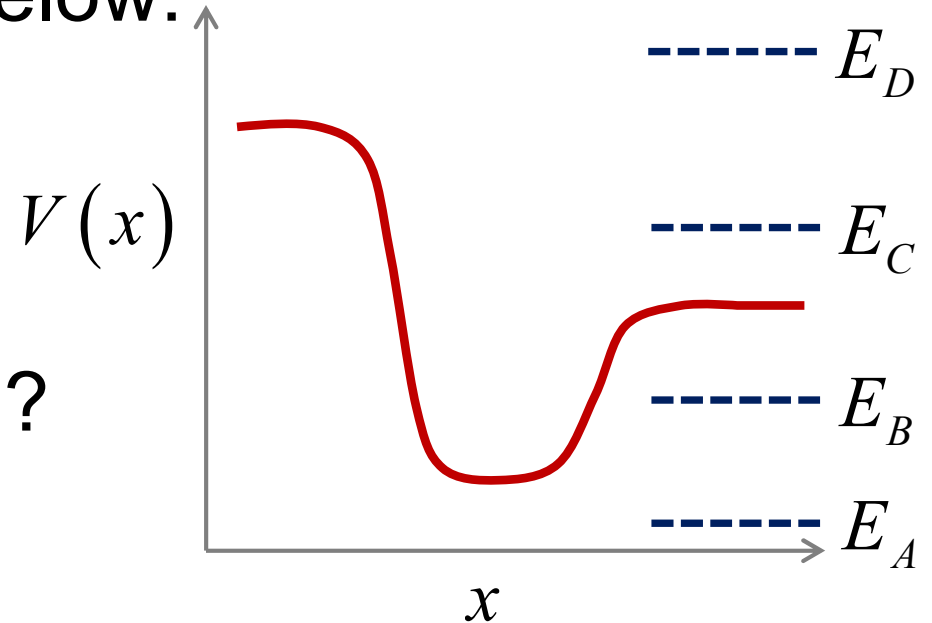
$$\int_x d^3x \psi_m^* \psi_n = \delta_{mn} = \begin{cases} 0 & \Leftarrow m \neq n \\ 1 & \Leftarrow m = n \end{cases}$$

Expectation value of operator  $A$  (e.g.  $p$ ,  $p^2/2m + V$ ) in state  $\psi_n$

$$\int_x d^3x \psi_n^* \hat{A} \psi_n$$

## Question 1

Consider the potential  $V(x)$  and the states  $A, B, C$  and  $D$  whose energies are shown below:



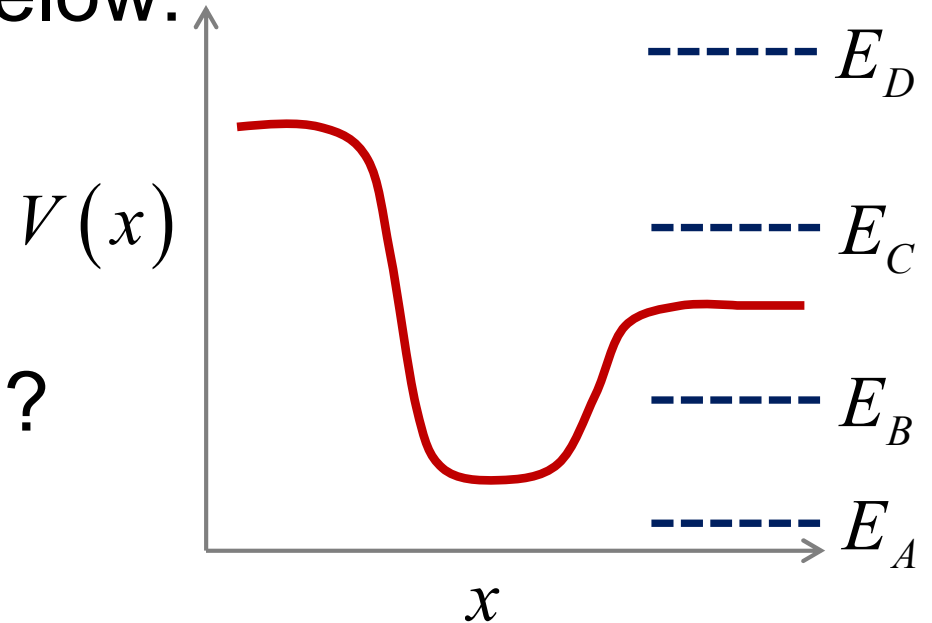
Which of the following is false?

