Introduction to the FDTD method

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FDTD (= finite-difference time-domain)

Principle

- 1. Start from Maxwell's equations
- 2. Replace all derivatives with finite-difference approximations
- 3. Done

Maxwell's curl equations

$$\nabla \times \overline{E} = -\mu \frac{\partial}{\partial t} \overline{H},$$

$$\nabla \times \overline{E} = -\mu \frac{\partial}{\partial t} \overline{H},$$

$$\nabla \times \overline{H} = \overline{J} + \sigma \overline{E} + \epsilon \frac{\partial}{\partial t} \overline{E},$$

$$\left(\overline{H} = \frac{1}{\mu}\overline{B}\right)$$

$$\begin{cases} \mu \frac{\partial}{\partial t} H_x = \frac{\partial}{\partial z} E_y - \frac{\partial}{\partial y} E_z \\ \mu \frac{\partial}{\partial t} H_y = \frac{\partial}{\partial x} E_z - \frac{\partial}{\partial z} E_x \\ \mu \frac{\partial}{\partial t} H_z = \frac{\partial}{\partial y} E_x - \frac{\partial}{\partial x} E_y \end{cases}$$

$$\begin{cases} \epsilon \frac{\partial}{\partial t} E_x = \frac{\partial}{\partial y} H_z - \frac{\partial}{\partial z} H_y - \sigma E_x \\ \epsilon \frac{\partial}{\partial t} E_y = \frac{\partial}{\partial z} H_x - \frac{\partial}{\partial x} H_z - \sigma E_y \\ \epsilon \frac{\partial}{\partial t} E_z = \frac{\partial}{\partial x} H_y - \frac{\partial}{\partial y} H_x - \sigma E_z \end{cases}$$

Central difference approximations

Notation

$$x = i \Delta x, \ y = j \Delta y, \ z = k \Delta z$$
 $t = n \Delta t$
$$F(x, y, z, t) = F(i \Delta x, j \Delta y, k \Delta z, n \Delta t) \Rightarrow F^{n}(i, j, k)$$

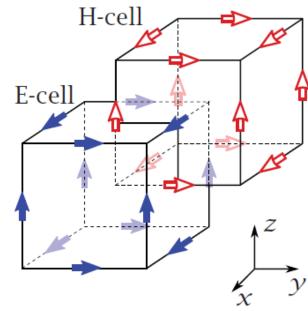
$$\frac{\partial F(x,y,z,t)}{\partial x} \approx \frac{F^n(i+\frac{1}{2},j,k) - F^n(i-\frac{1}{2},j,k)}{\Delta x} \qquad O(\Delta x^2)$$

$$\frac{\partial F(x,y,z,t)}{\partial y} \approx \frac{F^n(i,j+\frac{1}{2},k) - F^n(i,j-\frac{1}{2},k)}{\Delta y} \qquad O(\Delta y^2)$$

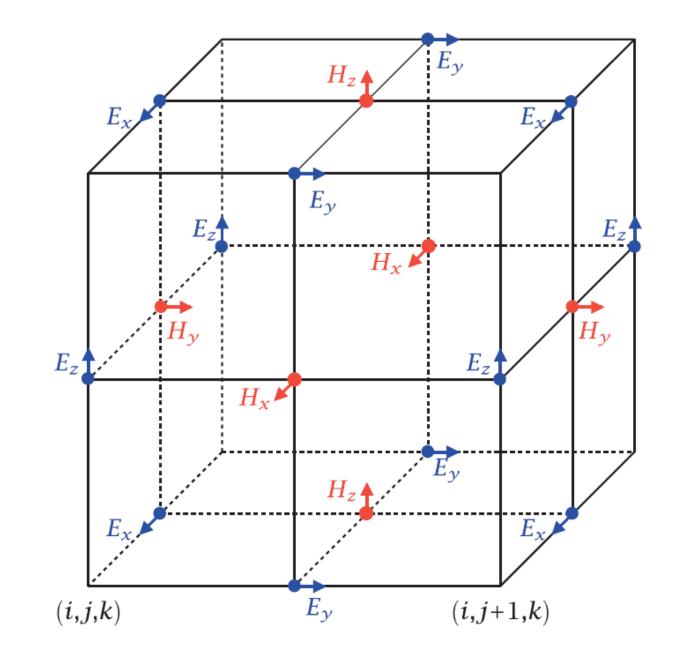
$$\frac{\partial F(x,y,z,t)}{\partial z} \approx \frac{F^n(i,j,k+\frac{1}{2}) - F^n(i,j,k-\frac{1}{2})}{\Delta z} \qquad O(\Delta z^2)$$

$$\frac{\partial F(x,y,z,t)}{\partial t} \approx \frac{F^{n+\frac{1}{2}}(i,j,k) - F^{n-\frac{1}{2}}(i,j,k)}{\Delta t} O(\Delta t^2)$$

Yee ce (Yee, 1966)

