

# E213 : Analysis of Decays of heavy vector boson $Z^0$

*A lab report written by*

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# Introduction

In this experiment, we would like to understand how data from a particle accelerator is analysed. For this, we consider the data from OPAL (Omni-Purpose Apparatus at LEP) experiment carried out at the Large electron-positron collider (LEP). In this particular report, we try to deduce different properties of the  $Z^0$  boson. For this, we first carry out event display analysis on a smaller set of data to understand how we can separate out different channels corresponding to this process. And after this, we carry out statistical analysis on the real world data to deduce important physical quantities, like the mass and decay width of  $Z^0$  and partial decay widths.

This report is divided into multiple sections. The first section discusses the requisite theoretical knowledge to understand this report. This is then followed by the section which discusses the solutions to pre-lab exercises, which are essential in understanding both the theoretical and experimental concepts. This is followed by analysis, where we explore and understand the data using various methods and finally discuss various physical quantities. This is then finally followed by conclusion, which summarises what had been done. The final parts of the report contain supplementary data and plots. All the data and source code for the analysis can be found here<sup>1</sup>.

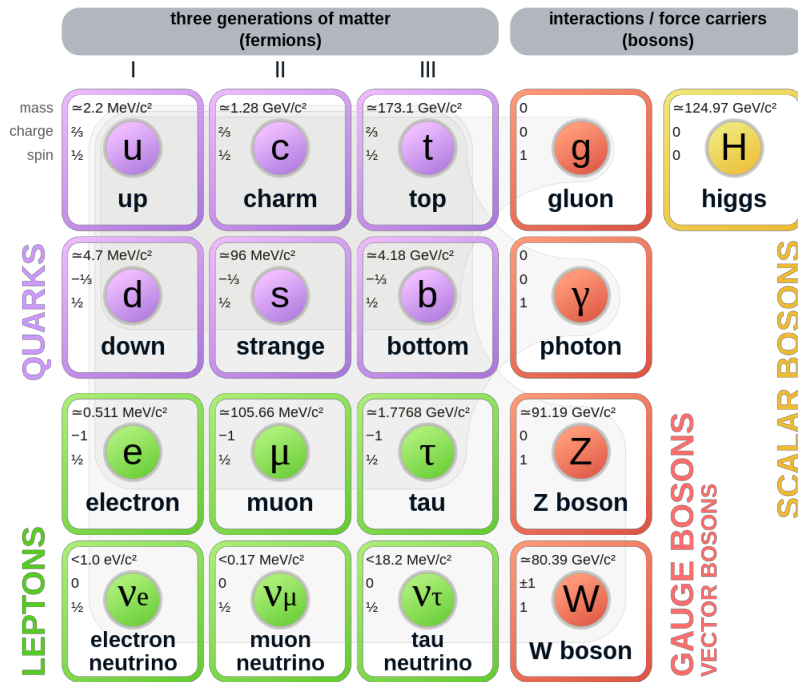
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<sup>1</sup><https://github.com/Mrunmoy-Jena/E213Lab>

# Prerequisite Knowledge

## 2.1 The Standard Model- A Brief Overview

The Standard Model of particle physics which has been one of the most successful and well-tested theories so far, provides the most fundamental description of nature by incorporating the elementary particles and their interactions. These elementary particles are categorized into two families: fermions (having half integer spins) which form matter as we know it, and bosons (with integer spins) serving as mediators of the three fundamental forces. While electromagnetic interactions are mediated by the photon ( $\gamma$ ); strong and weak interactions are mediated by gluons (g) and by  $W^\pm, Z^0$  bosons respectively. A fundamental particle that mediates gravitation has been only postulated theoretically, and is left out of the Standard Model, since the effects of gravity are too weak to play any important role in the realm of particle physics.



**Figure 2.1:** A chart showing the families of elementary particles in the Standard Model [1]

The fermions consists of three generations of quarks and leptons. The quarks have six flavours: up (u), down (d), charm (c), strange (s), top (t) and bottom (b). Similarly, the leptons consist of the electron (e), muon ( $\mu$ ) and tau ( $\tau$ ), each having its own associated charge-less and almost massless neutrino ( $\nu_e$ ,  $\nu_\mu$  and  $\nu_\tau$ ). Furthermore, each particle in the standard model has its own antiparticle. The quarks are able to form composite particles in either three quark combinations, called baryons ( $qqq/\bar{q}\bar{q}\bar{q}$ ) or a quark-antiquark pair, called a meson ( $q\bar{q}$ ). Mathematically, the elementary particles are described as elements of representations of certain symmetry groups. The gauge fields that couple to these particles (i.e. mediate the interactions) arise naturally as a consequence of invariance of their corresponding Lagrangian under

local phase transformations [2]. As such, the gauge symmetry that governs the Standard Model is given by:

$$SU(3)_{\text{Colour}} \times SU(2)_{\text{Left chiral}} \times U(1)_{\text{Y(Weak hypercharge)}}$$

## 2.2 The Unified Theory of Electromagnetic and Weak Interactions

Since the experiment deals with verifying some of the properties of the  $Z^0$  bosons, it is of interest to touch upon the theory of electroweak unification.

While electromagnetism and the theory of weak interactions were formulated separately, it was later on postulated that at higher energies ( $\sim 246$  GeV [3]), both these interactions would be unified into a single force. As such, the GSW(Glashow-Salam-Weinberg) electroweak model was developed in the 1960s to describe this unified force.

One finds that imposing the principle of local gauge invariance on the  $SU(2)_L$  symmetry group leads to the introduction of three gauge fields:  $W^{(1)}$ ,  $W^{(2)}$  and  $W^{(3)}$  (or  $W^0$  in some references) [2]. The physical  $W^+$  and  $W^-$  bosons (that mediate the weak charged current interaction) can be seen as the linear combinations:

$$W^\pm = \frac{1}{\sqrt{2}} \left( W^{(1)} \mp W^{(2)} \right) \quad (2.1)$$

However, the  $W^{(3)}$  field has no physical interpretation of its own. Therefore an additional symmetry, the  $U(1)_Y$  group is introduced. The field  $B$  (or equivalently  $Y^0$ ) arising as a consequence of this new symmetry, similarly does not have a physical meaning on its own. Rather, it was seen that linear combinations of the  $W^{(3)}$  ( $W^0$ ) and  $B$  ( $Y^0$ ) fields gives rise to the photon and the  $Z^0$  boson:

$$\begin{pmatrix} \gamma \\ Z^0 \end{pmatrix} = \begin{pmatrix} \cos \theta_W & \sin \theta_W \\ -\sin \theta_W & \cos \theta_W \end{pmatrix} \begin{pmatrix} B \\ W^{(3)} \end{pmatrix} \quad (2.2)$$

where  $\theta_W$  is the weak mixing/Weinberg angle

In addition to this, it is to be noted that the gauge fields  $W^{(1),(2),(3)}$  and  $B$  have to be massless, in order to respect gauge invariance under local  $SU(2)_L \times U(1)_Y$  gauge transformation. However, the physical gauge bosons  $W^\pm$  and  $Z^0$  are predicted to be massive, whereas the photon should remain massless. To explain this, the concept of electroweak spontaneous symmetry breaking was introduced. A massive scalar field (the Higgs field) is introduced, to which these bosons ( $W^\pm$ ,  $Z^0$ ) must couple to, in order to get their physical masses, while the photon does not interact with it [4]. The intricacies of the Higgs mechanism are not of immediate interest here, and can be understood from standard references [2, 5].

## 2.3 Physics Related to the $Z^0$ Resonance

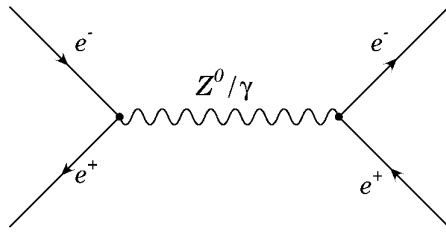
### 2.3.1 $e^+e^-$ Interactions

Before discussing the processes of interest involving  $Z^0$  production, we first list out the important ways in which  $e^+e^-$  pairs can interact [6]:

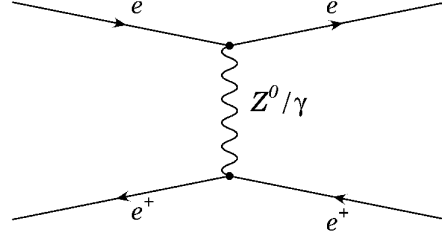
- $e^+e^- \rightarrow e^+e^-$ : Bhaba scattering (elastic scattering) of  $e^+e^-$  pairs

- $e^+e^- \rightarrow e^+e^-\gamma\gamma$ : two photon process;  $e^+e^-$  can scatter off of virtual photons, arising out of the incoming  $e^-$  and  $e^+$  themselves and these photons can then produce hadrons
- $e^+e^- \rightarrow f\bar{f}$ : Electron-positron pairs annihilate to produce a gauge boson ( $\gamma$  or  $Z^0$ ), which would in turn produce a fermion-antifermion ( $f\bar{f}$ ) pair. Here  $f\bar{f}$  are fermions other than  $e^+e^-$  since that process is already included under Bhaba scattering.
- $e^+e^- \rightarrow \gamma\gamma/\gamma\gamma\gamma$ : An electron-positron pair could also produce two or three photons

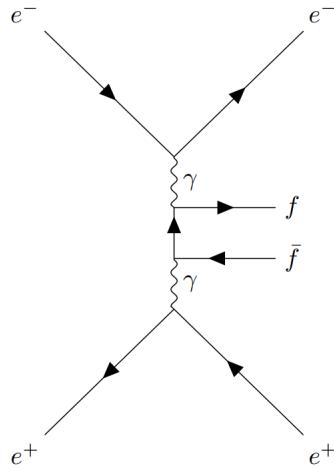
These  $e^+e^-$  interactions are displayed in Figure 2.2.



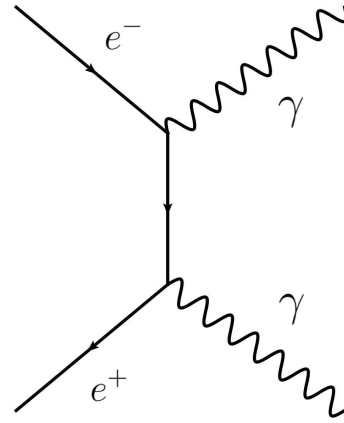
(a) s-channel Bhaba scattering



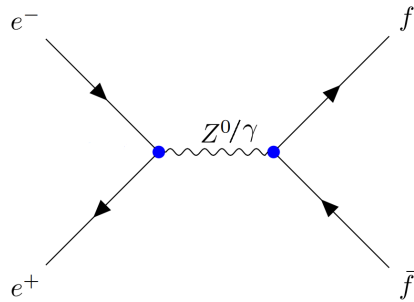
(b) t-channel Bhaba scattering



(c) Two photon process



(d) Production of two photons from  $e^+e^-$  pair



(e) Electron-positron annihilation to  $f\bar{f}$

**Figure 2.2:** Examples of  $e^+e^-$  interactions [6–8]

In case of the lowest order  $e^+e^- \rightarrow f\bar{f}$  process, in general, there are contributions to cross section from pure  $Z^0$ , pure  $\gamma$  as well as from  $\gamma - Z^0$  interference terms. However, near center of mass energies close to the  $Z^0$  resonance, the major contribution to the total cross section is from the pure  $Z^0$  term.

Additionally, it is easy to see that for the particular case of  $e^+e^- \rightarrow e^+e^-$  scattering (Bhaba scattering), there is a  $t$ -channel contribution in addition to the  $s$ -channel component. The cross sections of both these processes have different angular dependencies [6]

$$\left(\frac{d\sigma}{d\Omega}\right)_s \propto (1 + \cos^2 \theta) \quad (2.3)$$

$$\left(\frac{d\sigma}{d\Omega}\right)_t \propto (1 - \cos \theta)^{-2} \quad (2.4)$$

This means that while the  $s$  channel cross section has a major contribution at large angles (or small values of  $\cos \theta$ ), the  $t$ -channel contribution increases asymptotically at small angles (or large values of  $\cos \theta$ ). As shall be seen later on, removing the  $t$ -channel contribution is an essential step, since we are only concerned with finding the inherent forward-backward asymmetry associated with the  $s$ -channel processes.

### 2.3.2 Forward-Backward Asymmetry

When we consider the  $s$ -channel processes  $e^+e^- \rightarrow f\bar{f}$ , as discussed before, if the process is mediated purely by photons ( $\gamma$ ) then the differential cross section follows the  $(1 + \cos^2 \theta)$  dependence and it shows a symmetric dependence on the scattering angle  $\theta$ . However, for the same process mediated by  $Z^0$  bosons we find that there is an additional term contributing to the asymmetry:

$$\left(\frac{d\sigma}{d\Omega}\right)_{s(Z^0)} \propto a(1 + \cos^2 \theta) + 2b \cos \theta \quad (2.5)$$

From theory of electroweak interactions, this asymmetry (between the number of fermions produced in the forward direction,  $\theta > \pi/2$  and the backward direction,  $\theta < \pi/2$ ) can be understood to be arising from the fact that the  $Z^0$  does not couple equally to right handed and left handed fermions. This becomes more clear by looking at the form of the asymmetry term:

$$b = \left[(g_L^e)^2 - (g_R^e)^2\right] \left[\left(g_L^f\right)^2 - \left(g_R^f\right)^2\right] \quad (2.6)$$

where  $g_L^f$  and  $g_R^f$  signify the coupling of  $Z^0$  to left handed and right handed fermions respectively. It can be clearly seen that in case the these two couplings had been equal, the asymmetry term would have gone to zero [2].

