

## 0.1 Homework 3

### 0.1.1 Schwinger boson representation

A two-dimensional quantum harmonic oscillator contains two decoupled free bosons, whose annihilation operators can be represented as  $a$  and  $b$  respectively.  $a = \frac{1}{\sqrt{2}}(x + ip_x)$ ,  $b = \frac{1}{\sqrt{2}}(y + ip_y)$ . They satisfy the commutation relations  $[a, a^\dagger] = [b, b^\dagger] = 1$  and  $[a, b] = [a, b^\dagger] = 0$ . This system has  $U(2)$  symmetry, which includes an  $SU(2)$  subgroup. Let's explore how to construct the  $SU(2)$  representation using bosonic operators. Define  $S^x = \frac{1}{2}(a^\dagger b + b^\dagger a)$ ,  $S^z = \frac{1}{2}(a^\dagger a - b^\dagger b)$ .

1. Express  $S^y$  in terms of  $a$  and  $b$ . [Hint: Make  $\vec{S} \times \vec{S} = i\vec{S}$ ]

2. Prove that  $S^y$  is actually related to the angular momentum operator of the harmonic oscillator  $L = xp_y - yp_x$ , namely  $S^y = \frac{L}{2}$ .

□

3. Define the following set of states, where  $s = 0, 1/2, 1, \dots$ , and  $m = -s, -s+1, \dots, s-1, s$  (they are called the Schwinger boson representation),

$$|s, m\rangle = \frac{(a^\dagger)^{s+m}}{\sqrt{(s+m)!}} \frac{(b^\dagger)^{s-m}}{\sqrt{(s-m)!}} |\Omega\rangle$$

where  $|\Omega\rangle$  is the state annihilated by  $a$  and  $b$ , i.e.,  $a|\Omega\rangle = b|\Omega\rangle = 0$ . Prove that the state  $|s, m\rangle$  is indeed a simultaneous eigenstate of  $\vec{S}^2 = (S^x)^2 + (S^y)^2 + (S^z)^2$  and  $S^z$ , with eigenvalues  $s(s+1)$  and  $m$  respectively. [Hint: Use the particle number basis.]

□

### 0.1.2 1D tight-binding model

The Hamiltonian of a periodic tight-binding chain of length  $L$  is given by the following expression:

$$H_{\text{chain}} = -t \sum_{n=1}^L \left( \hat{a}_n^\dagger \hat{a}_{n+1} + \hat{a}_{n+1}^\dagger \hat{a}_n \right)$$

where  $t$  is the hopping matrix element between adjacent sites  $n$  and  $n+1$ ,  $\hat{a}_n^\dagger$  creates a fermion at site  $n$ , and the set of operators  $\{a_n^\dagger, a_n; n = 1, \dots, L\}$  satisfies the standard anticommutation relations:

$$\{a_n, a_{n'}^\dagger\} = \delta_{nn'}, \quad \{a_n, a_{n'}\} = 0, \quad \{a_n^\dagger, a_{n'}^\dagger\} = 0$$

We assume periodic boundary conditions, i.e., we consider  $a_{L+n}^\dagger = a_n^\dagger$ . The purpose of this problem is to prove that this Hamiltonian can be diagonalized by a linear transformation of the discrete Fourier transform form:

$$b_k^\dagger = \frac{1}{\sqrt{L}} \sum_{n=1}^L e^{ikn} a_n^\dagger$$

1. Let's require that  $b_k^\dagger$  remains invariant under any shift of the summation index  $n \rightarrow n + n'$  ("translation invariance"). Prove that this implies that the index  $k$  is quantized and determine the set of allowed  $k$  values. How many independent  $b_k^\dagger$  operators are there?

2. Verify that the set of  $b_k$  and  $b_k^\dagger$  operators also satisfies the above standard anticommutation relations. That is:

$$\{b_k, b_{k'}^\dagger\} = \delta_{kk'}, \quad \{b_k, b_{k'}\} = 0, \quad \{b_k^\dagger, b_{k'}^\dagger\} = 0$$

**Hint:** Use the identity  $\sum_{m=1}^L e^{i\frac{2\pi}{L}m} = 0$ .

3. Prove that the inverse transformation of the above has the form:

$$a_n^\dagger = \frac{1}{\sqrt{L}} \sum_k e^{-ikn} b_k^\dagger$$

where the sum is over the set of allowed  $k$  values determined in (a).

4. Show that  $b_k^\dagger$  is indeed a creation operator of a single-particle eigenstate of  $H_{\text{chain}}$  by proving that its commutator with the Hamiltonian has the form  $[H_{\text{chain}}, b_k^\dagger] = \varepsilon_k b_k^\dagger$ . Give the explicit expression for the corresponding eigenvalue  $\varepsilon_k$ .