# CHAPTER IV

### CAPACITANCE

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# 4.1 Capacitor

Two conductors of arbitrary shape, completely invlated from each other and their surroundings, form a capacitor. No matter what their shape, these conductors are called plates. When the capacitor is charged by connecting the plates to the opposite terminals of a battery, equal and opposite charges (say 4q and -q) appear on the

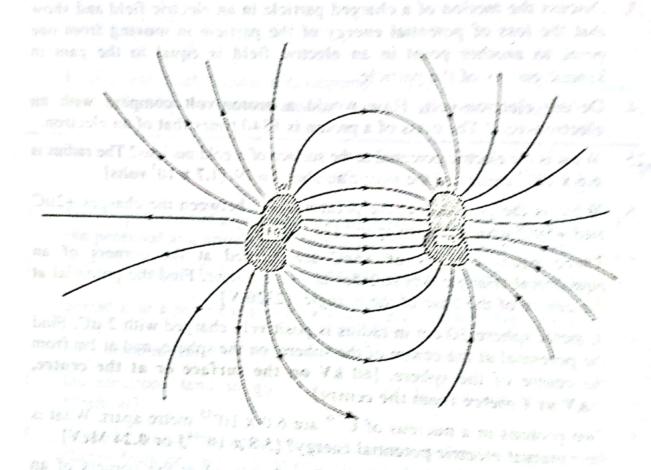


Fig. 4.1

two plates of the capacitor. By charge of a capacitor we mean the absolute value of the charge on either plate, the net charge on the capacitor being zero. The potential difference between the plates of the capacitor is the potential difference of the battery. Fig. 4.1 shows the general arrangement of a capacitor.

4.2 Capacitance

The charge q of a capacitor is found to be directly proportional to the potential difference between the plates. Or, The calculation of the electric field and the pore tight of forest

or, 
$$q = CV$$
 or,  $C = \frac{q}{V}$  or the probability of  $V$  of  $V$  or  $V$ 

The proportionality constant C is called the capacitance of the capacitor. Its value depends on

- (i) the geometry of each plate
- (ii) the spatial relationship between the plates and (iii)the medium in which the plates are immersed.

As can be seen from eqn. (4.1), the SI unit of capacitance is coulomb per volt. This unit occurs so often that it is given a special name - the farad.

1 farad (1F) = 1 coulomb per volt (1 C/V). So if a potential difference of 1 volt is needed to give a capacitor a charge of 1 coulomb, then the capacitance of the capacitor is said to be I farad.

The farad is a large unit. Submultiples of the farad, such as microfarad ( $1\mu F = 10^{-6}F$ ) and the picofarad ( $1pF = 10^{-12}F$ ) are more convenient units in practice. The later of t

of the potential difference between the plates is refried the

electric field E by the relation recomme

#### Calculation of capacitance

Once the geometry of a capacitor is known, its capacitance can be calculated. Since different capacitors have different plate geometries, it is wise to develop a general plan to simplify the process involved in the calculation. In brief, the plan is a large and

- (i) assume a charge q on the plates
- (ii) applying Gauss' law to calculate the electric field E between the plates in terms of the charge on the plates.
- (iii)knowing E, calculate the potential difference V between the the potential deference between the plates, we can set . self. =

charge q of a capacitor is  $\frac{q}{V} = \frac{q}{V}$  mind the parameter q is capacitor in q

The calculation of the electric field and the potential difference may be simplified by making certain assumptions. These are,

(a) calculating the electric field:

Gauss' law:

$$\in_0 \int \mathbf{E} \cdot d\mathbf{A} = \mathbf{q}$$
 stale does to (the geometry of (4.2)

VO = p .ao

Here q is the charge enclosed by the Gaussian surface, and the integral is carried out over that surface. Only those cases will be considered in which the Gaussian surface are such that whenever electric flux passes through it, the electric field E and dA will point in the same direction. Eqn. 4.2 then reduces to

I farad (IF) = I coulomb per volt (I C/V). So if a potenti

in which A is the area of that part of the Gaussian surface through which the flux passes. For convenience the Gaussian surface is so drawn that it completely encloses the charge on the positive plate.

