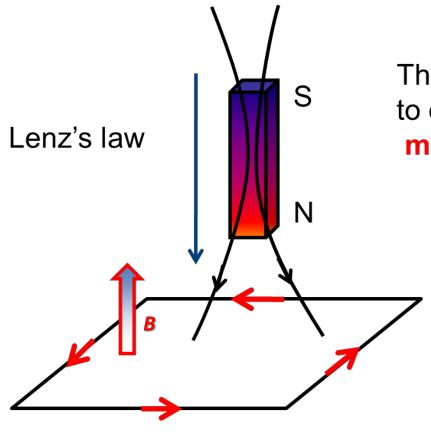
Faraday's Law

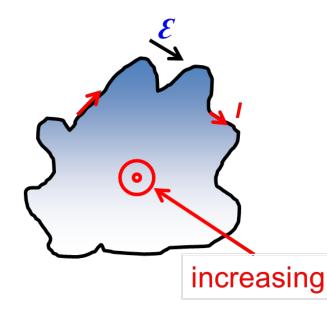


The induced current wants to oppose the change in magnetic flux



$$\mathcal{E} = -N \frac{d\Phi_{\scriptscriptstyle B}}{dt}$$

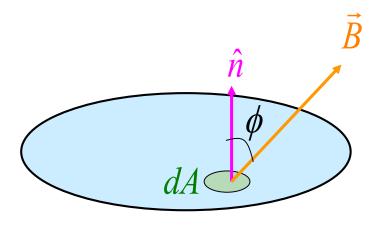
$$\iint \vec{E} \cdot d\vec{s} = -(N) \frac{d\Phi_B}{dt}$$



$$\Phi_B = \int B dA \cos \phi = \int \vec{B} \cdot d\vec{A}$$

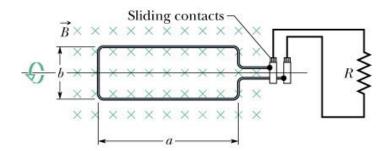
Review of Ch30

Faraday's Law

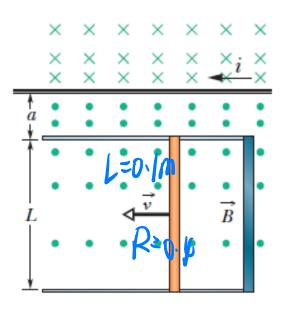


Methods for Changing Φ_R Through a Loop

- **1.** Change the magnitude of *B* within the loop.
- 2. Change either the total area of the coil or the portion of the area within the magnetic field.
- **3.** Change the angle ϕ between \vec{B} and \hat{n} .



The figure shows a rod of length L=10.0 cm that is forced to move at constant speed $\kappa=5.00$ m/s along horizontal rails./The rod, rails, and connecting strip at the right form a conducting loop. The rod has resistance $(0.400\,\Omega)$; the rest of the loop has negligible resistance. A current i=100 A through the long straight wire at distance a=10.0 mm from the loop sets up a magnetic field through the loop. Find:



- (a) the emf and current induced in the loop.
- (b) At what rate is thermal energy generated in the rod?
- (c) What is the magnitude of the force that must be applied to the rod to make it move at constant speed?

(d) At what rate does this force do work on the rod?

(a) We consider an infinitesimal horizontal strip of length x and width dr, parallel to the wire and a distance r from it; it has area A = x dr and the flux is

$$d\Phi_{B} = BdA = \frac{\mu_{0}i}{2\pi r} xdr \implies \Phi_{B} = \frac{\mu_{0}ix}{2\pi} \int_{a}^{a+L} \frac{dr}{r} = \frac{\mu_{0}ix}{2\pi} \ln\left(\frac{a+L}{a}\right).$$

According to Faraday's law the emf induced in the loop is

$$\varepsilon = \frac{d\Phi_B}{dt} = \frac{\mu_0 i}{2\pi} \frac{dx}{dt} \ln\left(\frac{a+L}{a}\right) = \frac{\mu_0 i v}{2\pi} \ln\left(\frac{a+L}{a}\right)$$

$$= \frac{\left(4\pi \times 10^{-7} \,\text{T} \cdot \text{m/A}\right) \left(100 \,\text{A}\right) \left(5.00 \,\text{m/s}\right)}{2\pi} \ln\left(\frac{1.00 \,\text{cm} + 10.0 \,\text{cm}}{1.00 \,\text{cm}}\right) = 2.40 \times 10^{-4} \,\text{V}.$$

By Ohm's law, the induced current is

$$i_{\ell} = \varepsilon / R = (2.40 \times 10^{-4} \text{ V}) / (0.400 \Omega) = 6.00 \times 10^{-4} \text{ A}.$$

Since the flux is increasing, the magnetic field produced by the induced current must be into the page in the region enclosed by the rod and rails. This means the current is clockwise.

