



Progress in MEOP based ^3He Polarization System

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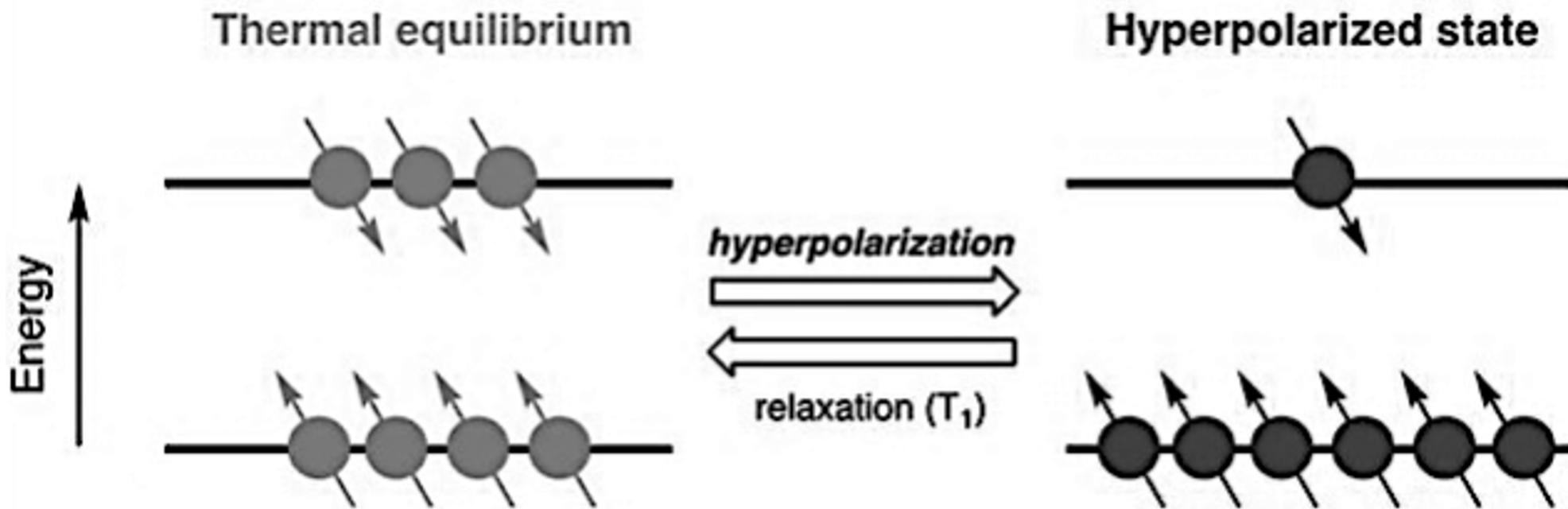
PNCMI 2025 Dongguan

Outline

- Introduction to ${}^3\text{He}$ Polarization
- The Basic Principle of MEOP
- Development of MEOP ${}^3\text{He}$ System
- Summary and Outlook

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^3He nuclei spin hyperpolarization



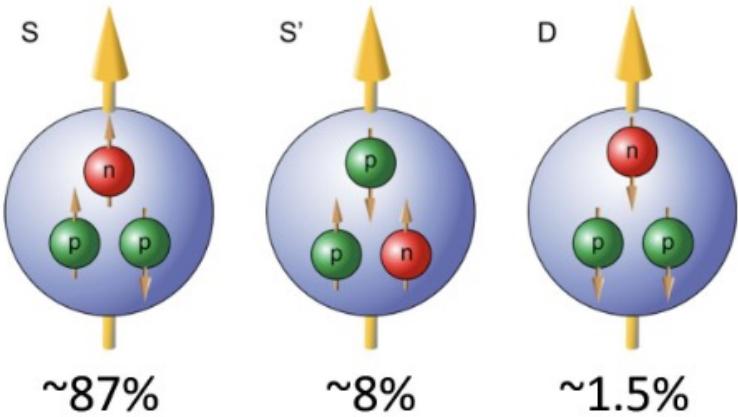
$$\text{Thermal equilibrium: } P_n = \frac{N_+ - N_-}{N_+ + N_-} \approx 10^{-5}$$

(For ^3He @ T=300 K, B=1 T)

$$\text{Hyperpolarized state: } P_n \approx 10^{-1}$$

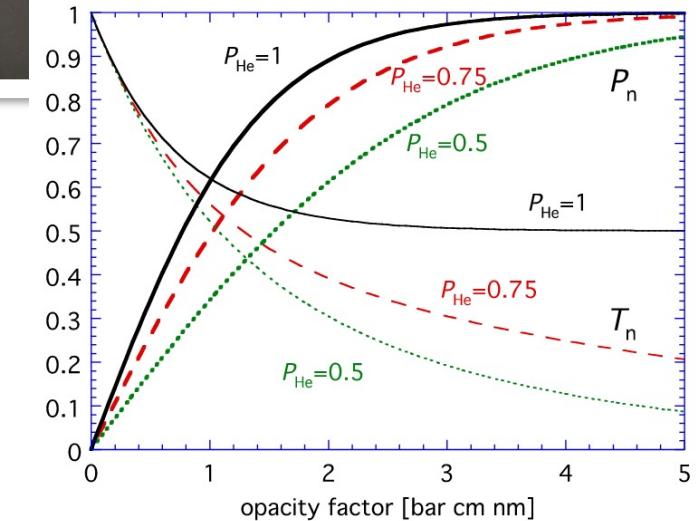
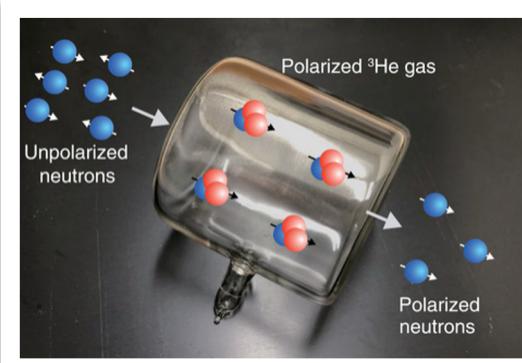
Application in nuclear and neutron physics

Polarized neutron target



$\sim 87\%$ of ${}^3\text{He}$ spin is carried by the unpaired neutron.