- a. State A cannot exist.
- b. State B is a bound state.
- c. State C is a bound state.
- d. State D is a scattering state.



## Question 1

Consider the potential  $V(x)$  and the states  $A, B, C$  and  $D$  whose energies are shown below:



Which of the following is false?

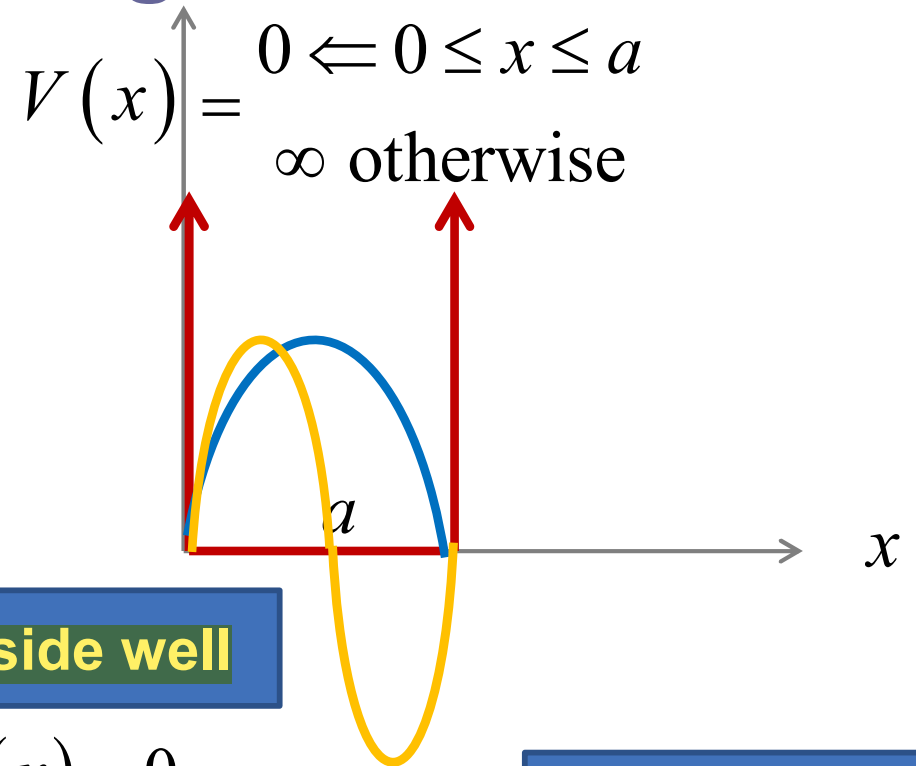
- a. State A cannot exist.
- b. State B is a bound state.
- c. State C is a bound state.
- d. State D is a scattering state.

# The infinite rectangular well

$$H\psi = E\psi$$

$$\Rightarrow \left( \frac{p^2}{2m} + V \right) \psi = E\psi$$

$$\Rightarrow \frac{-\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(x)\psi = E\psi$$



**Inside well**

$$\frac{-\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi \Rightarrow \frac{d^2\psi}{dx^2} = -k^2\psi$$

