

Coconut

## Chapter 15

# Oscillations

$\phi$  also called epoch

$(wt + \phi) \rightarrow$  phase

$\phi \rightarrow$  phase constant

$x_m =$  amplitude

$$x = x_m \cos(wt + \phi)$$

$$\frac{dx}{dt} = v = -w x_m \sin(wt + \phi)$$

$$\frac{dv}{dt} = a = -w^2 x_m \cos(wt + \phi)$$

$$v(t) = -w x(t)$$

$$a(t) = -w^2 x(t)$$

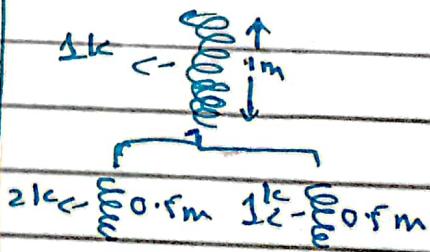
$$\omega = \sqrt{\frac{k}{m}}, \quad \omega = 2\pi f$$

$$2\pi f = \sqrt{\frac{k}{m}}$$

$$f = \frac{1}{2\pi} \sqrt{\frac{k}{m}} \Rightarrow T = 2\pi \sqrt{\frac{m}{k}}$$

$$\frac{v(t)}{x(t)} = \frac{-w x_m \sin \phi}{x_m \cos \phi} = -w \tan \phi$$

One spring with 'k' = 2 springs  
of each  $\frac{k}{2}$  in series



$\Rightarrow$  cutting the spring in half will double the spring constant.

$$F = -kx \rightarrow F_{max} = m a_{max}, \quad a_m = -w^2 x_m$$

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## Angular Frequency

$$x(t) = x_0 \cos(\omega t)$$

Cosine repeats when argument repeats by  $2\pi$

$$\omega t = \omega t + 2\pi$$

$$\omega t + \omega T = \omega t + 2\pi$$

$$\omega T = 2\pi$$

$$\omega = \frac{2\pi}{T} = \frac{2\pi}{f}$$

## Velocity

$$x(t) = x_m \cos(\omega t + \phi)$$

Taking derivative

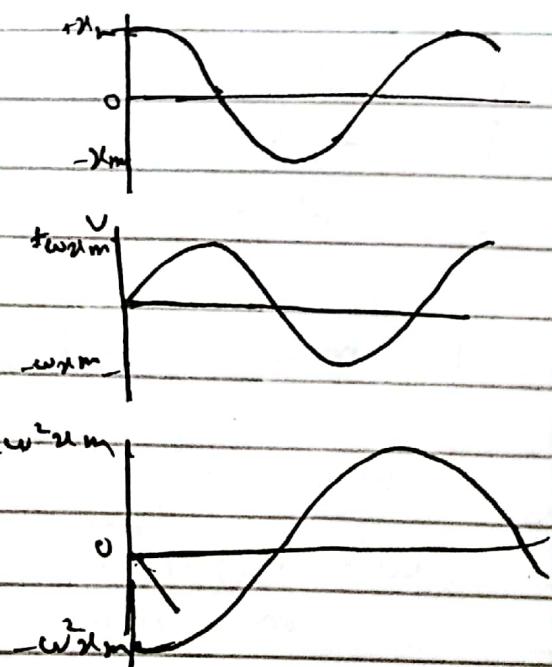
$$v(t) = \frac{dx(t)}{dt} = -x_m \omega \sin(\omega t + \phi)$$

## Acceleration

Phase shift =  $\pi/2$

$$a(t) = \frac{dv(t)}{t} = -\omega^2 x_m \cos(\omega t + \phi)$$

$$a(t) = -\omega^2 x(t)$$



## SHM EQUATION

## Force Law

$$F = ma = m(-\ddot{w}x) = -(m\omega^2)x$$

mass and  $\omega$  are constant

$$F = -kx$$

$$k = m\omega^2$$

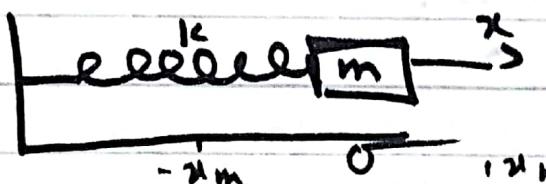
$$\omega^2 = \frac{k}{m} \Rightarrow \omega = \sqrt{\frac{k}{m}}$$

## MASS - SPRING SYSTEM

$$F \propto x$$

$$F \propto -x(t)$$

$$F = -kx(t) \quad (\text{Hooke's Law}) \quad ①$$



$$F = ma = m(-\omega^2 x(t)) \quad -②$$

$$k = mw^2$$

$$\omega = \sqrt{\frac{k}{m}} = \sqrt{\frac{m\omega^2}{m}} = \omega$$

$$\text{As } T = 2\pi \sqrt{\frac{m}{k}}$$

$$T = 2\pi \sqrt{\frac{m}{k}}$$

# TORSIONAL PENDULUM / ANGULAR SHO

$$F = -kx$$

Replace  $F \rightarrow T$  (Restoring torque)

$$x \rightarrow \theta$$

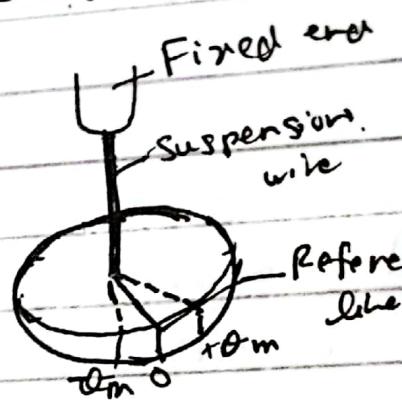
$$T = -k\theta(t)$$

$k$  is Greek  $\rightarrow k$  alpha

$\hookrightarrow$  called "Torsional Constant")