microfured (luF = 10 To and the piceform) (lpF = 10 T) are more

(b) Calculating the potential difference. The ni shift inclusives

The potential difference between the plates is related to the electric field E by the relation

Once the geometry of a capacitor is known, its capacitance can be (6.4) alated. Since different capacit 
$$\mathbf{Z}\mathbf{b}_{0}\cdot\mathbf{A}_{0}$$
  $\mathbf{z}=\mathbf{i}\mathbf{Y}_{1}\mathbf{z}\cdot\mathbf{A}_{0}\mathbf{Y}_{1}$  plate geometries, it is wise to develop a general plan to simplify the

the integral being evaluated along any path that starts on one plate and ends on the other. One should always choose a path that follows an electric field line from the positive plate to the negative plate as shown in Fig. 4.2, since the vectors  $\mathbf{E}$  and  $d\mathbf{S}$  point in the same direction along this path. It, therefore, follows that the quantity  $V_f - V_i$  is negative. Since we are looking for  $V_i$ , the absolute value of the potential difference between the plates, we can set  $V_f - V_i = -V$ . Eqn. 4.3 then becomes

in which the + and the - signs remind us that our path of integration starts on the positive plate and ends on the negative plate.

The electric field E between the plates is the sum of the fields due to the two plates i.e.,  $E = E_+ + E_-$  where  $E_+$  is the field due to charges on the positive plate while  $E_-$  is that due to charges on the negative plate. By Gauss' law both  $E_+$  and  $E_-$  are proportional to q so that E is also proportional to q. By eqn. 4.3. V is also proportional to q. This means that if q is doubled, E and V are also doubled. Because V is proportional to q, the ratio q/V is a constant and is independent of q.

#### (i) Capacitance of a parallel-plate capacitor

A parallel-plate capacitor formed of two parallel conducting plates of area A and separated by a distance d is shown in Fig. 4.2. If the plates are connected to the opposite terminals of a battery, then a charge +q appears on one plate and a charge -q on the other. If d is small enough compared to the plate dimensions, the electric field strength E between the plates will be uniform, which means that the lines of force will be parallel and evenly spaced. According to the laws of electro-magnetism, there should be some fringing or curving of the lines at the edges of the plates; for small enough d it can be neglected for the present purpose.

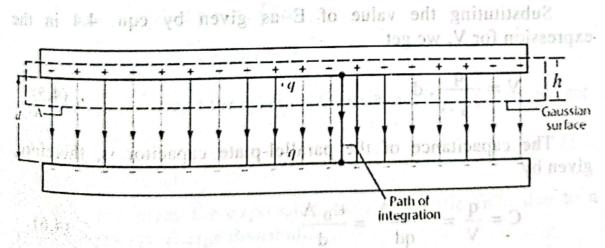


Fig. 4.2

Let us imagine a Gaussian surface of height h closed by plane caps of area A of the same shape and size of the capacitor plates. Because the electric field inside a conductor carrying a static charge is zero, the flux of E for the part of the Gaussian surface that lies inside the top capacitor plate is also zero. The flux E through the wall of the Gaussian surface is zero because, to the extent that the fringing of the lines of force can be neglected, E lies in the wall Thus the only part of the Gaussian surface which contributes to the electric flux is the Gaussian surface that lies between the plates Here E is constant and according to Gauss' law

one constant to q. This means that if q is doubled be and V are also seed because V is proportion 
$$\frac{p}{e} = A$$
.  $A = B$ .  $A = B$ .  $A = A$  constant is independent of q.

or,  $A = A = A$ 

(4.4)

in the sparated by a difference d is shown in Fig. 4.2.

The potential difference V between the plates can be obtain from eqn. 4.3. Or

The parameter 
$$\mathbf{V} = \mathbf{V} = \mathbf{E} \cdot \mathbf{A} \cdot \mathbf{I} = \mathbf{E} \cdot \mathbf{B} \cdot \mathbf{A} = \mathbf{E} \cdot \mathbf{A} \cdot \mathbf{B} \cdot \mathbf{A} = \mathbf{E} \cdot \mathbf{A} \cdot \mathbf{A} \cdot \mathbf{A} = \mathbf{A} \cdot \mathbf{A} \cdot$$

since E is constant and can be taken outside the integral and simply the plate separation d.