The forward backward asymmetry factor is given as the ratio:

$$\mathcal{A}_{fb} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = \frac{3b}{4a} \quad (2.7)$$

where  $\sigma_F$  and  $\sigma_B$  are the cross sections in forward and backward directions respectively.

It is found that at the resonance of  $Z^0$  boson,  $\mathcal{A}_{fb}$  simplifies to [6]:

$$\mathcal{A}_{fb}^f \approx 3 \left( \frac{g_V^f}{g_A^f} \right) = 1 - 4 \sin^2 \theta_W \quad (2.8)$$

Thus from the calculation of the forward-backward asymmetry, the ratio of  $g_V^f$  (vector coupling strength between  $Z^0$  and fermions) to  $g_A^f$  (axial vector coupling strength between  $Z^0$  and fermions) can be found, which in turn gives us the Weinberg (weak mixing) angle  $\theta_W$ .

### 2.3.3 Background Processes: Radiative Corrections

In order to test out the predictions of the Standard Model at a very high level of precision, higher order background processes need to be accounted for and as such the corrections made due to these processes are called 'radiative corrections'. The following are some of these processes that need to be accounted for [9]:

- **Initial state radiation (ISR):** The radiation of photons in the initial state lead to decrease in the centre of mass energy and thus affect the  $Z^0$  resonance peak parameters. This leads to a peak height reduction by upto 25 %, shift of peak to higher energy and increase in full width at half maximum (FWHM).
- **Final state radiation (FSR):** When photons or gluons are radiated in the final state, it is found that partial widths increase by the corresponding factors:

$$\Delta_{QED} = 1 + \frac{3}{4} \frac{Q_f^2 \alpha(m_Z^2)}{\pi} ; \Delta_{QCD} \approx 1 + \frac{\alpha_s(m_Z^2)}{\pi} \quad (2.9)$$

- **QED vacuum polarization:** The production of  $e^+e^-$  pairs from vacuum makes the QED coupling constant scale dependent:

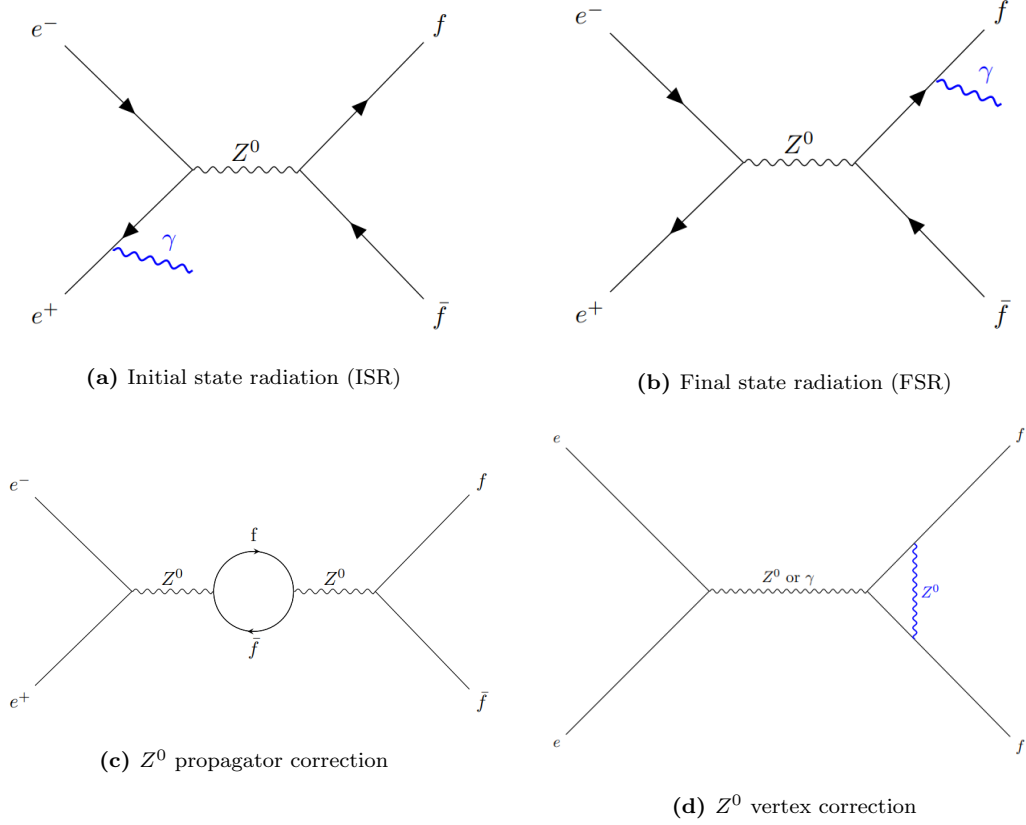
$$\alpha(0) \rightarrow \alpha(s) = \frac{\alpha_0}{1 - \Delta\alpha(s)} \quad (2.10)$$

It has been shown that the imaginary part of  $\Delta\alpha(s)$  has an influence on the  $\gamma/Z^0$  interference term.

- **Electroweak corrections:** Further contributions also arise from virtual processes such as higher order loops in  $Z^0$  propagator and vertex corrections.

Some instances of these radiative corrections are shown in Figure 2.3





**Figure 2.3:** Examples of radiative corrections [6]

### 2.3.4 Mass and Width of the $Z^0$ Boson

From the theory of electroweak interactions, it is known that the contribution the  $Z^0$  boson exchange propagator to the matrix element (in s-channel processes) can be given by:

$$\mathcal{M}_{Z^0} \propto \frac{g_{Z^0}^2}{q^2 - m_{Z^0}^2} = \frac{g_{Z^0}^2}{s - m_{Z^0}^2} \quad (2.11)$$

While one can see from this dependence that at low energies ( $\sqrt{s} \ll m_{Z^0}$ ), QED processes dominate and at high energies ( $\sqrt{s} \gg m_{Z^0}$ ), the  $Z^0$  exchange also becomes important; we see that around the  $Z^0$  resonance ( $\sqrt{s} \sim m_{Z^0}$ ), the propagator diverges. In order to correct for this, it is necessary to introduce a modification in the propagator for a decaying state. When an unstable particle has a decay rate  $\Gamma$ , its wavefunction gets modified to:

$$\psi \propto e^{-imt} \rightarrow e^{-imt} \quad (2.12)$$

which is equivalent to introducing an additional imaginary term in the mass:

$$m \rightarrow m - i\frac{\Gamma}{2} \quad (2.13)$$

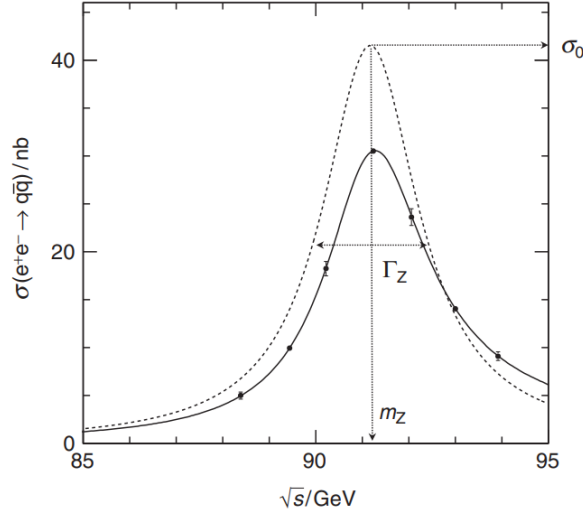
Accordingly, the  $Z^0$  propagator then changes to:

$$\frac{1}{s - m_{Z^0}^2} \rightarrow \frac{1}{s - m_{Z^0}^2} \quad (2.14)$$

On simplifying this result, taking the spin averaged matrix element and inserting the appropriate pre-factors, the complete form of the cross section in the process  $e^+e^- \rightarrow Z^0 \rightarrow f\bar{f}$  is given by:

$$\sigma_f(s) = \frac{12\pi}{M_Z^2} \frac{s\Gamma_e\Gamma_f}{(s - M_Z^2)^2 + \left(\frac{s\Gamma_Z}{M_Z}\right)^2} (\hbar^2 c^2), \quad (2.15)$$

The probability distribution that characterizes the dependence of this cross section of the centre of mass energy is called the Breit-Wigner distribution. This distribution is shown in Figure 2.4.



**Figure 2.4:** The Breit Wigner distribution of the cross section for  $e^+e^- \rightarrow q\bar{q}$  process at centre of mass energies around the  $Z^0$  resonance (the dashed and solid curves are the distribution after and before corrections for initial state radiation). [2]

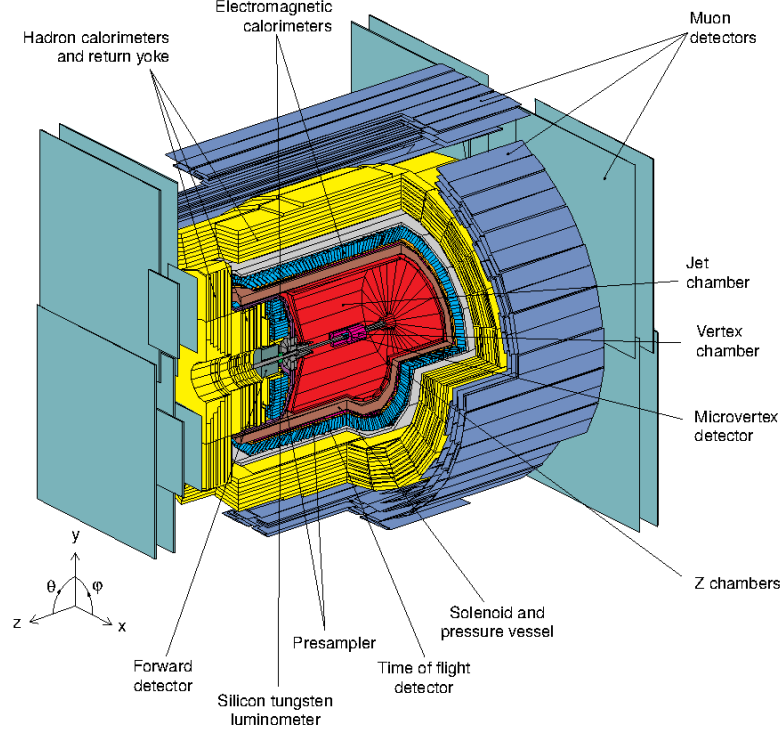
As can be seen from the figure, at the  $Z^0$  resonance, the physical parameters of the  $Z^0$  can easily be extracted. For example, the mass of the  $Z^0$  would simply be the value on the  $x$  axis where this peak occurs, its decay width can be found from the full width at half maximum (FWHM) of the curve, and reading off the  $y$  value at the resonance peak gives us the resonant cross section. This will in fact be one of the important tasks of the present experiment.

## 2.4 LEP Experiment and the OPAL Detector

The Large Electron Positron Collider (LEP) was built at CERN, with one of its major goals being making high precision measurements of the parameters of the  $Z^0$  boson, such as its mass, decay width and angular distributions of final state particles produced from  $Z^0$  decays. As such, for most of its period of operation from 1989 to 1995, it produced  $e^+e^-$  collisions at centre of mass energies very close to the  $Z^0$  resonance and about 17 million  $e^+e^- \rightarrow Z^0$  events were recorded in this duration [2]. The collisions took place at four different points in the circular collider, and hence four such detectors were used in this experiment: ALEPH (Apparatus for LEP PHysics), DELPHI (DEtector with Lepton, Photon and Hadron Identification), L3 (Third LEP experiment) and OPAL (Omni-Purpose Apparatus for LEP).

The components of the OPAL detector shall be described in further detail in this Section, since this experiment focuses on analysis of data collected by the OPAL.

### 2.4.1 The OPAL Detector



**Figure 2.5:** The OPAL Detector and its various components [10]

For analysis of events arising from this detector, a right handed coordinate system is followed. The origin is fixed at the point where  $e^+e^-$  collisions take place, the  $x$  axis points horizontally in the direction of the center of the LEP, while the  $z$  axis is defined to lie along along the beam path and the  $y$  axis is then just perpendicular to the  $x - z$  plane as usual.  $\theta$ , the polar angle is the angle between the  $y$  and  $z$  axes and the azimuthal angle  $\phi$  is measured counterclockwise between the  $x$  and  $y$  directions. The OPAL detector and this coordinate system is displayed in Figure 2.5.

Starting from the origin, as the  $e^+e^-$  collision products fly outwards, various detector components are encountered which serve to detect the type and various characteristics of the final state particles. These primary detector components are discussed here in brief [6, 11]:

- **Vertex detector:** This is a type of drift detector that surrounds the central beam pipe and plays a key role in locating vertices of short-lived decay products and also helps in improving the momentum resolution. It has dimensions of 1 m in length and measures 470 mm across (diameter), layered with 36 axial cells that measure positions in the  $r - \phi$  plane.
- **Jet chamber:** This is again a cylindrical drift chamber with a good spatial and track resolution, designed in a manner so as to efficiently record events associated with jets. Inside this chamber, there are 24 sectors, each of which contains 159 sensing wires, arranged parallel to the beam axis.

Beside the anode wires, there are alternating cathode wires positioned in between the potential and anode wires. The coordinates of the decay particles are obtained from parameters such as the positions of the wires and the drift time. Moreover, the signals from these wires determine the energy loss  $dE/dx$  of charged particles in this region.