Depends

- $\hookrightarrow$  length of wire
- $\hookrightarrow$  Diameter of wire
- $\hookrightarrow$  material of suspension wire



$$\text{As } T = 2\pi \sqrt{\frac{m}{k}}$$

Replace  $k \rightarrow k$ ;  $m \rightarrow I$  (rotational Inertia)

$$T = 2\pi \sqrt{\frac{I}{k}} \quad (\text{Torsional Pendulum})$$

We know  $T = I\alpha$

$$I\alpha = -k\theta(t) \Rightarrow \alpha = -k\theta(t)/I \quad \text{--- ①}$$

$$\text{also } \alpha = -\omega^2 \theta(t)$$

Comparing

$$\omega^2 = \frac{k}{I} \Rightarrow \omega = \sqrt{\frac{k}{I}}$$

$I = \frac{1}{2}MR^2$  for disk with rotational axis in middle of disk

# SIMPLE PENDULUM

$mg \cos \theta$  cancels  $T$

$$F_{rz} = mg \sin \theta$$

$$\tau = -L mg \sin \theta$$

$$\text{As } \tau = I\alpha$$

$$I\alpha = -L mg \sin \theta$$

$$\alpha = -\frac{L mg}{I} \sin \theta$$

For small Angle  $\sin \theta \approx \theta$

$$\alpha = -\frac{L mg}{I} \theta$$

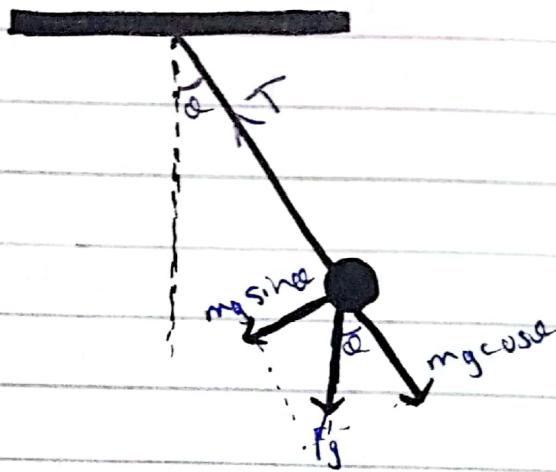
$$\alpha = -\omega^2 \theta$$

$$\omega^2 = \frac{L mg}{I}$$

$$\omega = \sqrt{\frac{L mg}{I}}$$

$$T = 2\pi \sqrt{\frac{I}{L mg}} \quad I = m L^2$$

$$T = 2\pi \sqrt{\frac{m L^2}{L mg}} = 2\pi \sqrt{\frac{L}{g}}$$



SHM is projection of uniform circular motion on diameter of circle in which the circular motion occurs.

## UNIFORM CIRCULAR MOTION AND SHM

using Projection method.

Displacement

$$\cos\theta = \frac{x}{x_m}$$

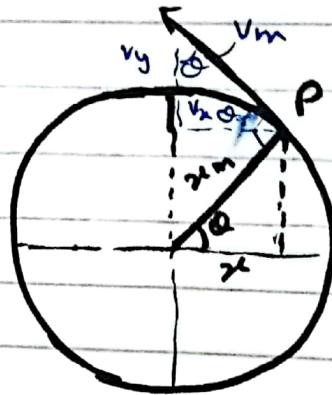
$$r = \frac{d}{t}$$

$$x(t) = x_m \cos\theta$$

$$\omega = \frac{\theta}{t}$$

$$x(t) = x_m \cos(\omega t + \phi)$$

$$\theta = \omega t$$



Velocity

$$\sin\theta = \frac{v(t)}{v_m}, \quad v \text{ is directed towards } -x\text{-axis}$$

$$\sin\theta = -V(t)/V_m$$

$$v(t) = -V_m \sin\theta$$

We know,  $V_m = \omega x_m$  and  $\theta = \omega t$

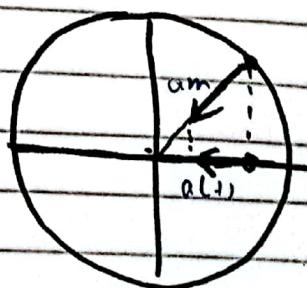
$$v(t) = -\omega x_m \sin(\omega t + \phi)$$

Acceleration

$$\cos\theta = -\frac{a(t)}{a_m} \quad (-ve \text{ in } x\text{-axis})$$

$$a(t) = -a_m \cos\theta$$

$$a(t) = -\omega^2 x_m \cos(\omega t + \phi)$$



## Electric Field Due to Electric Dipole

### Electric Dipole

Combination of two charges with equal magnitude and opposite sign placed at distance "d" apart is called an electric dipole.

### Electric Dipole Moment

The product of either charge and the displacement between them is called dipole moment

$$\vec{p} = q \vec{d}$$

From -ve to the charge.

### CASE-I

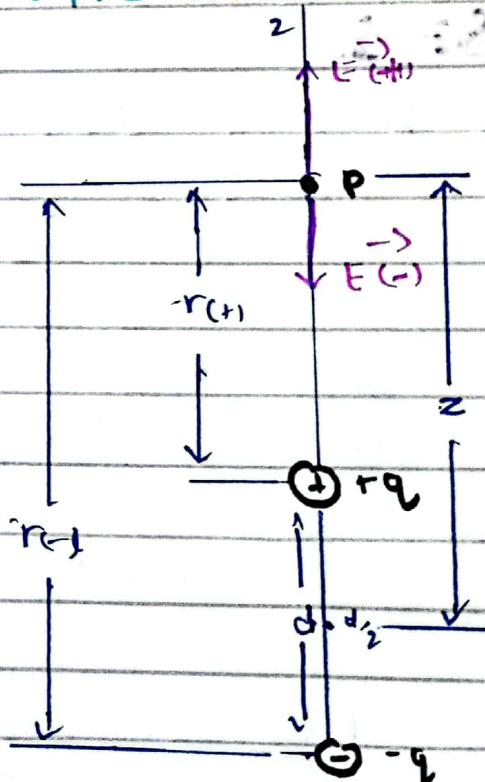
$$E_{\text{net}} = E(+)-E(-)$$

$$= \frac{1}{4\pi\epsilon_0} \frac{q}{r(+)} - \frac{1}{4\pi\epsilon_0} \frac{q}{r(-)}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{q}{(z-d_{12})^2} - \frac{1}{4\pi\epsilon_0} \frac{q}{(z+d_{12})^2}$$

$$= \frac{q}{4\pi\epsilon_0} \left[ \frac{1}{(z-d_{12})^2} - \frac{1}{(z+d_{12})^2} \right]$$

$$= \frac{q}{4\pi\epsilon_0} \left[ \frac{(z+d_{12})^2 - (z-d_{12})^2}{[(z-d_{12})(z+d_{12})]^2} \right]$$