Substituting the value of E as given by eqn. 4.4 in the expression for V, we get

$$V = \frac{q}{\epsilon_0 A} \cdot d \tag{4.5}$$

The capacitance of the parallel-plate capacitor is, therefore given by

$$C = \frac{q}{V} = \frac{q \in_0 A}{qd} = \frac{\in_0 A}{d}$$
 (4.6)

As can be seen from eqn. 4.6, the capacitance does indeed depend only on geometrical factors, namely, the plate area A and the plate separation d.

## (ii) Capacitance of a spherical capacitor

Fig. 4.3 shows a central cross-section of a capacitor that consists of two concentric spherical shells of radii a and b. As a Gaussian surface let us draw a sphere of radius r concentric with the two shells. Applying Gauss' law to this surface we obtain

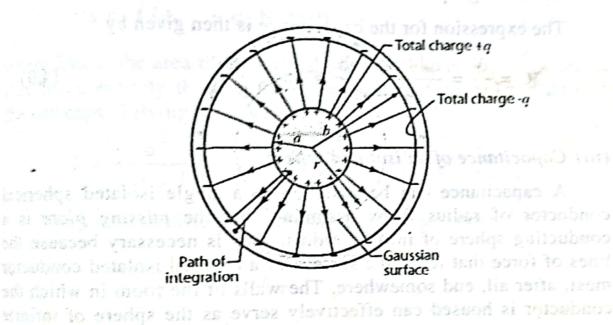


Fig. 4.3

Now the capacitance of a spherical capacitur is given by equ.

$$q = \epsilon_0 \oint E \cdot ds = \epsilon_0 \oint E ds$$

$$= \epsilon_0 E \oint ds = \epsilon_0 E (4\pi r^2)$$

where  $4\pi r^2$  is the area of the spherical Gaussian surface. Solving, we get

$$E = \frac{q}{4\pi\epsilon_0 r^2} \tag{4.7}$$

Eqn. 4.7 gives the expression for the electric field due to a uniform spherical charge distribution.

The expression for the potential difference between the two concentric spheres is given by

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$$\pi$$
 estimated  $\frac{q^{1}}{4\pi} \in \frac{q^{1}}{0}$ .  $\frac{d}{dr}$  drotsel from EA and the  $\frac{d}{dr}$  and  $\frac{d$ 

$$=\frac{q}{4\pi\epsilon_0}\int_a^b\frac{dr}{r^2}=\frac{q}{4\pi\epsilon_0}\left(\frac{1}{a}\frac{1}{a}\frac{1}{b}\right)$$

(8.4) of two concentric spherical shears 
$$a = 0$$
 and  $b = 0$  as  $a = 0$ .

(8.4) of two concentric spherical shears  $a = 0$ . As a concentric spherical shears  $a = 0$ . As a concentric sphere of tadays  $a = 0$ .

In deriving eqn. 4.8 we have used the fact that here ds = dr.

The expression for the capacitance is then given by

$$C = \frac{q}{V} = \frac{q(4\pi \in_0) ab}{q(b-a)} = 4\pi \in_0 \frac{ab}{b-a}$$
 (4.9)

### (iii) Capacitance of an isolated sphere

A capacitance can be assigned to a single isolated spherical conductor of radius R by assuming that the missing plate is a conducting sphere of infinite radius. This is necessary because the lines of force that leave the surface of a charged isolated conductor must, after all, end somewhere. The walls of the room in which the conductor is housed can effectively serve as the sphere of infinite radius. Fig. 4.3

Now the capacitance of a spherical capacitor as given by eqn. 4.9, is

$$C = 4\pi \in_{0} \frac{ab}{b-a} = 4\pi \in_{0} \frac{a}{1-a^{2}b^{2}} = ab \cdot 3 = ab$$

If the second sphere is of infinite radius, then b -> « Substituting R for a, we obtain

$$C = 4\pi \epsilon_0 R$$
 (4.10)

Eqn. 4.7 gives the expression for the electric field due to a (iv) Capacitance of a cylindrical capacitor grand harmand

Fig. 4.3, also serves to show, a cross-section of a cylindrical capacitor of length l formed by two co-axial cylinders of radii a and b. The length of the capacitor is assumed to be much greater than its radius, i.e., l >> b so that fringing of the lines of force at the ends radius, edge effect) can be ignored for the purpose of calculating the capacitance. Each plate contains a charge of magnitude q.