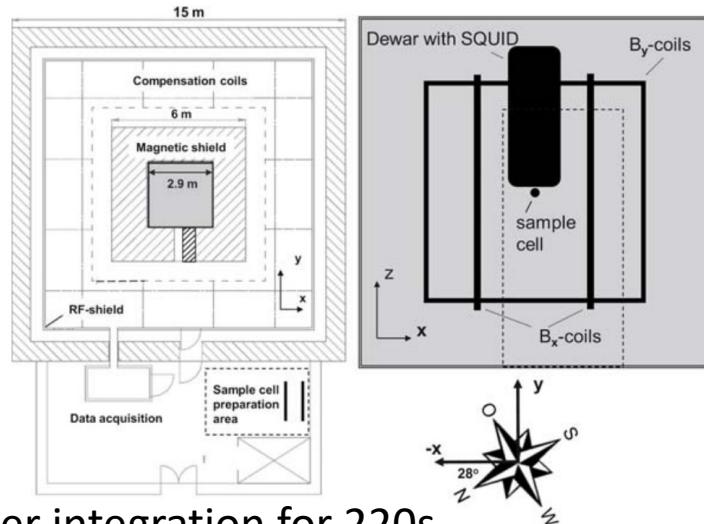
Neutron spin filter



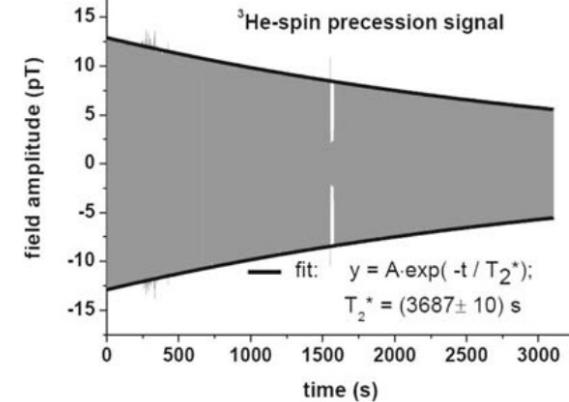
Strong spin-dependent neutron absorption cross-section.

Magnetic field metrology

Long coherence time enabling long-term measurements

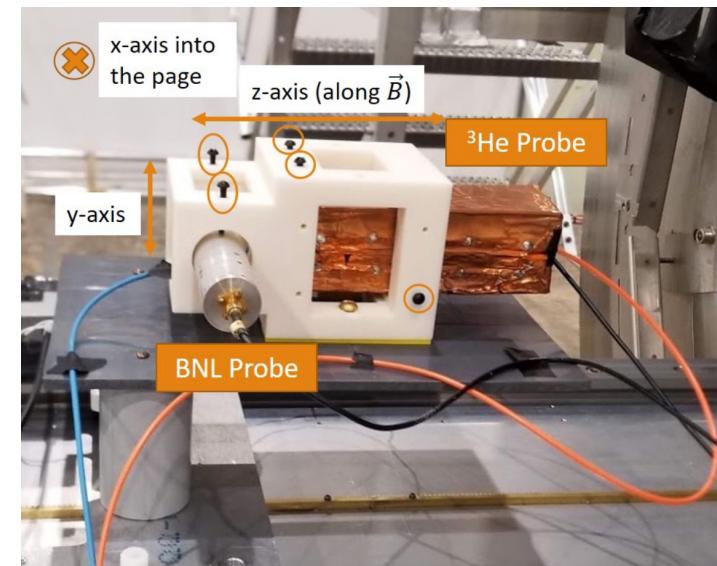
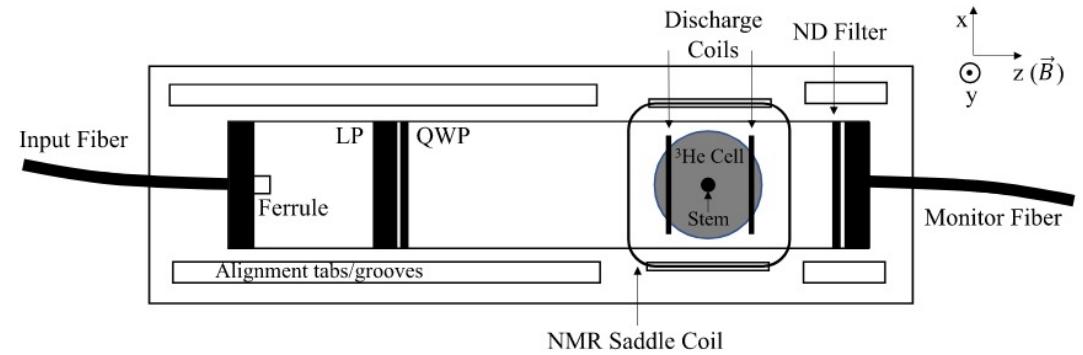


$\delta B \approx fT$ after integration for 220s



Gummel, C., et al. *The European Physical Journal D* 57 (2010): 303-320

High-field ($\sim T$) calibration with a precision of ppb



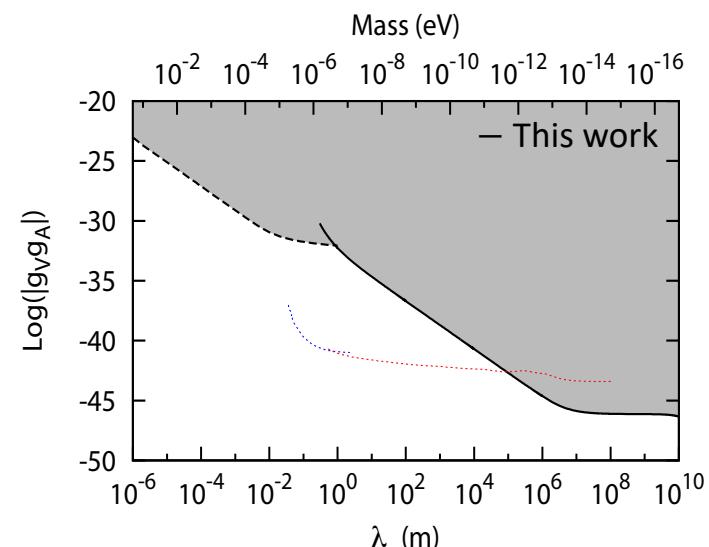
Farooq, Midhat, et al. *Physical Review Letters* 124.22 (2020): 223001.

Tests on fundamental physics

- Electric dipole moments (EDMs) searches
- Lorentz and CPT violation tests
- **Exotic spin-dependent interactions**
and axion dark matter searches.

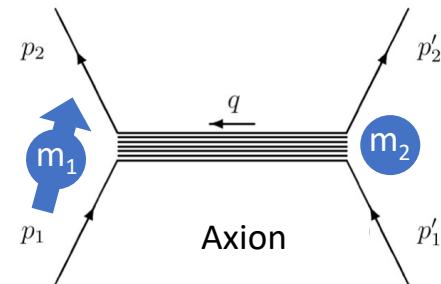
The fluctuated effective magnetic field B_{VA} causes the relaxation of ${}^3\text{He}$.

$T_1 \approx 347$ hrs[A. K. Petukhov, et. al. PRL 115, 182001 (2015)]

