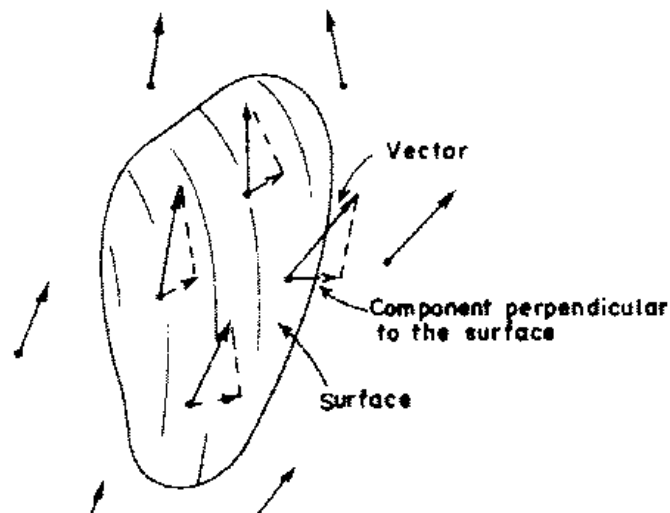


Chapter 2 Review on Vector Algebra and Vector Calculus

A *scalar* is specified by the magnitude of a numerical value, whereas a *vector* contains the information about the magnitude and the direction of a quantity. Vector analysis has to be invoked to properly describe the physical quantities in electromagnetics. For example, two important concepts, **flux** and **circulation**, frequently occur in Electromagnetics.

Flux: (average surface normal component)·(surface area)

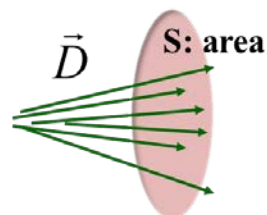
In plain text, a flux is the total “amount” of a vector quantity going out of a surface area. In math, flux is the surface integral of a vector normal to a surface.



E.g. Water flux \propto strength of water source



Electric Flux Φ_e

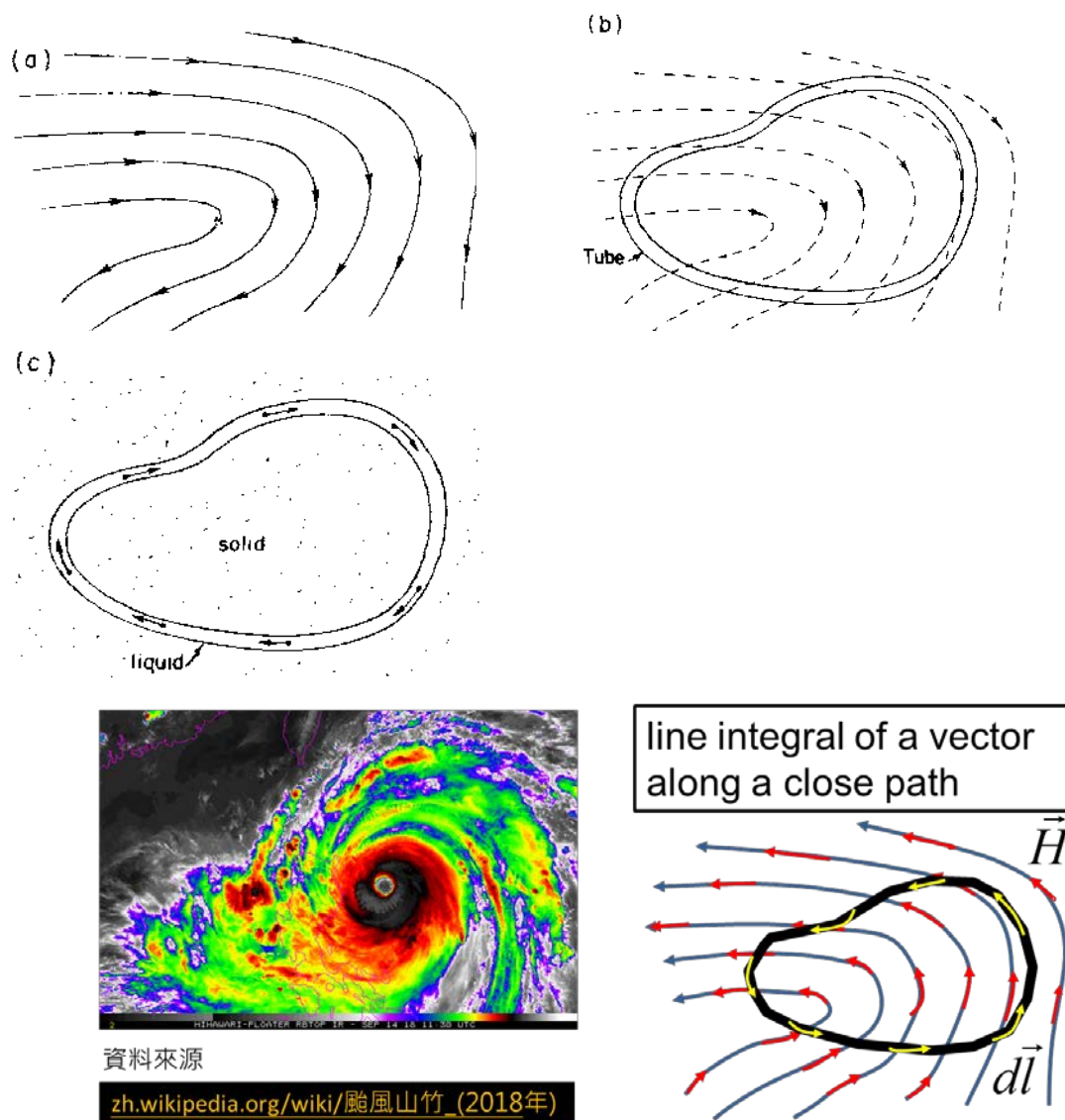


Circulation: (average tangential component)·(closed path)

CHAPTER 2 Review on Vector Algebra and Calculus

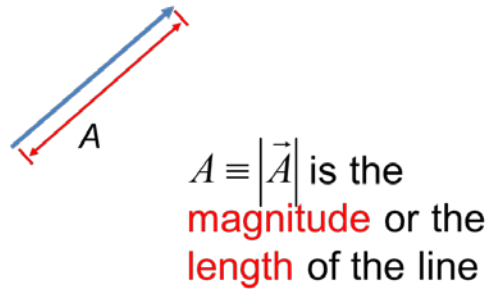
In plain text, circulation is the strength of a vector quantity going around a closed path. In math, circulation is the line integral of a vector along a closed path.

In the following we review the basic concept of vector algebra and calculus pertaining to the Cartesian, cylindrical, and spherical coordinates.