$$= \frac{q}{4\pi\epsilon_0} \left[ \frac{2z_d}{[z^2 - (d/2)^2]^2} \right] \quad \text{But } z > d/2$$

$$\Rightarrow \frac{q}{4\pi\epsilon_0} \left[ \frac{2z_d}{2^4} \right]$$

As  $p = q \times d \Rightarrow d = p/q$

$$= \frac{q}{4\pi\epsilon_0} \frac{2 \cdot p}{2^3 \cdot 4}$$

$E_{\text{net}} =$

$$\frac{p}{2\pi\epsilon_0 z^3}$$

## CASE - II

$$E = \kappa \frac{Q}{R^2} \quad (\text{General Equation})$$

$$E_1 = \kappa \frac{q}{R^2}, \quad E_2 = \kappa \frac{q}{R^2}$$

$$E_{1x} = \kappa \frac{q}{R^2} \cos\alpha$$

$$\cos\alpha = \frac{x}{R}, \quad R = \sqrt{x^2 + a^2}$$

$$E_{1x} = \kappa \frac{q}{R^2} \cos\alpha = \frac{\kappa q \cdot x}{(x^2 + a^2) \sqrt{x^2 + a^2}} = \frac{\kappa q x}{(x^2 + a^2)(\sqrt{x^2 + a^2})}$$

$$= \frac{\kappa q x}{(x^2 + a^2)^{3/2}}$$

$$\text{Similarly, } E_{2x} = \frac{\kappa q x}{(x^2 + a^2)^{3/2}}$$

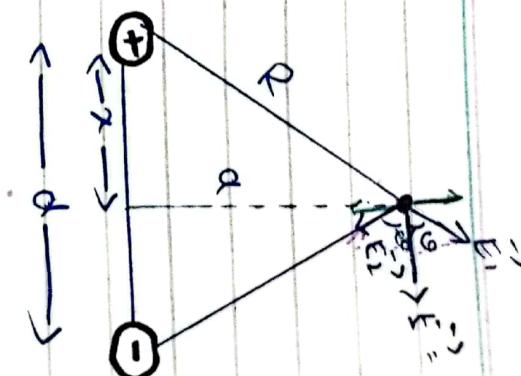
$$E_{\text{net}} = E_{1x} + E_{2x}$$

$$= \frac{2 \kappa q x}{(x^2 + a^2)^{3/2}} \quad \text{But } x = d_1 = \frac{2 \kappa q}{((d_2)^2 + a^2)^{3/2}} \cdot \frac{d_1}{2}$$

$$= \frac{\kappa q d_1}{(\alpha^2)^{3/2}} \quad (\alpha \gg d_2)$$

$$= \frac{\kappa q \frac{d_1}{\alpha^2}}{\alpha^3} \frac{\alpha^3}{q}$$

$$E_{\text{net}} = \frac{4\pi \epsilon_0 \alpha^3}{q}$$



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# Chapter 23 Gauss' Law

## Electric Flux

$$d\phi = \vec{E} \cdot d\vec{A}$$
$$\phi = \int \vec{E} \cdot d\vec{A}$$

( $d\vec{A}$  is patch vector area)  
(total flux)

$$\phi = \oint \vec{E} \cdot d\vec{A}$$

(net flux in closed surface (Gauss law))

$$\phi = (E \cos \alpha) A$$

## Gauss' Law And Coulomb's Law

Net flux  $\phi$  through a closed surface  $\phi = \frac{q}{\epsilon_0}$

$q$  is charge enclosed by surface,

$$\phi = \frac{Q}{\epsilon_0} \quad (Q = q_1 + q_2 + \dots + q_n)$$

$$\Rightarrow q = \phi \epsilon_0$$

$$\therefore \oint \vec{E} \cdot d\vec{A} = q_{\text{enc}} \quad (\text{Gauss Law})$$

When gaussian surface is spherical, circular

$\rightarrow E$  is constant

$$\cancel{\oint} \quad \cancel{E} \cdot E \oint dA = q$$



$$\text{Sum of patch areas of sphere} = 4\pi r^2$$

$$\epsilon_0 E (4\pi r^2) = q$$

$$E = \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r^2}$$

# Charged Isolated Conductor

The excess charge on isolated conductor will move entirely to the surface of the conductor. No excess charge will be found within the body of the conductor.

$$\text{Surface charge Density} = \sigma = \frac{Q}{A} = \frac{q}{A}$$

$$E = \frac{q}{4\pi r^2 \epsilon_0} = \frac{q}{A \epsilon_0}, \quad E = \frac{\sigma}{\epsilon_0} \quad (\text{For sphere only})$$

OR

$$\sigma = \frac{q}{A} \Rightarrow q = \sigma A$$

$$\text{Also } \epsilon_0 E f dA = q \Rightarrow \epsilon_0 E A = \sigma A$$

$$E = \frac{\sigma}{\epsilon_0}$$

# APPLICATIONS OF GAUSS LAW

## Cylindrical Symmetry

→ Positively charged Wire/Rod of Infinite Length.

$$\phi_{\text{total}} = \phi_{cs} + \phi_{\text{top}} + \phi_{\text{bottom}}$$

$$= E \cdot A \cos(0^\circ) + E \cdot A \cos(90^\circ) + E \cdot A \cos(90^\circ)$$

$$= EA_{cs}$$

$$\text{Curved surface area of cylinder} = 2\pi r L$$

$$\phi_{\text{total}} = E (2\pi r L)$$

$$\text{Linear charge density} = \lambda = \frac{q}{L}$$

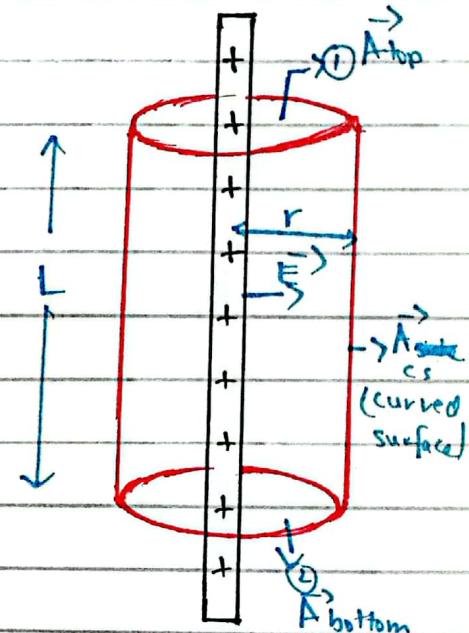
$$\Rightarrow \Phi_{\text{inside}} = \lambda \times L$$