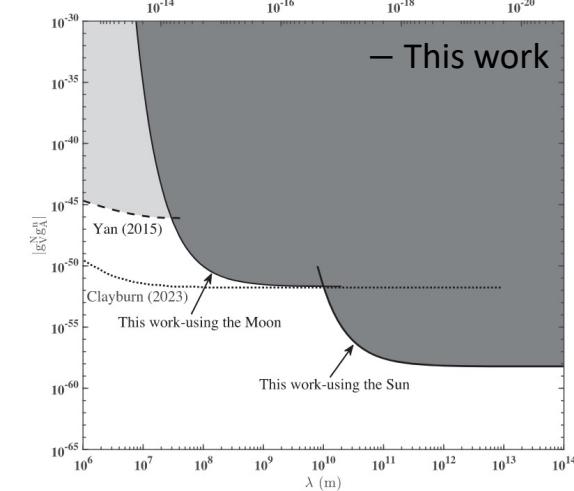
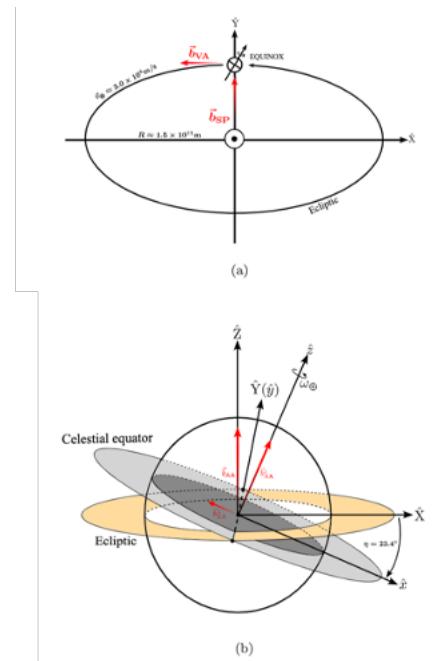


$$V_{VA} = \frac{\hbar g_V g_A}{2\pi} \frac{e^{-r/\lambda}}{r} \vec{\sigma} \cdot \vec{v}$$

$$= \hbar \gamma \vec{B}_{VA} \cdot \vec{\sigma}$$



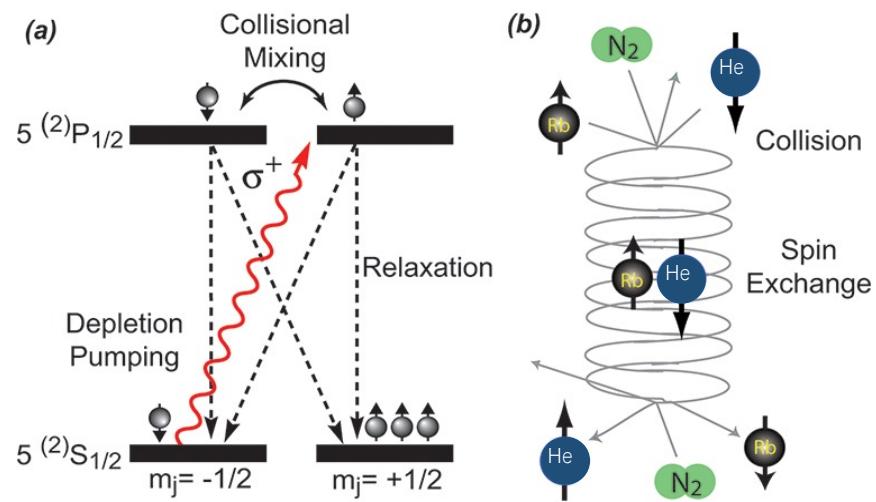
Measuring the effective magnetic field generated by the Sun and Moon using the ${}^3\text{He}-{}^{129}\text{Xe}$ comagnetometer.



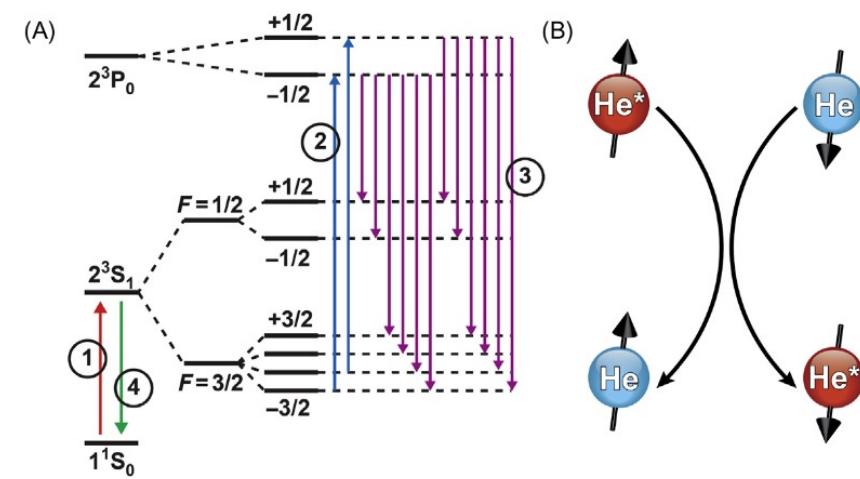
- Introduction to ${}^3\text{He}$ Polarization
- **The Basic Principle of MEOP**
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Indirect approaches of ${}^3\text{He}$ polarization

Spin Exchange Optical Pumping (SEOP)



Metastability-Exchange Optical Pumping (MEOP)



Spin Angular Momentum Transfer: photon \rightarrow electron \rightarrow nucleon

Comparison between these two approaches

- SEOP

1960: Bouchiat, M. Ab, et. al. *Physical review letters* 5.8 (1960): 373.

- Slow polarization production (several hours)
- Saturated ${}^3\text{He}$ Polarization (70%-80%)
- High pressure (0.5 to 13 bar)
- Gas cell heating required
- Easy for In-suit polarization
- Successfully applied to polarize ${}^{21}\text{Ne}$ 、 ${}^{129}\text{Xe}$ 、 ${}^{131}\text{Xe}$

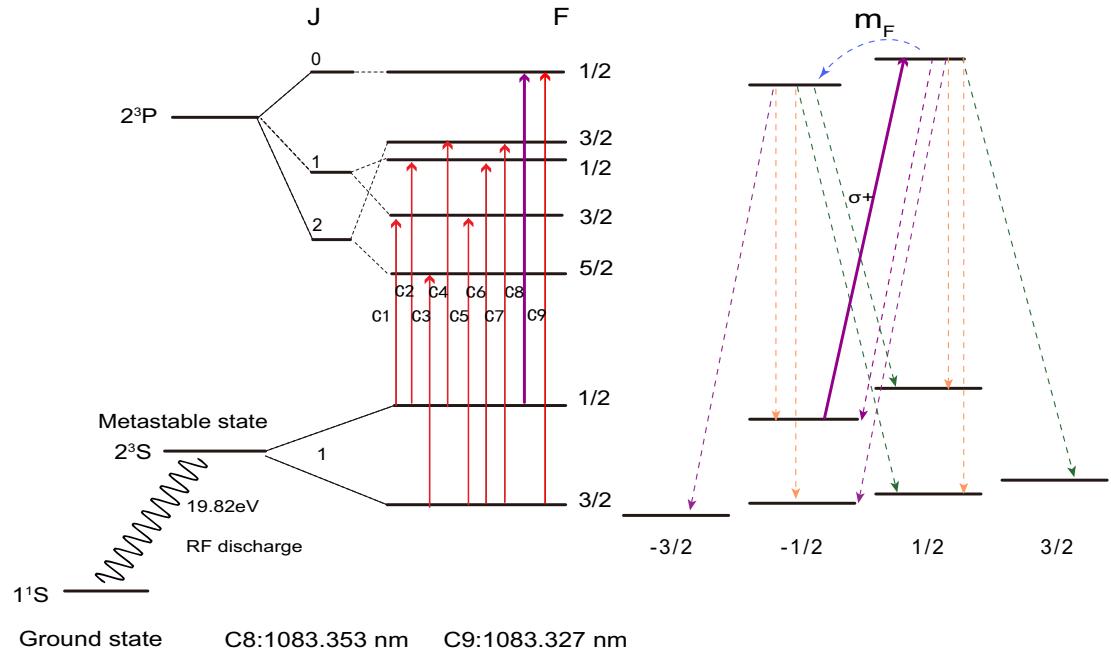
- MEOP

1962: Walters, G. K., et. al. *Physical Review Letters* 8.11 (1962): 439.

