

Compressible inviscid fluid flow solver using the MUSCLE-Hancock method and a HLLC Riemann solver.

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Abstract:

Work is work for some, but for some it is play.

Keywords: hydrodynamics

1 Introduction

For this research we develop a fluid flow solver for the solution of flow problems involving compressible inviscid ideal gases. The governing equations are the Euler equations defined as

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0 \quad (1.1)$$

$$\frac{\partial(\rho \mathbf{u})}{\partial t} + \nabla \cdot \{\rho \mathbf{u} \otimes \mathbf{u}\} + \nabla p = \mathbf{f} \quad (1.2)$$

$$\frac{\partial E}{\partial t} + \nabla \cdot [(E + p)\mathbf{u}] = q, \quad (1.3)$$

where ρ is the fluid density, $\mathbf{u} = [u_x, u_y, u_z] = [u, v, w]$ is the fluid velocity in cartesian coordinates, p is the fluid pressure, $\mathbf{f} = [f_x, f_y, f_z]$ is an arbitrary momentum-density source or sink, E is the material energy-density comprising kinetic energy-density, $\frac{1}{2}\rho||\mathbf{u}||^2$, and internal energy-density, ρe , such that $E = \frac{1}{2}\rho||\mathbf{u}||^2 + \rho e$, where e is the specific internal energy. The value q is an arbitrary energy-density source or sink.

The ideal gas law provides the closure relation

$$p = (\gamma - 1)\rho e \quad (1.4)$$

where γ is the ratio of the constant-pressure specific heat, c_p , to the constant-volume specific heat, c_v , i.e., $\gamma = \frac{c_p}{c_v}$, and is a material property.

1.1 Notation in preparation for numerical schemes

The conservation of mass-, momentum-, and energy equations can be written in the following form

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial}{\partial x} \mathbf{F}(\mathbf{U}) + \frac{\partial}{\partial y} \mathbf{G}(\mathbf{U}) + \frac{\partial}{\partial z} \mathbf{H}(\mathbf{U}) = \mathbf{Q} \quad (1.5)$$

where

$$\mathbf{U} = \begin{bmatrix} \rho \\ \rho u \\ \rho v \\ \rho w \\ E \end{bmatrix}, \quad \mathbf{F}(\mathbf{U}) = \begin{bmatrix} \rho u \\ \rho u u + p \\ \rho u v \\ \rho u w \\ u(E + p) \end{bmatrix}, \quad \mathbf{G}(\mathbf{U}) = \begin{bmatrix} \rho v \\ \rho v u \\ \rho v v + p \\ \rho v w \\ v(E + p) \end{bmatrix}, \quad \mathbf{H}(\mathbf{U}) = \begin{bmatrix} \rho w \\ \rho w u \\ \rho w v \\ \rho w w + p \\ w(E + p) \end{bmatrix}, \quad \text{and} \quad \mathbf{Q} = \begin{bmatrix} 0 \\ f_x \\ f_y \\ f_z \\ q \end{bmatrix}. \quad (1.6)$$

The \mathbf{U} vector is now a collection of the conserved variables, the \mathbf{F} , \mathbf{G} and \mathbf{H} vectors is representative of generic flux terms, and the \mathbf{Q} vector is a generic source term. We will be using this notation in the sections that follow.

1.2 General Finite Volume discretization

Consider the cell volume, in 3D, shown in Figure 1 below.

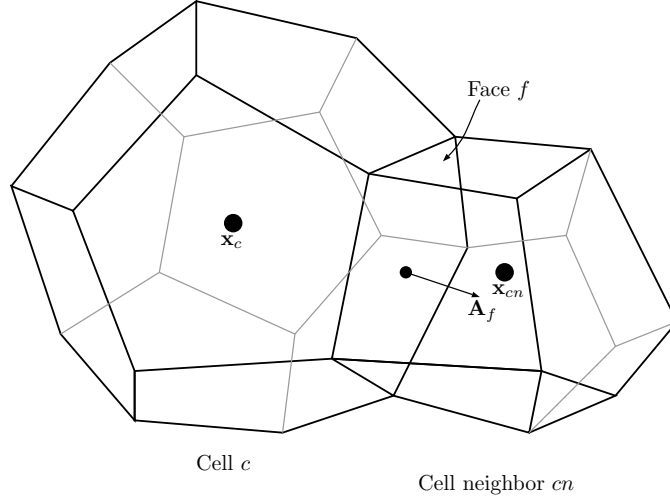


Figure 1: Schematic of a multidimensional cell.