- **z chambers:** As the name suggests, the  $z$  chambers serve the purpose of locating the  $z$  coordinates of decay particles when they leave the jet chamber. As such they are helpful in improving the resolutions of the polar angle and invariant mass distributions of charged particles. Consisting of 24 drift chambers, subdivided into 8 cells each, a high  $z$  resolution of 100-200  $\mu\text{m}$  is achieved by these chambers.
- **Solenoid:** The solenoid surrounds the detectors discussed so far (categorized under central detector) and generates a uniform magnetic field of 0.435 T along the beam direction. The magnetic field makes the charged particles follow a helical path in the central detectors and therefore is essential to measure their momenta.
- **Time of flight (TOF) system:** The role of the TOF system is to generate signals (light from a total of 160 scintillation chambers) that trigger the detector and to identify charged particles in the range of 0.6 to 2.5 GeV. Further it is useful in rejecting cosmic rays.
- **Electromagnetic calorimeter (ECAL):** This detector is made up of 9440 blocks of lead glass. The geometry of this detector is in the form of a barrel around the jet chamber with two endcaps, covering a total of 98 % of the solid angle. Particles such as photons, electrons and positrons deposit all of their energy in the ECAL either through Bremsstrahlung ( $e^+$  or  $e^-$  emitting a photon) or by pair production (a photon giving rise to an  $e^+e^-$  pair). Hadrons on the other hand, do deposit some of their energy here (in the form of electromagnetic showers), but mostly don't come to a stop in the ECAL. The reason for choosing lead glass for the ECAL is that the material provides an excellent resolution of:

$$\left(\frac{\sigma_E}{E}\right)_{\text{ECAL}} \approx \frac{5\%}{\sqrt{E}} \quad (2.16)$$

- **Hadron calorimeter (HCAL):** This detector has a similar geometry to that of that of ECAL and serves to detect mesons and baryons. In addition to ionisation energy loss, hadrons also undergo strong interaction with the nuclei that they encounter. As such there can be numerous decay products in the associated hadronic showers. Since the decay products travel significant distances between any two nuclear interactions, the HCAL occupies much more volume of the detector than ECAL and is made up of absorber material with a high density (iron slabs in this case). Because of the uncertainties in determining the amount of energy lost electromagnetically and in the energy loss in nuclear interactions, the energy resolution for the HCAL is poor compared to that of the ECAL:

$$\left(\frac{\sigma_E}{E}\right)_{\text{HCAL}} \approx \frac{120\%}{\sqrt{E}} \quad (2.17)$$

- **Muon detector:** Outside the HCAL, lie four layers of muon detectors. The muon detector provides a coverage of almost the entire solid angle of  $4\pi$  and has a barrel region and two endcap regions, similar to the ECAL and HCAL. The detection probability of muons in this detector is very high ( $\sim 100\%$ ) because almost all of the other decay products have either been absorbed in the ECAL or in the HCAL.

# Pre-Lab Exercises

In this chapter, we have provided our solutions to some theoretical questions that were needed to be solved before conducting the experiment. All the equations and standard values of parameters used to calculate the numerical results for these exercises are taken from [6], unless mentioned otherwise.

## 3.1 Calculation of partial decay widths for $Z^0 \rightarrow f\bar{f}$

In this exercise, we are required to calculate the partial decay widths of  $Z^0 \rightarrow f\bar{f}$ ; where  $f\bar{f}$  represent the following fermion-antifermion pairs : (i)  $e^+e^-$  (ii)  $\mu^+\mu^-$  (iii)  $\tau^+\tau^-$  (iv)  $q\bar{q}$ , where  $q$  represents all the flavours of quarks (except for t quark, because it is too heavy ( $M_t \approx 172.76$  GeV [12]) to be produced from  $Z^0$  decays). The partial decay widths have been calculated with the following formula:

$$\Gamma_f = \frac{N_c^f \sqrt{2}}{12\pi} G_F M_Z^3 \left( \left( g_V^f \right)^2 + \left( g_A^f \right)^2 \right) \quad (3.1)$$

where:

$N_c^f$ : colour factor, (1 for leptons, 3 for quarks)

$G_F = 1.16637 \times 10^{-5} \text{GeV}^{-2}$ , Fermi's constant

$M_Z = 91.182$  GeV, mass of  $Z^0$  boson

$g_V^f = I_3^f - 2Q_f \sin^2 \theta_W$ , vector coupling strength of  $Z^0$  to fermions

$g_A^f = I_3^f$ , axial-vector coupling strength of  $Z^0$  to fermions

$Q_f$ : electric charge of fermion  $f$

$I_3$ : third component of weak isospin

$\sin^2 \theta_W = 0.2312$ ,  $\theta_W$  is the Weinberg (weak-mixing) angle

Fermion	$Q_f$	$I_3^f$	$g_V^f$	$g_A^f$	$N_c^f$	$\Gamma_f^{(\text{calc})} / \text{MeV}$	$\Gamma_f^{(\text{ref})} / \text{MeV}$
$e^-, \mu^-, \tau^-$	-1	-0.5	-0.0376	-0.5	1	83.89	83.8
$u, c$	2/3	0.5	0.1917	0.5	3	285.34	299
$d, b, s$	-1/3	-0.5	-0.3459	-0.5	3	367.84	378
$\nu_e, \nu_\mu, \nu_\tau$	0	0.5	0.5	0.5	1	165.85	167.6

**Table 3.1:** Parameters  $Q_f, I_3^f, g_V^f, g_A^f, N_c^f$  for various fermion pairs and their partial decay widths

The calculated partial decay widths for the required fermion pairs have been listed under  $\Gamma_f^{(\text{calc})}$  in Table 3.1. Further, the reference [6] values of partial decay widths for the same fermion pairs are listed under  $\Gamma_f^{(\text{ref})}$  for comparison. We have also included the partial decay widths of the three neutrinos since they would be used in the solution to the next exercise.

One finds that calculated values of partial decay width for the lepton pairs are in close agreement with the literature value, deviating by about 0.1% to 1%. The slight deviation could be caused because the

$\gamma \rightarrow f\bar{f}$  term and interference terms have been neglected. In case of the quarks, the deviations from the reference values are higher ( $\sim 2.7\%$  to  $4.6\%$ ). This may be due to the fact that additionally, the effect of strong interactions have not been accounted for in our calculations.

## 3.2 Calculation of hadronic, leptonic and total decay widths and cross section

**Hadronic decay width:** The decay widths for the hadronic mode is given by the sum of the partial decay widths of the u,d,c,s and b quarks:

$$\Gamma_{had} = \Gamma_u + \Gamma_c + \Gamma_d + \Gamma_s + \Gamma_b = 2 \cdot \Gamma_{u,c} + 3 \cdot \Gamma_{d,s,b} = 1674.20 \text{ MeV}$$

**‘Charged’ decay width:** The charged leptons, e,  $\mu$ ,  $\tau$  will contribute to this decay width:

$$\Gamma_{charged \text{ leptonic}} = \Gamma_e + \Gamma_\mu + \Gamma_\tau = 3 \cdot \Gamma_{e,\mu,\tau} = 250.17 \text{ MeV}$$

**‘Neutral’ (invisible) decay width:** The uncharged leptons ( $\nu_e$ ,  $\nu_\mu$ ,  $\nu_\tau$ ) will contribute to this:

$$\Gamma_{neutral \text{ leptonic}} = \Gamma_{\nu_e} + \Gamma_{\nu_\mu} + \Gamma_{\nu_\tau} = 3 \cdot \Gamma_{\nu_e} = 497.55 \text{ MeV}$$

**Total  $Z^0$  decay width:** The total decay width for  $Z^0$  will just be the sum of the hadronic, charged leptonic and neutral leptonic decay widths:

$$\Gamma_{total} = \Gamma_{hadronic} + \Gamma_{charged \text{ leptonic}} + \Gamma_{neutral \text{ leptonic}} = 2421.92 \text{ MeV}$$

**Partial cross sections at maximum of resonance:** At resonance, the formula for calculating partial cross section for  $Z^0 \rightarrow f\bar{f}$  becomes:

$$\sigma_f^{peak} = \frac{12\pi\Gamma_e\Gamma_f}{M_Z^2\Gamma_Z^2} \quad (3.2)$$

The calculated partial cross sections for the different decay channels along with the respective decay widths are tabulated in Table 3.2.

Decay channel	Decay width / MeV	Partial cross section / $10^{-11} \text{ MeV}^{-2}$
Hadronic (u,d,c,s,b)	1674.20	10.79
Charged leptonic (e, $\mu$ , $\tau$ )	250.17	1.61
Neutral leptonic ( $\nu_e$ , $\nu_\mu$ , $\nu_\tau$ )	497.55	3.21
Total	2421.92	15.61

**Table 3.2:** Calculated decay widths and partial cross sections for different  $Z^0$  decay channels

### 3.3 Effect of additional generation on width of $Z^0$ resonance curve

In case it is possible for the  $Z^0$  to decay into an extra generation of light fermions (u,d,e, $\nu$ ), the total decay width will increase, and the new total decay width will be:

$$\Gamma_{total}^{(new)} = \Gamma_{total} + \Gamma_e + \Gamma_\nu + \Gamma_u + \Gamma_d = 3324.34 \text{ MeV}$$

The percentage increase in the width of the  $Z^0$  resonance curve will be:

$$\frac{\Gamma_{total}^{(new)} - \Gamma_{total}}{\Gamma_{total}} = \frac{902.42}{2421.92} \times 100\% = 37.26\%$$

### 3.4 Angular distributions of differential cross sections

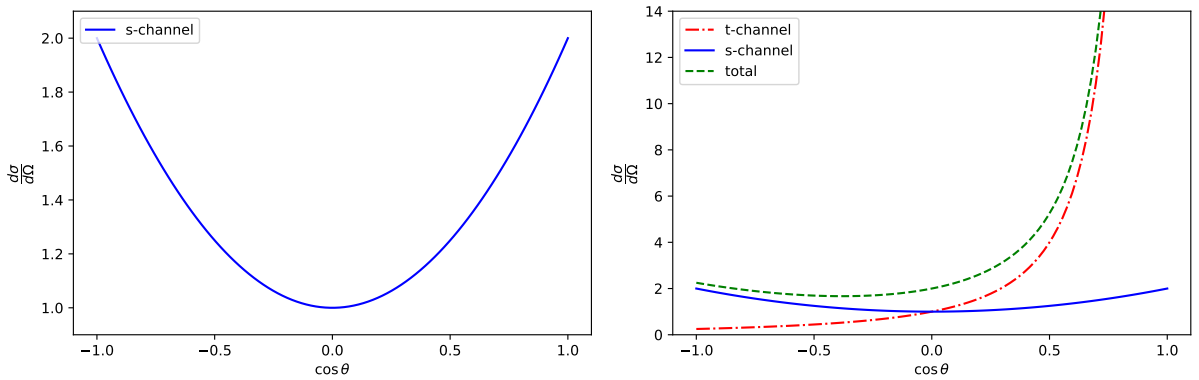
On studying the dependence of differential cross sections of  $e^+e^-$  processes on the azimuthal angle  $\theta$ , it is found that the s-channel has a symmetrical dependence on  $\theta$  (with an additional small asymmetric term, in the case of  $Z^0$  mediated process; which will be calculated later):

$$\left(\frac{d\sigma}{d\Omega}\right)_s \propto (1 + \cos^2 \theta) \quad (3.3)$$

For the case of a t-channel process, it is found that the differential cross section diverges quickly at small  $\theta$  values [6]:

$$\left(\frac{d\sigma}{d\Omega}\right)_t \propto (1 - \cos \theta)^{-2} \quad (3.4)$$

While the process  $e^+e^- \rightarrow \mu^-\mu^+$  can only take place through the s-channel; the  $e^+e^- \rightarrow e^+e^-$  process can occur through both s and t channels. The angular distributions of these two processes are shown in Figure 3.1. For the  $e^+e^- \rightarrow e^+e^-$  process, the contributions from the s and t channels are shown separately as well as their combined contribution to the differential cross section.



**Figure 3.1:** Angular distributions of  $e^+e^- \rightarrow \mu^+\mu^-$  (left) and  $e^+e^- \rightarrow e^+e^-$  (right)

### 3.5 Calculation of forward-backward asymmetry

We are required to calculate the forward-backward asymmetry factor,  $A_{FB}$  in the process  $e^+e^- \rightarrow \mu^+\mu^-$ . In order to do this, the following formula is used:

$$A_{FB} = \frac{3F_2}{4F_1} \quad (3.5)$$

where the parameters are given by:

$$F_1(s) = Q_f^2 - 2v_e v_f Q_f \Re(\chi) + (v_e^2 + a_e^2)(v_f^2 + a_f^2)|\chi|^2$$

$$F_1(s) = -2a_e a_f Q_f \Re(\chi) + 4v_e a_e v_f a_f |\chi|^2$$

$$v_f = \frac{g_V^f}{2 \sin \theta_W \cos \theta_W}$$

$$a_f = \frac{g_A^f}{2 \sin \theta_W \cos \theta_W}$$

$$\chi(s) = \frac{s}{\left( (s - M_Z^2) + i s \frac{\Gamma_Z}{M_Z} \right)}, \text{ Propagator}$$

$s$ : Square of centre-of-mass energy

The calculated asymmetry factors for various combinations of  $\sin^2 \theta_W$  and centre of mass energies,  $\sqrt{s}$  are given in Table 3.3.