As a Gaussian surface let us construct a coaxial cylinder of radius r and length l closed by end caps. Applying Gauss' law we then obtain

Isolated sphere
$$q = \epsilon_0 \oint E \cdot ds = \epsilon_0 \oint E ds$$

$$= \epsilon_0 \oint E \cdot ds = \epsilon_0 \oint E ds$$

$$= \epsilon_0 E \oint E ds = \epsilon_0 E (2\pi r l)^{-3\pi l}$$

where  $2\pi rl$  is the area of the curved part of the Gaussian surface, the flux being entirely through the cylindrical surface and not through the end caps. Solving for E we get

Example 4.1 A plane parallel cupacitor 
$$h = \frac{q}{2\pi \epsilon_0} = \frac{1}{2} p_{iotes}$$
 of radius  $r = 10.0$  cm, separated by a distance  $d = h_{iotes} = 10.0$ 

The potential difference between the plates is given by

$$V = \int_{r}^{r} E ds = \frac{q}{2\pi \epsilon_0 l} \int_{a}^{b} \frac{dr}{r}$$

$$V = \int_{r}^{r} E ds = \frac{q}{2\pi \epsilon_0 l} \int_{a}^{b} \frac{dr}{r}$$

q = CV where C is the capacitance of the capacitor and is given by  $\mathbf{d} = \mathbf{0} \cdot \mathbf{0}$ 

$$\frac{q}{2\pi\epsilon_0} \ln \frac{b}{a} \text{ and } A = 0$$

From the relation  $C = \frac{q}{V}$ , we then have

$$C = 2\pi \epsilon_0 \frac{\left(\frac{l^{(1)} O(1 \times 10^{-1})}{l^{(1)} O(1 \times 10^{-1})} \left(\frac{3.14 \times 10^{-1} O(1 \times 10^{-1})}{l^{(1)} O(1 \times 10^{-1})}\right)}{l^{(1)} O(1 \times 10^{-1})} = 0$$

$$(4.12)$$

As can be seen from eqn. 4.12 that, like a parallel-plate capacitor; the capacitance of a cylindrical capacitor depends only on geometrical factors, in this case l, b and a.

The capacitances of various capacitors derived in this section is summarized below.

respecting apply on his many control

108 Type of capacitor	Capacitance	Equation
	one E o d o mignist in	4.0
Spherical	$4\pi \in_0 \frac{ab}{b-a}$	
Isolated sphere	4π∈ <sub>0</sub> R	4.10
Cylindrical	$2\pi \in_0 \frac{l}{\ln(b/a)}$	4.12

It can be seen that every expression involves the constant emultiplied by a quantity that has the dimension of length.

Example 4.1 A plane-parallel capacitor has circular plates of radius r = 10.0 cm, separated by a distance d = 1.00mm. How much charge is stored on each plate when their electric potential difference has the value V = 100V?

Soln.

q = CV where C is the capacitance of the capacitor and is given by

$$C = \frac{\epsilon_0 A}{d}$$
 Here,  $A = \pi r^2 = \pi (0.1 \text{m})^2$   
= 3.14 x 10<sup>-2</sup>m<sup>2</sup>

even have  $\frac{1}{2} d = 1.00 \text{ mm} = 1 \times 10^{-3} \text{m}$ 

$$C = \frac{(8.85 \times 10^{-12} \text{ C}^2 / \text{N.m}^2) (3.14 \times 10^{-2} \text{ m}^2)}{1.00 \times 10^{-3} \text{ m}}$$

$$= 2.8 \times 10^{-10} \text{ F} = 280 \times 10^{-12} \text{ F}$$

$$= 280 \text{ pF}.$$

$$\therefore 0 = \text{CV} = (2.8 \times 10^{-10} \text{F}) (100 \text{ N})$$

$$\therefore q = CV = (2.8 \times 10^{-10} \text{F}) (100 \text{ V})$$

$$= 2.8 \times 10^{-8} \text{ coulomb}$$

$$= 28 \times 10^{-9} \text{ C} = 28 \text{ nC}.$$

Example 4.2 The area of each plate of an air-filled parallelplate capacitor is 1.1 × 10<sup>8</sup> metre<sup>2</sup>. What must be the separation between the plates, if the capacitance is to be 1.0 farad?  $= (4)(3.14)(8.85 \times 10^{-11} F/m)(6370 \times 10^{3} m)$ 