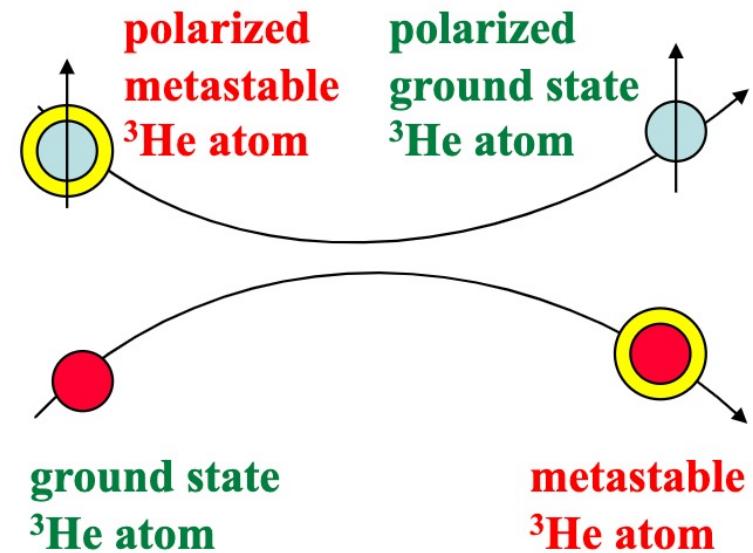
- Fast polarization production (several minutes)
- Saturated ${}^3\text{He}$ Polarization (70%-80%)
- Low pressure (1 mbar)
- Continuous RF discharge
- Need compression for application
- Exclusively applied to polarize Helium

Processes involved in low-field MEOP

- Optical Pumping



- Metastability-Exchange Collisions



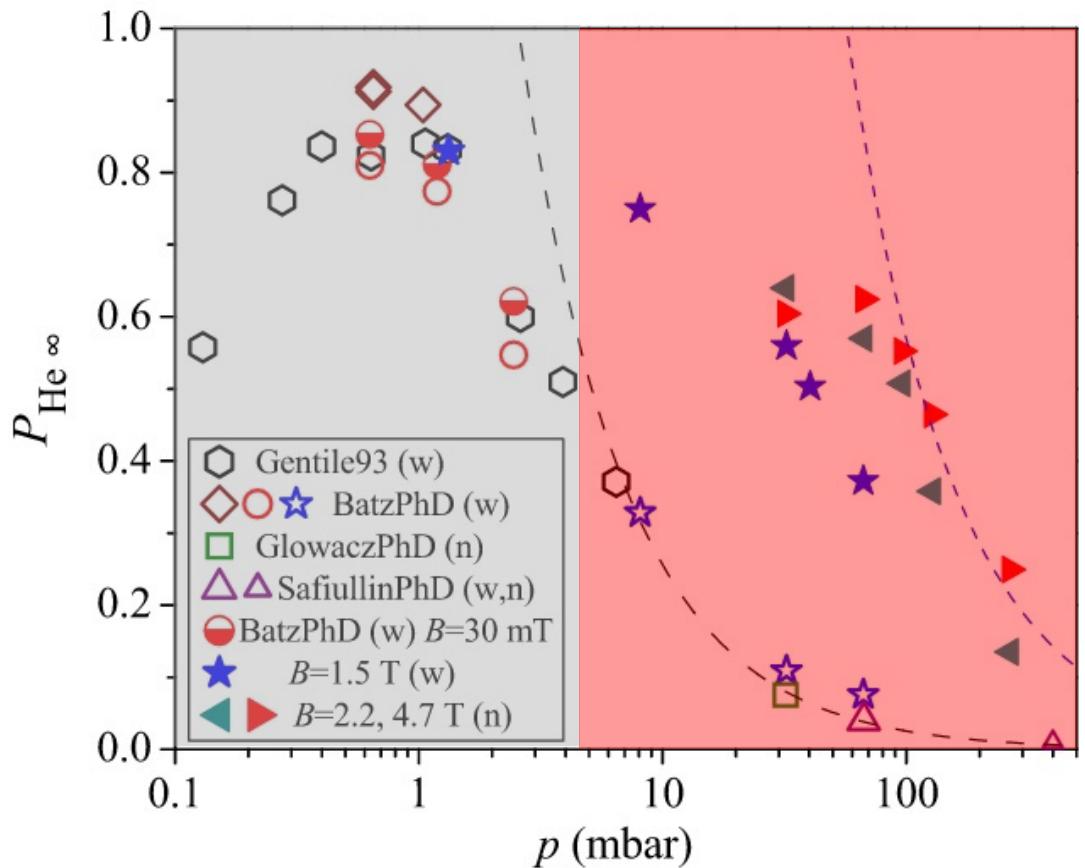
1. Striking the atoms into the metastable state by discharge,
2. The metastable state ^3He is polarized by the cycling optical transition between 2^3S-2^3P (e.g., C₈ or C₉ Component).

The excitation energy is exchanged in the collisions with nuclear spin unperturbed.

Recent Advancements in MEOP ^3He

Conditions of Efficient MEOP:

- low field ($\sim \text{mT}$), low pressure ($\sim \text{mbar}$)
- high field ($\sim \text{T}$), high pressure ($\sim 100 \text{ mbar}$)



MEOP ^3He without laser

PHYSICAL REVIEW A **98**, 063405 (2018)

Nuclear hyperpolarization of ^3He by magnetized plasmas

A. Maul,¹ P. Blümller,¹ P.-J. Nacher,² E. Otten,¹ G. Tastevin,² and W. Heil^{1,*}

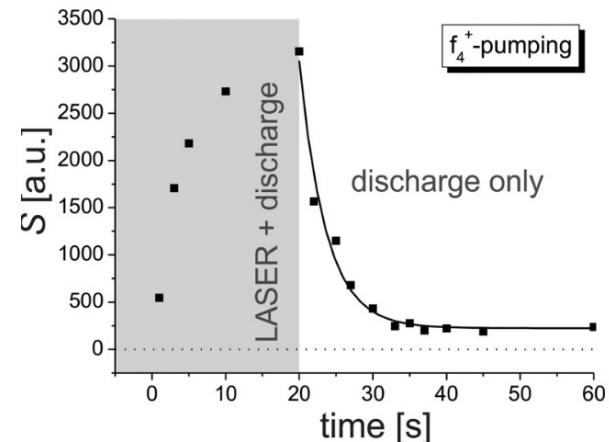
¹Institute of Physics, University of Mainz, 55128 Mainz, Germany

²Laboratoire Kastler Brossel, ENS-PSL University, Centre National de la Recherche Scientifique, Sorbonne Université, Collège de France, 24 rue Lhomond, 75005 Paris, France

(Received 20 June 2018; published 5 December 2018)

Polarization of Atoms in a Magnetized Plasma (PAMP)