$\sin^2 \theta_W \backslash \sqrt{s} / \text{GeV}$	89.225	91.225	93.225
0.21	-0.0937	0.0762	0.2317
0.23	-0.1639	0.0228	0.1965
0.25	-0.1948	0.0042	0.1906

**Table 3.3:** Forward backward asymmetry factors for different Weinberg angles ( $\theta_W$ ) and centre of mass energies ( $\sqrt{s}$ )



# Analysis

## 4.1 Part I: Analysis of Event Displays

The first part of the analysis of data consists of looking at event displays, which are generated using *GRope*. The main objective of this part is to understand how different channels affect different observed quantities and how we can come up with a filter to the quantities, so as to separate out different channels in a real world data. The filters are known as the cut criteria for this process. Every event display consists of an image and a list of quantities,

- Ctrk(N): Number of charged tracks
- Ctrk(SumP): Momentum of all charged tracks
- Ecal(SumE): Total energy in the electromagnetic calorimeter
- Hcal(sumE): Total energy in the hadronic calorimeter

We have eight different samples for this part - four *pure* samples and four *mixed* samples. Each of the four *pure* samples consists of only the events from a single channel, whereas the four *mixed* samples consists of events from all the channels. We would analyse *pure* samples first and come up with a cut criteria and test it on the *mixed* samples and see how it fares.

### 4.1.1 Analysis of $e^-e^+$ -channel Samples

The decay of  $Z^0$  boson into an electron-positron pair is one of the channels of this process. For this case, we expect to detect two charged tracks. Since the pair has negligible mass compared to  $Z^0$  boson, they carry away the center of mass energy, which implies  $SumP$  will have this value. By construction of the detectors, we expect all the energy to be deposited in the electromagnetic calorimeter, while the signals should be close to zero in other components. The event display of a typical event can be found in B.1, which shows two clear charged tracks and also shows that they deposit their energy in the electromagnetic calorimeter. By conservation of four-momenta, we should also expect to see the tracks exactly opposite to each other, which can also be seen here.

### 4.1.2 Analysis of $\mu^-\mu^+$ -channel Samples

The decay of  $Z^0$  boson into an muon-antimuon pair is another channel of this process. Just like the previous case, we expect to see two charged tracks and  $SumP$  value being close to the center of mass energy value. By construction of the detectors, we expect close to zero energy deposit in the calorimeters. But instead, we expect to see a signal in the muon chambers. The event display of a typical event can be found in B.2, which shows two clear charged tracks and signals in the muon chamber. We also see that conservation of four-momenta holds, since the tracks are exactly opposite to one another.

### 4.1.3 Analysis of $\tau^-\tau^+$ -channel Samples

The decay of  $Z^0$  boson into an tau-antitau pair is another channel of this process. Unlike its other leptonic cousins,  $\tau$  has a much shorter life, which results in interesting outcomes here. We expect to see a large number of charged tracks due to the decay, clustered exactly opposite to each other due to conservation of four-momenta. The dominant decay products are  $\mu$  and  $\pi$ , which can then be detected in the other components of the detectors. This channel can be classified by the number of charge tracks, called *prongs*. The event display of a *3-prong* event can be found in B.3. We can see the signals in muon chamber on one side and we can also see the deposit in electromagnetic calorimeter and hadronic calorimeter on the other side, which corresponds to the pions. We should also expect the *SumP* value to be less than the center of mass energy, since neutrinos are produced to conserve lepton number, which carry away some of the energy. Unlike the previous two cases, this channel doesn't have a very distinct filter property at the first sight.

### 4.1.4 Analysis of $q\bar{q}$ -channel Samples

The decay of  $Z^0$  boson into an quark-antiquark pair is another channel of this process. Since quarks cannot exist as free particles, they produce hadrons, which then decay into further particles that can be then detected at various components. This results in the so-called *jets*, which are much bigger clusters of charged tracks. Depending on the final products, we expect to notice signals in all of the components of the detectors. But one distinct property of this channel is the large number of charged tracks produced. The event display of a typical event can be found in B.4, which clearly shows the jets. Again, we expect the *SumP* value to be lesser than center of mass energy, since depending on the decay products, various neutrinos are produced to conserve lepton number, which carry away some of the energy.

### 4.1.5 Developing the cuts

With this, we now first extract the information of how different quantities behave for the different *pure* samples. The raw data can be found in the appendix A.1, A.2, A.3 and A.4. We plot histograms of all these quantities for different channels and identify *cuts* on the quantities to separate out the channels. The histograms are given in the fig 4.1. The bin sizes are chosen so that we are able to see the best possible separation of the different channels.

We see that  $Ctrk(N)$  clearly lets us distinguish  $q\bar{q}$  channel events from the other channels. Therefore, we require the value of  $Ctrk(N)$  to be greater than or equal to 8 to separate the hadronic channel and we require the value of  $Ctrk(N)$  to be lesser than 6 to separate all the leptonic channels. Note that this cut alone completely separates the hadronic channel events from the leptonic channel events. We see from the histogram for  $Ctrk(Sump)$  that we can separate muon channel and tau channel from each other by requiring  $Ctrk(Sump)$  greater than or equal to 75 for muon channel and lesser than or equal to 75 for tau channel. We do not impose cuts on any other channel for this quantity. Note that this separates only the muon and tau channel from each other, but these two events may still contain electron and hadronic channel events. This cut, when used with all the other cuts, ensures that we completely separate the channels from one another. We see from the histogram for  $Ecal(SumE)$  that it is possible to completely distinguish electron and muon events. This was expected from what we discussed previously. We require the value of  $Ecal(SumE)$  greater than or equal to 80 for electron channel and lesser than or equal to 10 for muon channel. We also impose a minor cut for tau channel -  $Ecal(SumE)$  less than or equal to 60 and for hadronic channel -  $Ecal(SumE)$  lies between 36 and 79. We see from the histogram



**Figure 4.1:** Histograms of  $Ctrk(N)$ ,  $Ctrk(Sump)$ ,  $Ecal(SumE)$  and  $Hcal(SumE)$  for the four channels, generated from *pure* samples.

for  $Hcal(SumE)$  that it is possible to completely separate electron channel from the other channel by requiring  $Hcal(SumE)$  to be less than 1. This quantity does not give us any other cut to separate out other channels. The final cuts are given in the table 4.1.

We now test the determined cuts on *mixed* samples. The raw data and the channel determined using

Channel	Ctrk(N)	Ctrk(Sump)	Ecal(SumE)	Hcal(SumE)
$e^-e^+$	(0,6)		$\geq 80$	$< 1$
$\mu^-\mu^+$	$< 6$	$\geq 75$	$\leq 10$	
$\tau^-\tau^+$	$< 6$	$\leq 75$	$\leq 60$	
$q\bar{q}$	$\geq 8$		[36,79]	

**Table 4.1:** Cuts determined from event display analysis

the cuts can be found in A.5. The event display was then checked visually if it matched the channel as determined by the cuts. We found that the all the events were determined accurately with our cuts. But we are not sure how perfect our cuts will be, when tested against real world data. For this purpose, we would like to refine our cuts before we carry out a statistical analysis. This is done in the first part of the next section.

## 4.2 Part II: Statistical Analysis of $Z^0$ Decays

In the previous subsection, we carried out an analysis based on event displays. This was possible because we did not have a lot of events to work with. Obviously, such an analysis would prove to be futile if we try to carry it out on a large set of data. Analysis of large sets of data would be the primary discussion of this subsection. We use the software, *ROOT* to analyse the data on a statistical basis. *ROOT* works with *.root* files, which contain all the information in a tree-like structure, called *ntuple*. The contents of

the *ntuple* are,

- RUN: Run number
- EVENT: Event number
- NCHARGED: Number of charged tracks
- PCHARGED: Total scalar sum of track momenta
- E\_ECAL: Total energy in electromagnetic calorimeter
- E\_HCAL: Total energy in hadronic calorimeter
- E\_LEP: LEP beam energy ( $= \sqrt{s}/2$ )
- COS\_THRU:  $\cos(\text{polar angle})$  between beam axis and thrust axis
- COS\_THET:  $\cos(\text{polar angle})$  between incoming positron and outgoing positive particle

We also have two categories of data, each containing a lot of *ntuples*,

- Monte Carlo (MC): These correspond to the “*pure*” events from the previous subsection. These are simulated - detector response to a calculated outgoing momentum four-vector for a specific process.
- Data: These correspond to the “*mixed*” events from the previous subsection. These are real life data recorded with the OPAL detector at specific energies around the  $Z^0$  resonance maximum.

*ROOT* can be used to set cuts on the above variables in a data and get histograms of different variables. This allows for a statistical analysis of the data. All the above information can be found in [6].

#### 4.2.1 Refining the cuts

As mentioned previously, we have two categories of data. There are four Monte Carlo (MC) data, one for each of the decay channels. There are six real world data and we will be using *data6.root* for our analysis in the following parts. Analysis of event displays already gave us an idea of what cuts to use, to extract the different channels. Our first step here will be to test our cuts on the MC data and check how they fare. And the second step will be to refine the cuts a little so that we are better able to extract different channels.

When it comes to  $e^-e^+$  final state decay channel, we would like to exclude t-channel events. This is because t-channel is possible only in the mode and for the sake of consistency, we would like to limit ourselves to only s-channel. From theory[6], we know that the t-channel dominates at large  $\cos\theta$ . By introducing a new cut to exclude events with large  $\cos\theta$ , we eliminate a most of the t-channel events. But this also means that we are eliminating some of the s-channel events. To account for this, we multiply the observed events after applying the modified cuts with correction factor. This correction factor is given by,

$$\delta = \frac{\int_{-1}^1 (1+x^2)dx}{\int_{-0.9}^{0.5} (1+x^2)dx} \approx 1.5829. \quad (4.1)$$

This factor is arrived at from theory, which gives the behaviour of s-channel as proportional to  $(1+\cos^2\theta)$  and the integral limits are correspond to the  $\cos\theta$  values which we exclude.

When it comes to  $\mu^-\mu^+$  final state decay channel, we observed a lot of events with *PCHARGED* equal

to exactly 0. These events are not physical. Hence we apply a cut to eliminate such events. In the cases of  $e^-e^+$  and  $\mu^-\mu^+$ , we also exclude  $\cos\theta$  values very close to 1 and  $-1$ , as the detector resolution is far from perfect close to the beam axis. To summarise, we have the following additions -  $\cos\theta \in [-0.9, 0.5]$ , to remove t-channel events in the case of  $e^-e^+$ ,  $PCHARGED > 0$ , to exclude unphysical events in the case of  $e^-e^+$  and  $\cos\theta \in [-0.9, 0.9]$ , to eliminate low resolution events in the cases of  $e^-e^+$  and  $\mu^-\mu^+$ . Interestingly, the channels  $\tau^-\tau^+$  and  $q\bar{q}$  did not require any additional global cuts. But, we did add the cut  $\cos\theta \in [-1, 1]$  to remove any unphysical events. Analysis done with *ROOT* on the  $e^-e^+$  MC can be found in the appendix, B.5, B.6, B.7, B.8 and B.9. Analysis on the other MC files and ultimately the data file are done similarly.

After testing our cuts, we arrive at the conclusion that we don't need any further refinements other than the additions mentioned above. With these additions, we go about analysing *data6.root*. The raw data of observed events using our cuts is given in the table 4.2 below. Note that the correction factor has not been applied to the  $e^-e^+$  event numbers but it is used in our latter calculations. The event numbers are exactly what we get after applying all the cuts appropriately.

MC Sample	Number of observed events				
	$e^-e^+$ cuts	$\mu^-\mu^+$ cuts	$\tau^-\tau^+$ cuts	$q\bar{q}$ cuts	Total (incl. global cuts, if any)
$e^-e^+$	18835	0	378	0	56720
$\mu^-\mu^+$	0	76209	8599	0	89887
$\tau^-\tau^+$	26	35	71131	135	79214
$q\bar{q}$	0	0	173	92164	98563

**Table 4.2:** Number of events with different cuts applied to each of the MC data.

#### 4.2.2 Efficiency Matrix

The efficiency matrix is a measure of how efficient the cuts are at extracting the different decay channels. If we consider the actual event numbers of the different channels as a  $4 \times 1$  matrix, the efficiency matrix will be a  $4 \times 4$  matrix and ideally, it should be a unit matrix. In this case, the event numbers observed matches the actual event numbers. But this is not the case usually. We should aim to achieve an efficiency matrix with diagonal elements as close to 1 as possible and the off-diagonal elements as close to 0 as possible. The efficiency matrix elements are given by[6],

$$\epsilon_{ij} = \frac{N_j^{i,cut}}{N_j^{j,all}}. \quad (4.2)$$

For example,  $\epsilon_{12}$  corresponds to  $e^-e^+$  cuts applied to  $\mu^-\mu^+$  events divided by the total observed  $\mu^-\mu^+$  events. Note that we construct the  $4 \times 1$  matrix in the following order -  $e^-e^+$  events,  $\mu^-\mu^+$  events,  $\tau^-\tau^+$  events and  $q\bar{q}$  events. With this efficiency matrix, one could extract the actual event numbers as follows,

$$N_{obs} = \epsilon N_{actual} \implies N_{actual} = \epsilon^{-1} N_{obs}. \quad (4.3)$$

This gives us an efficiency matrix,

$$\epsilon = \begin{pmatrix} 5.26 \times 10^{-1} & 0 & 3.28 \times 10^{-4} & 0 \\ 0 & 8.48 \times 10^{-1} & 4.42 \times 10^{-4} & 0 \\ 6.66 \times 10^{-3} & 9.57 \times 10^{-2} & 8.98 \times 10^{-1} & 1.76 \times 10^{-3} \\ 0 & 0 & 1.70 \times 10^{-3} & 9.35 \times 10^{-1} \end{pmatrix}. \quad (4.4)$$

