

FOURIER TRANSFORM FOR CONTINUOUS TIME APERIODIC SIGNAL

3.1 Fourier Transform for Continuous Time A-periodic Signal

Let us consider an a-periodic signal $x(t)$ with finite duration as shown in Fig-3.1. From a-periodic signal, we can create a periodic signal $x_p(t)$ with period T_p as shown in Fig-3.2. Clearly $x_p(t) = x(t)$ in the limit as $T_p \rightarrow \infty$ that is

$$x(t) = \lim_{T_p \rightarrow \infty} x_p(t) \text{-----} (27)$$

$$x_p(t) = \sum_{k=-\infty}^{\infty} c_k e^{j2\pi k F_0 t} \text{-----} (28)$$

$$C_k = \frac{1}{T_p} \int_{-\frac{T_p}{2}}^{\frac{T_p}{2}} x_p(t) e^{-j2\pi k F_0 t} dt \text{-----} (29)$$

$$C_k = \frac{1}{T_p} \int_0^{T_p} x(t) e^{-j2\pi k F_0 t} dt \text{-----} (30)$$

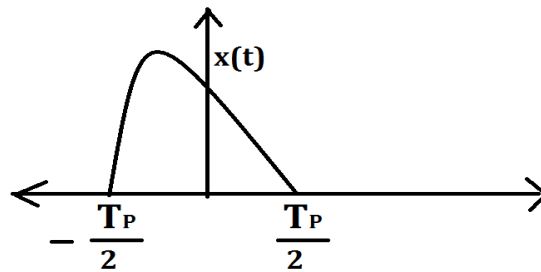


Fig-3.1 A-periodic signal $x(t)$

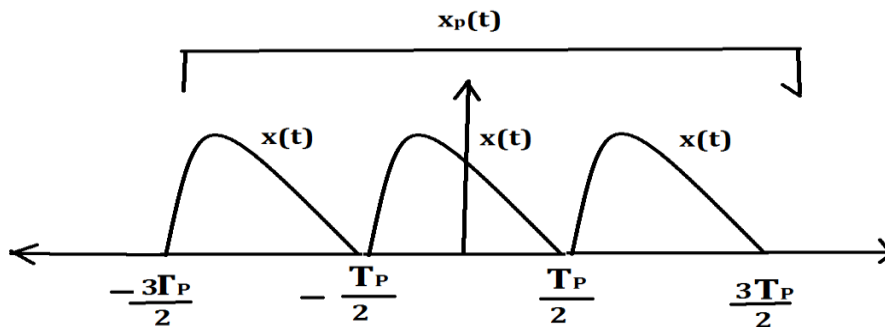


Fig-3.2: Periodic signal $x_p(t)$ constructed by three time repetition form of a-periodic signal $x(t)$ with period T_p

Since $x_p(t) = x(t)$ for $-\frac{T_p}{2} \leq t \leq \frac{T_p}{2}$ can be expressed as

$$C_k = \frac{1}{T_p} \int_{-\frac{T_p}{2}}^{\frac{T_p}{2}} x(t) e^{-j2\pi k F_0 t} dt \quad \text{-----} \quad (31)$$

It is also true that $x(t)=0$ for $|t| > \frac{T_p}{2}$, consequently the limits on the integral in (31) can be replaced by $-\infty$ & ∞ Hence,

$$C_k = \frac{1}{T_p} \int_{-\infty}^{\infty} x(t) e^{-j2\pi k F_0 t} dt \quad \text{-----} \quad (32)$$

Let us now define a function $X(F)$, called the Fourier transform of $x(t)$ as

$$X(F) = \int_{-\infty}^{\infty} x(t) e^{-j2\pi F t} dt \quad \text{-----} \quad (33)$$

$X(F)$ is the function of the continuous variable F if we compare (32)&(33) then Fourier coefficients C_K can be expressed in terms of $X(F)$ as

$$C_K = \frac{1}{T_p} X(kF_0) = X\left(\frac{K}{T_p}\right) \quad \text{-----} \quad (34)$$