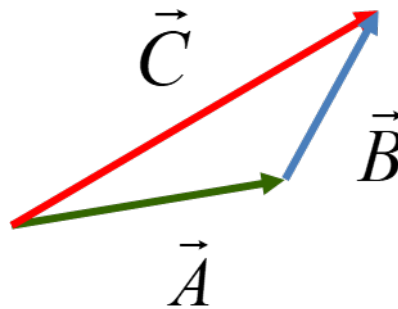


Basic Vector Algebra

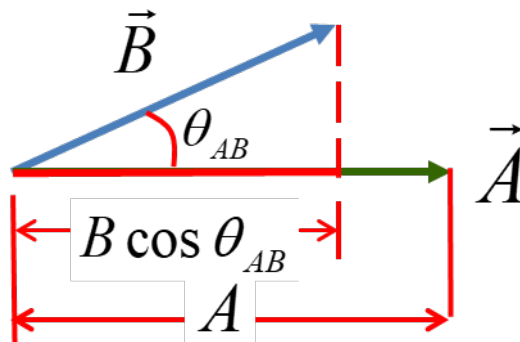
A Vector $\vec{A} = A\hat{a}_A$ has a magnitude of A pointing to the direction of \hat{a}_A where \hat{a}_A is a unit vector. A unit vector has a magnitude of 1.



Vector Sum $\vec{C} = \vec{A} + \vec{B}$, can be obtained by a head-to-tail rule shown below

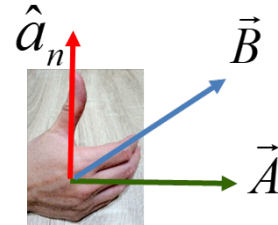
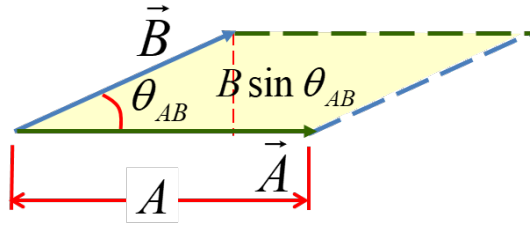


Scalar or Dot Product $\vec{A} \cdot \vec{B} = AB \cos \theta_{AB}$: multiplication of the projection of \vec{B} along \hat{a}_A and A .



Thus $A = \sqrt{\vec{A} \cdot \vec{A}}$ and $\vec{A} \cdot \vec{B} = \vec{B} \cdot \vec{A}$ (commutative)

Vector or Cross Product $\vec{A} \times \vec{B} = AB \sin \theta_{AB} \hat{a}_n$



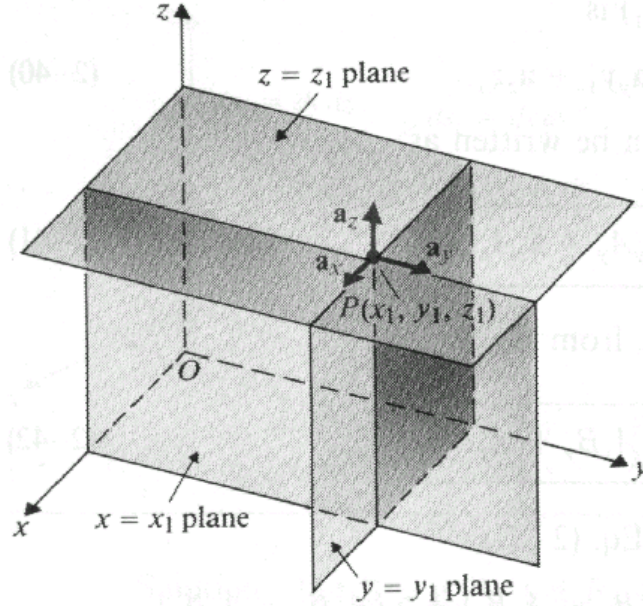
Note:

1. $|\vec{A} \times \vec{B}| = AB \sin \theta_{AB}$ is equal to the area of the parallelogram shown above.
2. The direction of the unit vector \hat{a}_n follows the so-called right-hand rule, with which you curl your right four fingers from the direction \vec{A} toward \vec{B} to find \hat{a}_n along the thumb direction. Therefore, the direction of the surface area formed by $\vec{A} \times \vec{B}$ is normal to the surface or perpendicular to \vec{A} and \vec{B} .
3. $\vec{A} \times \vec{B} = -\vec{B} \times \vec{A}$ (anti-commutative, surface areas are the same, but sense of direction is reversed)
4. $(\vec{A} \times \vec{B}) \times \vec{C} \neq \vec{A} \times (\vec{B} \times \vec{C})$
5. $(\vec{A} \times \vec{B}) \cdot \vec{C} = (\vec{B} \times \vec{C}) \cdot \vec{A} = (\vec{C} \times \vec{A}) \cdot \vec{B}$ is the volume of the parallelepiped formed by the three vectors $\vec{A}, \vec{B}, \vec{C}$.

Orthogonal Coordinate Systems

Cartesian Coordinates (x, y, z) : suitable for problems with rectangular

symmetry. The three coordinates are within the range of $-\infty \leq x \leq \infty$, $-\infty \leq y \leq \infty$, $-\infty \leq z \leq \infty$.



A general expression of a vector:

$$\vec{A} = \hat{a}_x A_x + \hat{a}_y A_y + \hat{a}_z A_z$$

The three unit vectors, $\hat{a}_x, \hat{a}_y, \hat{a}_z$ are arranged such that

$$\hat{a}_x \times \hat{a}_y = \hat{a}_z, \quad \hat{a}_y \times \hat{a}_z = \hat{a}_x, \quad \hat{a}_z \times \hat{a}_x = \hat{a}_y, \quad \hat{a}_x \cdot \hat{a}_y = 0,$$

$$\hat{a}_y \cdot \hat{a}_z = 0, \quad \hat{a}_z \cdot \hat{a}_x = 0$$

A differential length:

$$d\vec{l} = \hat{a}_x dx + \hat{a}_y dy + \hat{a}_z dz = d\vec{l}_x + d\vec{l}_y + d\vec{l}_z$$

A differential surface:

$$\begin{aligned} d\vec{s} &= \hat{a}_x ds_x + \hat{a}_y ds_y + \hat{a}_z ds_z \\ &= d\vec{l}_y \times d\vec{l}_z + d\vec{l}_z \times d\vec{l}_x + d\vec{l}_x \times d\vec{l}_y, \end{aligned}$$

where $ds_x = \left| d\vec{l}_y \times d\vec{l}_z \right| = dydz$

$$ds_y = \left| d\vec{l}_z \times d\vec{l}_x \right| = dzdx$$

$$ds_z = \left| d\vec{l}_x \times d\vec{l}_y \right| = dxdy$$

*Note that a surface has a direction normal to the surface.