Using Gauss' Law

$$\phi = \frac{Q_{\text{inside}}}{\epsilon_0}$$

$$E (2\pi r L) = \frac{\lambda L}{\epsilon_0}$$

$$E = \frac{\lambda}{2\pi r \epsilon_0}$$



Alternate

$$\epsilon_0 \oint \vec{E} \cdot d\vec{A} = q_{\text{enc}}$$

$$0 + EA_{cs} + 0 = \frac{q_{\text{enc}}}{\epsilon_0}$$

$$E (2\pi r L) = \frac{\lambda L}{\epsilon_0}$$

$$E = \frac{\lambda}{2\pi r \epsilon_0}$$

→ Magnitude of E-field at a distance r, direction is  $\perp$  to line of charge

# PLANNER SYMMETRY

## Nonconducting Sheet

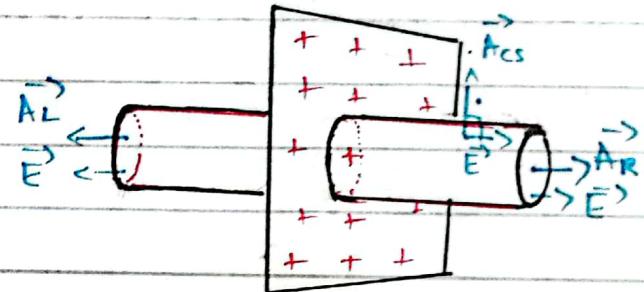
→ Infinite Sheet of Charge

$$\phi_{\text{total}} = \phi_p + \phi_c + \phi_{cs}$$

$$\rightarrow EA \cos 0^\circ + EA \cos 0^\circ + EA_{cs} \cos 90^\circ \\ = EA + EA$$

$$\phi_{\text{total}} = 2EA$$

$$\text{Gauss Law}, \quad \phi = \frac{q}{\epsilon_0}$$

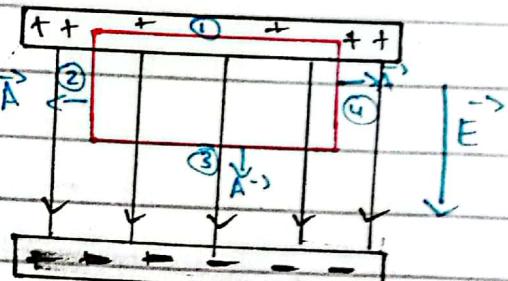


$$2EA = \frac{q}{\epsilon_0} \Rightarrow E = \frac{q}{2A\epsilon_0} \Rightarrow E = \frac{\sigma}{2\epsilon_0} \quad (q/A = \sigma)$$

## Conducting Sheet

→ Two parallel sheets / Parallel plate conductor

$$\phi_{\text{total}} = \phi_1 + \phi_2 + \phi_3 + \phi_4$$



→ E inside conductor is always zero

$$= 0 + EA \cos 90^\circ + EA \cos 0^\circ + EA \cos 90^\circ$$

$$\phi_{\text{total}} = EA$$

Gauss' Law

$$\phi = \frac{q}{\epsilon_0}$$

$$EA = \frac{q}{\epsilon_0} \Rightarrow E = \frac{q}{A\epsilon_0}, \quad \frac{q}{A} = \sigma$$

$$E = \frac{\sigma}{\epsilon_0}$$

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# Spherical Symmetry

- 'r' is radius of gaussian surface
- 'R' is radius of spherical object/surface / Shell
- Excess charge is always uniformly distributed on shell surface

## SHELL THEOREM

The charges are in layers of ~~the~~ shell, not a solid sphere.

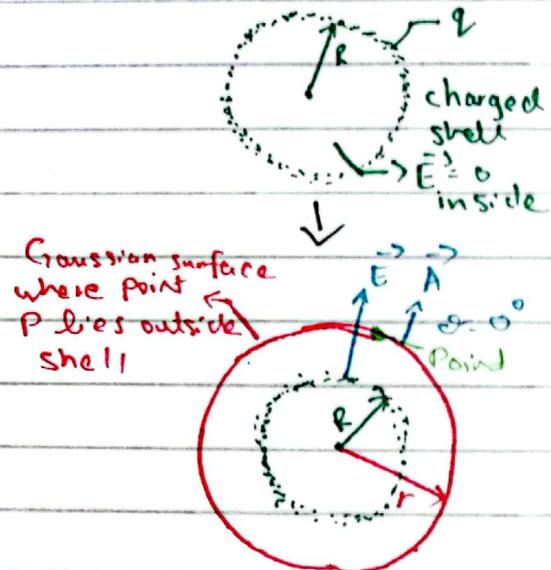
### CASE-I Point is Outside the Shell

→ Draw circular gaussian surface outside the shell, equidistant from the shell

$$\rightarrow r > R$$

Gauss Law,

$$E \cdot \oint \vec{E} \cdot d\vec{A} = q_{enc}$$



Electric field intensity is constant because of spherical shape

At every point, angle between  $\vec{E}$  and patch element  $d\vec{A}$  is  $0^\circ$

$$\oint E d\cos 0^\circ = \oint E dA = E \oint dA = E (4\pi r^2)$$

Gauss Law,  $E_0 \phi = q$

$$E_0 [E(4\pi r^2)] = q$$

$$E = \frac{1}{4\pi \epsilon_0} \cdot \frac{q}{r^2} = \frac{kq}{r^2} \quad (\text{Also the expression of } \vec{E} \text{ of point charge})$$

→ Charged shell behaves as point charge for all points outside the shell

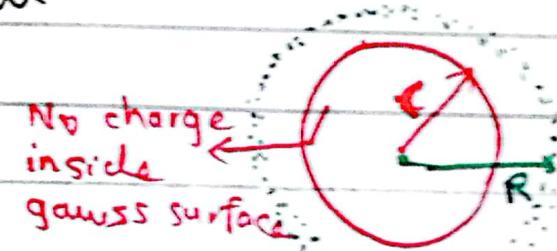
## CASE-II Point P is Inside the Shell

→ Draw gaussian surface inside the shell equidistant from all sides.

$$\rightarrow r < R$$

Gauss Law

$$\epsilon_0 \oint \vec{E} \cdot d\vec{A} = q_{enc}$$



But no charge is enclosed by the gaussian surface

$$\therefore q_{enc} = 0$$

$$\epsilon_0 \oint \vec{E} \cdot d\vec{A} = 0$$

$$\vec{E} \oint d\vec{A} = 0$$

~~Either~~ As area can never be zero,  $\oint d\vec{A} \neq 0$

$$\therefore E = 0$$

→ Electric field inside the shell of charge is zero and any charge placed inside shell will experience no force.