Those Fourier coefficient as sample of $X(F)$ taken at multiples of F_0 and scaled by F_0 (multiplied by $\frac{1}{T_p}$). Substitution for C_K from (34) in to (28) yields

$$x_p(t) = \frac{1}{T_p} \sum_{k=-\infty}^{\infty} X\left(\frac{K}{T_p}\right) e^{j2\pi k F_0 t} \quad \text{-----} \quad (35)$$

If T_p is infinity then we write $\Delta F = \frac{1}{T_p}$ with substitution, equation (35) becomes as

$$x_p(t) = \sum_{k=-\infty}^{\infty} X(k\Delta F) e^{j2\pi k \Delta F t} \Delta F \quad \text{-----} \quad (36)$$

$$\text{Now } \lim_{T_p \rightarrow \infty} x_p(t) = \lim_{\Delta F \rightarrow 0} \sum_{k=-\infty}^{\infty} X(k\Delta F) e^{j2\pi k \Delta F t} \Delta F \quad \text{-----} \quad (37)$$

Apply the equation-(23) on equation -(33) and we get following equation as below

$$x(t) = \lim_{P \rightarrow \infty} x_P(t) = \lim_{\Delta F \rightarrow 0} \sum_{k=-\infty}^{\infty} X(k\Delta F) e^{j2\pi k\Delta F t} \quad \Delta F \text{ ----- (38)}$$

if $\Delta F \rightarrow 0$ then $k\Delta F$ also act as cautious variable (although k is integer). We are replacing $k\Delta F$ by F which is act as continuous variable of frequency. So equation (33) and (34) can be written as given form

$$X(F) = \int_{-\infty}^{\infty} x(t) e^{-j2\pi F t} dt \text{ ----- (39)}$$

$$x(t) = \int_{-\infty}^{\infty} X(F) e^{j2\pi F t} dF \text{ ----- (40)}$$

Where $X(F)$ is the Fourier transform of $x(t)$.

Summary 3(A): Fourier Transform for Continuous Time Aperiodic Signal

$$\begin{aligned} x(t) &= \int_{-\infty}^{\infty} X(F) e^{j2\pi F t} dF \\ X(F) &= \int_{-\infty}^{\infty} x(t) e^{-j2\pi F t} dt \end{aligned}$$

3.2 Energy Density and Power Density Spectrum for Continuous Time Aperiodic Signal

A continuous time aperiodic signal has infinite energy which is given as

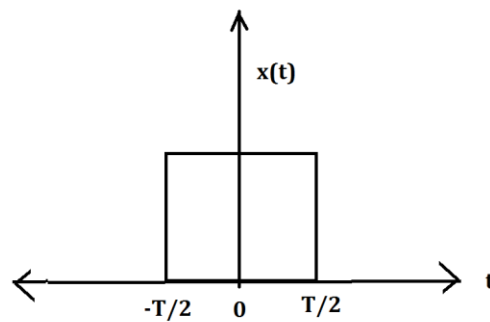
$$E_X = \int_{-\infty}^{+\infty} [x(t)]^2 dt = \int_{-\infty}^{+\infty} [X(F)]^2 dF \text{ ----- (41)}$$

The $[X(F)]^2$ is represented the energy density spectrum of continuous time aperiodic signal

In this case, we can't calculate power density spectrum of aperiodic signal due to unknown period.

Example-05: Determine the Fourier transform and energy density spectrum of rectangular pulse define as

$$x(t) = \begin{cases} 0, & \frac{T}{2} < |t| \\ 1, & \frac{T}{2} \geq |t| \end{cases}$$



Solution:

We know that from summery-3(A) (see page -19) for Continuous Time A-periodic Signal

$$X(F) = \int_{-\infty}^{\infty} x(t) e^{-j2\pi Ft} dt$$

$$X(F) = \int_{-\infty}^{-T/2} x(t) e^{-j2\pi Ft} dt + \int_{-T/2}^{T/2} x(t) e^{-j2\pi Ft} dt + \int_{T/2}^{\infty} x(t) e^{-j2\pi Ft} dt$$

$$X(F) = \int_{-\infty}^{-T/2} 0 e^{-j2\pi Ft} dt + \int_{-T/2}^{T/2} 1 e^{-j2\pi Ft} dt + \int_{T/2}^{\infty} 0 e^{-j2\pi Ft} dt$$

$$X(F) = \int_{-T/2}^{T/2} 1 e^{-j2\pi Ft} dt$$

$$X(F) = \frac{\sin(\pi FT)}{\pi F}$$

$$X(F) = T \frac{\sin(\pi FT)}{\pi FT}$$

$$\text{Or } X(F) = T \text{sinc}(\pi FT)$$

Fourier transform for given rectangular pulse is $T \frac{\sin(\pi FT)}{\pi FT}$

So energy density spectrum for given rectangular pulse is $X(F)^2 = [T \frac{\sin(\pi FT)}{\pi FT}]^2$

The graphical representation of energy density spectrum for given rectangular pulse as shown in fig-(3.2)

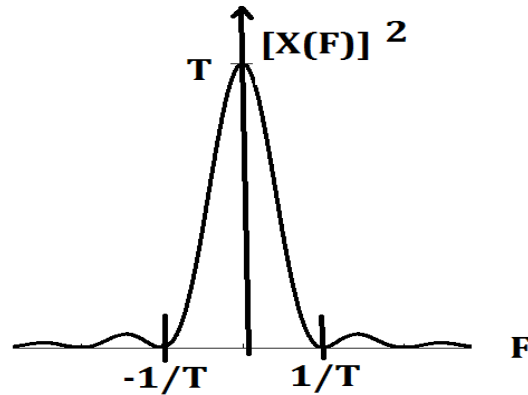


Fig-3.3: Energy density spectrum (EDS) for given rectangular pulse

Properties of the Fourier transform

The Fourier transform has many useful properties that make calculations easier and also help thinking about the structure of signals and the action of systems on signals.

✓ **Linearity:**

The Fourier transform is **linear**: if

$$x_1(t) \leftrightarrow X_1(\omega) \quad \text{and} \quad x_2(t) \leftrightarrow X_2(\omega),$$

then

$$c_1x_1(t) + c_2x_2(t) \leftrightarrow c_1X_1(\omega) + c_2X_2(\omega)$$

for any two numbers c_1 and c_2 .

Proof: obvious –

$$\begin{aligned}\mathcal{F}[c_1x_1(t) + c_2x_2(t)] &= \int_{-\infty}^{\infty} [c_1x_1(t) + c_2x_2(t)] e^{-j\omega t} dt \\ &= c_1 \int_{-\infty}^{\infty} x_1(t) e^{-j\omega t} dt + c_2 \int_{-\infty}^{\infty} x_2(t) e^{-j\omega t} dt \\ &= c_1X_1(\omega) + c_2X_2(\omega)\end{aligned}$$

Linearity

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for any two numbers c_1 and c_2 .

Proof: obvious –

$$\begin{aligned} \mathcal{F}[c_1 x_1(t) + c_2 x_2(t)] &= \int_{-\infty}^{\infty} [c_1 x_1(t) + c_2 x_2(t)] e^{-j\omega t} dt \\ &= c_1 \int_{-\infty}^{\infty} x_1(t) e^{-j\omega t} dt + c_2 \int_{-\infty}^{\infty} x_2(t) e^{-j\omega t} dt \\ &= c_1 X_1(\omega) + c_2 X_2(\omega) \end{aligned}$$

Time Shifting:

If $x(t) \leftrightarrow X(\omega)$, then $x(t - c) \leftrightarrow X(\omega)e^{-j\omega c}$ for any constant c .