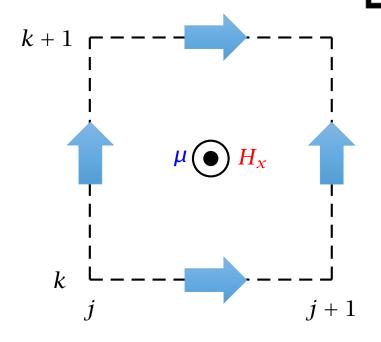


- → Electric field
- **➡** Magnetic field



Derivation of FDTD update equations

$$\mu \frac{\partial}{\partial t} H_x = \frac{\partial}{\partial z} E_y - \frac{\partial}{\partial y} E_z$$



$$\mu = \mu(i, j + \frac{1}{2}, k + \frac{1}{2})$$

$$\mu = \mu(i, j + \frac{1}{2}, k + \frac{1}{2})$$

$$\frac{\partial}{\partial t} H_x^n(i, j + \frac{1}{2}, k + \frac{1}{2}) \approx \frac{H_x^{n + \frac{1}{2}}(i, j + \frac{1}{2}, k + \frac{1}{2}) - H_x^{n - \frac{1}{2}}(i, j + \frac{1}{2}, k + \frac{1}{2})}{\Delta t}$$

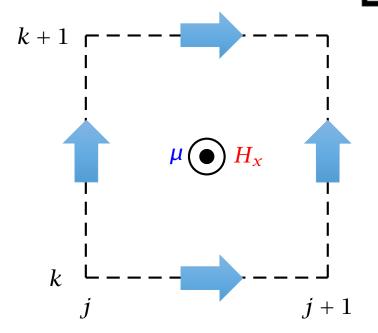
$$\frac{\partial}{\partial z} E_y^n(i, j + \frac{1}{2}, k + \frac{1}{2}) \approx \frac{E_y^n(i, j + \frac{1}{2}, k + 1) - E_y^n(i, j + \frac{1}{2}, k)}{\Delta z}$$

$$k \int_{j}^{k} ----\int_{j+1}^{l} \frac{\partial}{\partial z} E_{y}^{n}(i,j+\frac{1}{2},k+\frac{1}{2}) \approx \frac{E_{y}^{n}(i,j+\frac{1}{2},k+1) - E_{y}^{n}(i,j+\frac{1}{2},k)}{\Delta z}$$

$$\frac{\partial}{\partial y} E_z^n(i,j+\frac{1}{2},k+\frac{1}{2}) \approx \frac{E_z^n(i,j+1,k+\frac{1}{2}) - E_z^n(i,j,k+\frac{1}{2})}{\Delta y}$$

Derivation of FDTD update equations

$$\mu \frac{\partial}{\partial t} H_x = \frac{\partial}{\partial z} E_y - \frac{\partial}{\partial y} E_z$$



$$\mu(i,j+\frac{1}{2},k+\frac{1}{2})\underbrace{\frac{H_x^{n+\frac{1}{2}}(i,j+\frac{1}{2},k+\frac{1}{2}) - H_x^{n-\frac{1}{2}}(i,j+\frac{1}{2},k+\frac{1}{2})}_{\Delta t}}_{\Delta t}$$

$$(i,j+\frac{1}{2},k+1) - E_y^n(i,j+\frac{1}{2},k) \qquad E_z^n(i,j+1,k+\frac{1}{2}) - E_z^n(i,j,k+\frac{1}{2})$$

$$=\frac{E_y^n(i,j+\frac{1}{2},k+1)-E_y^n(i,j+\frac{1}{2},k)}{\Delta z}-\frac{E_z^n(i,j+1,k+\frac{1}{2})-E_z^n(i,j,k+\frac{1}{2})}{\Delta y}$$

FDTD update equations

$$k+1 + \frac{\Delta t}{\mu(i,j+\frac{1}{2},k+\frac{1}{2})\Delta z} + \frac{\Delta t}{\mu(i,j+\frac{1}{2},k+\frac{1}{2})\Delta y} + \frac{\Delta t}{\mu(i,j+\frac{1}{2},k+\frac{1}{2})\Delta y} = H_y^{n-\frac{1}{2}}(i+\frac{1}{2},j,k+\frac{1}{2})$$

$$H_x^{n+\frac{1}{2}}(i,j+\frac{1}{2},k+\frac{1}{2}) = H_x^{n-\frac{1}{2}}(i,j+\frac{1}{2},k+\frac{1}{2})$$

$$+\frac{\Delta t}{\mu(i,j+\frac{1}{2},k+\frac{1}{2})\Delta z}\left[E_{y}^{n}(i,j+\frac{1}{2},k+1)-E_{y}^{n}(i,j+\frac{1}{2},k)\right]$$

$$+\frac{\Delta t}{\mu(i,j+\frac{1}{2},k+\frac{1}{2})\Delta y}\left[E_z^n(i,j,k+\frac{1}{2})-E_z^n(i,j+1,k+\frac{1}{2})\right]$$

$$H_{\mathcal{Y}}^{n+\frac{1}{2}}(i+\frac{1}{2},j,k+\frac{1}{2}) = H_{\mathcal{Y}}^{n-\frac{1}{2}}(i+\frac{1}{2},j,k+\frac{1}{2})$$

$$+\frac{\Delta t}{\mu(i+\frac{1}{2},j,k+\frac{1}{2})\Delta x}\left[E_z^n(i+1,j,k+\frac{1}{2})-E_z^n(i,j,k+\frac{1}{2})\right]$$

$$+\frac{\Delta t}{\mu(i+\frac{1}{2},j,k+\frac{1}{2})\Delta z}\left[E_x^n(i+\frac{1}{2},j,k)-E_x^n(i+\frac{1}{2},j,k+1)\right]$$