A differential volume:

$$dv = dxdydz = (d\vec{l}_x \times d\vec{l}_y) \cdot d\vec{l}_z = (d\vec{l}_y \times d\vec{l}_z) \cdot d\vec{l}_x = (d\vec{l}_z \times d\vec{l}_x) \cdot d\vec{l}_y$$

Given $\vec{A} = \hat{a}_x A_x + \hat{a}_y A_y + \hat{a}_z A_z$ and $\vec{B} = \hat{a}_x B_x + \hat{a}_y B_y + \hat{a}_z B_z$

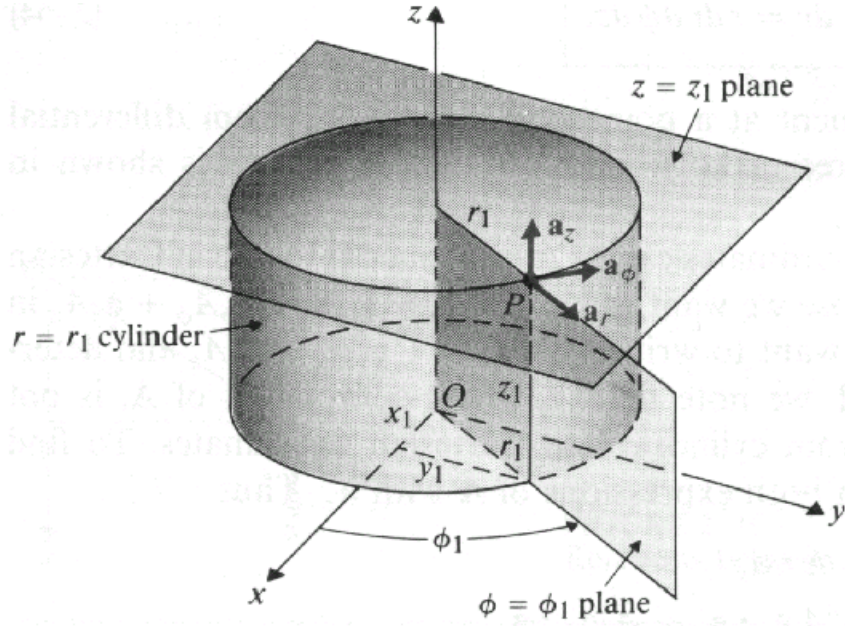
Vector scalar product:

$$\vec{A} \cdot \vec{B} = A_x B_x + A_y B_y + A_z B_z$$

Vector cross product

$$\vec{A} \times \vec{B} = \begin{vmatrix} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix}$$

Cylindrical Coordinates (r, ϕ, z): suitable for problems with cylindrical symmetry.



Radius distance: $0 \leq r \leq \infty$, azimuthal angle: $0 \leq \phi < 2\pi$, longitudinal coordinate: $-\infty \leq z \leq \infty$. A general expression of a vector:

$$\vec{A} = \hat{a}_r A_r + \hat{a}_\phi A_\phi + \hat{a}_z A_z$$

The three unit vectors, $\hat{a}_r, \hat{a}_\phi, \hat{a}_z$ are arranged such that

$$\hat{a}_r \times \hat{a}_\phi = \hat{a}_z, \quad \hat{a}_\phi \times \hat{a}_z = \hat{a}_r, \quad \hat{a}_z \times \hat{a}_r = \hat{a}_\phi, \quad \hat{a}_r \cdot \hat{a}_\phi = 0,$$

$$\hat{a}_\phi \cdot \hat{a}_z = 0, \quad \hat{a}_z \cdot \hat{a}_r = 0$$

A differential length:

$$d\vec{l} = \hat{a}_r dr + \hat{a}_\phi r d\phi + \hat{a}_z dz = d\vec{l}_r + d\vec{l}_\phi + d\vec{l}_z$$

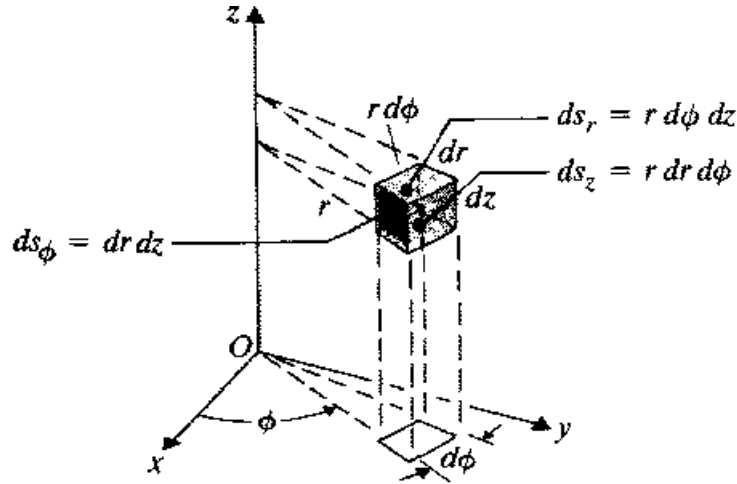
A differential surface:

$$\begin{aligned} d\vec{s} &= \hat{a}_r ds_r + \hat{a}_\phi ds_\phi + \hat{a}_z ds_z \\ &= d\vec{l}_\phi \times d\vec{l}_z + d\vec{l}_z \times d\vec{l}_r + d\vec{l}_r \times d\vec{l}_\phi, \end{aligned}$$

where $ds_r = r d\phi dz$, $ds_\phi = dr dz$, $ds_z = r dr d\phi$.

A differential volume:

$$dv = r dr d\phi dz = (\vec{dl}_r \times \vec{dl}_\phi) \cdot \vec{dl}_z = (\vec{dl}_\phi \times \vec{dl}_z) \cdot \vec{dl}_r = (\vec{dl}_z \times \vec{dl}_r) \cdot \vec{dl}_\phi$$



Given $\vec{A} = \hat{a}_r A_r + \hat{a}_\phi A_\phi + \hat{a}_z A_z$ and $\vec{B} = \hat{a}_r B_r + \hat{a}_\phi B_\phi + \hat{a}_z B_z$

Vector scalar product:

$$\vec{A} \cdot \vec{B} = A_r B_r + A_\phi B_\phi + A_z B_z$$

Vector cross product:

$$\vec{A} \times \vec{B} = \begin{vmatrix} \hat{a}_r & \hat{a}_\phi & \hat{a}_z \\ A_r & A_\phi & A_z \\ B_r & B_\phi & B_z \end{vmatrix}$$

Coordinate Transformation between the Cartesian and Cylindrical coordinates

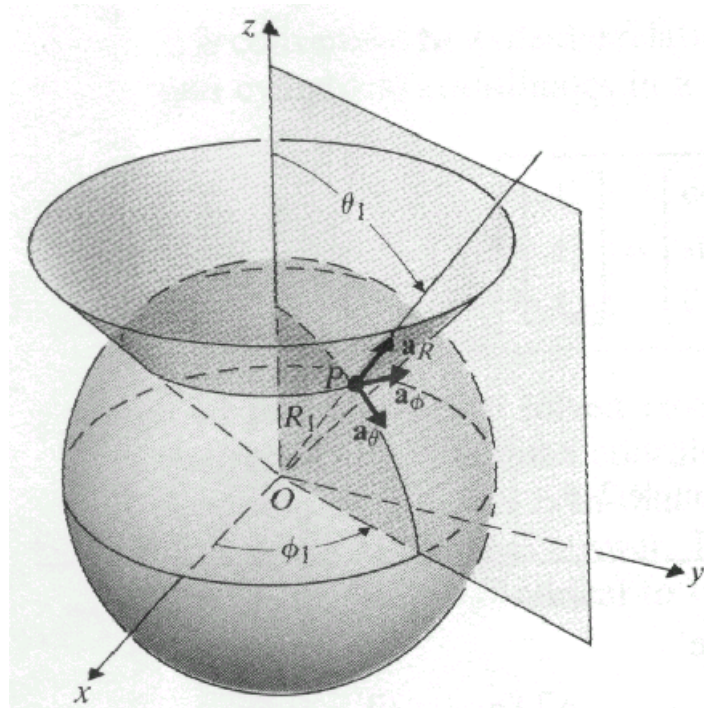
$$\begin{bmatrix} A_x \\ A_y \\ A_z \end{bmatrix} = \begin{bmatrix} \cos \phi & -\sin \phi & 0 \\ \sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} A_r \\ A_\phi \\ A_z \end{bmatrix}$$

and

$$x = r \cos \phi, y = r \sin \phi, z = z$$

$$r = \sqrt{x^2 + y^2}, \phi = \tan^{-1} \frac{y}{x}, z = z$$

Spherical Coordinates (R, θ, ϕ): suitable for problems with spherical symmetry. Radius distance: $0 \leq R \leq \infty$, colatitude angle: $0 \leq \theta \leq \pi$, azimuthal angle: $0 \leq \phi < 2\pi$.