Proof:

$$\begin{aligned} \mathcal{F}[x(t - c)] &= \int_{-\infty}^{\infty} x(t - c) e^{-j\omega t} dt \\ &= \int_{-\infty}^{\infty} x(t) e^{-j\omega(t+c)} dt \\ &= e^{-j\omega c} \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt \\ &= X(\omega) e^{-j\omega c}. \end{aligned}$$

✓ **Frequency Shifting:**

If $x(t) \leftrightarrow X(\omega)$, then $x(t)e^{j\omega_0 t} \leftrightarrow X(\omega - \omega_0)$ for any **real** ω_0 .

Proof:

$$\begin{aligned}\mathcal{F}[x(t)e^{j\omega_0 t}] &= \int_{-\infty}^{\infty} x(t)e^{j\omega_0 t} e^{-j\omega t} dt \\ &= \int_{-\infty}^{\infty} x(t)e^{-j(\omega - \omega_0)t} dt \\ &= X(\omega - \omega_0).\end{aligned}$$

Multiplication by a complex exponential

If $x(t) \leftrightarrow X(\omega)$, then $x(t)e^{j\omega_0 t} \leftrightarrow X(\omega - \omega_0)$ for any **real** ω_0 .

Proof:

$$\begin{aligned}\mathcal{F}[x(t)e^{j\omega_0 t}] &= \int_{-\infty}^{\infty} x(t)e^{j\omega_0 t} e^{-j\omega t} dt \\ &= \int_{-\infty}^{\infty} x(t)e^{-j(\omega - \omega_0)t} dt \\ &= X(\omega - \omega_0).\end{aligned}$$

Time Reversal

$$f(-t) \xleftrightarrow{F} F(-j\omega)$$

Pf)

$$\begin{aligned} F[f(-t)] &= \int_{-\infty}^{\infty} f(-t) e^{-j\omega t} dt = \int_{t=-\infty}^{t=\infty} f(-t) e^{-j\omega t} dt \\ &= \int_{-t=-\infty}^{-t=\infty} f(t) e^{j\omega t} d(-t) = \int_{-t=-\infty}^{-t=\infty} f(t) e^{j\omega t} (-dt) \\ &= - \int_{t=\infty}^{t=-\infty} f(t) e^{j\omega t} dt = \int_{t=-\infty}^{t=\infty} f(t) e^{j\omega t} dt \\ &= \int_{-\infty}^{\infty} f(t) e^{j\omega t} dt = F(-j\omega) \end{aligned}$$

✓ **The Modulation Theorem:**

If $x(t) \leftrightarrow X(\omega)$, then $x(t) \cos(\omega_0 t) \leftrightarrow \frac{1}{2}[X(\omega + \omega_0) + X(\omega - \omega_0)]$.

Proof: use linearity and the last property to get

$$\begin{aligned}\mathcal{F}[x(t) \cos(\omega_0 t)] &= \mathcal{F}\left[\frac{1}{2}x(t) (e^{j\omega_0 t} + e^{-j\omega_0 t})\right] \\ &= \frac{1}{2}\mathcal{F}[x(t)e^{j\omega_0 t}] + \frac{1}{2}\mathcal{F}[x(t)e^{-j\omega_0 t}] \\ &= \frac{1}{2}[X(\omega - \omega_0) + X(\omega + \omega_0)].\end{aligned}$$

Multiplication by a cosine

If $x(t) \leftrightarrow X(\omega)$, then $x(t) \cos(\omega_0 t) \leftrightarrow \frac{1}{2}[X(\omega + \omega_0) + X(\omega - \omega_0)]$.