We first apply a spatial integration of Eq. (1.5) over the finite volume of a cell,

$$\int_V \left(\frac{\partial \mathbf{U}}{\partial t} + \nabla \cdot \mathcal{F}(\mathbf{U}) \right) dV = \int_V \mathbf{Q} dV \quad (1.7)$$

where $\mathcal{F}(\mathbf{U}) = [\mathbf{F}(\mathbf{U}), \mathbf{G}(\mathbf{U}), \mathbf{H}(\mathbf{U})]$. Next, using Gauss's divergence theorem, allows us to write

$$\int_V \frac{\partial \mathbf{U}}{\partial t} dV + \int_S \mathbf{n} \cdot \mathcal{F} dA = \int_V \mathbf{Q} dV. \quad (1.8)$$

Now, using cell c as the volume of integration, and assuming \mathbf{U} and \mathbf{Q} constant over the cell, with values \mathbf{U}_c and \mathbf{Q}_c , the above equation becomes

$$V_c \frac{\partial \mathbf{U}_c}{\partial t} + \sum_{f=0}^{N_{f,c}-1} \mathbf{A}_f \cdot \mathcal{F}_f = V_c \mathbf{Q}_c. \quad (1.9)$$

where V_c is the volume of the cell, $N_{f,c}$ is the number of faces for cell c , \mathbf{A}_f is the area-vector of face f , which is the product of the face area, A_f , and the face normal, \mathbf{n}_f (i.e., $\mathbf{A}_f = A_f \mathbf{n}_f$), and $\mathcal{F}_f = \mathcal{F}(\mathbf{U}_f)$ is the face flux vector. The treatment of the \mathcal{F}_f term is the topic of the MUSCL-Hancock method which we detail in section 2.

1.3 Multidimensional transformation of interface fluxes

Most of the Riemann solver schemes presented in [1] are for one dimensional geometries. A transformation technique is prescribed in [1] that allows one to seamlessly use the one dimensional formulations.

Given the arbitrary face area vector, $\mathbf{A}_f = \mathbf{n}_f A_f$, we first seek a rotation matrix, $R_{\hat{\mathbf{i}}}$, to rotate any vector about an axis $\mathbf{a}_{\hat{\mathbf{i}}}$ such that it is aligned with $\hat{\mathbf{i}}$. To determine $R_{\hat{\mathbf{i}}}$ we first set the rotation axis as

$$\mathbf{a}_{\hat{\mathbf{i}}} = \begin{cases} \frac{\mathbf{n}_f \times \hat{\mathbf{i}}}{\|\mathbf{n}_f \times \hat{\mathbf{i}}\|}, & \text{if } \mathbf{n}_f \cdot \hat{\mathbf{i}} < 1 - \epsilon \\ \hat{\mathbf{j}}, & \text{if } \mathbf{n}_f \cdot \hat{\mathbf{i}} \geq 1 - \epsilon \end{cases} \quad (1.10)$$

and the angle of rotation, $\theta_{\hat{\mathbf{i}}}$, as

$$\theta_{\hat{\mathbf{i}}} = \arccos(\mathbf{n}_f \cdot \hat{\mathbf{i}}). \quad (1.11)$$