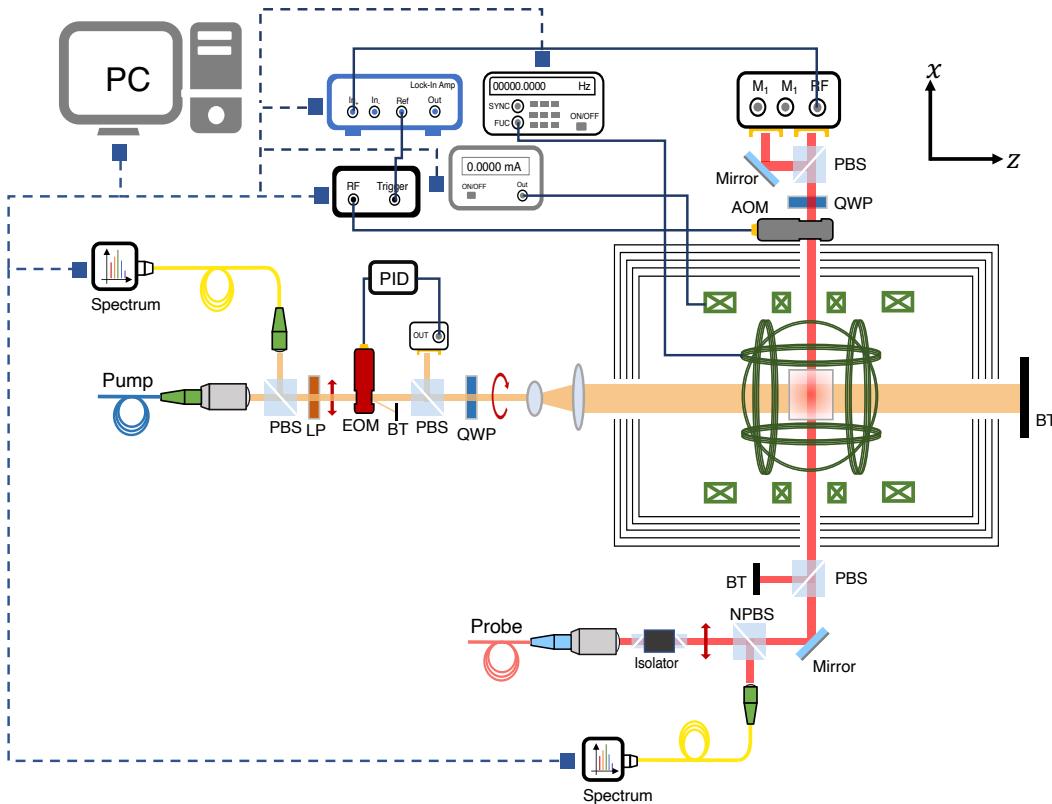
With an rf excitation of 100 MHz and magnetic field $B_0=4.7 \text{ T}$, polarizations in the 1 to 9% are achieved in mbar range without laser.



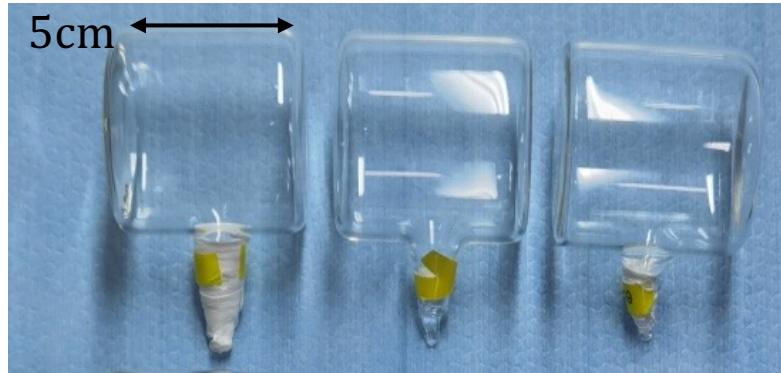
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Experimental Setup

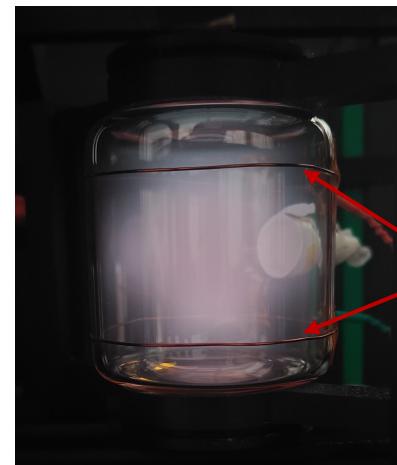
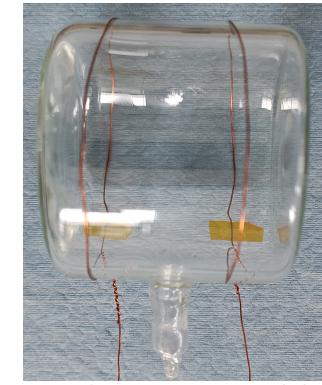
Schematic of experimental setup



Homemade Duran glass cell



Homemade GE180 glass cell



Discharge ring

The guiding magnetic field

The effect of B_0 gradient on the relaxation time

$$\frac{1}{T_1} \approx \frac{8R^4\gamma^2}{175D}(|\nabla B_x|^2 + |\nabla B_y|^2)$$

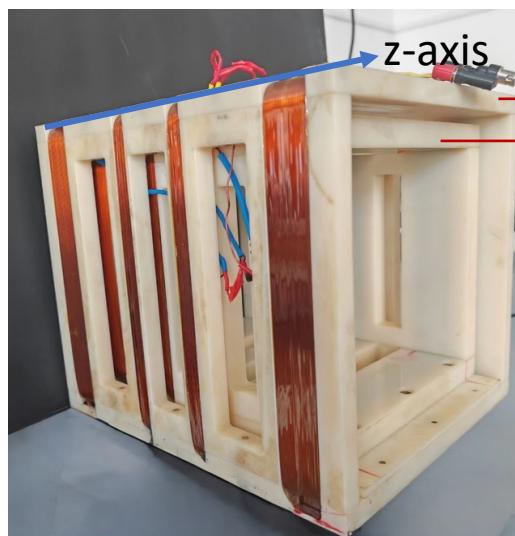
Low pressure $\frac{1}{T_2} \approx \frac{4R^4\gamma^2}{175D}(|\nabla B_x|^2 + |\nabla B_y|^2 + 2|\nabla B_z|^2)$

Cates, G. D., et. al. *Physical review A* 37.8 (1988): 2877.