$$\Rightarrow \psi = A \sin(kx) + B \cos(kx), \quad k = \sqrt{2mE/\hbar^2}$$

**Outside well**

$$\psi(x) = 0$$

**Normalization**

$$\int_{-\infty}^{\infty} dx |\psi_n|^2 = 1$$

$$\int_{-\infty}^{\infty} dx |A|^2 \sin^2\left(\frac{n\pi x}{a}\right) = 1$$

$$\Rightarrow A = \sqrt{2/a}$$

**Boundary condition:  
ψ is continuous**

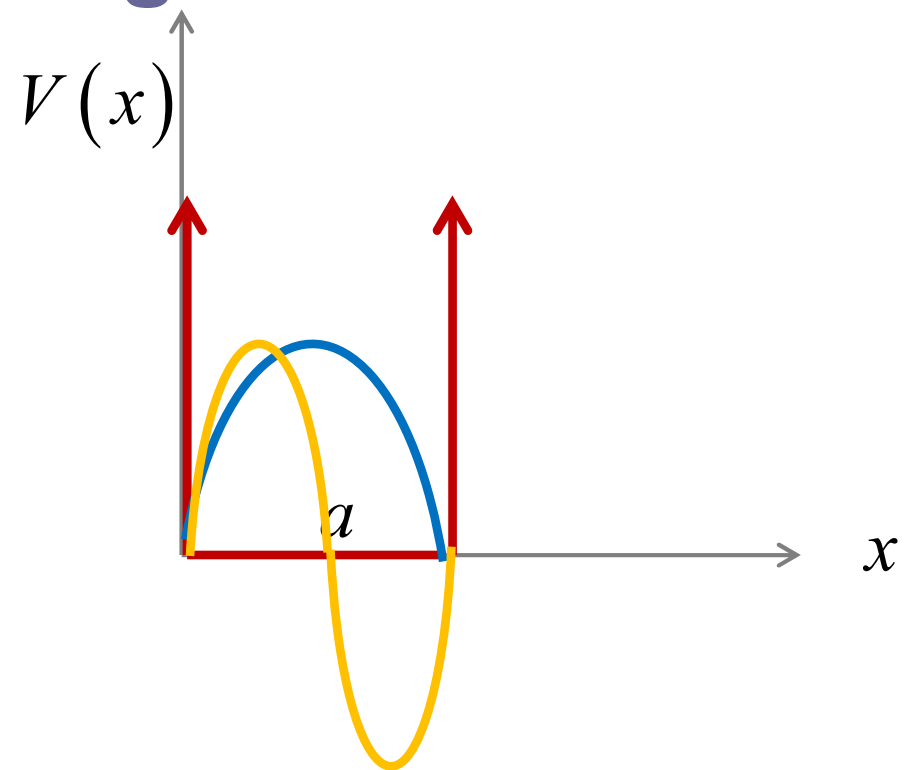
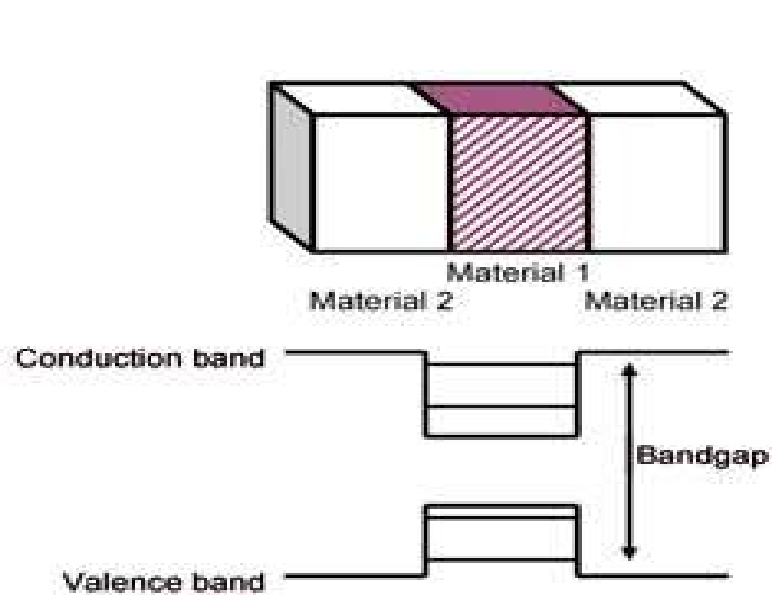
$$\psi(0) = 0 \quad \psi(a) = 0$$

$$\Rightarrow B = 0 \quad \Rightarrow ka = n\pi$$

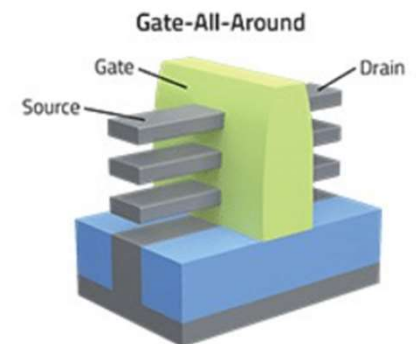
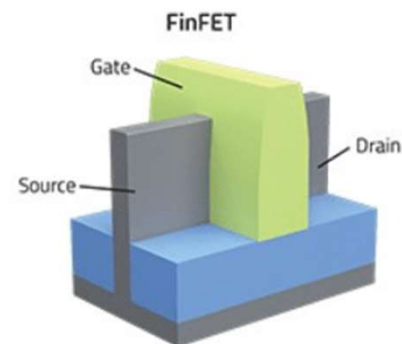
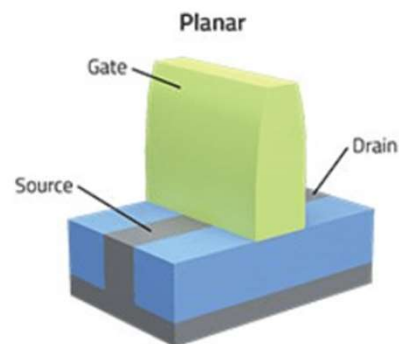
$$\psi_n = A \sin\left(\frac{n\pi x}{a}\right)$$