The three unit vectors, $\hat{a}_R, \hat{a}_\theta, \hat{a}_\phi$ are arranged such that

$$\hat{a}_R \times \hat{a}_\theta = \hat{a}_\phi, \quad \hat{a}_\theta \times \hat{a}_\phi = \hat{a}_R, \quad \hat{a}_\phi \times \hat{a}_R = \hat{a}_\theta, \quad \hat{a}_R \cdot \hat{a}_\theta = 0,$$

$$\hat{a}_\theta \cdot \hat{a}_\phi = 0, \quad \hat{a}_\phi \cdot \hat{a}_R = 0$$

A general expression of a vector

$$\vec{A} = \hat{a}_R A_R + \hat{a}_\theta A_\theta + \hat{a}_\phi A_\phi$$

A differential length:

$$d\vec{l} = \hat{a}_R dR + \hat{a}_\theta R d\theta + \hat{a}_\phi R \sin \theta d\phi = d\vec{l}_R + d\vec{l}_\theta + d\vec{l}_\phi$$

differential surface:

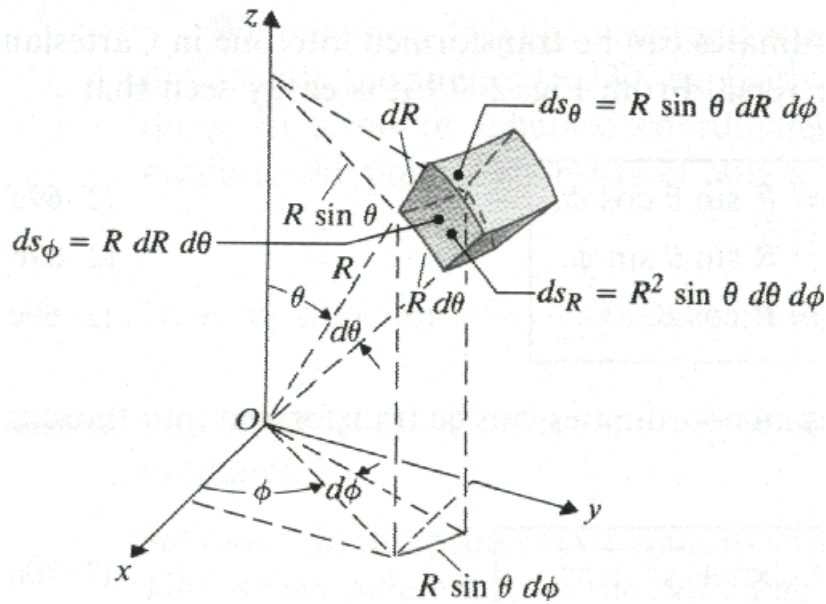
$$d\vec{s} = d\vec{s}_R + d\vec{s}_\theta + d\vec{s}_\phi,$$

where

$$ds_R = R^2 \sin \theta d\theta d\phi, \quad ds_\theta = R \sin \theta dR d\phi, \quad ds_\phi = R dR d\theta$$

A differential volume:

$$dv = R^2 \sin \theta dR d\theta d\phi$$



Given $\vec{A} = \hat{a}_R A_R + \hat{a}_\theta A_\theta + \hat{a}_\phi A_\phi$ and $\vec{B} = \hat{a}_R B_R + \hat{a}_\theta B_\theta + \hat{a}_\phi B_\phi$

Vector dot product:

$$\vec{A} \cdot \vec{B} = A_R B_R + A_\theta B_\theta + A_\phi B_\phi$$

Vector cross product:

$$\vec{A} \times \vec{B} = \begin{vmatrix} \hat{a}_R & \hat{a}_\theta & \hat{a}_\phi \\ A_R & A_\theta & A_\phi \\ B_R & B_\theta & B_\phi \end{vmatrix}$$

Coordinate Transformation between the Cartesian and spherical coordinates

$$\begin{bmatrix} A_x \\ A_y \\ A_z \end{bmatrix} = \begin{bmatrix} \sin \theta \cos \phi & \cos \theta \cos \phi & -\sin \phi \\ \sin \theta \sin \phi & \cos \theta \sin \phi & \cos \phi \\ \cos \theta & -\sin \theta & 0 \end{bmatrix} \begin{bmatrix} A_R \\ A_\theta \\ A_\phi \end{bmatrix}$$

$$x = R \sin \theta \cos \phi, y = R \sin \theta \sin \phi, z = R \cos \theta$$

$$R = \sqrt{x^2 + y^2 + z^2}, \theta = \tan^{-1} \frac{\sqrt{x^2 + y^2}}{z}$$

$$\phi = \tan^{-1} \frac{y}{x}$$

General Form

In an orthogonal coordinate system, the three unit vectors can be expressed by $\hat{a}_{u_1}, \hat{a}_{u_2}, \hat{a}_{u_3}$ are arranged such that

$$\hat{a}_{u_1} \times \hat{a}_{u_2} = \hat{a}_{u_3}, \quad \hat{a}_{u_2} \times \hat{a}_{u_3} = \hat{a}_{u_1}, \quad \hat{a}_{u_3} \times \hat{a}_{u_1} = \hat{a}_{u_2},$$

$$\hat{a}_{u_i} \cdot \hat{a}_{u_j} = 0 \text{ for } i, j = 1, 2, 3$$

A general expression of a vector

$$\vec{A} = \hat{a}_{u_1} A_{u_1} + \hat{a}_{u_2} A_{u_2} + \hat{a}_{u_3} A_{u_3}$$

A differential length:

$$d\vec{l} = \hat{a}_{u_1} h_1 du_1 + \hat{a}_{u_2} h_2 du_2 + \hat{a}_{u_3} h_3 du_3 = d\vec{l}_{u_1} + d\vec{l}_{u_2} + d\vec{l}_{u_3}$$

where h_1, h_2, h_3 are called *metric coefficients*.