(b) Thermal energy is being generated at the rate

$$P = i_{\ell}^{2} R = (6.00 \times 10^{-4} \text{ A})^{2} (0.400 \Omega) = 1.44 \times 10^{-7} \text{ W}.$$

(c) The force of the external agent must have the same magnitude as the magnetic force and must be in the opposite direction.

$$dF_B = i_{\ell}B dr = (\mu_0 i_{\ell} i / 2\pi r) dr.$$

$$F_{B} = \frac{\mu_{0}i_{\ell}i}{2\pi} \int_{a}^{a+L} \frac{dr}{r} = \frac{\mu_{0}i_{\ell}i}{2\pi} \ln\left(\frac{a+L}{a}\right)$$

$$= \frac{\left(4\pi \times 10^{-7} \,\mathrm{T \cdot m/A}\right) \left(6.00 \times 10^{-4} \,\mathrm{A}\right) \left(100 \,\mathrm{A}\right)}{2\pi} \ln\left(\frac{1.00 \,\mathrm{cm} + 10.0 \,\mathrm{cm}}{1.00 \,\mathrm{cm}}\right)$$

$$= 2.87 \times 10^{-8} \,\mathrm{N}. \qquad \text{toward the right}$$

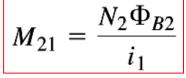
The external agent must therefore apply a force of 2.87×10^{-8} N, to the left.

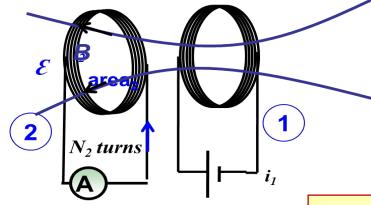
(d) The external agent does work at the rate

$$P = F v = (2.87 \times 10^{-8} \text{ N})(5.00 \text{ m/s}) = 1.44 \times 10^{-7} \text{ W}.$$

All the energy supplied by the agent is converted to thermal energy.

Faraday's law & Inductance





$$N_2\Phi_{B2} = M_{21}i_1$$

$$\mathcal{E}_2 = -N_2 \frac{d\Phi_{B2}}{dt}$$

$$\mathcal{E}_2 = -M_{21} \frac{di_1}{dt}$$

$$M = \frac{N_2 \Phi_{B2}}{i_1} = \frac{N_1 \Phi_{B1}}{i_2}$$

M : Mutual inductance;

$$L = \frac{N\Phi_B}{i}$$

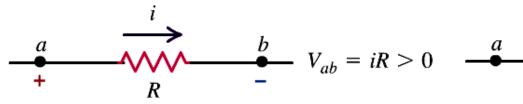
$$N\frac{d\Phi_B}{dt} = L\frac{di}{dt}$$

$$\mathcal{E} = -L\frac{di}{dt}$$

L self inductance

Potential difference across an inductor

(a) Resistor with current *i* flowing from *a* to *b*: potential drops from *a* to *b*.



(b) Inductor with *constant* current *i* flowing from *a* to *b*: no potential difference.

i constant:
$$di/dt = 0$$

$$0000$$

$$E = 0$$

$$V_{ab} = L\frac{di}{dt} = 0$$

(c) Inductor with *increasing* current *i* flowing from *a* to *b*: potential drops from *a* to *b*.

$$\begin{array}{c|c}
 & i \text{ increasing: } di \middle| dt > 0 \\
\hline
 & b \\
\hline
 & V_{ab} = L \frac{di}{dt} > 0
\end{array}$$

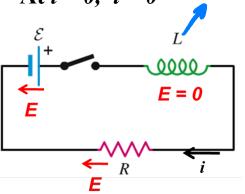
把它视为专派

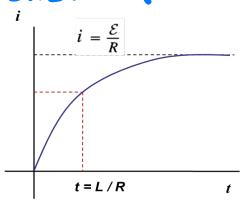
(d) Inductor with decreasing current i flowing from a to b: potential increases from a to b.

$$\frac{a}{\underbrace{-}} \underbrace{\frac{decreasing: di|dt < 0}{b}}_{\underbrace{+}} V_{ab} = L \frac{di}{dt} < 0$$