Given a cut, whether or not a particular event passes it can be modelled with a binomial distribution, just like modelling a coin toss. In the limit when such an “experiment” is conducted on large sample size, the probability mass function of the binomial distribution can be approximated by a normal distribution. In which case, the standard deviation is given by,

$$\Delta\epsilon_{ij} = \sqrt{\frac{\epsilon_{ij}(1 - \epsilon_{ij})}{N}}, \quad (4.5)$$

where  $\epsilon_{ij}$  is a particular element of the efficiency matrix and  $N$  is the total events corresponding to that matrix element. The above treatment can be found in any introductory book on probability theory, for example Feller[13]. This gives us the standard deviation in the efficiency matrix elements,

$$\Delta\epsilon = \begin{pmatrix} 2.10 \times 10^{-3} & 0 & 6.44 \times 10^{-5} & 0 \\ 0 & 1.12 \times 10^{-3} & 7.47 \times 10^{-5} & 0 \\ 3.41 \times 10^{-4} & 9.81 \times 10^{-4} & 1.08 \times 10^{-3} & 1.33 \times 10^{-4} \\ 0 & 0 & 1.46 \times 10^{-4} & 7.85 \times 10^{-5} \end{pmatrix}. \quad (4.6)$$

For our calculations in the following sections, we require  $N_{actual}$ . Therefore, we invert the matrix. The calculations for error in the inverse matrix elements is not straightforward. We refer to [14], which gives the error in the inverse matrix element to be,

$$(\Delta\epsilon^{-1})_{ij}^2 = \sum_{\alpha=1}^4 \sum_{\beta=1}^4 (\epsilon^{-1})_{i\alpha}^2 (\Delta\epsilon)_{\alpha\beta}^2 (\epsilon^{-1})_{\beta j}^2. \quad (4.7)$$

With this we get the following inverse efficiency matrix and the error in the inverse efficiency matrix elements,

$$\epsilon^{-1} = \begin{pmatrix} 1.90 & 7.85 \times 10^{-5} & -6.95 \times 10^{-4} & 1.30 \times 10^{-6} \\ 7.36 \times 10^{-6} & 1.18 & -5.80 \times 10^{-4} & 1.09 \times 10^{-6} \\ -1.41 \times 10^{-2} & -1.26 \times 10^{-1} & 1.11 & -2.09 \times 10^{-3} \\ 2.57 \times 10^{-5} & 2.29 \times 10^{-4} & -2.03 \times 10^{-3} & 1.07 \end{pmatrix} \pm \begin{pmatrix} 7.59 \times 10^{-3} & 1.54 \times 10^{-5} & 1.36 \times 10^{-4} & 2.75 \times 10^{-7} \\ 1.30 \times 10^{-6} & 1.67 \times 10^{-3} & 9.81 \times 10^{-5} & 2.02 \times 10^{-7} \\ 7.26 \times 10^{-4} & 1.31 \times 10^{-3} & 1.33 \times 10^{-3} & 1.59 \times 10^{-4} \\ 2.58 \times 10^{-6} & 1.98 \times 10^{-5} & 1.75 \times 10^{-4} & 8.98 \times 10^{-4} \end{pmatrix} \quad (4.8)$$

### 4.2.3 Cross Sections

With the above inverse efficiency matrix, we can now calculate the actual event numbers and hence, the cross sections corresponding to different modes. The actual event numbers,  $N_{actual}$ , will simply be

$\epsilon^{-1}N_{obs}$ , as explained earlier. The error in actual event numbers is then,

$$\Delta N_{obs,i} = \sqrt{\sum_{j=1}^4 N_{obs,j}^2 (\Delta \epsilon_{ij}^{-1})^2 + \sum_{j=1}^4 (\epsilon_{ij}^{-1})^2 N_{obs,j}}. \quad (4.9)$$

The events are assumed to obey Poisson statistics and hence,  $\Delta N_{obs}$  is taken as  $\sqrt{N_{obs}}$ . With the actual event numbers, we can calculate the cross section as,

$$\sigma = \frac{N_{actual}}{\int \mathcal{L} dt} + \text{cf}(\text{radiation}), \quad (4.10)$$

with the errors,

$$\Delta \sigma = \sqrt{\frac{(\Delta N_{actual})^2}{(\int \mathcal{L} dt)^2} + \frac{N_{actual}^2 (\Delta \int \mathcal{L} dt)^2}{(\int \mathcal{L} dt)^4}}. \quad (4.11)$$

where  $\int \mathcal{L} dt$  is the integrated luminosity and  $\text{cf}(\text{radiation})$  is the radiative correction factors. We use the integrated luminosity values and the radiative correction factors from [6]. The actual event numbers, integrated luminosity values and the radiative correction factor can be found in the tables A.6, A.7 in the appendix and the calculated cross sections for different  $\sqrt{s}$  values can be found in the table 4.3.

$\sqrt{s}$ [GeV]	$\sigma_{ee}$ [nb]	$\sigma_{mm}$ [nb]	$\sigma_{tt}$ [nb]	$\sigma_{qq}$ [nb]
88.47	$0.39 \pm 0.03$	$0.30 \pm 0.02$	$0.47 \pm 0.03$	$7.19 \pm 0.10$
89.46	$0.82 \pm 0.04$	$0.65 \pm 0.03$	$0.72 \pm 0.03$	$14.20 \pm 0.14$
90.22	$1.26 \pm 0.04$	$1.16 \pm 0.03$	$1.12 \pm 0.03$	$25.70 \pm 0.21$
91.22	$1.69 \pm 0.02$	$1.82 \pm 0.02$	$1.75 \pm 0.02$	$40.75 \pm 0.22$
91.97	$1.12 \pm 0.05$	$1.32 \pm 0.04$	$1.21 \pm 0.04$	$28.96 \pm 0.27$
92.96	$0.45 \pm 0.04$	$0.55 \pm 0.03$	$0.66 \pm 0.04$	$13.67 \pm 0.20$
93.71	$0.28 \pm 0.03$	$0.34 \pm 0.02$	$0.40 \pm 0.03$	$8.20 \pm 0.13$

**Table 4.3:** Calculated cross section values for different  $\sqrt{s}$  values.

#### 4.2.4 Forward Backward Asymmetry and Weak mixing angle

As discussed earlier in Section 2.3.2, the forward-backward asymmetry arises as a result of the different couplings of the  $Z^0$  to left handed and right handed fermions, and the weak mixing (Weinberg) angle can be extracted from this factor. We consider the  $\mu^- \mu^+$  decay channel as we don't have the problem of t-channel like in the  $e^- e^+$  decay mode. We also get two clear tracks in the case of  $\mu^- \mu^+$  decay mode. We can calculate the asymmetry by measure the event numbers for  $\cos\theta > 0$  and  $\cos\theta < 0$ . The forward backward asymmetry then is,

$$A_{fb} = \frac{N_+ - N_-}{N_+ + N_-} + \text{cf}(A_{fb}), \quad (4.12)$$

where  $N_+$  is the event numbers with  $\cos\theta > 0$  and  $N_-$  is the event numbers with  $\cos\theta < 0$ ,  $\text{cf}(A_{fb})$  is the  $A_{fb}$  correction factor, which can be found in the table A.7. The error in  $A_{fb}$  is given by,

$$\Delta A_{fb} = \sqrt{N_- \left( \frac{2N_+}{(N_+ + N_-)^2} \right)^2 + N_+ \left( \frac{2N_-}{(N_+ + N_-)^2} \right)^2}. \quad (4.13)$$

The calculated  $A_{fb}$  can be found in table 4.4.

With this, we can calculate the weak mixing angle. For the data events, we consider the  $A_{fb}$  value at

$\sqrt{s}$ [GeV]	$N_+$	$N_-$	$A_{fb}$
$\mu^- \mu^+$ data			
88.47	51	69	$-0.128 \pm 0.090$
89.46	148	156	$-0.007 \pm 0.057$
90.22	272	319	$-0.063 \pm 0.041$
91.22	4297	4384	$0.008 \pm 0.011$
91.97	387	380	$0.039 \pm 0.036$
92.96	173	123	$0.231 \pm 0.057$
93.71	189	150	$0.209 \pm 0.054$
$\mu^- \mu^+$ MC			
91.22	37807	38402	$0.010 \pm 0.014$

**Table 4.4:** Forward and Backward event numbers for  $\mu^- \mu^+$  mode data and MC events with the calculated  $A_{fb}$  values.

$Z^0$  resonance energy and calculate the weak mixing angle (using Eq. 2.8) as,

$$\sin^2 \theta_W = 0.2369 \pm 0.0085, \quad (4.14)$$

and for the MC event, we calculate the weak mixing angle as,

$$\sin^2 \theta_W = 0.2352 \pm 0.0098. \quad (4.15)$$

The literature value for weak-mixing angle is  $\sin^2 \theta_W = 0.23122 \pm 0.00003$  [3]. This value is within one standard deviation of both the results calculated with the  $\mu^- \mu^+$  MC and the data file.

#### 4.2.5 Lepton Universality

Lepton universality states that, since the masses of leptons are much smaller compared to that of the  $Z^0$  boson, the cross section of all the three leptonic decay modes must be the same, at  $Z^0$  resonance. From the table 4.3, we can see that at  $Z^0$  resonance ( $= 91.22 \text{ GeV}$ ),

$$\begin{aligned} \sigma_e &= 1.69 \pm 0.02 \text{ nb} \\ \sigma_\mu &= 1.82 \pm 0.02 \text{ nb} \\ \sigma_\tau &= 1.75 \pm 0.02 \text{ nb} \end{aligned} \quad (4.16)$$

We see that even though the values are almost equal, they are not exactly equal. The theoretical value calculated with values from [6] gives  $\sigma_f = 1.99 \text{ nb}$ . We see that the values calculated from the data are several standard deviations away from the theoretical value, even though they match up to the order and in fact, quite close. We postulate that this is due to the fact that we don't have the best possible cut criteria. That is, we are undercounting the event numbers to some extent, which results in the lower cross section values. And since the factor by which we undercount differs between different modes, we get cross sections that slightly differ from each other. This problem could be overcome by substantially increasing



the number of events. We calculate the ratios of the hadronic to the leptonic modes as follow,

$$\begin{aligned}\frac{\sigma_{had}}{\sigma_e} &= 24.18 \pm 0.30 \\ \frac{\sigma_{had}}{\sigma_\mu} &= 22.43 \pm 0.24 \\ \frac{\sigma_{had}}{\sigma_\tau} &= 23.28 \pm 0.26\end{aligned}\tag{4.17}$$

The ratios above directly reflect the deviations we had in the cross section values and a little greater than the theoretical value. The literature values are given by  $\sigma_{had}/\sigma_e = 20.804 \pm 0.050$ ,  $\sigma_{had}/\sigma_\mu = 20.785 \pm 0.033$  and  $\sigma_{had}/\sigma_\tau = 20.764 \pm 0.045$  respectively [15]. Again, we are several orders of standard deviation away. But, both they are of the same order of magnitude and comparable. This means that the problem in our data is not fundamental and could be overcome by improving our counts.

#### 4.2.6 Breit-Wigner Fit of Cross Section

Now we extract the all important quantities from our data. To do this, we fit our cross section values against a Breit-Wigner curve of the form,

$$\sigma_f(s) = \frac{12\pi}{M_Z^2} \frac{s\Gamma_e\Gamma_f}{(s - M_Z^2)^2 + \left(\frac{s\Gamma_Z}{M_Z}\right)^2} (\hbar^2 c^2),\tag{4.18}$$

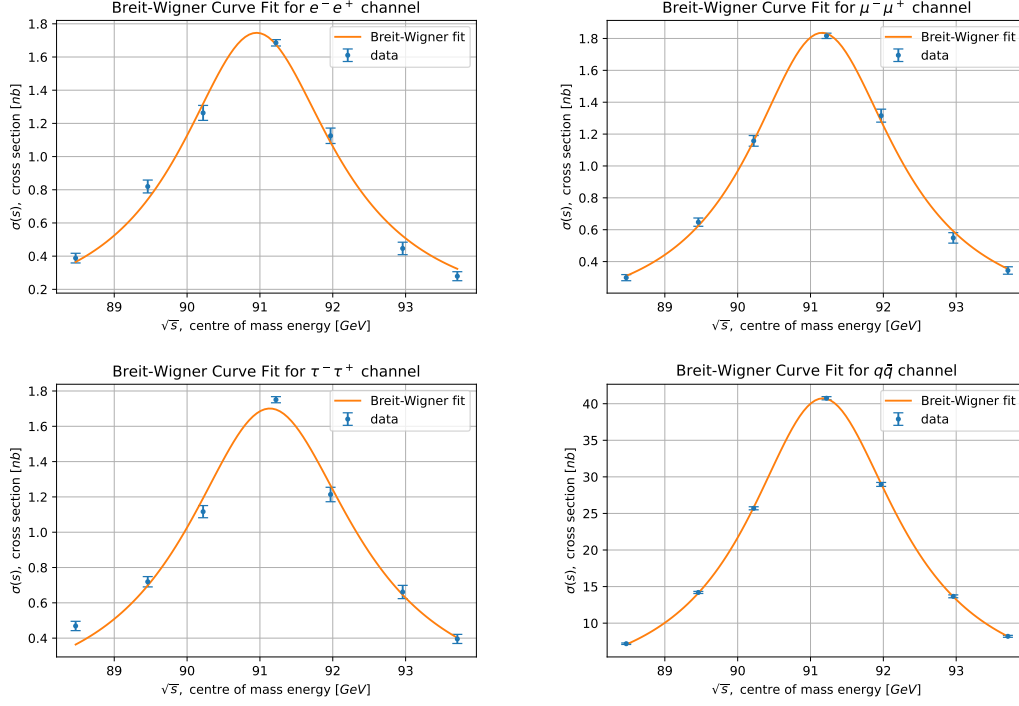
which has three parameters,  $M_Z$ ,  $\Gamma_Z$  and  $\Gamma_e\Gamma_f$ . The factor  $\hbar^2 c^2 = 3.893793719 \times 10^5 \text{nb GeV}^2$  is used to convert the cross section values to SI units from natural units. We used *scipy.optimize* module for *Python* to fit the function against our data. The plots are given in 4.2. The fit parameters obtained along with the reduced  $\chi^2$  values for the fit can be found in the table 4.5.