FDTD update equations

$$H_z^{n+\frac{1}{2}}(i+\frac{1}{2},j+\frac{1}{2},k) = H_z^{n-\frac{1}{2}}(i+\frac{1}{2},j+\frac{1}{2},k)$$

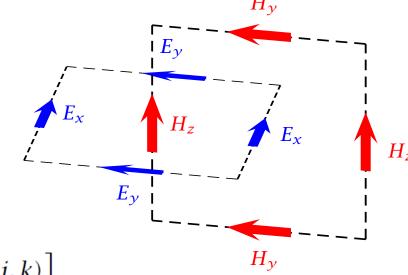
$$+\frac{\Delta t}{\mu(i+\frac{1}{2},j+\frac{1}{2},k)\Delta y}\left[E_x^n(i+\frac{1}{2},j+1,k)-E_x^n(i+\frac{1}{2},j,k)\right]$$

$$+\frac{\Delta t}{\mu(i+\frac{1}{2},j+\frac{1}{2},k)\Delta x}\left[E_{y}^{n}(i,j+\frac{1}{2},k)-E_{y}^{n}(i+1,j+\frac{1}{2},k)\right]$$

$$E_{x}^{n+1}(i+\frac{1}{2},j,k) = \frac{2\epsilon(i+\frac{1}{2},j,k) - \sigma(i+\frac{1}{2},j,k)\Delta t}{2\epsilon(i+\frac{1}{2},j,k) + \sigma(i+\frac{1}{2},j,k)\Delta t} E_{x}^{n}(i+\frac{1}{2},j,k)$$

$$+\frac{2\Delta t}{[2\epsilon(i+\frac{1}{2},j,k)+\sigma(i+\frac{1}{2},j,k)\Delta t]\Delta y}\left[H_z^{n+\frac{1}{2}}(i+\frac{1}{2},j+\frac{1}{2},k)-H_z^{n+\frac{1}{2}}(i+\frac{1}{2},j-\frac{1}{2},k)\right]$$

$$+\frac{2\Delta t}{[2\epsilon(i+\frac{1}{2},j,k)+\sigma(i+\frac{1}{2},j,k)\Delta t]\Delta z}\left[H_{y}^{n+\frac{1}{2}}(i+\frac{1}{2},j,k-\frac{1}{2})-H_{y}^{n+\frac{1}{2}}(i+\frac{1}{2},j,k+\frac{1}{2})\right]$$

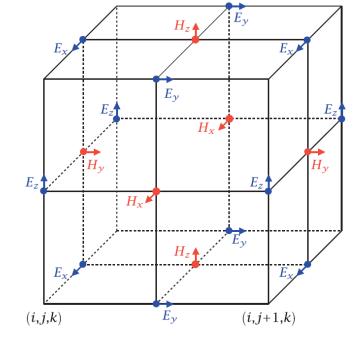


FDTD update equations

$$E_{y}^{n+1}(i,j+\frac{1}{2},k) = \frac{2\epsilon(i,j+\frac{1}{2},k) - \sigma(i,j+\frac{1}{2},k)\Delta t}{2\epsilon(i,j+\frac{1}{2},k) + \sigma(i,j+\frac{1}{2},k)\Delta t} E_{y}^{n}(i,j+\frac{1}{2},k)$$

$$+\frac{2\Delta t}{[2\epsilon(i,j+\frac{1}{2},k)+\sigma(i,j+\frac{1}{2},k)\Delta t]\Delta z}\left[H_x^{n+\frac{1}{2}}(i,j+\frac{1}{2},k+\frac{1}{2})-H_x^{n+\frac{1}{2}}(i,j+\frac{1}{2},k-\frac{1}{2})\right]$$

$$+\frac{2\Delta t}{[2\epsilon(i,j+\frac{1}{2},k)+\sigma(i,j+\frac{1}{2},k)\Delta t]\Delta x}\left[H_z^{n+\frac{1}{2}}(i-\frac{1}{2},j+\frac{1}{2},k)-H_z^{n+\frac{1}{2}}(i+\frac{1}{2},j+\frac{1}{2},k)\right]$$