A differential surface:

$$d\vec{s} = d\vec{s}_{u_1} + d\vec{s}_{u_2} + d\vec{s}_{u_3} ,$$

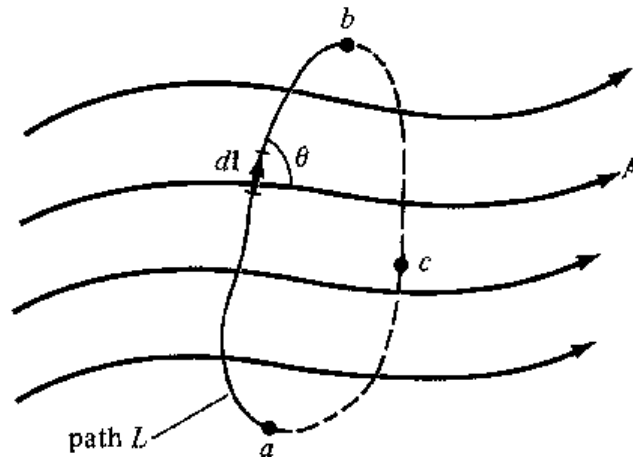
coordinate system relations	Cartesian Coor. (x, y, z)	Cylindrical Coor. (r, ϕ , z)	Spherical Coor. (R, θ , ϕ)
Base vectors \hat{a}_{u_1} \hat{a}_{u_2} \hat{a}_{u_3}	\hat{a}_x \hat{a}_y \hat{a}_z	\hat{a}_r \hat{a}_ϕ \hat{a}_z	\hat{a}_R \hat{a}_θ \hat{a}_ϕ
Metric coefficient h_1 h_2 h_3	1 1 1	1 r 1	1 R $R \sin \theta$

Vector Calculus

Line Integral

$$\int_L \vec{F} \cdot d\vec{l} = \int_a^b F \cos \theta dl : \text{ the integral of vector } \vec{F} \text{ along path } L.$$

$$\oint_L \vec{F} \cdot d\vec{l} : \text{ integral of } \vec{F} \text{ along a closed path } L. \text{ Note that a closed path defines a surface.}$$

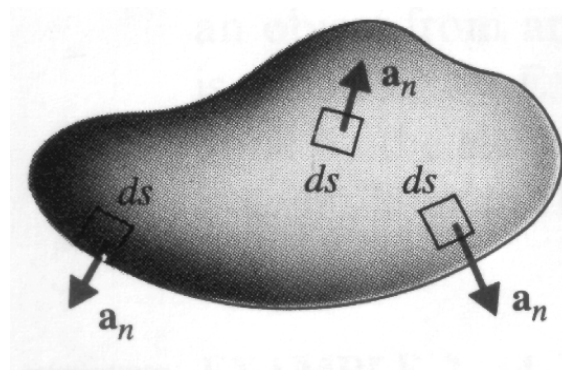


Surface Integral

$$\int_S \vec{A} \cdot d\vec{s} = \int_S A \cos \theta ds = \int_S \vec{A} \cdot \hat{a}_n ds : \text{the integral of } \vec{A} \text{ across an}$$

open surface S , or the *flux* of \vec{A} through S .

$\oint_S \vec{A} \cdot d\vec{s}$: the integral of \vec{A} across a closed surface S . Note that a closed surface defines a volume. The direction of a closed surface is defined outward a volume.

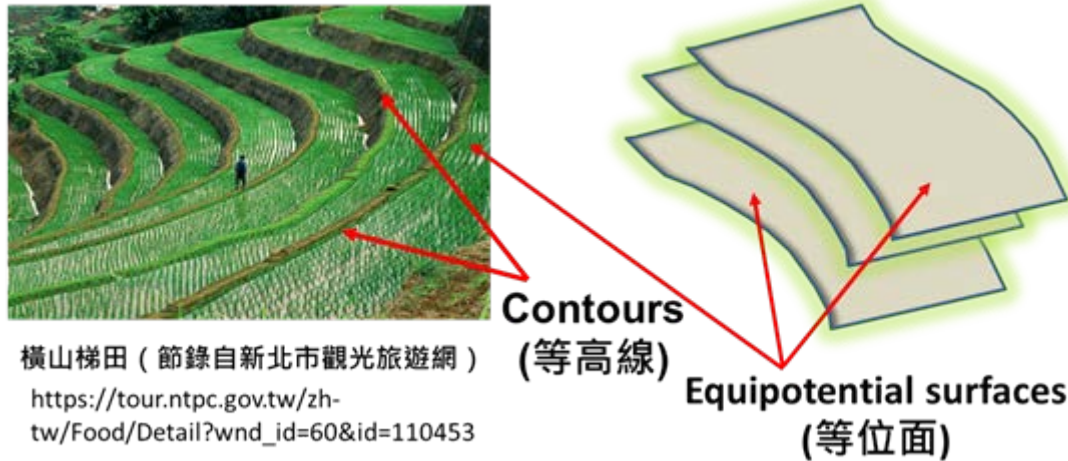


Volume Integral

$$\int_V \rho_v dv : \text{the integral of the function } \rho_v \text{ over a volume } V.$$

Gradient of a Scalar : a vector having a magnitude equal to the

maximum rate of change of a scalar in space, and a direction along the maximum change.



Consider the directional derivative of a scalar function V along an arbitrary direction \hat{a}_l

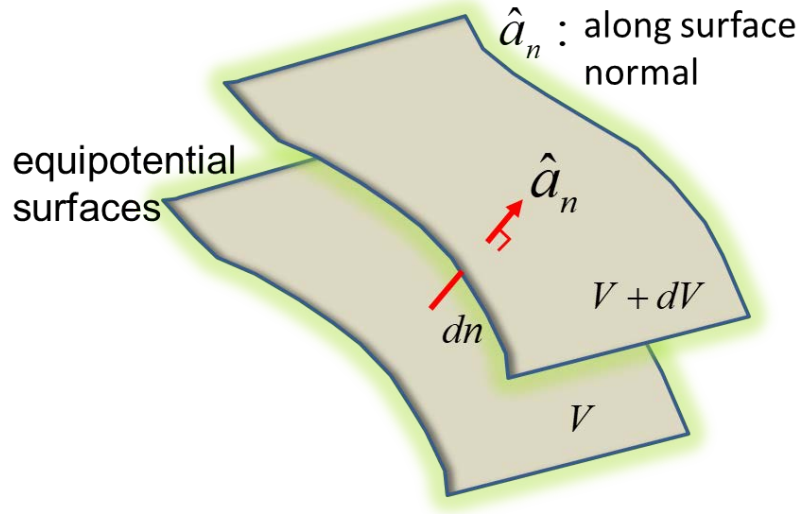
$$\frac{dV}{dl} = \frac{dV}{dn} \frac{dn}{dl} = \frac{dV}{dn} \cos \alpha = \frac{dV}{dn} \hat{a}_n \cdot \hat{a}_l,$$

where \hat{a}_n is a unit vector normal to the $V = \text{constant}$ surface. Refer to the following plot, $\hat{a}_n \cdot \hat{a}_l = \cos \alpha$ has a value less than 1.