We have three parameters and seven data points. This gives us 4 degrees of freedom for our fit. The

Channel	$M_Z$ [GeV]	$\Gamma_Z$ [GeV]	$\Gamma_e\Gamma_f$ [GeV <sup>2</sup> ]	$\chi_{red}^2$
$e^-e^+$	$90.973 \pm 0.047$	$2.591 \pm 0.158$	$6.615 \times 10^{-3} \pm 6.29 \times 10^{-4}$	3.47
$\mu^-\mu^+$	$91.117 \pm 0.021$	$2.462 \pm 0.060$	$6.306 \times 10^{-3} \pm 2.46 \times 10^{-4}$	1.10
$\tau^-\tau^+$	$91.159 \pm 0.058$	$2.825 \pm 0.179$	$7.690 \times 10^{-3} \pm 7.85 \times 10^{-4}$	8.34
$q\bar{q}$	$91.187 \pm 0.004$	$2.515 \pm 0.010$	$0.1460 \pm 0.0009$	0.58

**Table 4.5:** Fit parameters and reduced  $\chi^2$  values for the data

$\mu^-\mu^+$  fit gives the best reduced  $\chi^2$  value. The reduced  $\chi^2$  value for  $q\bar{q}$  is less than 1, which implies that we are over-fitting the data or the error variances are overestimated. The reduced  $\chi^2$  value for  $e^-e^+$  is greater than 1, which implies that we are under-fitting the data or the error variances are underestimated. Whereas, the reduced  $\chi^2$  value for  $\tau^-\tau^+$  is much greater than 1, which implies that we don't have a good fit[16]. But, none of the fits have a corresponding *p - valule* less than 0.05, which means they're not significant enough to be considered seriously different from previous literature results[17]. One major issue with using reduced  $\chi^2$  statistic test here is that we have only a limited number of points. Therefore, we don't have a lot of freedom to fit the data with. This ultimately results in a bad value, even though the fit is seemingly nice giving decent values for the parameters. If we had such a good data with more data points, we should obtain a better reduced  $\chi^2$  value. Our opinion is that the residuals give a better picture on whether we have a good fit, in our case. And all the residuals here are very small compared to the absolute values that we fit, which means that we have a decent fit for our data. For a better



**Figure 4.2:** Breit-Wigner fits for the four different decay channels

quantitative analysis of goodness of fit, we suggest methods like *bootstrapping*. Reduced  $\chi^2$  statistic test is not considered meaningful for limited number of points by some sources[18].

The fit parameters directly give us the values of mass of the  $Z^0$  boson,  $M_Z$  and the decay width of the  $Z^0$  boson,  $\Gamma_Z$ . The mean value of the quantities are,

$$\begin{aligned} M_Z &= 91.123 \pm 0.020 \text{ GeV} \\ \Gamma_Z &= 2.598 \pm 0.062 \text{ GeV}. \end{aligned} \tag{4.19}$$

The literature values are  $M_Z = 91.1876 \pm 0.0021 \text{ GeV}$  and  $\Gamma_Z = 2.4952 \pm 0.0023 \text{ GeV}$  [15]. The literature value for  $M_Z$  lies within  $4\sigma$  of our calculated value, whereas the literature value for  $\Gamma_Z$  lies within  $2\sigma$  of our calculated value. We also note that the literature value for  $M_Z$  lies within one standard deviation of the value obtained from the  $q\bar{q}$  fit, which had the most counts. This hints us that we could improve our accuracy of the results by significantly increasing the number of events analysed.

#### 4.2.7 Partial Width of Different Channels and Number of Light Neutrino Generations

The third fit parameter doesn't directly give us the partial decay width of different channels. It gives the product of the partial decay width of  $e^-$  mode and the partial decay width of  $f$  mode, the final state fermion in consideration. In the case of  $e^-e^+$  channel, this reduces to  $\Gamma_e^2$ , which lets us calculate the partial decay width of the  $e^-$  mode. This can then be used to calculate the partial width of other channels. This is given in the table 4.6.

The literature values for  $\Gamma_f$  are taken from [15]. The literature values lie within one standard deviation of the calculated partial width, which is impressive. But we also note that the precision of our results is

Channel	$\Gamma_f$ [MeV]	$\Gamma_f$ (lit.) [MeV]
$e^-e^+$	$81.33 \pm 5.47$	$83.91 \pm 0.12$
$\mu^-\mu^+$	$77.53 \pm 6.03$	$83.99 \pm 0.18$
$\tau^-\tau^+$	$94.56 \pm 11.55$	$84.08 \pm 0.22$
$q\bar{q}$	$1795.52 \pm 121.31$	$1744.4 \pm 2.0$

**Table 4.6:** Partial decay widths of different channels and corresponding literature values

not as good as the literature value precision. How precise the value of  $\Gamma_f$  is, ultimately depends on the precision of the cross section values, which can be improved only by observing more events.

With all the decay widths in hand, we can now determine the number of generations of light neutrinos. We use the value for  $\Gamma_\nu$  from [6], which is  $\Gamma_\nu = 167.6 MeV$ . With this, we can calculate the number of neutrino generations using,

$$n_\nu = \frac{\Gamma_Z - \Gamma_e - \Gamma_\mu - \Gamma_\tau - \Gamma_q}{\Gamma_\nu}. \quad (4.20)$$

This gives the number of neutrino generations as,

$$n_\nu = 3.28 \pm 0.45. \quad (4.21)$$

This tells us that number of neutrino generations are 3. And within one standard deviation, our result absolutely excludes the possibility of fourth neutrino generation.

#### 4.2.8 Discussion

One of the main systematic error in this experiment would be due to the systematic errors of the OPAL detector. The systematic error associated with this is not mentioned in the manual for this experiment [6]. For example, we do not have information about uncertainties associated with different components of the detector. An analysis including all those factors will be beyond the scope of this experiment, in our opinion. This uncertainty also impacts the main quantity we measured in this experiment,  $N$ . In our analysis, we had assumed a statistical error of  $\sqrt{N}$ , but we have no knowledge of other systematic errors on this quantity.

And for some of the values taken from the manual for this experiment, like radiative corrections, there was no mention of any uncertainty in the quantity. Ideally, we should identify and add uncertainties in these quantities to improve our accuracy.

# Conclusion

Our objective to understand how analysis is carried out with data from a particle collider was achieved in this experiment. We started with event display analysis, which let us develop a cut criteria to separate out different channels involved in the decay of  $Z^0$ . This was then further refined with the help of *Monte-Carlo* simulation events. Once we had the final cuts, we went about calculating the efficiency matrix, which was used to determine the actual counts. This was then used to calculate the cross sections. The cross section values at  $Z^0$  resonance was then used to verify lepton universality. We noted that even though the cross sections are about the same for the three channels, they are several standard deviations away from the literature value. The measurement of forward-backward asymmetry in the muon channel helped us determine the weak mixing angle. The literature value lies within  $1\sigma$  of the value obtained here. We then fit the cross sections of different channels at different centre of mass energies to *Breit-Wigner curve*, which enabled us to determine the mass and decay width of  $Z^0$  and the partial decay widths of the different channels. The literature value for the mass of  $Z^0$  lies within  $4\sigma$  of the calculated value and the literature value for the decay width of  $Z^0$  lies within  $2\sigma$  of the calculated value, which is impressive. The literature value for partial decay widths of all the channels lie within  $1\sigma$  of the calculated values. We also determined the maximum possible light neutrino generations and eliminated the possibility of a fourth light neutrino within  $1\sigma$ . All this analysis ultimately relied on the measurement of  $N$ , which implies that we could improve both the accuracy and precision of the results by significantly increasing  $N$ . Systematic errors, if known quantitatively, would also improve our results.

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# Data

**Table A.1:** Event display analysis data for  $e^-e^+$  channel (pure data).

EVENT	Ctrk(N)	Ctrk(Sump)	Ecal(SumE)	Hcal(SumE)
1	2	50.9	82.6	0
2	2	91.9	90	0
3	3	82.5	92.3	0
4	2	80.9	86.8	0
5	2	38.1	89.5	0
6	2	83.8	87.5	0
7	2	87.4	93.2	0
8	2	69.3	90.7	0
9	2	86.1	89.4	0.5
10	2	90.3	90.6	0
11	2	92.1	88.5	0.5
12	3	81.7	91.6	0
13	2	89.6	92.5	0
14	2	61.1	89.2	0
15	3	88.4	89.1	0
16	2	90.9	90.5	0.3
17	2	64.6	88.8	0
18	2	95.6	96.2	0
19	2	93	90.8	0
20	2	94.1	89.2	0

**Table A.2:** Event display analysis data for  $\mu^-\mu^+$  channel (pure data).

EVENT	Ctrk(N)	Ctrk(Sump)	Ecal(SumE)	Hcal(SumE)
1	2	90.1	1.6	7
2	2	93	1.6	8.7
3	2	96.8	2	0
4	2	89.1	2.3	8.5
5	2	90.25	1.5	7.2
6	2	91.8	1.8	4.3
7	2	86.3	3.7	3.3
8	2	99.2	1.3	2.9
9	2	88.2	1.6	3
10	2	90.9	1.3	6.7
11	2	95.6	2.5	6.1
12	2	75.3	3.1	6.8
13	2	85.2	5.8	4.4
14	2	98.6	3.6	5.7
15	2	86.8	1.9	7.9
16	2	98	1.9	2
17	2	108.3	2	8.5
18	2	92.4	3.6	6.7
19	2	92	1.9	22.6
20	2	92.6	3.6	5.7

**Table A.3:** Event display analysis data for  $\tau^-\tau^+$  channel (pure data).

EVENT	Ctrk(N)	Ctrk(Sump)	Ecal(SumE)	Hcal(SumE)
1	5	74	51.1	10.2
2	2	46.5	17.3	8.2
3	2	30.8	1.6	6.3
4	2	29.5	10.2	4.1
5	2	33.1	1.5	10.6
6	2	24.4	12.4	11.7
7	4	36	16.1	5.7
8	2	41.3	11.1	20
9	2	49.7	5.2	20.3
10	2	33.4	23.6	6.9
11	2	14.1	3.3	6.3
12	2	19.7	15.9	3.8
13	2	26.8	16.5	3.4
14	5	23.4	27	17.1
15	2	23.8	29.4	3.6
16	2	39	18.9	4.4
17	2	24.1	46.5	7.3
18	5	38.5	28.5	0
19	2	35.3	51.8	2.3
20	2	17.8	2.5	5

**Table A.4:** Event display analysis data for  $q\bar{q}$  channel(pure data).

EVENT	Ctrk(N)	Ctrk(Sump)	Ecal(SumE)	Hcal(SumE)
1	15	37.7	37	14.1
2	17	39.2	66.8	9.9
3	46	64.6	53	13
4	36	45.3	53.2	7.7
5	41	59.9	53.2	13.8
6	9	21.9	65.2	8.8
7	16	55.9	50.4	24.3
8	30	38.1	68.3	13.8
9	22	34.4	75.5	6.2
10	36	51.2	62.3	5.5
11	23	63.1	56	17.2
12	23	59	60.6	8.5
13	26	62.2	67.2	20.4
14	30	43.3	71.7	4.3
15	40	47.8	61.4	5.7
16	19	67.9	52.1	10.6
17	14	52.1	61	4.4
18	29	82.6	53.8	16.4

**Table A.5:** Event display analysis for mixed data

EVENT	Ctrk(N)	Ctrk(Sump)	Ecal(SumE)	Hcal(SumE)	Channel
1	19	39.5	44.3	15.6	$q\bar{q}$
2	36	42.8	57.1	12.5	$q\bar{q}$
3	2	95.7	93.4	0	$e^-e^+$
4	2	90.8	1.4	4.1	$\mu^-\mu^+$
5	4	36.5	35.8	10.8	$\tau^-\tau^+$
6	2	97	2.2	8.9	$\mu^-\mu^+$
7	68	42.9	48.5	6.2	$q\bar{q}$
8	5	35	40.8	3.3	$\tau^-\tau^+$
9	21	75.8	45.8	21	$q\bar{q}$
10	2	95.2	1.3	7.9	$\mu^-\mu^+$
11	2	22.7	34.4	0	$\tau^-\tau^+$
12	4	44.3	37.8	2.6	$\tau^-\tau^+$
13	21	53.1	36.2	22.9	$q\bar{q}$
14	2	89.5	92	0	$e^-e^+$
15	2	89.1	89.7	0	$e^-e^+$
16	2	4.1	4.4	0	$\tau^-\tau^+$
17	2	87.8	1.4	4.3	$\mu^-\mu^+$
18	2	75.3	90	0	$e^-e^+$
19	2	93.7	1.6	6.8	$\mu^-\mu^+$
20	2	67.1	93.6	0	$e^-e^+$

**Table A.6:** Observed event numbers and actual events numbers for different  $\sqrt{s}$  values.

$\sqrt{s}$ [GeV]	$e^-e^+$	$\mu^-\mu^+$	$\tau^-\tau^+$	$q\bar{q}$
	$N_{obs}$			
88.47	106	120	251	3281
89.46	261	304	425	7413
90.22	415	591	693	14709
91.22	4839	8581	10178	221068
91.97	393	767	863	18728
92.96	150	296	427	8100
93.71	178	339	459	8635
	$N_{actual}$			
88.47	$202 \pm 20$	$141 \pm 13$	$256 \pm 18$	$3508 \pm 61$
89.46	$496 \pm 31$	$358 \pm 21$	$416 \pm 23$	$7927 \pm 92$
90.22	$789 \pm 39$	$697 \pm 29$	$661 \pm 30$	$15729 \pm 130$
91.22	$920 \pm 137$	$1023 \pm 111$	$971 \pm 120$	$236399 \pm 541$
91.97	$747 \pm 38$	$904 \pm 33$	$820 \pm 33$	$20026 \pm 147$
92.96	$285 \pm 23$	$349 \pm 20$	$419 \pm 23$	$8662 \pm 96$
93.71	$338 \pm 25$	$400 \pm 22$	$448 \pm 24$	$9234 \pm 100$

**Table A.7:** Integrated luminosity values for *data6.root* and radiative and  $A_{FB}$  corrections for different  $\sqrt{s}$  values[6].