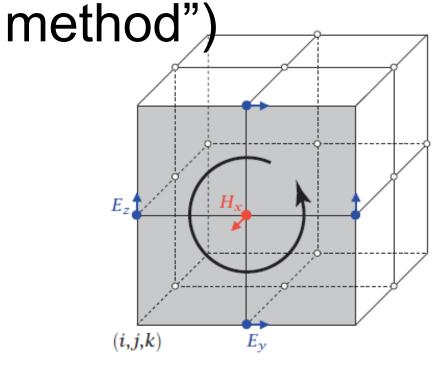


$$\frac{E_z^{n+1}(i,j,k+\frac{1}{2})}{2\epsilon(i,j,k+\frac{1}{2}) + \sigma(i,j,k+\frac{1}{2})\Delta t} E_z^n(i,j,k+\frac{1}{2})$$

$$+\frac{2\Delta t}{[2\epsilon(i,j,k+\frac{1}{2})+\sigma(i,j,k+\frac{1}{2})\Delta t]\Delta x}\left[H_{y}^{n+\frac{1}{2}}(i+\frac{1}{2},j,k+\frac{1}{2})-H_{y}^{n+\frac{1}{2}}(i-\frac{1}{2},j,k+\frac{1}{2})\right]$$

$$+\frac{2\Delta t}{[2\epsilon(i,j,k+\frac{1}{2})+\sigma(i,j,k+\frac{1}{2})\Delta t]\Delta y}\left[H_{x}^{n+\frac{1}{2}}(i,j-\frac{1}{2},k+\frac{1}{2})-H_{x}^{n+\frac{1}{2}}(i,j+\frac{1}{2},k+\frac{1}{2})\right]$$

Derivation of FDTD update equations from integral form of Maxwell's equations ("FIT



$$\nabla \times \overline{E} = -\mu \frac{\partial}{\partial t} \overline{H}$$

$$\downarrow \qquad \qquad \downarrow$$

$$\int_{S} \nabla \times \overline{E} \cdot \overline{dS} = -\int_{S} \mu \frac{\partial}{\partial t} \overline{H} \cdot \overline{dS}$$

$$\downarrow \qquad \qquad \downarrow$$

$$\int_{\partial S} \overline{E} \cdot \overline{dc} = -\int_{S} \mu \frac{\partial}{\partial t} \overline{H} \cdot \overline{dS}$$

Use the mid-ordinate numerical integration method

$$\int_{2\mathcal{E}} \overline{E} \cdot \overline{dc} \approx E_y^n(i, j + \frac{1}{2}, k) \Delta y + E_z^n(i, j + 1, k + \frac{1}{2}) \Delta z - E_y^n(i, j + \frac{1}{2}, k + 1) \Delta y - E_z^n(i, j, k + \frac{1}{2}) \Delta z$$

$$-\int\limits_{S}\mu\frac{\partial}{\partial t}\overline{H}\cdot\overline{dS}\approx-\mu(i,j+\tfrac{1}{2},k+\tfrac{1}{2})\frac{\left[H_{x}^{n+\tfrac{1}{2}}(i,j+\tfrac{1}{2},k+\tfrac{1}{2})-H_{x}^{n-\tfrac{1}{2}}(i,j+\tfrac{1}{2},k+\tfrac{1}{2})\right]}{\Delta t}\Delta y\Delta z$$

Example: 1D FDTD

$$\mu \frac{\partial}{\partial t} H_y = \frac{\partial}{\partial x} E_z$$
$$\epsilon \frac{\partial}{\partial t} E_z = \frac{\partial}{\partial x} H_y$$

$$H_y^{n+\frac{1}{2}}(i+\frac{1}{2}) = H_y^{n-\frac{1}{2}}(i+\frac{1}{2}) + \frac{\Delta t}{\mu(i+\frac{1}{2})\Delta x} \left(E_z^n(i+1) - E_z^n(i)\right)$$

$$E_z^{n+1}(i) = E_z(i)^n + \frac{\Delta t}{\epsilon(i)\Delta x} \left(H_y^{n+\frac{1}{2}}(i+\frac{1}{2}) - H_y^{n+\frac{1}{2}}(i-\frac{1}{2}) \right)$$

Implementation in MATLAB

$$\begin{split} H_y^{n+\frac{1}{2}}(i+\frac{1}{2}) &= H_y^{n-\frac{1}{2}}(i+\frac{1}{2}) + \frac{\Delta t}{\mu(i+\frac{1}{2})\Delta x} \left(E_z^n(i+1) - E_z^n(i) \right) \\ E_z^{n+1}(i) &= E_z(i)^n + \frac{\Delta t}{\epsilon(i)\Delta x} \left(H_y^{n+\frac{1}{2}}(i+\frac{1}{2}) - H_y^{n+\frac{1}{2}}(i-\frac{1}{2}) \right) \end{split}$$

```
% Magnetic field update equation
Hy(1:K-1) = Hy(1:K-1) + Db(1:K-1) .* (Ez(2:K) - Ez(1:K-1));
% Electric field update equation
Ez(2:K-1) = Ez(2:K-1) + Cb(2:K-1) .* (Hy(2:K-1) - Hy(1:K-2));
X coordinate 0
                    1.0
                                            3.0
                                                                              6.0
                                2.0
                                                       4.0
                                                                   5.0
H field indexing
F field
indexing
```

