

Notes for Qballs With Phaseshift

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1 ODE for Qballs

Q - balls are a class of spherically symmetric solutions in 3 + 1 dimensions arising in a family of field theories.

Here, we examine the modeling of a 2 + 1 dimension case of these such solutions. In previous results, we already know that the Q ball is governed by the following equation:

$$\ddot{\phi} - \nabla^2 \phi + \phi - A|\phi| \phi + B|\phi^2| \phi = 0 \quad (1)$$

where $\phi = \phi(r, t)$, a radially symmetric solution that is varying with respect to time. Previously, work has been done to seek solutions in form of

$$\phi(r, t) = \sigma(r)e^{i\omega t} \quad (2)$$

and generate numerical approximations accordingly. We notice that we can add a phase shift factor, which motivates the following proposal:

$$\phi(r, t) = \sigma(r)e^{i(\omega t + k)} \quad (3)$$

where k is the phase shift. Then if we substitute equation (3) into equation (1), we still obtain the same equation governing $\sigma(r)$:

$$\nabla^2 \sigma(r) + (\omega^2 - 1)\sigma(r) + A\sigma(r)^2 + B\sigma(r)^3 = 0 \quad (4)$$

This means that there are no change to the ordinary differential equations governing the Q balls with the additional phase change factor. We know that this equation is:

2 Boost

But normally, we are interested at Q balls moving at relativistic speeds, in which we need to consider the Lorentz Transformations. These are normally called boost.

We use

$$\phi'(t', x', y') = \sigma(x', y')e^{i(\omega t + k)} = \phi'_1 + i\phi'_2 \quad (5)$$

to denote the Q ball field in the rest frame. Notice then,

$$\phi'_1(t', x', y') = \sigma(x', y') \cos(\omega t' + k) \quad (6)$$

and

$$\phi'_2(t', x', y') = \sigma(x', y') \sin(\omega t' + k) \quad (7)$$

and the corresponding time derivatives π'_1 , π'_2 of ϕ'_1 and ϕ'_2 correspondingly are

$$\pi'_1(t', x', y') = \frac{\partial \phi'_1}{\partial t'} \frac{\partial t'}{\partial t} + \frac{\partial \phi'_1}{\partial x'} \frac{\partial x'}{\partial t} + \frac{\partial \phi'_1}{\partial y'} \frac{\partial y'}{\partial t} \quad (8)$$

$$\pi'_2(t', x', y') = \frac{\partial \phi'_2}{\partial t'} \frac{\partial t'}{\partial t} + \frac{\partial \phi'_2}{\partial x'} \frac{\partial x'}{\partial t} + \frac{\partial \phi'_2}{\partial y'} \frac{\partial y'}{\partial t} \quad (9)$$

Now we apply the lorentz transformations $x = \gamma(x' - vt')$, $y' = y$, and $t = \gamma(t' - vx')$

$$\phi_1(t, x, y) = \sigma(\gamma(x + vt), y) \cos(\omega\gamma(t + vx) + k) \quad (10)$$

$$\phi_2(t, x, y) = \sigma(\gamma(x + vt), y) \sin(\omega\gamma(t + vx) + k) \quad (11)$$

$$(12)$$

and for π'_1 and π'_2 , we have

$$\pi_1(t, x, y) = \gamma \frac{\partial \phi'_1}{\partial t'} + \gamma v \frac{\partial \phi'_2}{\partial x'} \quad (13)$$

$$= -\omega\phi_2 + \gamma v \frac{\partial \sigma(\gamma(x + vt), y)}{\partial x} \cos(\omega\gamma(t + vx) + k) \quad (14)$$

$$\pi_2(t, x, y) = \gamma \frac{\partial \phi'_1}{\partial t'} + \gamma v \frac{\partial \phi'_2}{\partial x'} \quad (15)$$

$$= \omega\phi_1 + \gamma v \frac{\partial \sigma(\gamma(x + vt), y)}{\partial x} \sin(\omega\gamma(t + vx) + k) \quad (16)$$

At $t = 0$, we then have all together:

$$\phi_1(t, x, y) = \sigma(\gamma x, y) \cos(\omega\gamma vx + k) \quad (17)$$

$$\phi_2(t, x, y) = \sigma(\gamma x, y) \sin(\omega\gamma vx + k) \quad (18)$$

$$\pi_1(t, x, y) = -\omega\phi_2 + \gamma v \frac{\partial \sigma(\gamma x, y)}{\partial x} \cos(\omega\gamma vx + k) \quad (19)$$

$$\pi_2(t, x, y) = \omega\phi_1 + \gamma v \frac{\partial \sigma(\gamma x, y)}{\partial x} \sin(\omega\gamma vx + k) \quad (20)$$

And this can be calculated numerically by the program. Notice the partial derivatives can be done using the np.grad function.

3 Anti-Qball

Now we want to consider an anti Qball, which takes the form

$$\phi(r, t) = f(r)e^{-i\omega t} \quad (21)$$

Now, with this form, we substitute into

$$\ddot{\phi} - \nabla^2 \phi + \phi - A|\phi| \phi + B|\phi^2| \phi = 0 \quad (22)$$

And we obtain

$$\nabla^2 \sigma(r) + (\omega^2 - 1)\sigma(r) + A\sigma(r)^2 + B\sigma(r)^3 = 0 \quad (23)$$

which is the same governing differential equation as before. Thus this implementation of the anti Q ball should not be too hard to complete. We will hopefully have some time to edit the code this week.