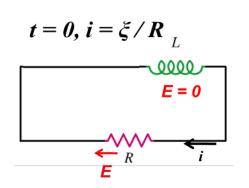
RL Circuit

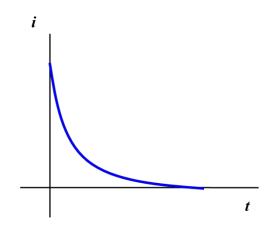






$$0 + iR - \mathscr{E} = -L \frac{di}{dt}$$
$$i = \frac{\mathscr{E}}{R} (1 - e^{-Rt/L})$$

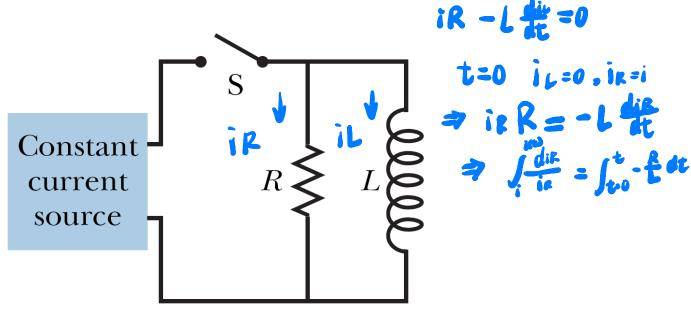




$$L\frac{di}{dt} + iR - \mathcal{E} = 0$$

$$i = \frac{\mathscr{E}}{R} e^{-t/\tau_L}$$

As shown in below figure, after switch S is closed at time t=0, the emf of the source is automatically adjusted to maintain a constant current i through S. (a) Find the current through the inductor as a function of time. (b) At what time is the current through the resistor equal to the current through the inductor?



(a) We assume i is from left to right through the closed switch. We let i_1 be the current in the resistor and take it to be downward. Let i_2 be the current in the inductor, also assumed downward. The junction rule gives $i = i_1 + i_2$ and the loop rule gives $i_1R - L(di_2/dt) = 0$. According to the junction rule, $(di_1/dt) = -(di_2/dt)$. We substitute into the loop equation to obtain $L\frac{di_1}{dt} + i_1R = 0.$

This equation is similar to Eq. 30-44, and its solution is the function given as Eq. 30-45:

$$i_1=i_0e^{-Rt/L},$$

where i_0 is the current through the resistor at t = 0, just after the switch is closed. Now just after the switch is closed, the inductor prevents the rapid build-up of current in its branch, so at that moment $i_2 = 0$ and $i_1 = i$. Thus $i_0 = i$, so

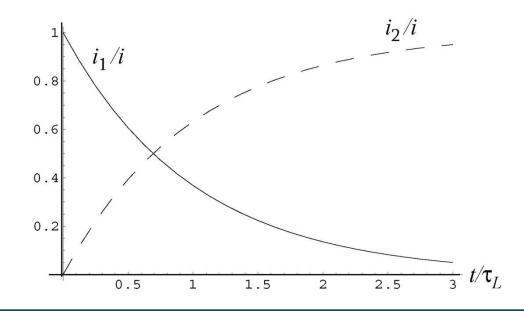
$$i_1 = ie^{-Rt/L}, \quad i_2 = i - i_1 = i(1 - e^{-Rt/L}).$$

(b) When
$$i_2 = i_1$$
, $e^{-Rt/L} = 1 - e^{-Rt/L} \implies e^{-Rt/L} = \frac{1}{2}$.

Taking the natural logarithm of both sides we obtain

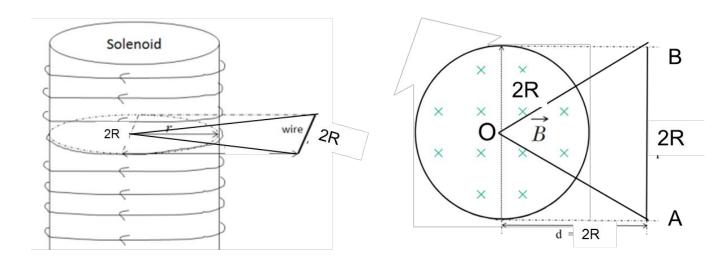
$$\left(\frac{Rt}{L}\right) = \ln 2 \implies t = \frac{L}{R} \ln 2.$$

A plot of $i_{1,2}/i$ as a function of t/τ_L is shown below.



An ideal long solenoid has 50 turns/cm, radius R = 5 cm and length l = 52 cm. An ideal battery of $\varepsilon = 10$ V and resistor of $R_1 = 10$ Ω are in series. A triangle wire loop, whose one vertex is at the center of one coil, is put in the plane of that coil as shown in figures.