Contents

- Principle of FDTD
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- Basic properties
 - Stability
 - Dispersion
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- Advantages and weaknesses
- Examples

Stability

Courant-Friedrichs-Lewy (CFL) condition:

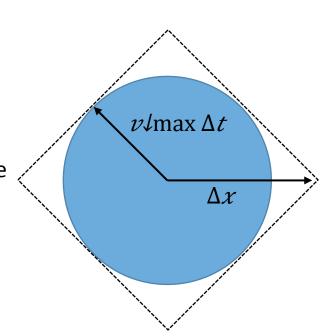
$$\Delta t \le \frac{1}{v_{\text{max}}} \frac{1}{\sqrt{\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2} + \frac{1}{\Delta z^2}}}$$

$$v_{\max} = \left\{ \frac{1}{\sqrt{\epsilon \mu}} \right\}_{\max}$$

In 3D:

$$\Delta x = \Delta y = \Delta z = \Delta$$
$$\Delta t \le \frac{\Delta}{\sqrt{3}v_{\text{max}}}$$

Numerical domain of dependence must include analytical domain of dependence



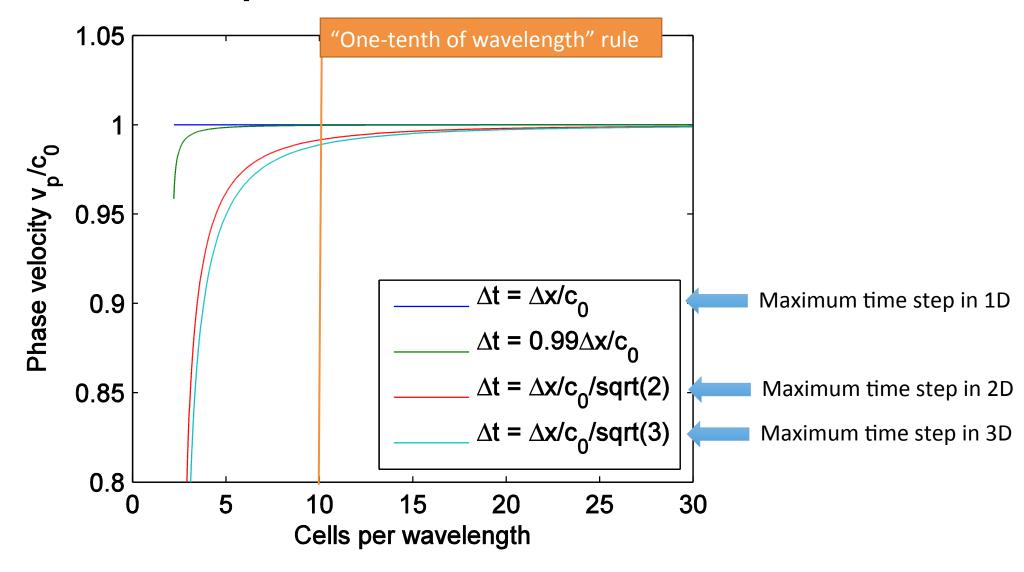
Numerical dispersion

Accurate formula
$$\frac{\omega^2}{c^2} = k_x^2 + k_y^2 + k_z^2$$

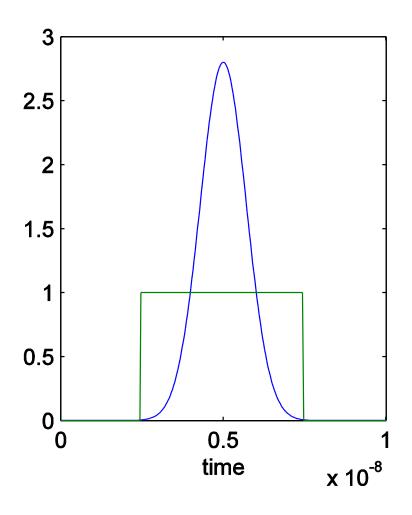
$$\frac{\sin^2(\omega\Delta t/2)}{(c\Delta t)^2} = \frac{\sin^2(\tilde{k}_x\Delta x/2)}{(\Delta x)^2} + \frac{\sin^2(\tilde{k}_y\Delta y/2)}{(\Delta y)^2} + \frac{\sin^2(\tilde{k}_z\Delta z/2)}{(\Delta z)^2}$$

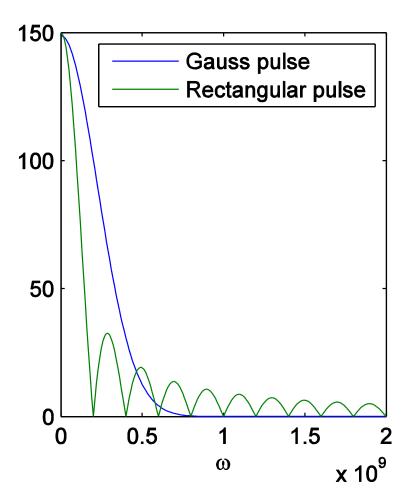
Numerical phase velocity
$$v_p=rac{\omega}{ ilde{k}}$$