$\sqrt{s}$ [GeV]	$\int \mathcal{L} dt [\text{nb}]^{-1}$	Rad. Correction [nb]		$A_{FB}$ Correction
		Hadronic	Leptonic	
88.47	$675.9 \pm 5.7$	+2.0	+0.09	0.021512
89.46	$800.8 \pm 6.6$	+4.3	+0.20	0.019262
90.22	$873.7 \pm 7.1$	+7.7	+0.36	0.016713
91.22	$7893.5 \pm 54.3$	+10.8	+0.52	0.018293
91.97	$825.3 \pm 6.9$	+4.7	+0.22	0.030286
92.96	$624.6 \pm 5.5$	-0.2	-0.01	0.062196
93.71	$942.2 \pm 7.7$	-1.6	-0.08	0.093850



# Plots

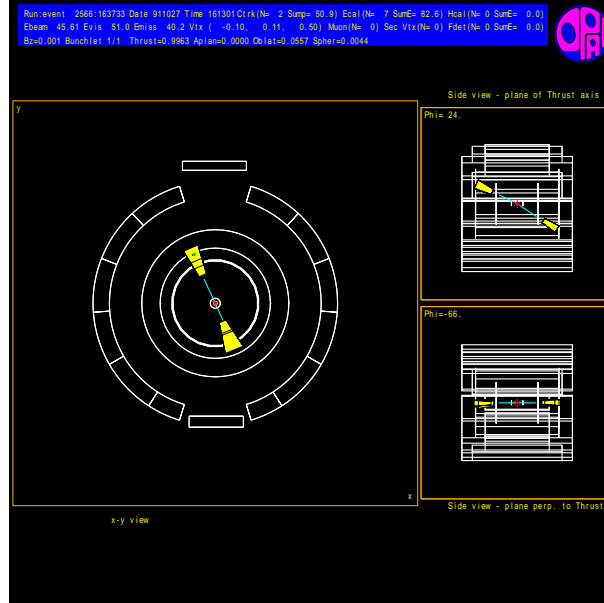


Figure B.1: Event display of  $e^-e^+$ -channel sample

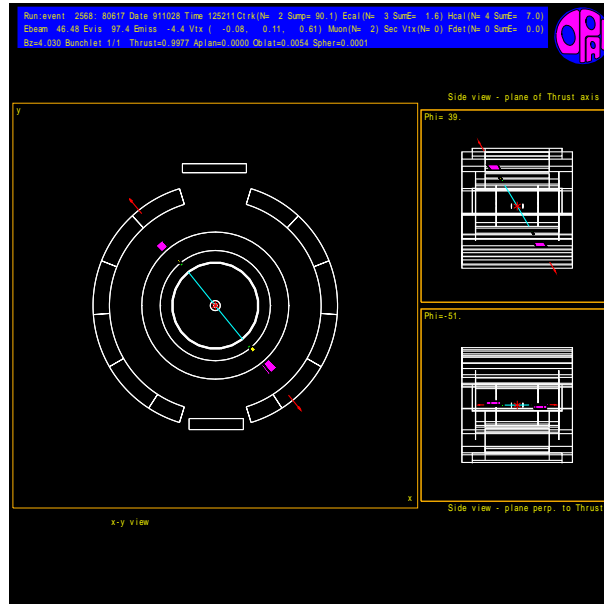
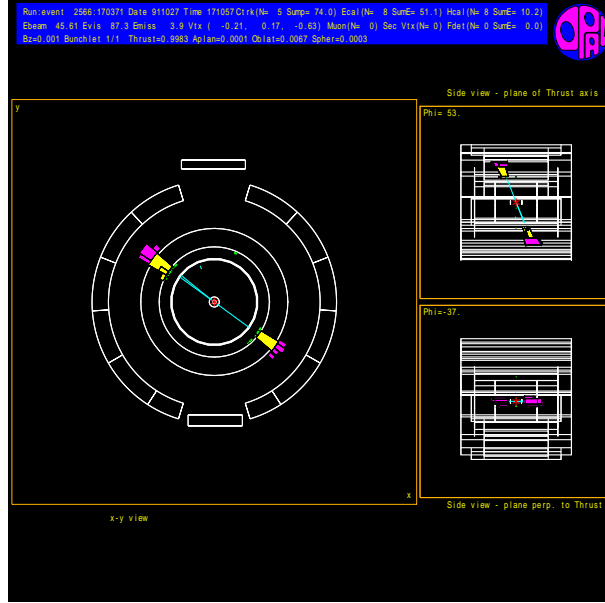
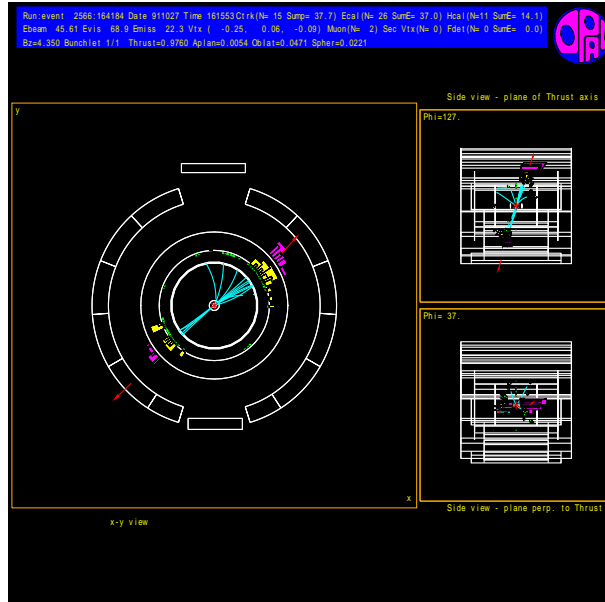


Figure B.2: Event display of  $\mu^-\mu^+$ -channel sample



**Figure B.3:** Event display of  $\tau^- \tau^+$ -channel sample



**Figure B.4:** Event display of  $q\bar{q}$ -channel sample

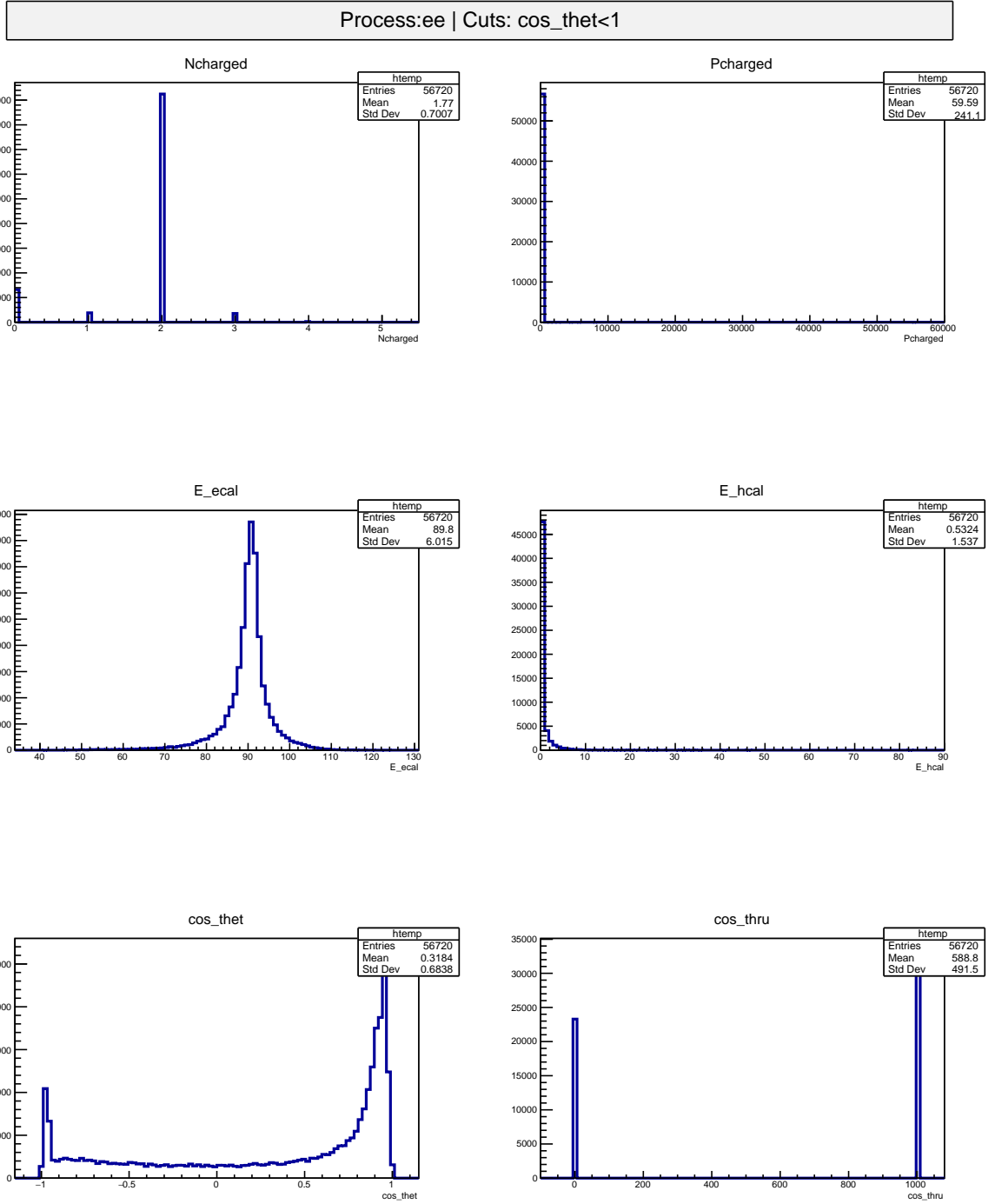


Figure B.5: Global cuts applied to  $e^-e^+$  MC

Process:ee | Cuts: Ncharged>0 && Ncharged<6 && E\_eal>80 && E\_hcal<1 && cos\_thet>-0.9 && cos\_thet<0.5