$$E_n = \frac{\hbar^2 n^2 \pi^2}{2ma^2}$$

# The infinite rectangular well



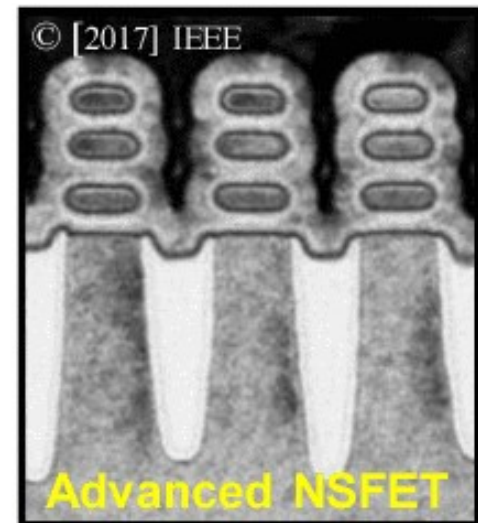
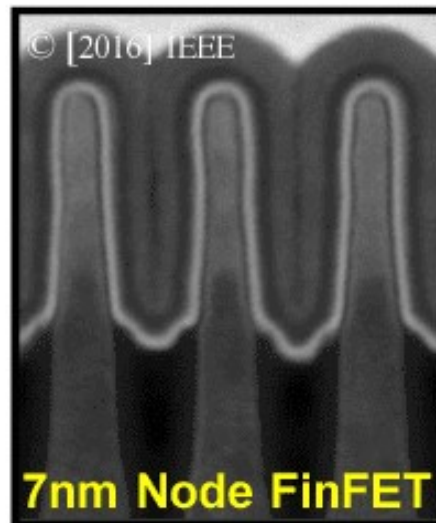
**Quantization or Quantum Confinement: MOSFET, QW Laser...**



## Question 2

Consider electrons at the bottom of the conduction band in a semiconductor fin (see figure) at room temperature. Assume  $m^* = m_0$  for this semiconductor. The fin thickness for which one needs to consider quantum confinement of electrons in the fin is less than about:

- a. 10 Å
- b. 100 Å
- c. 1000 Å
- d. Never



## Question 2

Consider electrons at the bottom of the conduction band in a semiconductor fin (see figure) at room temperature. Assume  $m^* = m_0$  for this semiconductor. The fin thickness for which one needs to consider quantum confinement of electrons in the fin is less than about:

- a. 10 Å
- b. 100 Å
- c. 1000 Å
- d. Never

$$\begin{aligned} E - E_C &= \frac{p^2}{2m^*} \sim k_B T \\ \Rightarrow p &= \sqrt{2m^* k_B T} \\ p &= \hbar k = \frac{h}{\lambda} \\ \Rightarrow \lambda &= \frac{h}{\sqrt{2m^* k_B T}} = 7.6 \text{ nm} \end{aligned}$$