Numerical dispersion



Spectra





Anisotropic dispersion in 2D and 3D

$$\Delta x = \Delta y = \Delta \qquad c\Delta t = \frac{\Delta}{2} \le \frac{\Delta}{\sqrt{3}}$$

$$\frac{\lambda}{\Delta} = 20$$

$$\lambda/\Delta = 10$$

$$\frac{v_p}{c} = \frac{0.98}{0.95}$$

$$\frac{0.96}{0.95}$$

$$0.00 \qquad 0.25 \qquad 0.5 \qquad 0.75 \qquad 1.00 \qquad 1.25 \qquad 1.50$$

$$\lambda/\Delta = 5$$

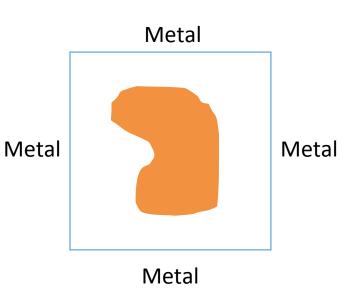
Example: 2D FDTD

- 2D FDTD and anisotropic dispersion
 - example3a.m

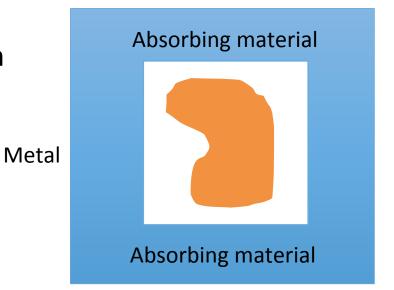
- Be careful with dielectric materials
 - wavelength is shorter => finer cell size is needed
 - example3b.m

Absorbing boundary conditions

- In the basic form, FDTD can only model boxes, with ideally conducting walls
 - How to terminate the computation domain to model free space?
- FDTD is very good at modelling different materials:
 - Berenger 1994: Make the walls of the box from an unphysical material that absorbs anything!
- ABCs are essential for any FDTD code
 - Example4.m, example5a, 5b, 5c



Metal



Metal

Metal

Contents

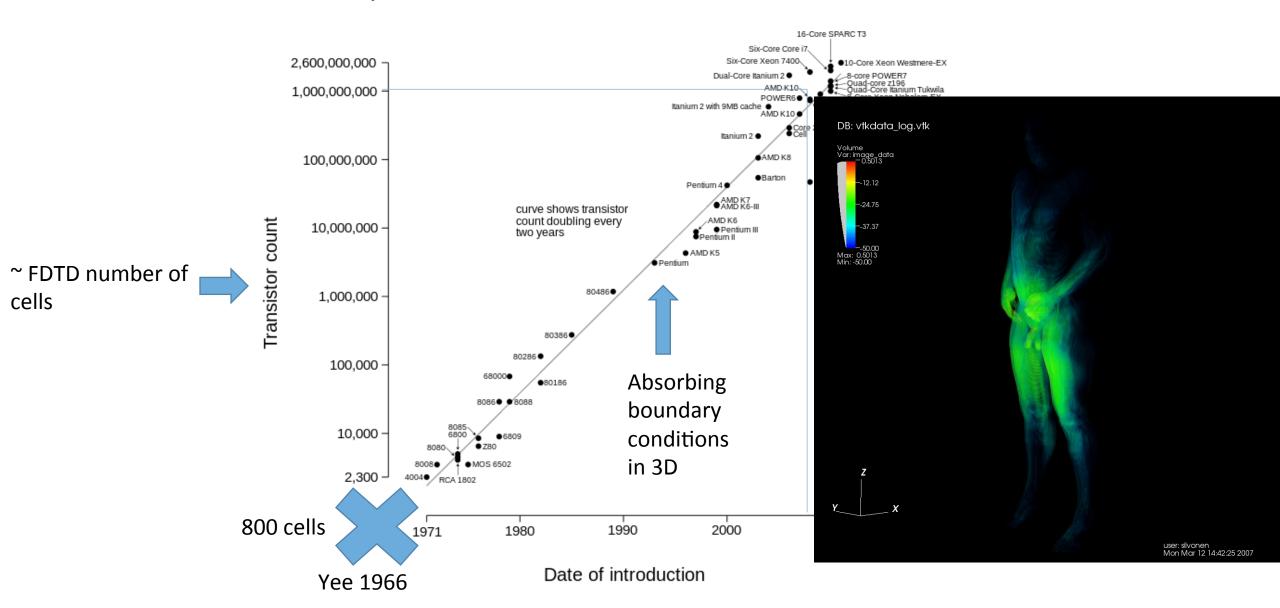
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Advantages of FDTD (1)

- Simple equations
- Can be parallelized easily
- Scales linearly with number of unknowns
- No need to solve equation systems

=> good for very large problems

Microprocessor Transistor Counts 1971-2011 & Moore's Law



Advantages of FDTD (2)

- Any kind of time-domain sources
- Geometry and boundary conditions are taken into account automatically. Any shape can be modeled easily
- Different media can be modelled naturally: non-linearity, inhomogeneity, anisotropy, complex geometry (metamaterials)