# ELECTRIC FIELD DUE TO SPHERICAL SHELL OF CHARGE

The charges are in solid spherical shell, not just a layer

## CASE-I Point is Outside Spherical Charge

Draw gaussian surface outside charged sphere

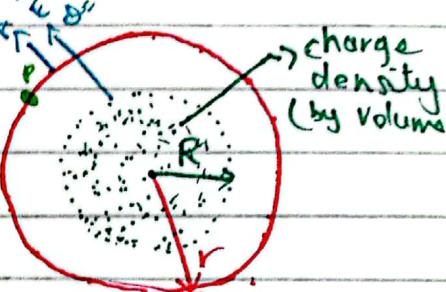
$$r > R$$

Since charge is uniformly distributed,

$$\rho = \frac{\text{Total charge}}{\text{Volume}} = \frac{q}{V}$$

$$\rho = \frac{q}{\frac{4}{3}\pi R^3} \quad (\text{Take } R, \text{ not } r)$$

$$q = \rho \left[ \frac{4}{3}\pi R^3 \right]$$



Gauss' Law

$$\oint \vec{E} \cdot d\vec{A}$$

$\theta = 0^\circ$  (At every point)

$$= \oint E dA \cos 0^\circ = \oint EdA = E \oint dA = E \oint 4\pi r^2$$

Gauss Law

$$\epsilon_0 E = q$$

$$\epsilon_0 [E(4\pi r^2)] = q$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \Rightarrow \frac{1}{4\pi\epsilon_0 r^2} \cdot \rho \frac{4}{3}\pi R^3 = \frac{\rho R^3}{3\epsilon_0 r^2}$$

$\rightarrow$  Spherical charge distribution at distance 'r' has same value as if charges were concentrated at centre of sphere

## CASE-II Point is Inside Sphere of Charges

Draw gaussian surface inside charged sphere

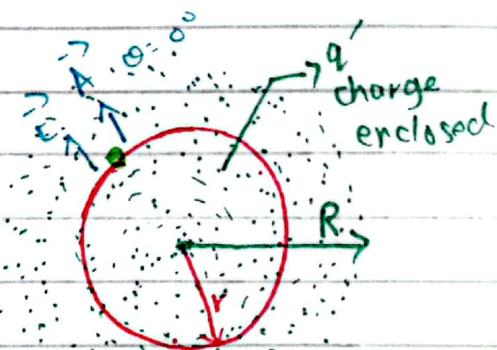
$$r < R$$

Charge is uniformly distributed

$$\rho = \frac{\text{Total charge}}{\text{Volume}} = \frac{q}{V}$$

$$\rho = \frac{q}{\frac{4}{3}\pi R^3}$$

$$q = \rho \left[ \frac{4}{3}\pi R^3 \right] - \textcircled{1}$$



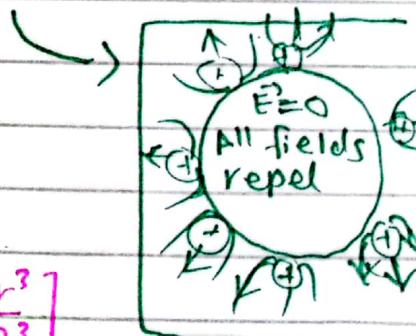
- Charges outside gaussian surface does not set  $\vec{E}$  inside surface

- Only enclosed charges setup a field  $\vec{E}$

$$\Rightarrow q' = \rho \left[ \frac{4}{3}\pi r^3 \right] - \textcircled{2} \quad (\text{Take } r, \text{ not } R)$$

Discarding both \textcircled{1} and \textcircled{2}

$$\frac{q'}{q} = \frac{\rho \left[ \frac{4}{3}\pi R^3 \right]}{\rho \left[ \frac{4}{3}\pi r^3 \right]} \Rightarrow \frac{q'}{q} = \frac{R^3}{r^3} \Rightarrow q' = q \left[ \frac{r^3}{R^3} \right]$$



Gauss' Law,  ~~$\oint \vec{E} \cdot d\vec{l} = \frac{q}{\epsilon_0}$~~   $\oint \vec{E} \cdot d\vec{l} = E \oint dA = E(4\pi r^2) \quad (a = a')$

$$\epsilon_0 \phi = q$$

$$\epsilon_0 [E(4\pi r^2)] = q'$$

$$E = \frac{1}{4\pi\epsilon_0 r^2} \cdot \frac{q'}{r^2} = \frac{1}{4\pi\epsilon_0 r^2} \cdot \frac{1}{r^2} \cdot \left[ \frac{r^3}{R^3} \right] q, \quad E = \frac{q}{4\pi\epsilon_0 r^3}, \quad E = \frac{q}{4\pi\epsilon_0 r^3}$$

$\rightarrow$  When  $r = 0, E = 0$

$$\rightarrow \text{When } r = R \Rightarrow E = \frac{1}{4\pi\epsilon_0 r^2} \cdot \frac{q}{r^2} \quad (E_{\text{max}})$$

# Ch 28 Magnetic Field

$$\vec{F}_B = q \vec{v} \times \vec{B}$$

$$F_B = qvB \sin\theta$$

Unit of Magnetic Field = 1 Tesla = 1 T =  $\frac{nI}{A \cdot m}$  (SI-unit)

\*  $1 T = 10^4 G$  (gauss)

$$V = \sqrt{\frac{2K}{m}} = \sqrt{\frac{2qV}{m}}, \quad I = \frac{q}{A}, \quad J = neV_d$$

(ne is carrier charge density)  
(n = number of charges per unit volume)

$$n = \frac{BI}{Vle}, \quad L = \frac{A}{d}$$

$$V = VBd$$

(Potential) ↓  
(velocity)