We then apply Rodrigues's formula as detailed in appendix A to obtain $R_{\hat{\mathbf{i}}}$, the associated rotation matrix. With this matrix in hand we can form the general transformation matrix, $T_{\hat{\mathbf{i}}}$, defined as

$$T_{\hat{\mathbf{i}}} = \text{diag}(1, R, 1) = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & R_{00} & R_{01} & R_{02} & 0 \\ 0 & R_{10} & R_{11} & R_{12} & 0 \\ 0 & R_{20} & R_{21} & R_{22} & 0 \\ 0 & 0 & 0 & 0 & 1 \end{bmatrix} \quad (1.12)$$

which can be used to define

$$\hat{\mathbf{F}}(\mathbf{U}) = \mathbf{F}(T_{\hat{\mathbf{i}}}\mathbf{U}) \quad (1.13)$$

such that

$$\mathbf{n}_f \cdot \mathcal{F}_f = \mathbf{n} \cdot [\mathbf{F}(\mathbf{U}_f), \mathbf{G}(\mathbf{U}_f), \mathbf{H}(\mathbf{U}_f)] = T_{\hat{\mathbf{i}}}^{-1} \hat{\mathbf{F}}_f \quad (1.14)$$

where we also define

$$\mathbf{F}^*(\mathbf{U}) = T_{\hat{\mathbf{i}}}^{-1} \hat{\mathbf{F}}(\mathbf{U}) \quad (1.15)$$

allowing us to write Eq. (1.9) as

$$\frac{\partial \mathbf{U}_c}{\partial t} + \frac{1}{V_c} \sum_{f=0}^{N_{f,c}-1} A_f \mathbf{F}_f^* = \mathbf{Q}_c. \quad (1.16)$$

In this form any Riemann solver can simply be supplied with \mathbf{U} for the cells on either side of an interface, and the face normal, \mathbf{n}_f , in order to use the classical one dimensional formulations contained in [1]. The given Riemann solver will then produce the appropriate value for \mathbf{F}_f^* .

2 MUSCL-Hancock Method (MHM)

The MUSCL-Hancock method is conceptually simple. It prescribes how the interface fluxes are to be computed at the beginning of a time step and then how to supply the necessary inputs to a Riemann-solver for the computation of the preserved variables at the end of the timestep.

MUSCL stands for Monotone Upstream-centered Scheme for Conservation Laws. The scheme as modified by S. Hancock gives the method its name. We will refer to this method by using the abbreviation *MHM*.

A single time step using MHM involves four basic steps.

2.1 Step 1 - Compute the maximum timestep

Generally, the time step size, Δt , needs to be limited in order to ensure numerical stability. The Courant-Friedrichs-Lewy (CFL) condition is the ratio

$$\frac{\|\mathbf{u}\| \Delta t}{L_c} \leq CFL, \quad (2.1)$$

where L_c is the characteristic length of cell c and the value of CFL is subject to the relevant numerical scheme under consideration. For explicit schemes the value of CFL is generally less than 1. When $CFL = 1$ the above expression indicates that the distance traveled by a unit of information, traveling at a velocity \mathbf{u} , over a time period Δt , cannot exceed the length of a cell.

By specifying the value of CFL the CFL-limited time step size for cell c , $\Delta t_{hydro,c}$, is then

$$\Delta t_{hydro,c} = \frac{CFL L_c}{\|\mathbf{u}\|} \quad (2.2)$$

where $\|\mathbf{u}\|$ has not yet been resolved. One option for \mathbf{u} is to use the maximum possible wave speed based on the velocity and sound speed in cell c , \mathbf{u}_c and a_c respectively, as

$$\|\mathbf{u}\| = \|\mathbf{u}_c\| + a_c, \quad (2.3)$$

where

$$a_c = \sqrt{\frac{\gamma_c p_c}{\rho_c}}. \quad (2.4)$$

The simulation time-step size-limit, Δt_{hydro} , is then

$$\Delta t_{hydro} = \min_c \left(\Delta t_{hydro,c} \right). \quad (2.5)$$

For finer control of the simulation the user may also specify a maximum time step size, Δt_{max} . Finally, the simulation time step size, Δt , is then

$$\Delta t = \max \left(\Delta t_{max}, \Delta t_{hydro} \right). \quad (2.6)$$

2.2 Step 2 - Estimate the gradient $\nabla \mathbf{U}$

Provided that an orthogonal mesh is used, the gradient can be estimated in each cell c from