Therefore $dV/dn \geq dV/dl \Rightarrow$ the gradient of a scalar can be defined as

$$\nabla V \equiv \hat{a}_n \frac{dV}{dn},$$

and the directional derivative can be rewritten as $\frac{dV}{dl} = (\nabla V) \cdot \hat{a}_l$



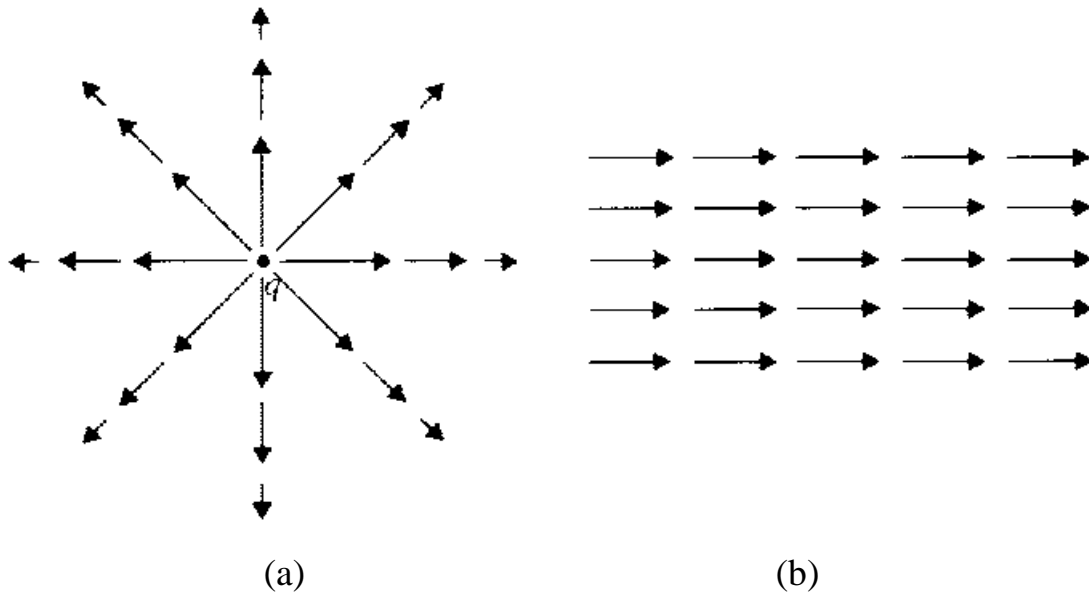
It follows that $dV = (\nabla V) \cdot d\vec{l}$ or

$$\begin{aligned}
 dV &= \frac{\partial V}{\partial l_{u_1}} dl_{u_1} + \frac{\partial V}{\partial l_{u_2}} dl_{u_2} + \frac{\partial V}{\partial l_{u_3}} dl_{u_3} \\
 &= \left(\frac{\partial V}{\partial l_{u_1}} \hat{a}_{u_1} + \frac{\partial V}{\partial l_{u_2}} \hat{a}_{u_2} + \frac{\partial V}{\partial l_{u_3}} \hat{a}_{u_3} \right) \cdot (dl_{u_1} \hat{a}_{u_1} + dl_{u_2} \hat{a}_{u_2} + dl_{u_3} \hat{a}_{u_3}) \\
 &= \left(\frac{\partial V}{\partial l_{u_1}} \hat{a}_{u_1} + \frac{\partial V}{\partial l_{u_2}} \hat{a}_{u_2} + \frac{\partial V}{\partial l_{u_3}} \hat{a}_{u_3} \right) \cdot d\vec{l} \\
 \Rightarrow \nabla V &= \left(\frac{\partial V}{\partial l_{u_1}} \hat{a}_{u_1} + \frac{\partial V}{\partial l_{u_2}} \hat{a}_{u_2} + \frac{\partial V}{\partial l_{u_3}} \hat{a}_{u_3} \right)
 \end{aligned}$$

In the Cartesian coordinate system, the gradient of a scalar V can be found to be

$$\nabla V = \frac{\partial V}{\partial x} \hat{a}_x + \frac{\partial V}{\partial y} \hat{a}_y + \frac{\partial V}{\partial z} \hat{a}_z$$

Divergence of a Vector Field



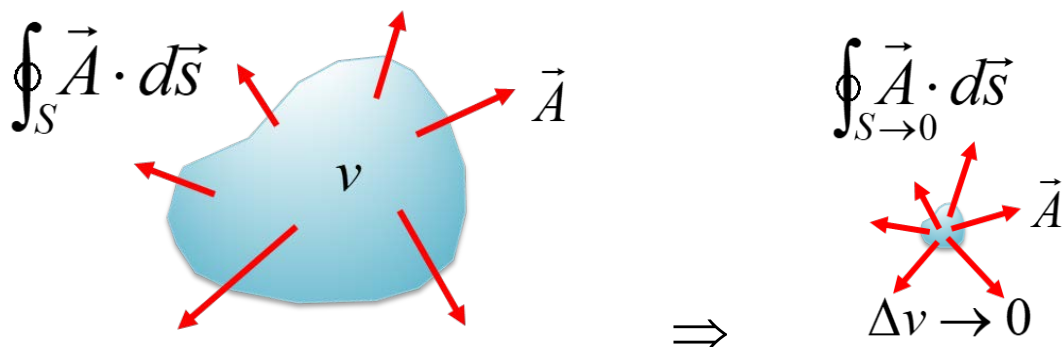
(a) A net outward flux surrounding q .

(b) An uniform flux of fields to the right.

Divergence of \vec{A} : $\nabla \cdot \vec{A} \equiv \lim_{\Delta v \rightarrow 0} \frac{\oint_S \vec{A} \cdot d\vec{s}}{\Delta v}$ is a scalar equal to the

net outward flux of \vec{A} per unit volume at a point in space. Therefore

$\nabla \cdot \vec{A}$ is a point function that describes the aforementioned physical quantity at a point location.



$$\Rightarrow \text{Divergence Theorem} \quad \int_V \nabla \cdot \vec{A} dv \equiv \oint_S \vec{A} \cdot d\vec{s}$$

Divergence Theorem

Partition V into many small v_i surrounded by small surfaces s_i

$\int_V \nabla \cdot \vec{A} dv \equiv \oint_S \vec{A} \cdot d\vec{s}$
 S : surface enclosing V

$\int_V \nabla \cdot \vec{A} dv \equiv \sum_i \lim_{\Delta v \rightarrow 0} \frac{\oint_{s_i} \vec{A}_i \cdot d\vec{s}}{\Delta v} \Delta v$

$= \oint_S \vec{A} \cdot d\vec{s}$

fluxes are summed to zero

Uncompensated surface

• In x, y, z coordinate system, $\nabla \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$

In the Cartesian coordinate system, the divergence of a vector

$\vec{A} = \hat{a}_x A_x + \hat{a}_y A_y + \hat{a}_z A_z$ is expressed by

$$\nabla \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$$

For a vector satisfying $\nabla \cdot \vec{B} = 0$, vector \vec{B} is said to be solenoidal.

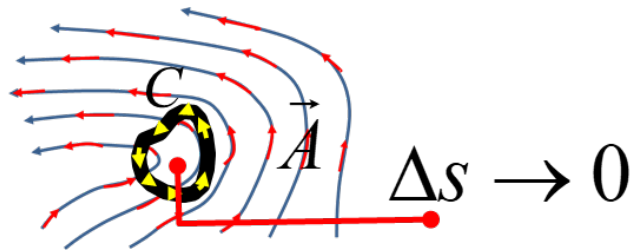
uniform flux of fields to the right

$-Bds + Bds$

$\oint_S \vec{B} \cdot d\vec{s} = 0$

$\nabla \cdot \vec{B} = 0$ (solenoidal field)

Curl of a Vector Field



$$\nabla \times \vec{A} \equiv \lim_{\Delta S \rightarrow 0} \frac{\hat{a}_n \oint_C \vec{A} \cdot d\vec{l}}{\Delta S} : \text{maximum net circulation of } \vec{A} \text{ per}$$

unit area at a point in space. The direction of $\nabla \times \vec{A}$ is chosen to be the surface normal direction of the infinitesimal area ΔS with which the net circulation is a maximum. Therefore $\nabla \times \vec{A}$ is also a point function that describes the aforementioned physical quantity at a point location.