\*  $F_B$  is deflecting Force, does not do work  
 $F_E$  does work

$$F = qE, \quad F = qVR \sin\theta, \quad F = ILB \sin\theta, \quad F = \frac{mv^2}{r}, \quad F = mg$$

$$V = \frac{IB}{Ine}, \quad n = \frac{IB}{Vle}, \quad E = V_d B$$

\*  $1 eV = 1.6 \times 10^{-19} J$

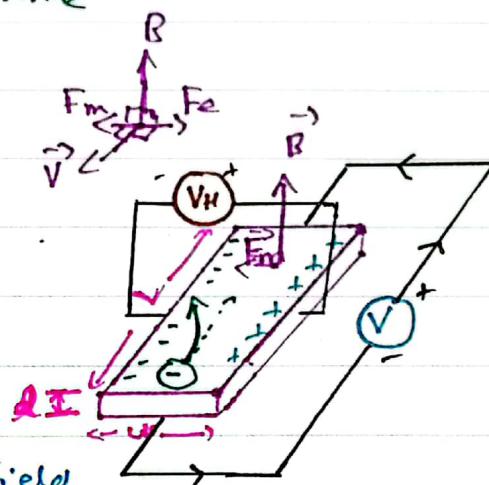


# THE HALL EFFECT

Hall Effect is production of Potential Difference (Hall voltage) across current carrying conductor (in presence of magnetic field), perpendicular to both current and magnetic field.

In Simple Words : Every current is automatically associated with magnetic field.

"When potential difference is applied across length of conductor, Hall voltage is produced across width of conductor."



$$F_E = F_B$$

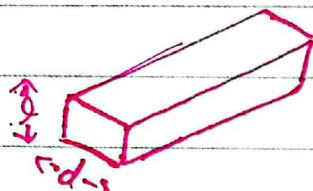
$$eE = eV_d B \quad \therefore (qE = qV_d B)$$

$$E = V_d B \quad \text{--- (1)}$$

$$\rightarrow J = neV_d$$

$$V_d = \frac{J}{ne} = \frac{I}{Ane}$$

$$V_d = \frac{I}{ldne} \quad \Rightarrow (A = ld)$$



$t$  = thickness  
 $d$  = width

$$\text{Also } E = \frac{V}{d}$$

\*  $n$  = number density of charge carriers  
(number per unit volume)

$$\rightarrow (1) \quad \frac{V}{d} = \frac{I}{ldne} B$$

$$V = \frac{IB}{lne} \quad \Rightarrow \quad n = \frac{IB}{Vle}$$



# CIRCULATING CHARGED PARTICLE

- \*  $\vec{F}_B$  is ~~the~~ only deflecting force, does zero work
- \*  $\vec{F}_e$  does work on particle

$$F_B = F_c \quad (\text{centripetal Force})$$

## Radius

$$qVB = \frac{mv^2}{r}$$

$$r = \frac{mv}{qB}$$

## Time Period

$$T = \frac{\text{circumference}}{\text{velocity}} = \frac{2\pi r}{v} = \frac{2\pi}{v} \left( \frac{mv}{qB} \right)$$

$$T = \frac{2\pi m}{qB}$$

## Frequency

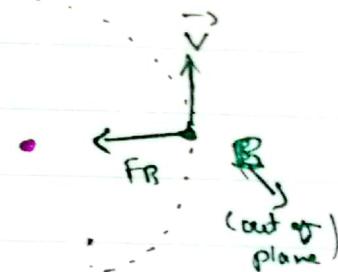
$$f = \frac{1}{T} = \frac{qB}{2\pi m}$$

## Angular Frequency

$$\omega = 2\pi f$$

$$\omega = 2\pi \left( \frac{qB}{2\pi m} \right)$$

$$\omega = \frac{qB}{m}$$



# Magnetic Force on Current Carrying Wire

We know  $F = qV_d B$  since - ①

$$\Rightarrow q = It = I \frac{L}{V_d}$$

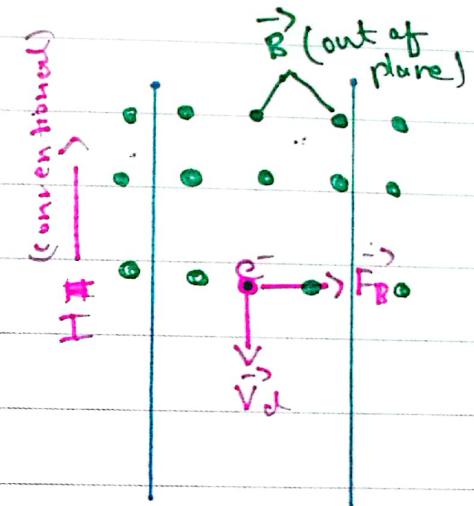
$$\textcircled{1} \Rightarrow F = I \frac{L}{V_d} V_d B \text{ since}$$

$$F = ILB \sin\theta$$

$$\theta = 90^\circ$$

$$F = ILB$$

$$\boxed{\vec{F}_B = I \vec{L} \times \vec{B}}$$



\* If wire is bent, we can break in small pieces and differentiate

$$d\vec{F}_B = I d\vec{L} \times \vec{B}$$

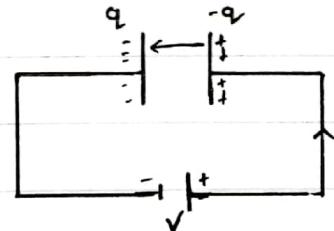
# Ch # 5 CAPACITANCE

→ Ability to store charge

$$Q \propto V \quad (Q \text{ is charge on one plate, don't sum both charges})$$

$$Q = CV$$

$$C = \frac{Q}{V} \quad (\text{Farad} = F = CV')$$



## Parallel Plate Capacitor

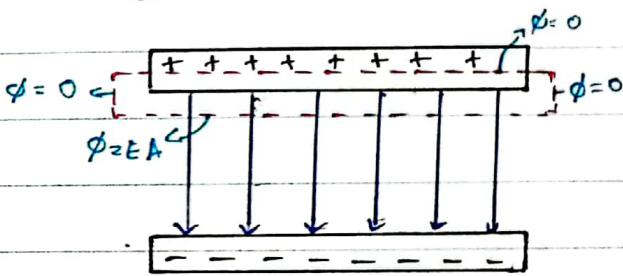
→ First we find Potential Difference by integrating  $\vec{E}$  with  $d\vec{s}$  (path length)

