Step4: Apply B.c. &1.c. to delete unsatisfying terms From B.C. (D&@ (I(0)=0, I(L)=0), we know I(x) should be sinusoidal type: B.c.020 I(x) = c/coskx + casinkx => I(x) = Ansin Zx, (K= Z), n=1,2,3,...  $T(t) = C_3 \cos akt + c_4 \sin akt \Rightarrow T(t) = B_n \cos \frac{n\pi a}{L} + C_4 \sin \frac{n\pi a}{L} + C_5 \sin \frac{n\pi a}{L} + C_6 \sin \frac{n\pi$ u(x,t) = I(x) T(t) = (An sin = x) (Bn as Tt + Cn sin = t) n=1,2,3... So from B.C. Ol Q, we obtain an infinite number of solutions (one for each n) as u(x,t) = (Dn cos Lt + Ensin Lt) sin Lx (the coefficients are combined as Dn. Enfor simplicity) Step 5: Use superposition and proper conditions to get the solution as a Fourier sevies Although none of hexit) obtained in step 4 satisfies the 1.C. 324 it is possible to combine them as an infinite series that does. > U(x,t) = \( D\_n cos \( \frac{1}{L} t + E\_n \sin \( \frac{1}{L} t \) \( \sin \frac{1}{L} \times \) From 1.c. 3 (u(x,0) = f(x))

From 1.c. (3) (u(x, 0) = f(x))  $u(x, 0) = \sum_{n=1}^{\infty} D_n \sin \frac{nz}{L} x = f(x)$  (Note: f(x) is expanded by a founier sine series) Dn can be found as how we find bn coefficient  $Dn = \frac{2}{L} \int_0^L f(x) sf(n) \frac{dx}{L} x dx$ From 1.c. (4)  $\left(\frac{\partial u}{\partial x}\right|_{x=0}^{\infty} g(x)$ 

From 1. c. (a)  $\left(\frac{\partial u}{\partial t}\Big|_{t=0} = g(x)\right)$   $\frac{\partial u}{\partial t} = \sum_{n=1}^{\infty} \left(-D_n \frac{n\pi a}{L} \sin \frac{n\pi a}{L} t + E_n \frac{n\pi a}{L} \cos \frac{n\pi a}{L} t\right) \sin \frac{n\pi}{L} \times$ At t=0  $\frac{\partial u}{\partial t}\Big|_{t=0} = \sum_{n=1}^{\infty} \left(E_n \frac{n\pi a}{L}\right) \sin \frac{n\pi}{L} \times = g(x) \quad \text{(Note: g(x) is expanded by a former sine series)}$ 

 $E_{n} \frac{nza}{L} = \frac{2}{L} \int_{0}^{L} g(x) \sin \frac{nz}{L} x dx \Rightarrow E_{n} = \frac{2}{nza} \int_{0}^{L} g(x) \sin \frac{nz}{L} x dx$ 

Final solution is u(x,t) = \(\sigma\) (Du cos \(\frac{12a}{2}\) t+ En sin \(\frac{12a}{2}\) sin \(\frac{12a}{2}\) x, where  $D_n = \frac{2}{2} \int_0^1 f(x) \sin(\frac{12a}{2}x) dx, \quad E_n = \frac{2}{n2a} \int_0^1 g(x) \sin(\frac{12a}{2}x) dx$ 

Remarks:

- 1) When using sov to solve PDEs, The trick of assuming solution as the
- 2 In the examples, we repeatedly encountered the BVP

Such Brp actually belongs to a more general category of Brps, called

# Sturm - Lionville boundary-value publem

General form:

This general form can also be rewritten in a "self-adjoint form"

subject to B.c.

A, 
$$y(a) + B$$
,  $y'(a) = 0$   
 $A_2 y(b) + B_2 y'(cb) = 0$ 

Sturm-Lionville equation subject to the B.c. is referred to the "Sturm-Lionville boundary-value problem"

Properties of Starm-Lionville Bup:

There exists an infinite number of

- For each eigenvalue, there is one
- Eigenfunctions are s

Remarks:

- O Any homogeneous 2nd order linear DE can be rewritten
- 2 Because the eigenfunctions form an orthogonal set, a given function fix) can be expanded by

Many DEs in engineering/physics are and can be rewritten as a Stam-Liouville equation. When solving these DEs with B.C., they all have

Example I: Legendre's equation (1-x2)y"-2xy+n(u+1)y=0,+(xx)

- 1) has the
- 2) eigen values
- 3) eigenfunctions

4) A function f(x) can be expanded in terms of Legeldre polynomials as:  $f(x) = P_n(x)$ 

Practice: Use Legendre polynomials to expand a function fex)

## Example I : Hermite's equation y'-2xy'+2ny =0

- i) has the
- 2) eigenvalues
- 3) eigenfunctions
- 4) A function can be expanded in terms of Hermite polynomials

Example II: BVP we encountered when solving Laplace's eq 2 wave eq

- 1) has the
- 2) eigenvalues
- 3) eigenfunctions
- 4) A function fex) can be expanded in terms of

More examples,

We already learned (in ch 6) how to find the solution (eigenfunction) for some important special DEs. In the following, we will do more practices on solving the most basic yet important S-L DVP:

 $\frac{E_{XI}}{y''+\lambda y=0}$ , y'(0)=0, y'(1)=0

Eyz: Same with Ex, but with B.c. y(0) = 0, y(1)+y(1)=0