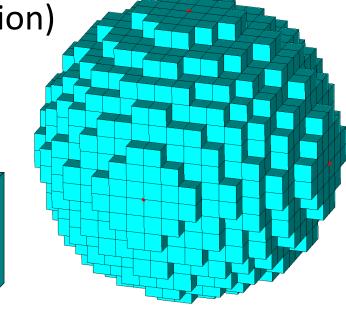
Weaknesses of FDTD (1)

Not good for slow phenomena (huge number of time steps needed)
 Example: 1 mm grid resolution -> time step = 1.9e-12 s
 Phenomenon lasting 1 ms -> Number of time steps = 5e8

• Curved shapes are problematic (staircase approximation)

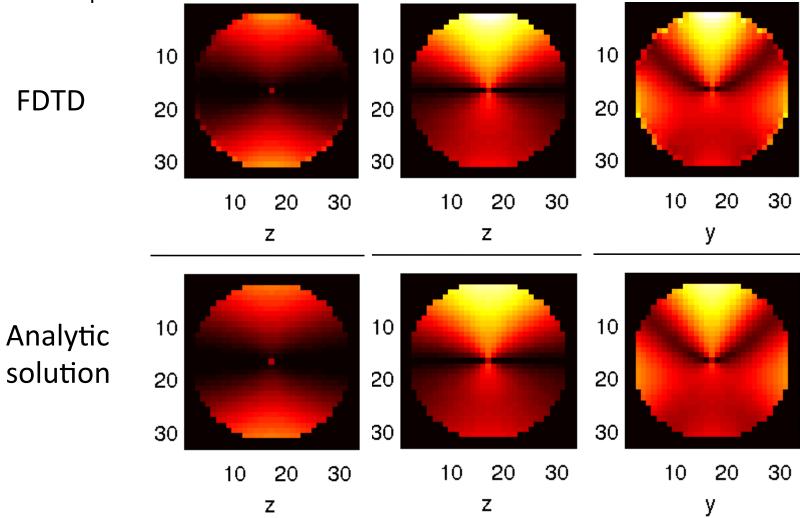
• High permittivity medium requires a fine grid

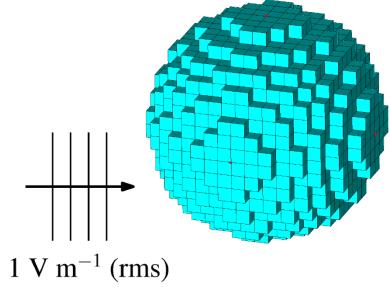
One-tenth of wavelength rule



Comparison with analytic solution

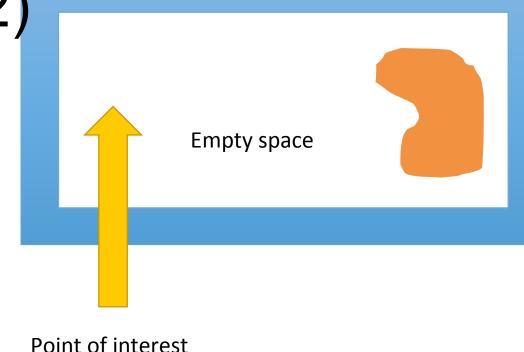
Radial component of the electric field in a sphere

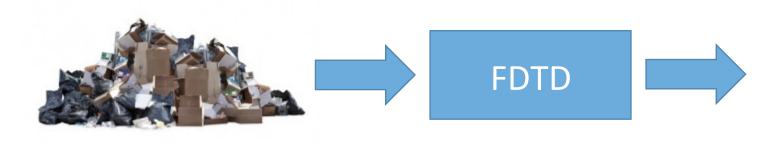




Weaknesses of FDTD (2)

- Computation domain must be finite
 - Absorbing boundary conditions
- Entire computational domain needs to be gridded (also empty space)
- Results depend on the choice of coordinate axes
- Error control







Error in FDTD

- Truncation error from difference approximation (time and space)
- Dispersion error, numerical anisotropy
- Unphysical Poynting theorem, conservation of energy
- Floating point (round-off) error
- Staircase approximation error
- Absorbing boundary condition error
- Modelling dielectric/lossy materials etc.