→ To find  $\vec{E}$ , take gaussian surface

$$\oint \vec{E} \cdot d\vec{A} = q$$

$$\Rightarrow E = \frac{q}{A \epsilon_0} \quad (\text{Derived in ch 23})$$

$$\Rightarrow q = EA \epsilon_0$$



→ Now calculating Potential Difference

→ Integration path will be -ve to the plate

$$\Delta V = V_f - V_i$$

$$V_f - V_i = - \int \vec{E} \cdot d\vec{s}$$

$$V = \int \vec{E} \cdot d\vec{s} = E \int \vec{E} \cdot d\vec{s} = E |d\vec{s}|$$

$V = E d$

$$C = \frac{Q}{V} = \frac{EA\epsilon_0}{Ed}$$

$$C = \frac{A\epsilon_0}{d}$$

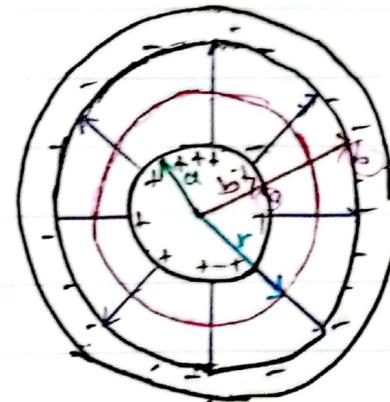


# Cylindrical Capacitance

$$q = \epsilon_0 A E \quad (\text{Derived in Ch 23})$$

$$q = \epsilon_0 E (2\pi r L) \quad \begin{array}{l} \text{area of gaussian surface} \\ \text{height} \end{array}$$

$$E = \frac{q}{2\pi\epsilon_0 L r}$$



$$V = - \int E ds = - \frac{q}{2\pi\epsilon_0 L} \int_b^0 \frac{dr}{r} =$$

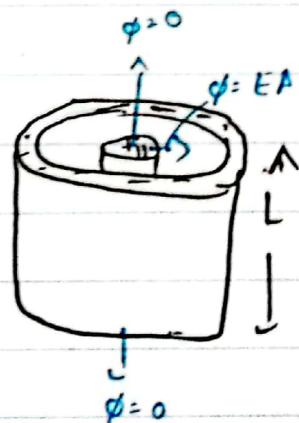
$$V = \frac{-q}{2\pi\epsilon_0 L} [\ln r]_b^0 = \frac{q}{2\pi\epsilon_0 L} [\ln r]_a^b \quad (ds = -dr)$$

$$V = \frac{q}{2\pi\epsilon_0 L} [\ln(b) - \ln(a)]$$

$$V = \frac{q}{2\pi\epsilon_0 L} \ln\left(\frac{b}{a}\right)$$

$$\Rightarrow C = \frac{q}{V} = q \times \frac{2\pi\epsilon_0 L}{q \times \ln(b/a)}$$

$$C = \frac{2\pi\epsilon_0 L}{\ln(b/a)}$$



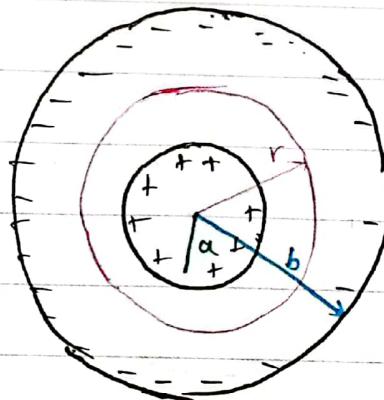
# Spherical Capacitor

$$q = \epsilon_0 EA$$

$$q = \epsilon_0 E (4\pi r^2)$$

Surface area of gaussian surface

$$V = \int_{-}^{+} E ds = - \int_{-}^{+} Edr = \epsilon_0 \int_{-}^{+} Edr \quad (ds = -dr)$$



$$V = \int_a^b \frac{q}{4\pi r^2 \epsilon_0} dr = \frac{q}{4\pi \epsilon_0} \left[ \frac{1}{r} \right]_a^b$$

$$V = \frac{q}{4\pi \epsilon_0} \left| -\frac{1}{r} \right|_a^b = -\frac{q}{4\pi \epsilon_0} \left[ \frac{1}{b} - \frac{1}{a} \right] = -\frac{q}{4\pi \epsilon_0} \frac{(a-b)}{ab}$$

$$V = \frac{q}{4\pi \epsilon_0} \left( \frac{b-a}{ab} \right)$$

$$C = \frac{q}{V} = q \times \frac{4\pi \epsilon_0 (ab)}{q (b-a)}$$

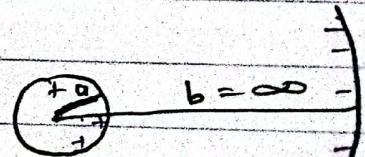
$$C = 4\pi \epsilon_0 \left( \frac{ab}{b-a} \right)$$

## Isolated Sphere

$$C = \frac{4\pi \epsilon_0}{\frac{1}{a} - \frac{1}{b-\infty}} = \frac{4\pi \epsilon_0}{\frac{1}{a} - 0} = 4\pi \epsilon_0 a$$

$$a = R$$

$$C = 4\pi \epsilon_0 R$$





## Capacitors in Parallel

- Capacitance increases
- Charge is distributed across capacitors
- Voltage at each capacitor remains same

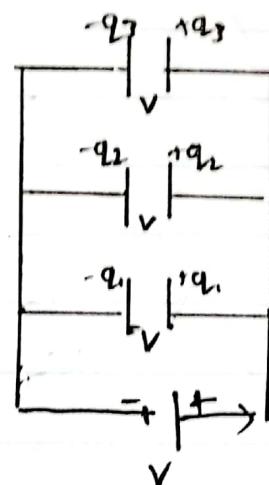
$$q_1 = C_1 V, q_2 = C_2 V, q_3 = C_3 V$$

$$Q = q_1 + q_2 + q_3 = (C_1 + C_2 + C_3)V$$

$$Q = C_{eq} V$$

~~$$C_{eq} = \frac{Q}{V}$$~~

$$C_{eq} = C_1 + C_2 + C_3 + \dots + C_n$$

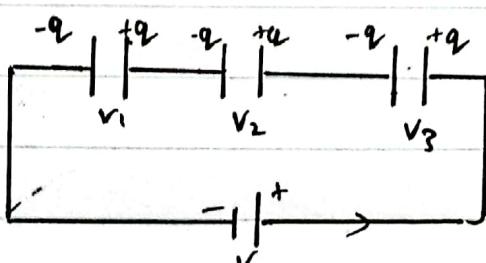


## Capacitance in Series

- Capacitance Decreases
- Charge remains same
- Voltage is distributed

$$V_1 = \frac{q}{C_1}, V_2 = \frac{q}{C_2}, V_3 = \frac{q}{C_3}$$

$$V = V_1 + V_2 + V_3 = q \left( \frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3} \right)$$



$$\frac{V}{q} = \left( \frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3} \right)$$

$$\frac{1}{C_{eq}} = \frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3} + \dots + \frac{1}{C_n}$$