Proof: use linearity and the last property to get

$$\begin{aligned}\mathcal{F}[x(t) \cos(\omega_0 t)] &= \mathcal{F}\left[\frac{1}{2}x(t) (e^{j\omega_0 t} + e^{-j\omega_0 t})\right] \\ &= \frac{1}{2}\mathcal{F}[x(t)e^{j\omega_0 t}] + \frac{1}{2}\mathcal{F}[x(t)e^{-j\omega_0 t}] \\ &= \frac{1}{2}[X(\omega - \omega_0) + X(\omega + \omega_0)].\end{aligned}$$

If $x(t) \leftrightarrow X(\omega)$ and $v(t) \leftrightarrow V(\omega)$, then

$$x(t) \star v(t) \leftrightarrow X(\omega)V(\omega)$$

Proof:

$$\begin{aligned}\mathcal{F}[x(t) \star v(t)] &= \int_{-\infty}^{\infty} [x(t) \star v(t)] e^{-j\omega t} dt \\&= \int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} x(\lambda) v(t - \lambda) d\lambda \right) e^{-j\omega t} dt \\&= \int_{-\infty}^{\infty} x(\lambda) \underbrace{\left(\int_{-\infty}^{\infty} v(t - \lambda) e^{-j\omega t} dt \right)}_{\mathcal{F}[v(t-\lambda)]} d\lambda \\&= \int_{-\infty}^{\infty} x(\lambda) V(\omega) e^{-j\omega \lambda} d\lambda \\&= V(\omega) \int_{-\infty}^{\infty} x(\lambda) e^{-j\omega \lambda} d\lambda \\&= X(\omega) V(\omega).\end{aligned}$$

Let $x(t)$ and $v(t)$ be real-valued signals. Then

$$\int_{-\infty}^{\infty} x(t)v(t)dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} \overline{X(\omega)}V(\omega)d\omega$$

Proof:

$$\begin{aligned} \int_{-\infty}^{\infty} x(t)v(t)dt &= \int_{-\infty}^{\infty} x(t) \left(\frac{1}{2\pi} \int_{-\infty}^{\infty} V(\omega)e^{j\omega t}d\omega \right) dt \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} V(\omega) \left(\int_{-\infty}^{\infty} x(t)e^{j\omega t}dt \right) d\omega \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} V(\omega)X(-\omega)d\omega \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \overline{X(\omega)}V(\omega)d\omega, \end{aligned}$$

where we used the fact that, since $x(t)$ is real,

$$\overline{X(\omega)} = \int_{-\infty}^{\infty} x(t)e^{j\omega t}dt = X(-\omega).$$

An important consequence of Parseval's theorem is that

$$\int_{-\infty}^{\infty} x^2(t)dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |X(\omega)|^2d\omega.$$

In other words, signal energy can be computed both in time domain and in frequency domain (up to a factor of $1/2\pi$).

Fourier Transform--Delta Function

The Fourier transform of the delta function is given by

$$\begin{aligned} \mathcal{F}_x [\delta(x - x_0)](k) &= \int_{-\infty}^{\infty} \delta(x - x_0) e^{-2\pi i k x} dx \\ &= e^{-2\pi i k x_0}. \end{aligned}$$

Fourier Transform--Cosine

$$\begin{aligned}\mathcal{F}_x [\cos (2 \pi k_0 x)](k) &= \int_{-\infty}^{\infty} e^{-2 \pi i k x} \left(\frac{e^{2 \pi i k_0 x} + e^{-2 \pi i k_0 x}}{2} \right) d x \\&= \frac{1}{2} \int_{-\infty}^{\infty} [e^{-2 \pi i (k-k_0) x} + e^{-2 \pi i (k+k_0) x}] d x \\&= \frac{1}{2} [\delta (k - k_0) + \delta (k + k_0)],\end{aligned}$$

where $\delta (x)$ is the [delta function](#).

Fourier Transform--Sine

$$\begin{aligned}\mathcal{F}_x [\sin (2 \pi k_0 x)](k) &= \int_{-\infty}^{\infty} e^{-2 \pi i k x} \left(\frac{e^{2 \pi i k_0 x} - e^{-2 \pi i k_0 x}}{2 i} \right) d x \\&= \frac{1}{2} i \int_{-\infty}^{\infty} [-e^{-2 \pi i (k-k_0) x} + e^{-2 \pi i (k+k_0) x}] d x \\&= \frac{1}{2} i [\delta (k + k_0) - \delta (k - k_0)],\end{aligned}$$

where $\delta (x)$ is the [delta function](#).