# The delta-function barrier

## Delta-function

$$\delta(x) = 0 \Leftarrow x \neq 0$$

$$\int_{-\infty}^{+\infty} \delta(x) dx = 1$$

$$\Rightarrow \int_{-\infty}^{+\infty} f(x+a) \delta(x) dx = f(a)$$

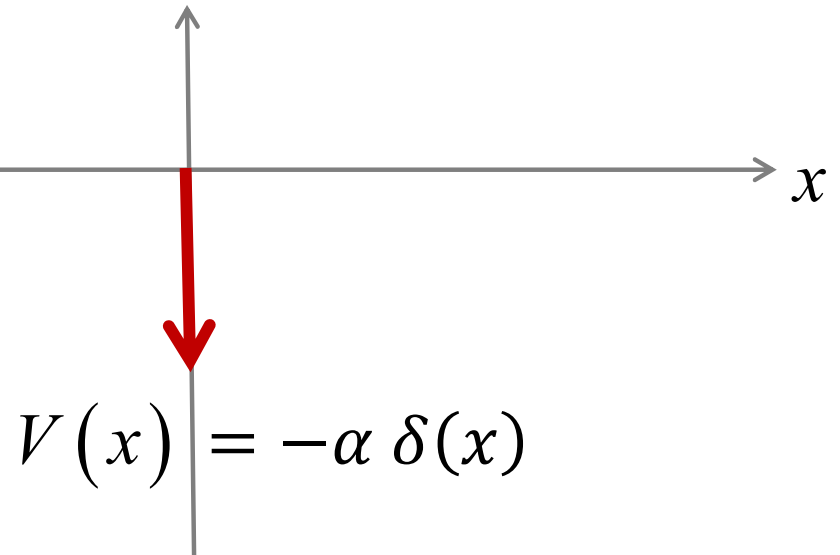
$$x < 0$$

$$\frac{-\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi \Rightarrow \frac{d^2\psi}{dx^2} = \kappa^2\psi$$

$$\kappa = \left( \frac{2m(-E)}{\hbar^2} \right)$$

$$x > 0$$

$$\psi = C_+ e^{-\kappa x}$$



**Boundary condition:**  
 **$\psi$  is continuous**

$$C_- = C_+ = \psi(0)$$

# The delta-function barrier

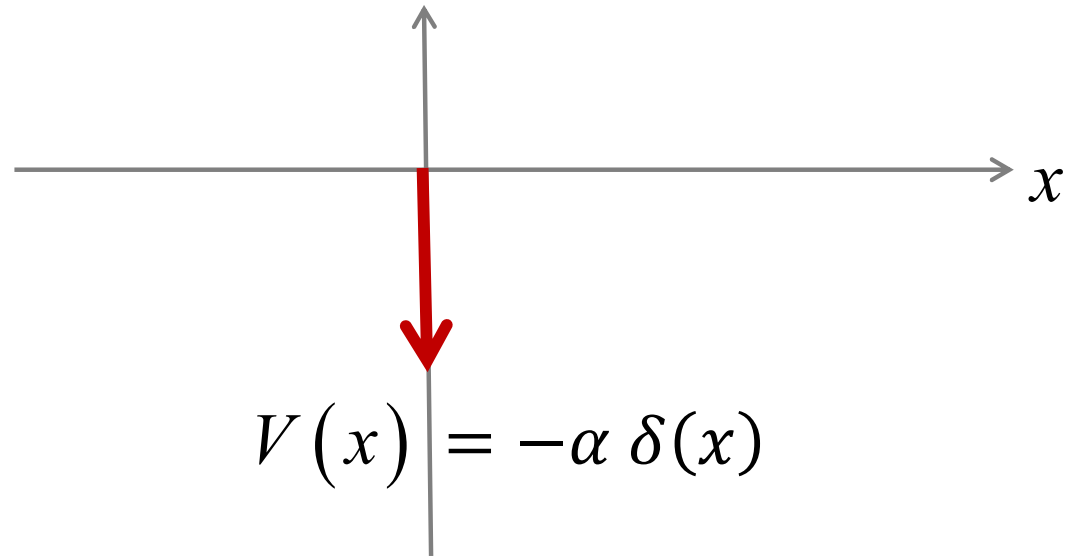
$$x < 0$$

$$\psi = \psi(0)e^{+\kappa x}$$

$$x > 0$$

$$\psi = \psi(0)e^{-\kappa x}$$

**Boundary condition:**  
 $\psi'$  is continuous  
 unless  $V$  is  $\infty$



$$\int_{0^-}^{0^+} dx \left[ \frac{-\hbar^2}{2m} \frac{d^2\psi}{dx^2} - \alpha \cdot \delta(x)\psi \right] = \int_{0^-}^{0^+} dx \cdot E\psi$$

$$\Rightarrow \frac{-\hbar^2}{2m} \left[ \frac{d\psi}{dx} \right]_{0^-}^{0^+} - \alpha \cdot \psi(0) = 0 \Rightarrow \Delta\psi'(0) = -\frac{2m\alpha}{\hbar^2} \psi(0)$$

$$\Delta\psi'(0) = -\kappa\psi(0) - \kappa\psi(0) = -2\kappa\psi(0) = -\frac{2m\alpha}{\hbar^2} \psi(0) \Rightarrow \kappa = \frac{m\alpha}{\hbar^2}$$

$$\Rightarrow -\frac{2mE}{\hbar^2} = \left( \frac{m\alpha}{\hbar^2} \right)^2 \Rightarrow E = -\frac{m\alpha^2}{2\hbar^2}$$

# Finis

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