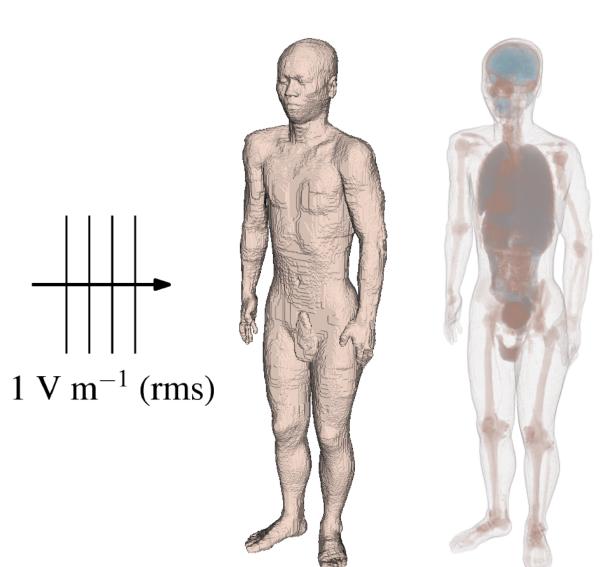
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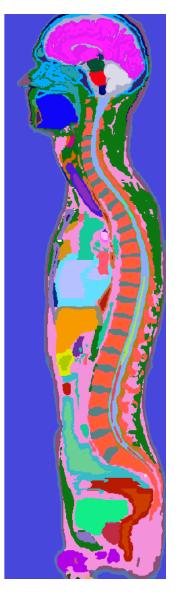
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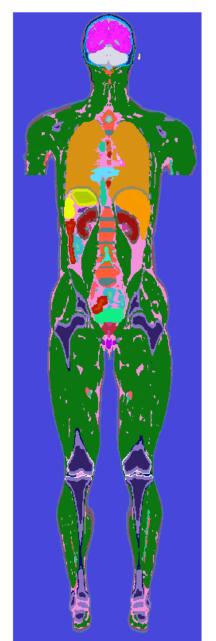
Applications of FDTD

- Radar cross section and scattering
- Metamaterials
- Antenna analysis example 7a
- Electronic component design
- Electromagnetic compatibility (EMC)
- Waveguides, resonators, filters example 7b
- Human exposure to EM waves

Human exposure to EM wav



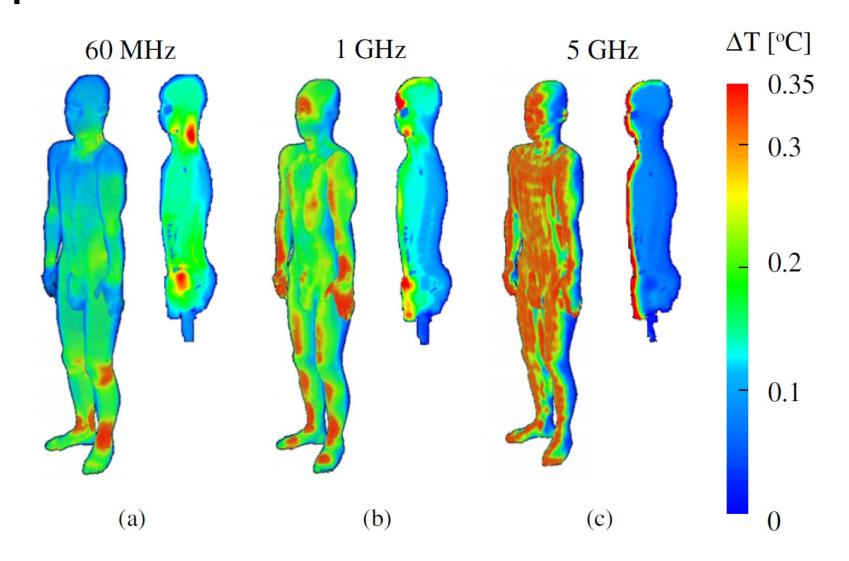




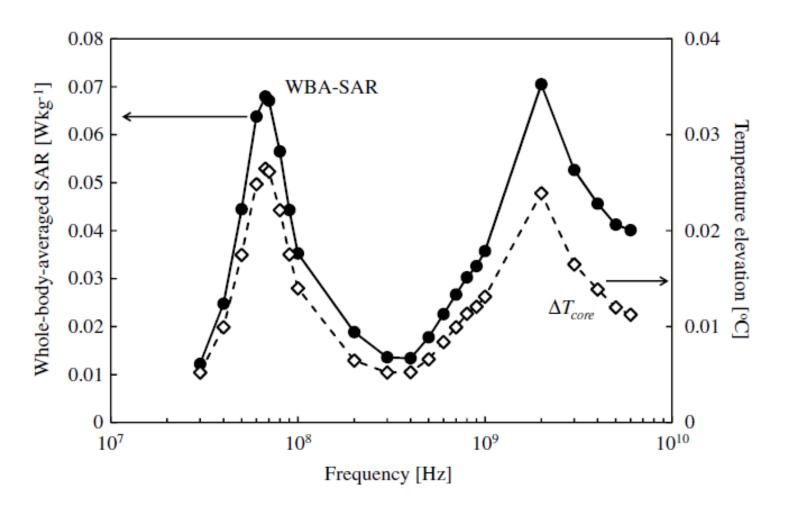
Energy absorption dΒ 0.00 -90.0 300 MHz 900 MHz 2450 MHz 70 MHz 1 m 33 cm 12 cm 4.3 m

Temperature rise

Absorbed power = 0.4 W/kg



Power absorption versus temperature rise



When to use FDTD?

- Use FDTD first
- Use FDTD for making animations
- Use FDTD for large heterogeneous geometries
- Use FDTD to model many things simultaneously

- Don't use FDTD at low frequencies
- Don't use FDTD with too large cell size
- Don't use FDTD if you need 99.9% accuracy