$$C = \frac{\epsilon_0 \text{ A}}{d}$$
Here,  $X = 1.1 \times 10^8 \text{m}^2$ 
or,  $d = \frac{\epsilon_0 \text{ A}}{C}$ 

$$= \frac{(8.85 \times 10^{-12} \text{ F/m})(1.1 \times 10^8 \text{ m}^2)}{1.0 \text{ F}}$$

$$= 9.735 \times 10^{-4} \text{ m}$$

$$= 9.735 \times 10^{-1} \text{ mm} = 0.9735 \text{ mm}.$$

Example 4/3 How much charge is stored in a capacitor consisting of two concentric spheres of radii 30 and 31 cm if the potential difference is 500V? dignel tinu and enastinged and

Soln.
$$q = CV \text{ (where } Here, a = 30 \text{ cm} = 0.30 \text{ m}$$

$$C = 4\pi \epsilon_0 \frac{ab}{(b-a)}$$

$$= \frac{(4) (3.14) (8.85 \times 10^{-12} \text{ F/m}) (0.30 \text{ m}) (0.31 \text{ m})}{(0.31-0.30) \text{m}}$$

$$= \frac{10.34 \times 10^{-12}}{0.01} \text{ F} = 1034 \times 10^{-12} \text{ F}$$

$$= 1.034 \times 10^{-9} \text{ F}$$

$$= 1.034 \text{ nF}$$

$$q = CV = (1.03 \times 10^{-9} \text{ F}) (500 \text{ V}) = 517 \text{ nC}.$$

in a centar's way, ity equivalent curacipance Example 4.4 What is the capacitance of the Earth, viewed as an isolated conducting sphere of radius 6370 km?

Example had the area of each place of an air filled alors distinct

Soln.

$$C = 4\pi \epsilon_0 R$$

$$= (4) (3.14) (8.85 \times 10^{-12} F/m) (6370 \times 10^3 m)$$

$$= 7.08 \times 10^{-6} = 708 \ \mu F.$$

Example 4.5 The space between the conductors of a long coaxial cable, used to transmit TV signals, has an inner radius a = 0.15mm and an outer radius b = 2.1 mm. What is the capacitance per unit length of this cable? [UIXI ] [m] [ ] OIXE8.8

Soln.

The capacitance of a coaxial cable is given by (eqn. 4.12)

$$C = 2\pi \in_0 \frac{l}{\ln(b/a)}$$
 where *l* is the length of the cable.

Hence capacitance per unit length is

Hence capacitance per unit length is white at some this form

$$\frac{C}{l_{m}} = \frac{2\pi \epsilon_{0}}{\ln(b/a)}$$
Here,  $b = 2.1 \text{ mm}$ 

$$= 2.1 \times 10^{-3} \text{m}$$

$$= \frac{(2) (3.14) (8.85 \times 10^{-12} \text{ F/m})}{\ln \left(\frac{2.1 \times 10^{-3} \text{m}}{0.15 \times 10^{-3} \text{m}}\right)}$$

$$= \frac{(2) (3.14) (8.85 \times 10^{-12} \text{ F/m})}{\ln \left(\frac{2.1 \times 10^{-3} \text{m}}{0.15 \times 10^{-3} \text{m}}\right)}$$

$$= \frac{55.578 \times 10^{-12}}{2.64} \text{ F/m}$$

$$= 21 \times 10^{-12} \text{F/m} = 21 \text{pF}.$$

# 4.3 Capacitors in series and parallel

In analyzing electric circuits, very often it is desirable to know the equivalent capacitance of two or more capacitors that are connected in a certain way. By equivalent capacitance is meant the capacitance of a single capacitor that can be substituted for the combination with no change in the operation of the rest of the circuit. With such a replacement, the circuit can be simplified, 50