\Rightarrow **Stoke's Theorem** $\int_S (\nabla \times \vec{A}) \cdot d\vec{s} \equiv \oint_C \vec{A} \cdot d\vec{l}$, where C is the path surrounding the surface S .

Stokes' Theorem $\int_S (\nabla \times \vec{A}) \cdot d\vec{s} \equiv \oint_C \vec{A} \cdot d\vec{l}$
 C : path surrounding surface S .

Partition S into many small s_i surrounded by small path c_i

$$\int_S \nabla \times \vec{A} \cdot d\vec{s} \rightarrow \sum_i \int_{s_i} \nabla \times \vec{A}_i \cdot d\vec{s} \equiv \sum_i \lim_{\Delta s_i \rightarrow 0} \frac{\oint_{c_i} \vec{A}_i \cdot d\vec{l}}{\Delta s_i} \Delta s_i$$

Line integrals are summed to zero

Uncompensated path

$$= \oint_C \vec{A} \cdot d\vec{l}$$

For a vector satisfying $\nabla \times \vec{A} = 0$, vector \vec{A} is said to be irrotational.

In the Cartesian coordinate system, the curl of a vector

$\vec{A} = \hat{a}_x A_x + \hat{a}_y A_y + \hat{a}_z A_z$ is expressed by

$$\begin{aligned}\nabla \times \vec{A} &= \begin{vmatrix} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{vmatrix} \\ &= \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}\right)\hat{a}_x + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x}\right)\hat{a}_y + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y}\right)\hat{a}_z\end{aligned}$$

Laplacian operator of a scalar field $\nabla^2 V \equiv \nabla \cdot (\nabla V)$

In the xyz coordinate system, the expression is given by

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$

Laplacian operator of a vector field

$$\nabla^2 \vec{A} \equiv \nabla(\nabla \cdot \vec{A}) - \nabla \times \nabla \times \vec{A}$$

In the xyz coordinate system, the expression is given by

$$\nabla^2 \vec{A} = \hat{a}_x \nabla^2 A_x + \hat{a}_y \nabla^2 A_y + \hat{a}_z \nabla^2 A_z$$

Two Null Identities

$\nabla \times (\nabla V) = 0$, no net circulation around a vector normal to an equipotential surface.

$\nabla \cdot (\nabla \times \vec{A}) = 0$, no net outward flux around the maximum circulation of a vector.

A Quick Reference

$$\begin{aligned}\nabla V &= \hat{a}_{u_1} \frac{\partial \mathcal{V}}{h_1 \partial u_1} + \hat{a}_{u_2} \frac{\partial \mathcal{V}}{h_2 \partial u_2} + \hat{a}_{u_3} \frac{\partial \mathcal{V}}{h_3 \partial u_3} \\ \nabla \cdot \vec{A} &= \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (h_2 h_3 A_1) + \frac{\partial}{\partial u_2} (h_1 h_3 A_2) + \frac{\partial}{\partial u_3} (h_1 h_2 A_3) \right] \\ \nabla \times \vec{A} &= \frac{1}{h_1 h_2 h_3} \begin{vmatrix} \hat{a}_{u_1} h_1 & \hat{a}_{u_2} h_2 & \hat{a}_{u_3} h_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ h_1 A_1 & h_2 A_2 & h_3 A_3 \end{vmatrix}\end{aligned}$$

Helmholtz's Theorem:

A vector field is determined to within an additive constant if both its divergence and its curl are specified everywhere.

A similar statement is “if two vector fields have the same values for their curl, divergence, and surface dot products at a boundary, the two vectors fields are the same within an additive constant.”

In mathematics, the *Helmholtz's Theorem* is equivalent to: If

$$(i) \nabla \cdot \vec{A} = \nabla \cdot \vec{B}, (ii) \nabla \times \vec{A} = \nabla \times \vec{B}, \text{ and } (iii) \vec{A} \cdot d\vec{s} = \vec{B} \cdot d\vec{s}$$

on the surface surrounding the volume in question, then $\vec{A} = \vec{B} +$ a constant vector.

Lemma: Green's first identity

The divergence theory, $\int_V \nabla \cdot \vec{G} dv = \oint_S \vec{G} \cdot d\vec{s}$, where \vec{G} is any

vector field in space. Define $\vec{G} = \phi \nabla \psi$ and substitute into

$$\int_V \nabla \cdot \vec{A} dv = \oint_S \vec{A} \cdot d\vec{s}, \text{ with } \nabla \cdot (\phi \nabla \psi) = \nabla \phi \cdot \nabla \psi + \phi \nabla^2 \psi.$$

One has

$\int_V (\nabla \phi \cdot \nabla \psi + \phi \nabla^2 \psi) dv = \oint_S \phi \nabla \psi \cdot d\vec{s}$, which is called *Green's first identity*. For the special case $\phi = \psi$, Green's first identity becomes

$$\int_V |\nabla \phi|^2 dv + \int_V \phi \nabla^2 \phi dv = \oint_S \phi \nabla \phi \cdot d\vec{s} \quad (*)$$

Proof of Helmholtz's Theorem:

Define $\vec{C} = \vec{A} - \vec{B}$.

a. $\nabla \times \vec{C} = \nabla \times \vec{A} - \nabla \times \vec{B} = 0 \Rightarrow \vec{C} = -\nabla \phi$ (the minus sign is immaterial)

b. $\nabla \cdot \vec{C} = \nabla \cdot \vec{A} - \nabla \cdot \vec{B} = 0 \Rightarrow \nabla \cdot \vec{C} = -\nabla^2 \phi = 0$

c. $\vec{C} \cdot d\vec{s} = \vec{A} \cdot d\vec{s} - \vec{B} \cdot d\vec{s} = 0 \Rightarrow \vec{C} \cdot d\vec{s} = -\nabla \phi \cdot d\vec{s} = 0$

Substitute a-c into (*) $\Rightarrow \int_V |\nabla \phi|^2 dv = 0$. But $|\nabla \phi|^2 > 0$ and therefore $\vec{C} = \vec{A} - \vec{B} = -\nabla \phi = 0 \Rightarrow \vec{A} = \vec{B}$. Note that if one adds a constant to $\vec{C} = -\nabla \phi$, it does not change the final answer. We thus prove the Helmholtz theorem. A simple statement for the theorem is that there exists a unique solution for the vector field in a space if its divergence, curl, and boundary values are specified uniquely. Therefore in electromagnetics the divergence and curl of a field are specified by so-called Maxwell's equations and the boundary condition of a field is specified in a real-world problem.

The following two are extracted from S. Ramo, J. Whinnery, and T. van Duzer, *Fields and Waves I Communication Electronics*, John Wiley & Sons.