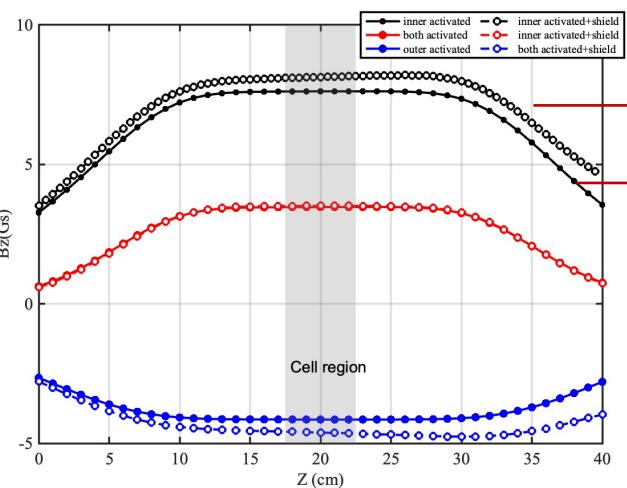
Dipole-free Merritt Coil:

Zero dipole moment outside the coil

Minimize the shielding effect on the field uniformity



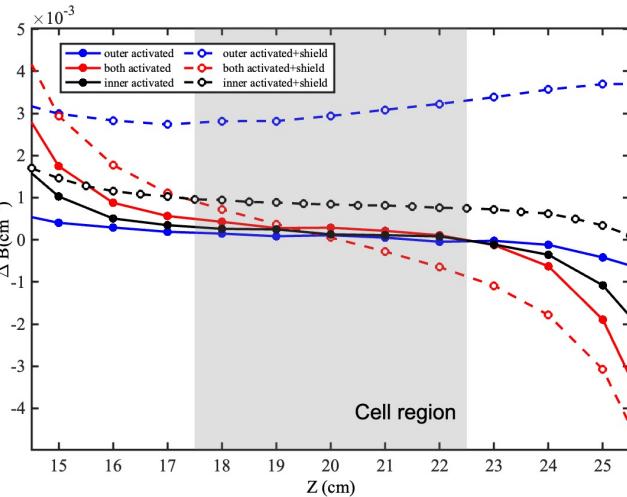
The B_z Component along z-axis



Merritt coil in the shield
Merritt coil without shield

The field of Merritt coil is asymmetry in the shield;

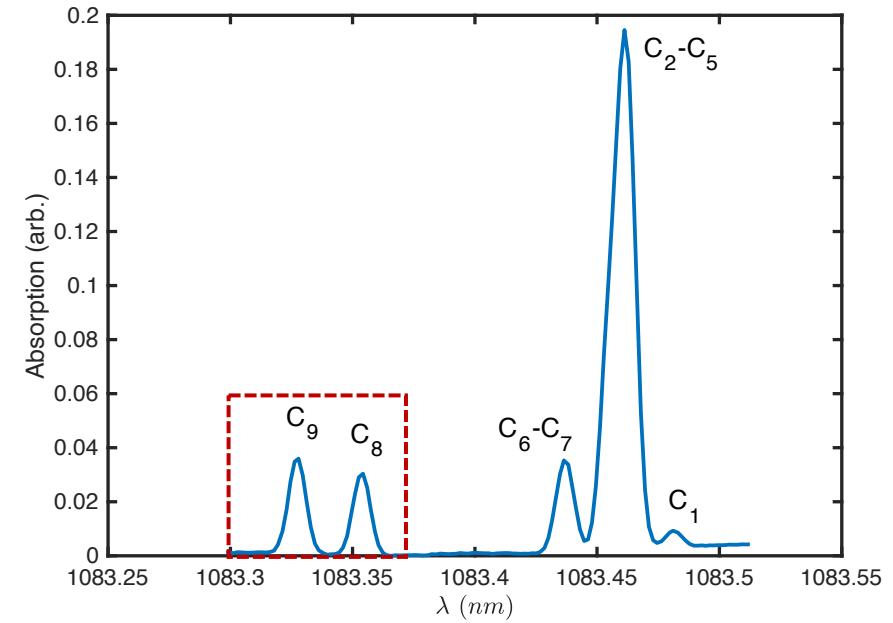
The gradient $\Delta B/B_0$ along z-axis



The dipole-free Merritt coil restores symmetry of the field.

Pump Laser

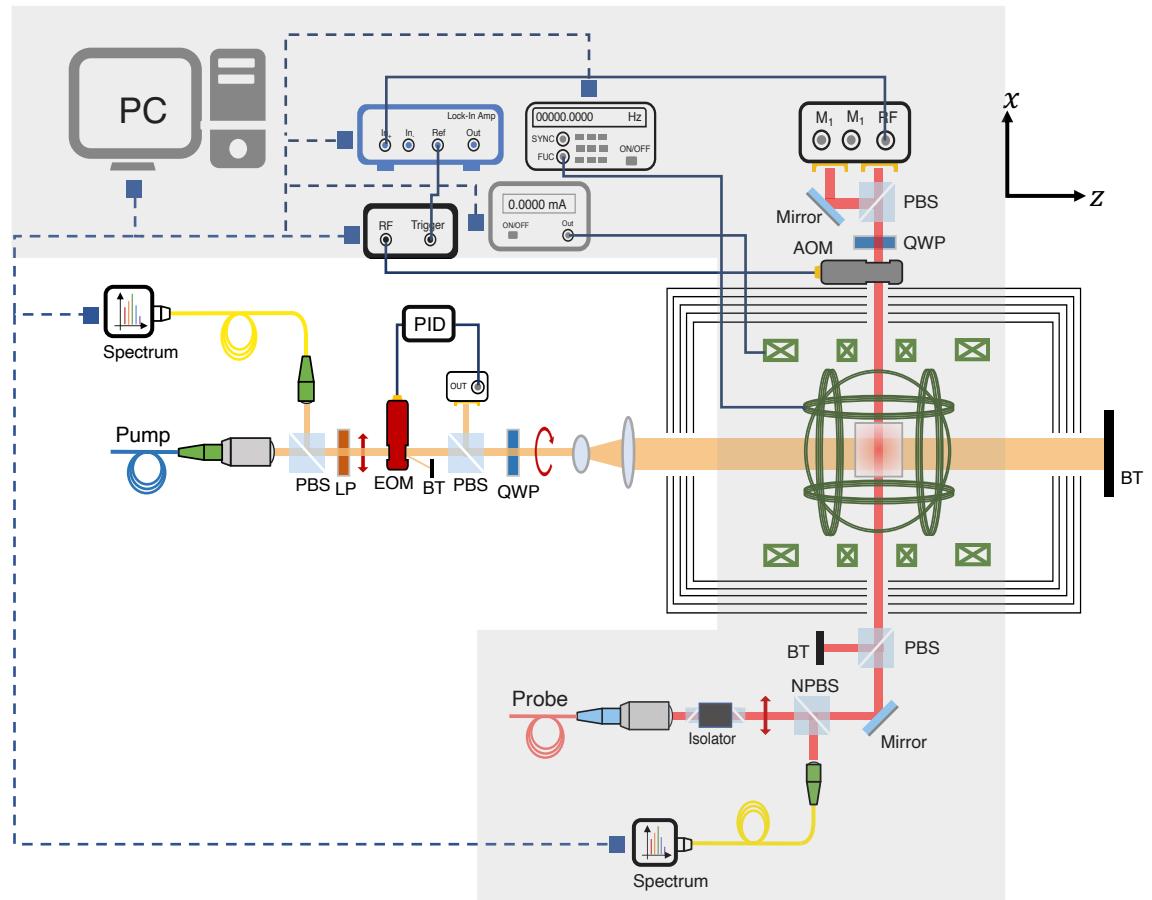
- Pumping laser: Keopsys Ytterbium Laser
- Power: ~ 2 W
- center wavelength: locked to transition line C_8
- linewidth: 2 GHz (matches the natural linewidth of ${}^3\text{He}$)



Experimental Spectrum of the 1083nm transition lines

Measurement and acquisition system

Detection Laser: DBR laser with power of about 1mW.