- a) What is the mutual inductance for this arrangement?
- If the triangle wire loop is replaced by only the straight wire AB, at the same location. **5ms** after the solenoid connected to the battery,
- b) Find the induced E(r) in the plane shown in the left figure
- c) What is the induced emf in the straight wire AB



For *RL DC* circuit, the increasing current in the loop is:

$$i = \frac{\varepsilon}{R_1} (1 - e^{-R_1 t/L})$$

The uniform B field inside the solenoid is : $B = \mu_0 ni$

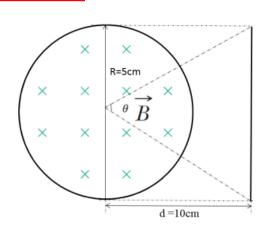
By using the definition of mutual inductance, we can get the mutual inductance for arrangement of the loop and solenoid:

$$M = \frac{\Phi_{tri}}{i_s} = \frac{BA_1}{i_s} = \frac{\mu_0 n A_1 i_s}{i_s} = \mu_0 n A_1 \qquad M = \frac{N_2 \Phi_{B2}}{i_1} = \frac{N_1 \Phi_{B1}}{i_2}$$

$$M = \frac{N_2 \Phi_{B2}}{i_1} = \frac{N_1 \Phi_{B1}}{i_2}$$

$$A_1 = \frac{1}{2}R^2\theta = \frac{1}{2}(5cm)^2(0.93) = 1.16 \times 10^{-3} m^2$$

$$M = \mu_0 n A_1 = 4\pi \times 10^{-7} \times (50 \times 10^2) \times (1.16 \times 10^{-3}) = 7.29 \times 10^{-6} H$$



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According to the symmetry of the B field distribution, we can imagine that the E field must also have the same symmetry. So we can choose the closed loop shown in the figure to calculate the induced electric field:

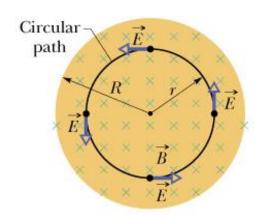
$$\oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi_B}{dt} = \pi r^2 \frac{dB}{dt} \text{ for } r < R \text{ or } \pi R^2 \frac{dB}{dt} \text{ for } r > R$$

$$\oint \vec{E} \cdot d\vec{l} = \oint E dl \cos 0^{\circ} = E \oint dl = 2\pi r E$$

$$E = \frac{r}{2} \frac{dB}{dt} = \frac{r}{2} \mu_0 n \frac{di}{dt} = Ar \ r < R, \ A = \frac{1}{2} \mu_0 n \frac{di}{dt}$$



$$E = \frac{R^2}{2r} \mu_0 n \frac{di}{dt} = \frac{E_1}{r} \quad r > R, \qquad E_1 = \frac{R^2}{2} \mu_0 n \frac{di}{dt}$$

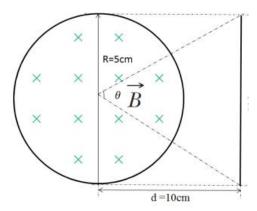


The self-inductance of the solenoid is:

$$L = \mu_0 n^2 A l = 4\pi \times 10^{-7} \times (50 \times 10^2)^2 \times \pi (5 \times 10^{-2})^2 \times 0.52 m = 0.13 H$$

Method 1:

Because the induced E field has no radial component, so the induced emf in the straight wire exactly equals the induced emf in the triangle loop, so we have:



$$\frac{di}{dt} = \frac{\varepsilon}{R} \frac{R}{L} e^{-Rt/L} = \frac{\varepsilon}{L} e^{-Rt/L} = \frac{10}{0.13} e^{-(10)(0.005)/0.13} = 52.36 \text{ A/s}$$

$$\varepsilon_{in} = A_1 \mu_0 n \frac{di}{dt} = 1.16 \times 10^{-3} \times 4\pi \times 10^{-7} \times 50 \times 10^2 \times 52.36 = 3.82 \times 10^{-4} V$$

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Method 2:

We have known the induced electric field distribution, so we do the following integral directly:

$$\varepsilon = \int_{i}^{f} \vec{E} \cdot d\vec{l} = \int_{i}^{f} \frac{E_{1}}{r} dl \cos \phi$$

$$= \int_{i}^{f} \frac{E_{1} \cos \phi}{2R} 2R \sec^{2} \phi \cos \phi d\phi$$

$$= \int_{i}^{f} E_{1} d\phi = E_{1} \theta = \frac{R^{2}}{2} \theta \mu_{0} n \frac{di}{dt}$$

$$= A_{1} \mu_{0} n \frac{di}{dt}$$