CHAPTER 2 Review on Vector Algebra and Calculus

VECTOR DIFFERENTIAL OPERATIONS

Rectangular Coordinates (x, y, z)

$$\begin{aligned}\nabla\Phi &= \hat{\mathbf{x}} \frac{\partial\Phi}{\partial x} + \hat{\mathbf{y}} \frac{\partial\Phi}{\partial y} + \hat{\mathbf{z}} \frac{\partial\Phi}{\partial z} \\ \nabla \cdot \mathbf{D} &= \frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} + \frac{\partial D_z}{\partial z} \\ \nabla \times \mathbf{H} &= \hat{\mathbf{x}} \left(\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} \right) + \hat{\mathbf{y}} \left(\frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} \right) + \hat{\mathbf{z}} \left(\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right) \\ \nabla^2\Phi &= \frac{\partial^2\Phi}{\partial x^2} + \frac{\partial^2\Phi}{\partial y^2} + \frac{\partial^2\Phi}{\partial z^2} \\ \nabla^2\mathbf{A} &= \hat{\mathbf{x}}\nabla^2 A_x + \hat{\mathbf{y}}\nabla^2 A_y + \hat{\mathbf{z}}\nabla^2 A_z\end{aligned}$$

Cylindrical Coordinates (r, ϕ, z)

$$\begin{aligned}\nabla\Phi &= \hat{\mathbf{r}} \frac{\partial\Phi}{\partial r} + \hat{\boldsymbol{\phi}} \frac{1}{r} \frac{\partial\Phi}{\partial\phi} + \hat{\mathbf{z}} \frac{\partial\Phi}{\partial z} \\ \nabla \cdot \mathbf{D} &= \frac{1}{r} \frac{\partial}{\partial r} (rD_r) + \frac{1}{r} \frac{\partial D_\phi}{\partial\phi} + \frac{\partial D_z}{\partial z} \\ \nabla \times \mathbf{H} &= \hat{\mathbf{r}} \left[\frac{1}{r} \frac{\partial H_z}{\partial\phi} - \frac{\partial H_\phi}{\partial z} \right] + \hat{\boldsymbol{\phi}} \left[\frac{\partial H_r}{\partial z} - \frac{\partial H_z}{\partial r} \right] + \hat{\mathbf{z}} \left[\frac{1}{r} \frac{\partial(rH_\phi)}{\partial r} - \frac{1}{r} \frac{\partial H_r}{\partial\phi} \right] \\ \nabla^2\Phi &= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial\Phi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2\Phi}{\partial\phi^2} + \frac{\partial^2\Phi}{\partial z^2} \\ \nabla^2\mathbf{A} &= \hat{\mathbf{r}} \left(\nabla^2 A_r - \frac{2}{r^2} \frac{\partial A_\phi}{\partial\phi} - \frac{A_r}{r^2} \right) + \hat{\boldsymbol{\phi}} \left(\nabla^2 A_\phi + \frac{2}{r^2} \frac{\partial A_r}{\partial\phi} - \frac{A_\phi}{r^2} \right) + \hat{\mathbf{z}} (\nabla^2 A_z)\end{aligned}$$

Spherical Coordinates (r, θ, ϕ)

$$\begin{aligned}\nabla\Phi &= \hat{\mathbf{r}} \frac{\partial\Phi}{\partial r} + \hat{\boldsymbol{\theta}} \frac{1}{r} \frac{\partial\Phi}{\partial\theta} + \frac{\hat{\boldsymbol{\phi}}}{r \sin\theta} \frac{\partial\Phi}{\partial\phi} \\ \nabla \cdot \mathbf{D} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 D_r) + \frac{1}{r \sin\theta} \frac{\partial}{\partial\theta} (\sin\theta D_\theta) + \frac{1}{r \sin\theta} \frac{\partial D_\phi}{\partial\phi} \\ \nabla \times \mathbf{H} &= \frac{\hat{\mathbf{r}}}{r \sin\theta} \left[\frac{\partial}{\partial\theta} (H_\phi \sin\theta) - \frac{\partial H_\theta}{\partial\phi} \right] \\ &\quad + \frac{\hat{\boldsymbol{\theta}}}{r} \left[\frac{1}{\sin\theta} \frac{\partial H_r}{\partial\phi} - \frac{\partial}{\partial r} (rH_\phi) \right] + \frac{\hat{\boldsymbol{\phi}}}{r} \left[\frac{\partial}{\partial r} (rH_\theta) - \frac{\partial H_r}{\partial\theta} \right] \\ \nabla^2\Phi &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial\Phi}{\partial r} \right) + \frac{1}{r^2 \sin\theta} \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial\Phi}{\partial\theta} \right) + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2\Phi}{\partial\phi^2} \\ \nabla^2\mathbf{A} &= \hat{\mathbf{r}} \left[\nabla^2 A_r - \frac{2}{r^2} \left(A_r + \cot\theta A_\theta + \csc\theta \frac{\partial A_\phi}{\partial\phi} + \frac{\partial A_\theta}{\partial\theta} \right) \right] \\ &\quad + \hat{\boldsymbol{\theta}} \left[\nabla^2 A_\theta - \frac{1}{r^2} \left(\csc^2\theta A_\theta - 2 \frac{\partial A_r}{\partial\theta} + 2 \cot\theta \csc\theta \frac{\partial A_\phi}{\partial\phi} \right) \right] \\ &\quad + \hat{\boldsymbol{\phi}} \left[\nabla^2 A_\phi - \frac{1}{r^2} \left(\csc^2\theta A_\phi - 2 \csc\theta \frac{\partial A_r}{\partial\phi} - 2 \cot\theta \csc\theta \frac{\partial A_\theta}{\partial\phi} \right) \right]\end{aligned}$$

VECTOR FORMULAS

$$\nabla(\Phi + \psi) = \nabla\Phi + \nabla\psi$$

$$\nabla \cdot (\mathbf{A} + \mathbf{B}) = \nabla \cdot \mathbf{A} + \nabla \cdot \mathbf{B}$$

$$\nabla \times (\mathbf{A} + \mathbf{B}) = \nabla \times \mathbf{A} + \nabla \times \mathbf{B}$$

$$\nabla(\Phi\psi) = \Phi\nabla\psi + \psi\nabla\Phi$$

$$\nabla \cdot (\psi\mathbf{A}) = \mathbf{A} \cdot \nabla\psi + \psi\nabla \cdot \mathbf{A}$$

$$\nabla \cdot (\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot \nabla \times \mathbf{A} - \mathbf{A} \cdot \nabla \times \mathbf{B}$$

$$\nabla \times (\Phi\mathbf{A}) = \nabla\Phi \times \mathbf{A} + \Phi\nabla \times \mathbf{A}$$

$$\nabla \times (\mathbf{A} \times \mathbf{B}) = \mathbf{A}\nabla \cdot \mathbf{B} - \mathbf{B}\nabla \cdot \mathbf{A} + (\mathbf{B} \cdot \nabla)\mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B}$$

$$\nabla \cdot \nabla\Phi = \nabla^2\Phi$$

$$\nabla \cdot \nabla \times \mathbf{A} = 0$$

$$\nabla \times \nabla\Phi = 0$$

$$\nabla \times \nabla \times \mathbf{A} = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2\mathbf{A}$$

$$\nabla(\mathbf{A} \cdot \mathbf{B}) = (\mathbf{A} \cdot \nabla)\mathbf{B} + (\mathbf{B} \cdot \nabla)\mathbf{A} + \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A})$$

$$\mathbf{A} \cdot \mathbf{B} \times \mathbf{C} = \mathbf{B} \cdot \mathbf{C} \times \mathbf{A} = \mathbf{C} \cdot \mathbf{A} \times \mathbf{B}$$

$$\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) = \mathbf{B}(\mathbf{A} \cdot \mathbf{C}) - \mathbf{C}(\mathbf{A} \cdot \mathbf{B})$$