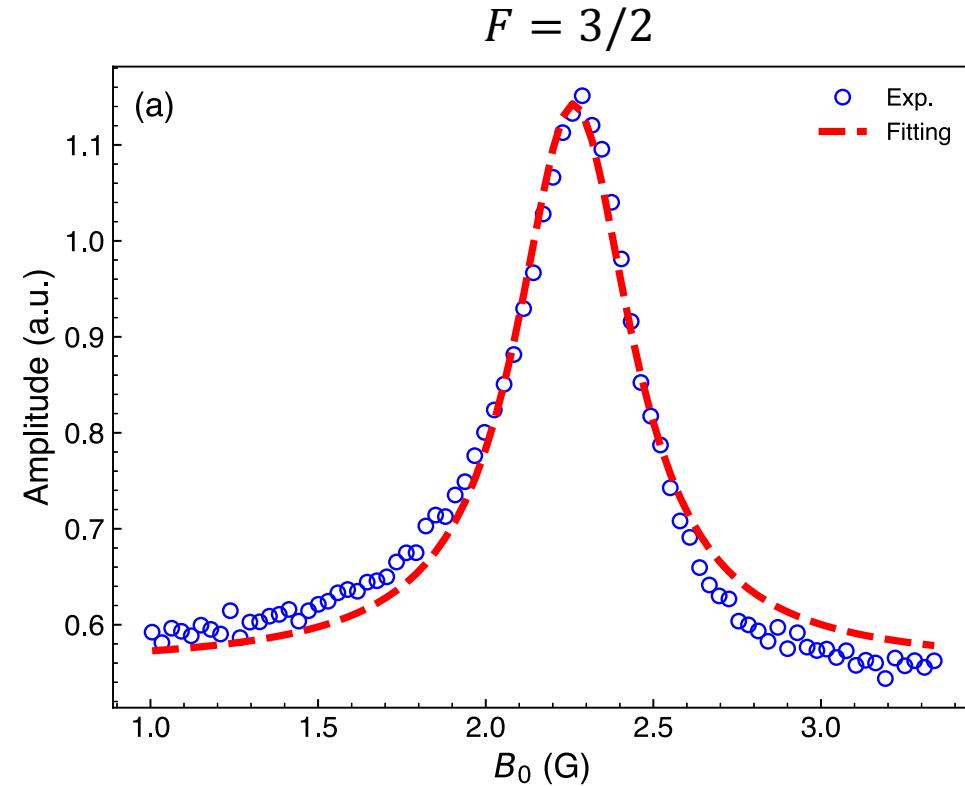
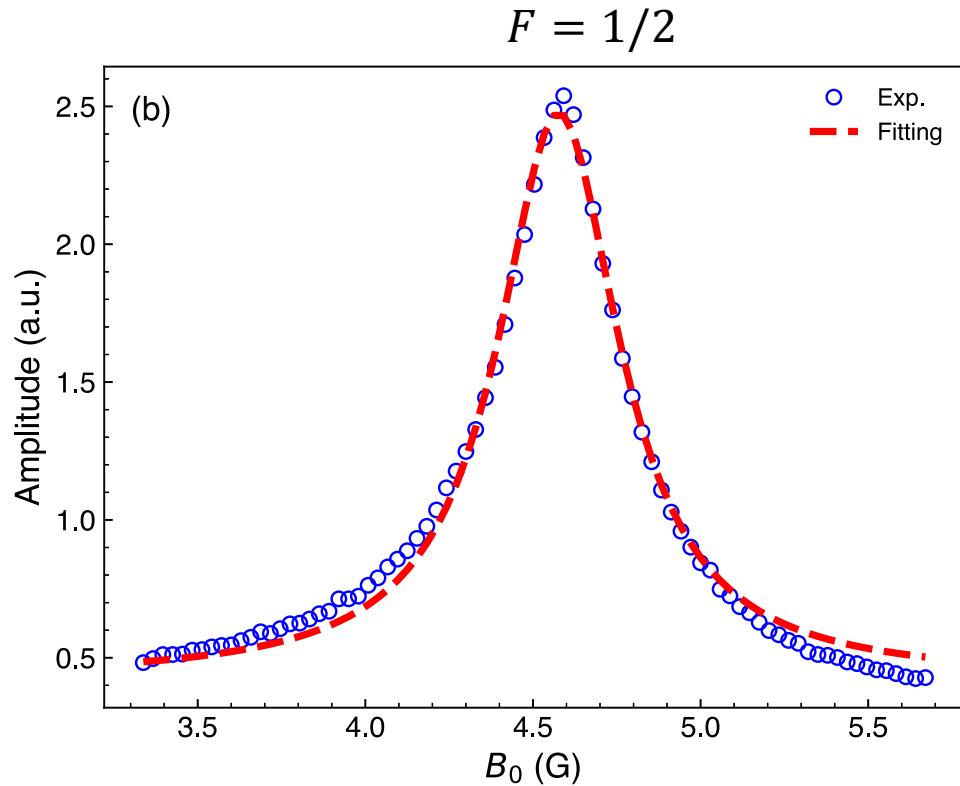
1. Optical polarimetry:

The wavelength of the laser is set to the transition line C₉; the power is modulated by an optical chopper and the absorption signal is extracted by lock-in method.

2. Magnetic resonance:

The wavelength of the laser is detuned, and the Faraday rotation signal is detected via homodyne detection.

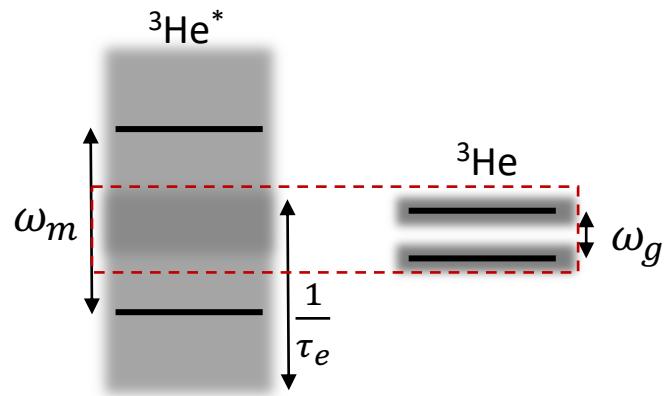
Magnetic resonance of the metastable states



The metastability-exchange collisions rate derived are
 $1/\tau_e = 1.02 \pm 0.02 \times 10^6 \text{ s}^{-1}$