# Ch 29 Magnetic Fields due to Currents

## Magnetic Field Outside a Current Carrying Wire

According to Amperes Law

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{\text{enclosed}}$$

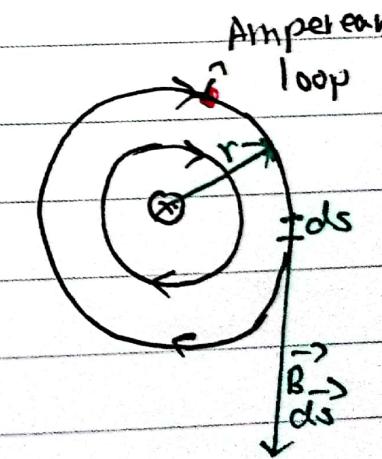
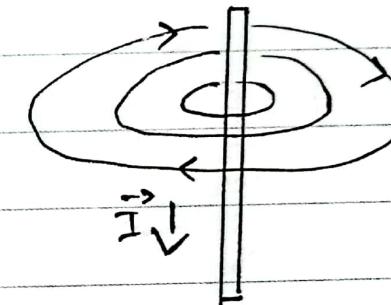
$d\vec{s}$  = circumference (small part)

$$B \oint d\vec{s} = \mu_0 I_{\text{enc}} \cdot \cos 0^\circ$$

$$\theta = 0^\circ, \oint d\vec{s} = 2\pi r$$

$$B(2\pi r) = \mu_0 I \quad (\text{All current is enclosed in loop})$$

$$B = \frac{\mu_0 I}{2\pi r} \quad (\mu_0 = 4\pi \times 10^{-7} \text{ NA}^{-2} \text{ or Tm/A})$$



# Magnetic Field Inside a Current Carrying Wire

According to Ampere's Law

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{enc}$$

$$B(2\pi r) = \mu_0 I_{enc}$$

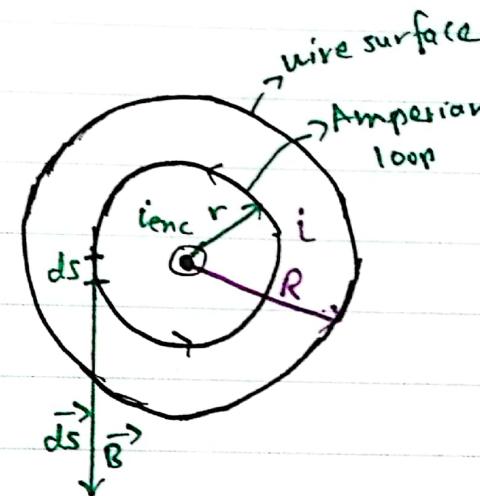
Now,

$$\frac{I_{enc}}{I_t} = \frac{\pi r^2}{\pi R^2} \quad (I_t = \text{total current})$$

$$I_{enc} = I_t \frac{r^2}{R^2}$$

$$\Rightarrow B(2\pi r) = \mu_0 I_t \frac{r^2}{R^2}$$

$$B = \left( \frac{\mu_0 I}{2\pi R^2} \right) \cdot r$$



# In Solenoids

Ampere's Law

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{\text{enc}}$$

$$\oint \vec{B} \cdot d\vec{s} =$$

$$= \int_a^b B \cdot ds \cos 90^\circ + \int_b^c B \cdot ds \cos 90^\circ + \int_c^d B \cdot ds \cos 90^\circ + \int_d^a B \cdot ds \cos 90^\circ$$

$$= \int_a^b B ds \cos 90^\circ + \int_b^c B ds \cos 90^\circ + 0 + \int_d^a B ds \cos 90^\circ$$

$$= B \int_a^b ds = B h$$

$$\text{Now, } I_{\text{enc}} = n h I \quad (n: \frac{\text{number of turns}}{\text{total length}}, h = \text{length})$$

$$\Rightarrow B h = \mu_0 I_{\text{enc}}$$

$$B h = \mu_0 n h I$$

$$\therefore B = \mu_0 n I$$

# Toroids

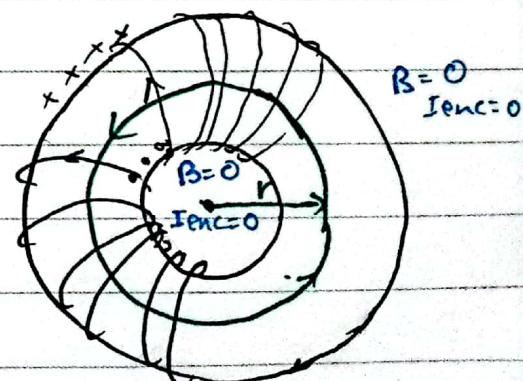
$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{\text{enc}}$$

$$B \oint ds = \mu_0 I_{\text{enc}}$$

$$B (2\pi r l) = \mu_0 N I$$

$$(I_{\text{enc}} = N I) \quad N = \text{Number of turns}$$

$$B = \frac{\mu_0 N I}{2\pi r}$$



$$\boxed{2\pi r = \text{total length}}$$

$$\therefore \frac{N}{2\pi r} = n$$

# Force Between Two Parallel Wires

Field of  $a$  ( $\vec{B}_a$ ) produces force  
on  $b$  ( $\vec{F}_{ba}$ )

$$\vec{B}_a = \frac{\mu_0 i_a}{2\pi d} \hat{z}$$

$$\vec{F}_{ba} = \mu_0 i_b L \vec{i}_a \times \vec{B}_a \\ = \frac{\mu_0 L i_a i_b}{2\pi d} \hat{y}$$

$$F_{ba} = \frac{\mu_0 L i_a i_b}{2\pi d}$$