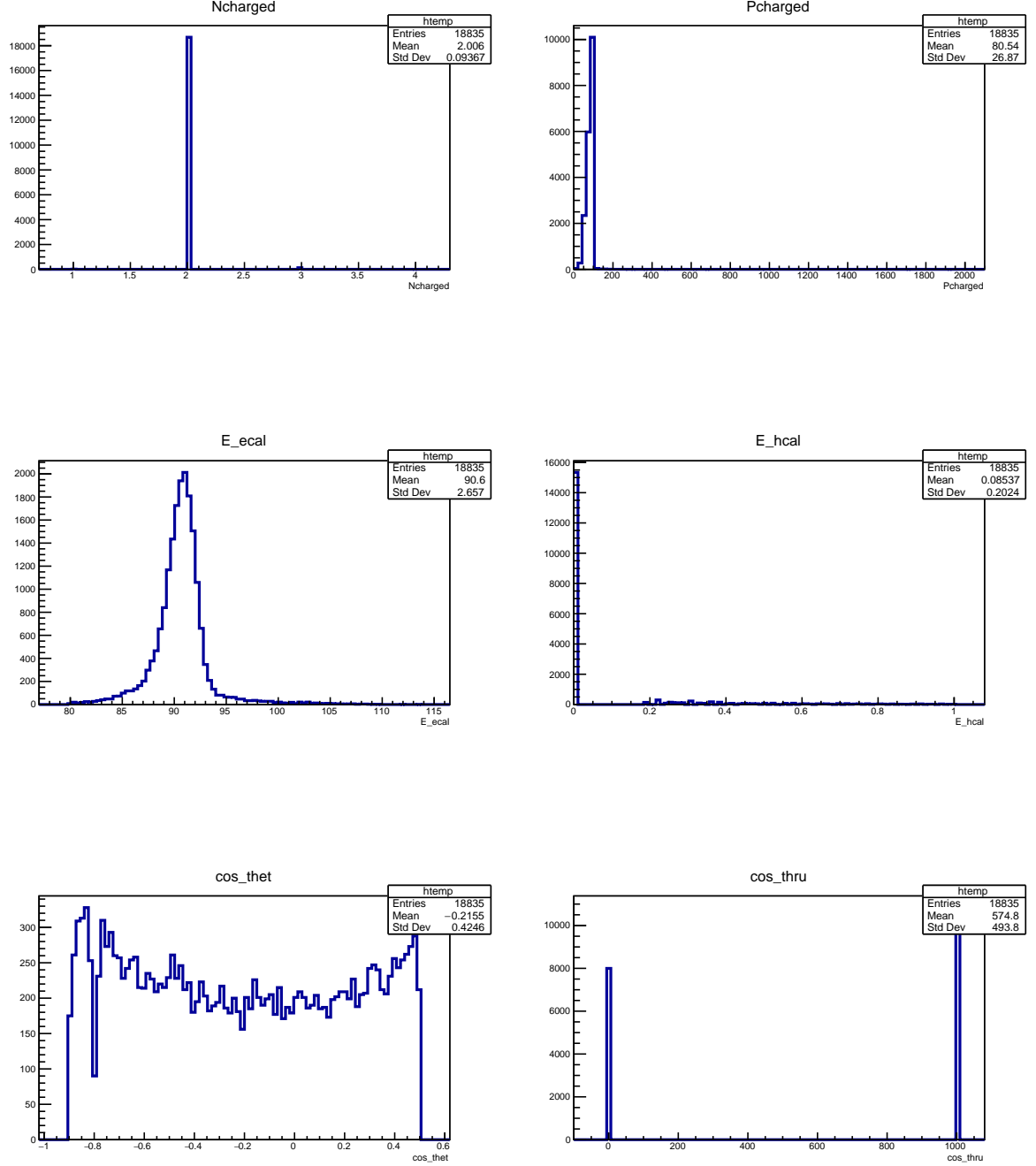
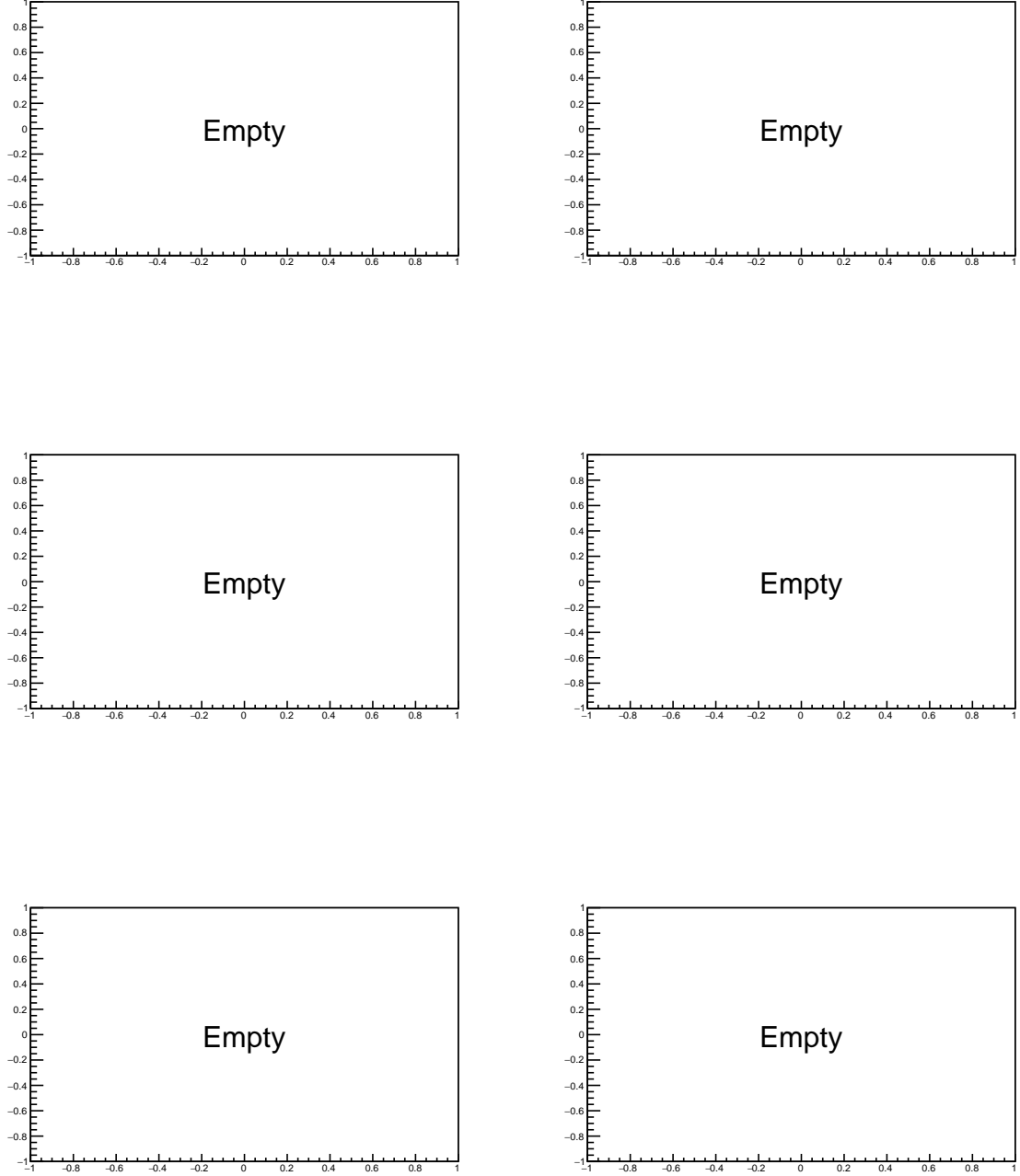


Figure B.6:  $e^-e^+$  cuts applied to  $e^-e^+$  MC

Process:ee | Cuts: Ncharged<6 && Pcharged>75 && E\_eal<10 && cos\_thet>-0.9 && cos\_thet<0.9



**Figure B.7:**  $\mu^-\mu^+$  cuts applied to  $e^-e^+$  MC

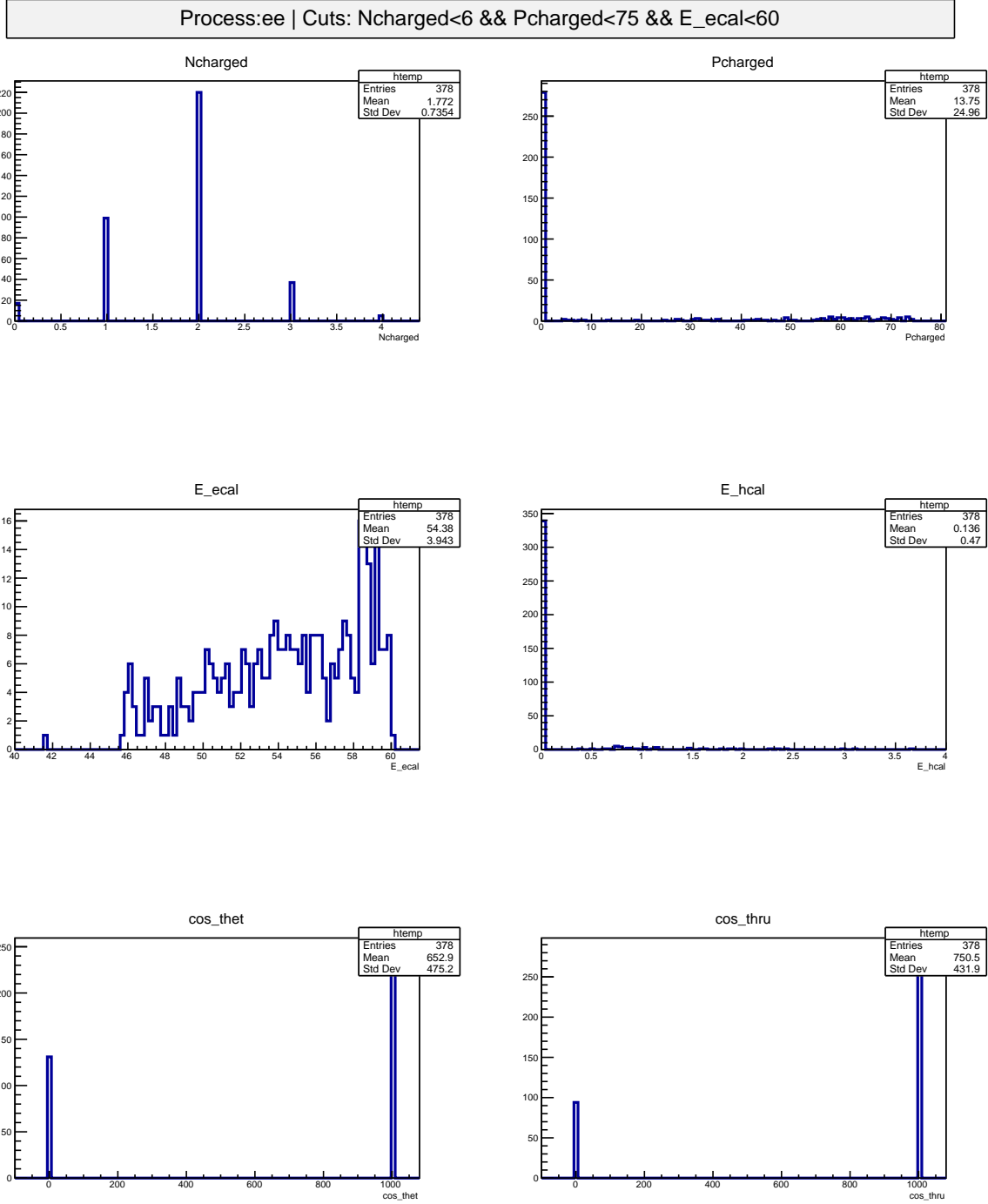
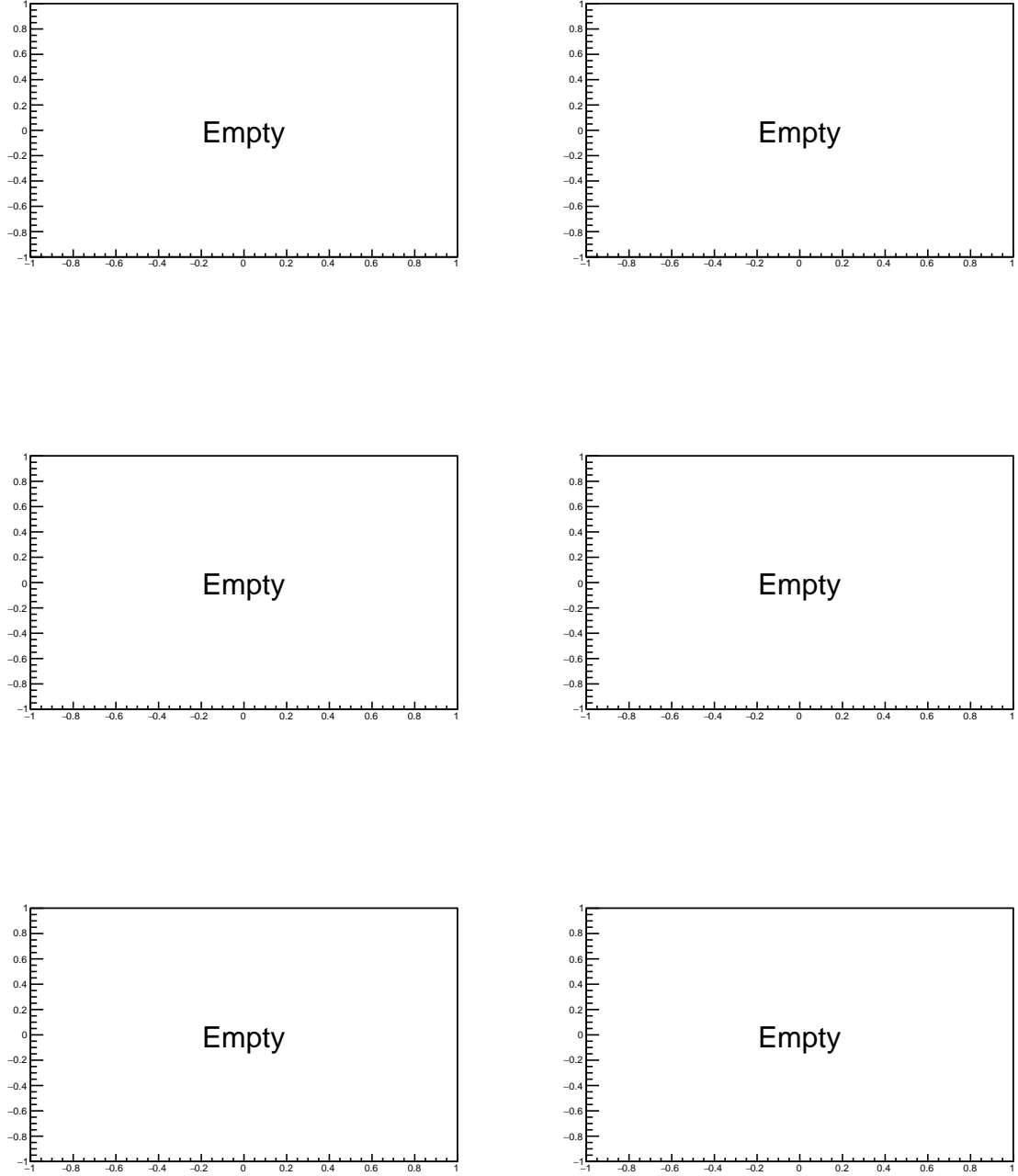


Figure B.8:  $\tau^-\tau^+$  cuts applied to  $e^-e^+$  MC

Process:ee | Cuts: Ncharged>8 && E\_ecal>36 && E\_ecal<79



**Figure B.9:**  $q\bar{q}$  cuts applied to  $e^-e^+$  MC

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