Magnetic Resonance Measurement

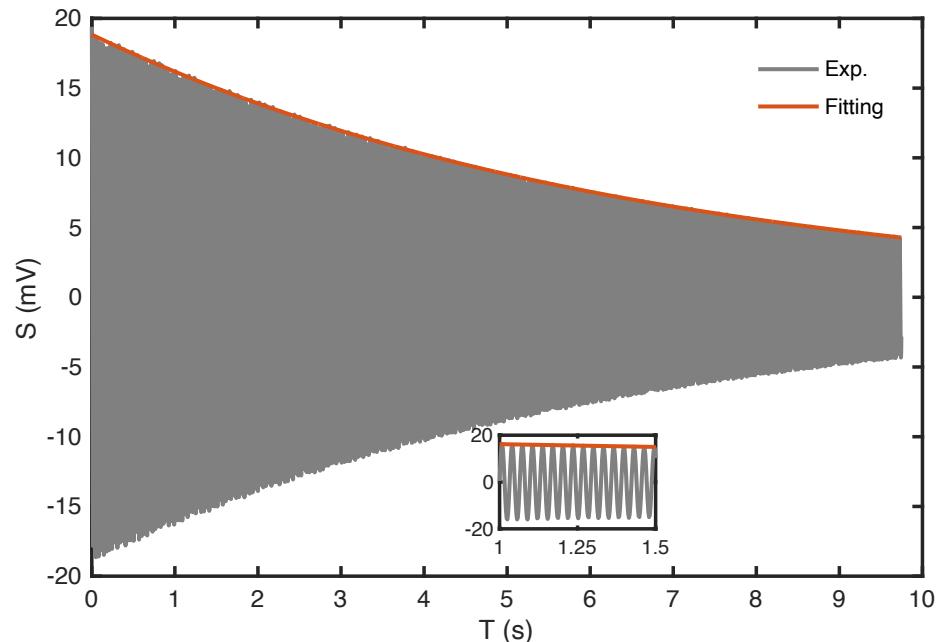
The coherence of the ground state is transferred to the metastable state by metastability-exchange collisions.



If $(\omega_m \ll \frac{1}{\tau_e})$, the metastable state and ground state is coupled.

Optical Detection of NMR

Magnetic resonance of the ground state



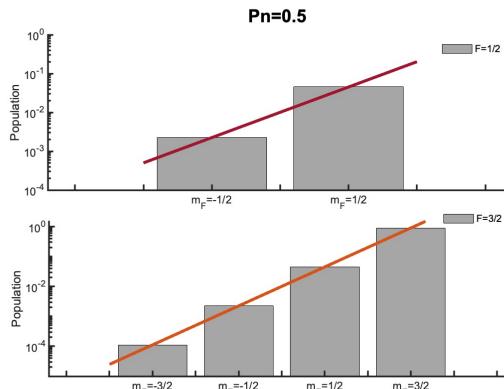
Optical detected FID Signal

Larmor frequency $\omega_0 = 30.7429 \pm 0.00002$ Hz
Transverse relaxation time $T_2 = 6.58 \pm 0.01$ s

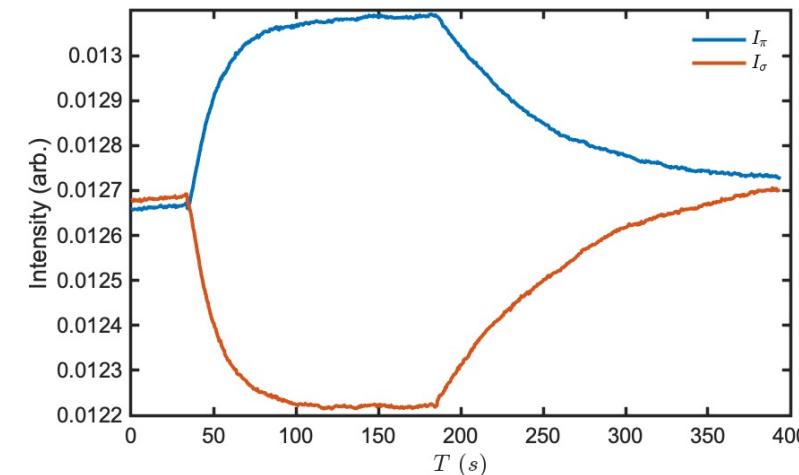
Optical Polarimetry

A spin-temperature distribution of the metastable state is decided by the ground ${}^3\text{He}$ polarization.

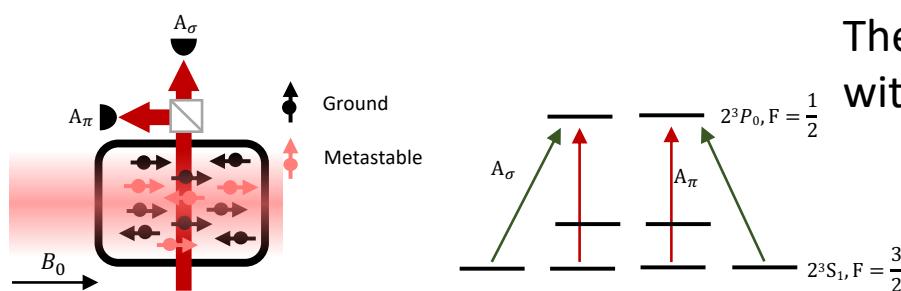
$$N_{m_F} \propto e^{\beta m_F}, \quad \beta = \ln \frac{1 + P_n}{1 - P_n}$$



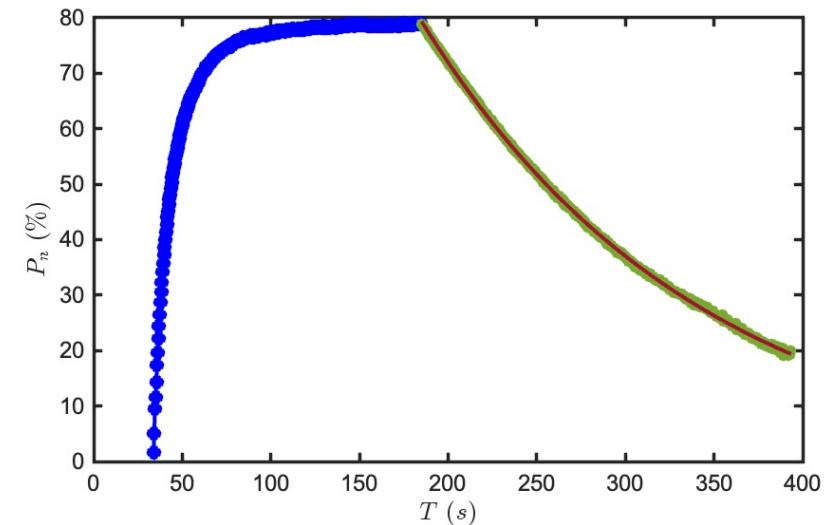
Spin-temperature distribution



The absorption of the probe light with different polarization



$$P_n = \sqrt{\frac{1 - A_\pi/A_\sigma}{1 + 2A_\pi/A_\sigma}}$$



The polarization buildup and decay

Saturated Polarization $P_n = 79 \pm 1 \%$
 Longitudinal relaxation time $T_1 = 158 \pm 2 \text{ s}$

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Summary

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- (1). Applications of polarized ^3He , particularly in fundamental physics
- (2). The basic of the MEOP
- (4). Progress in developing the low-field MEOP ^3He system, including
 - a. A dipole-free coil system for obtaining uniform field in the shield,
 - b. The magnetic resonance measurement in the system,
 - c. Optical polarimetry of the nuclear polarization.

- Outlook

The metastability exchange collisions may be harnessed to squeeze the noble gas nuclear spin.

A. Serafin, et. al. Physical review letters, 127, 013601 (2021)

Thanks for watching

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