$$\{\nabla \mathbf{U}\}_c^n \approx \frac{1}{V_c} \sum_{f=0}^{N_{f,c}-1} \left\{ \mathbf{A}_f \otimes \left(\frac{\|\mathbf{x}_f - \mathbf{x}_c\|}{\|\mathbf{x}_{cn} - \mathbf{x}_c\|} \mathbf{U}_c^n + \frac{\|\mathbf{x}_c - \mathbf{x}_f\|}{\|\mathbf{x}_{cn} - \mathbf{x}_c\|} \mathbf{U}_{cn}^n \right) \right\}, \quad (2.7)$$

where V_c is the volume of cell c , $N_{f,c}$ is the number of faces for cell c , A_f the face area-vector of face f (i.e., $\mathbf{A}_f = A_f \mathbf{n}_f$), \mathbf{x}_{cn} is the centroid of the neighboring cell cn at face f , and finally \mathbf{U}_{cn} is the finite volume cell-constant value of \mathbf{U} for cell cn . For non-orthogonal meshes the gradient can be corrected as shown in [2].

Limiting:

The cell-wise gradients computed in Eq. (2.7) requires limiting to avoid numerical oscillations. The chosen limiting scheme, to ensure specific properties when coupled with radiation transport (i.e., asymptotic diffusion limit), is the double minmod limiter, as prescribed in [3]. The definition of the double minmod limiter in [3] is either misleading or incorrectly defined as it implies that all gradients are limited to ≥ 0 , therefore we define this limiter in detail here based on the cited literature of [3].

The **general vector-based minmod limiter**, for an M amount of vectors with each vector having N elements, is defined as

$$\text{minmod}(\mathbf{U}^0, \dots, \mathbf{U}^{M-1}) = \begin{bmatrix} \text{minmod}(U_0^0, \dots, U_0^{M-1}) \\ \vdots \\ \text{minmod}(U_{N-1}^0, \dots, U_{N-1}^{M-1}) \end{bmatrix}, \quad (2.8)$$

where U_n^m denotes the n -th entry of the m -th vector. The **general scalar-based minmod limiter**, for M amount of scalar elements, is defined as

$$\text{minmod}(c_0, \dots, c_{M-1}) = \begin{cases} 0, & \text{if } \text{sign}(a_0) \neq \text{sign}(a_m) \text{ for any } m \\ \min(a_0, \dots, a_{M-1}), & \text{if } a_m > 0 \text{ for all } m \\ \max(a_0, \dots, a_{M-1}), & \text{if } a_m < 0 \text{ for all } m \end{cases}. \quad (2.9)$$

These two general limiters are used to define the **double minmod limiter** for the gradient of cell c , which has M amount of neighbor cells.

$$\text{double minmod limited } \{\nabla \mathbf{U}\}_c^n = \text{minmod} \left(\{\nabla \mathbf{U}\}_c^n, \alpha \left\{ \mathbf{x}_{ccn,x} \otimes \frac{\mathbf{U}_{cn} - \mathbf{U}_c}{\|\mathbf{x}_{ccn}\|^2} \right\} \text{ for all } cn \in [0, M-1] \right) \quad (2.10)$$

where $\mathbf{x}_{ccn} = \mathbf{x}_{cn} - \mathbf{x}_c$ and $\alpha = 2$ denoting the “double”. When $\alpha = 0$ the scheme reduces to the unlimited scheme and if $\alpha = 1$ the scheme is the standard minmod limiter.

2.3 Step 3 - Advance the conserved variables by half a time step

Advance the cell-centered values over half a time step as

$$\mathbf{U}_c^{n+\frac{1}{2}} = \mathbf{U}_c^n - \frac{\frac{1}{2}\Delta t^n}{V_c} \sum_{f=0}^{N_{f,c}-1} \left(A_f \mathbf{F}_f^{*n} \right) + \frac{1}{2} \Delta t^n \mathbf{Q}. \quad (2.11)$$

where $\mathbf{F}_f^{*n} = \mathbf{F}^*(\mathbf{U}_f^n)$. \mathbf{U}_f^n is extrapolated from \mathbf{U}_c^n as

$$\mathbf{U}_f^n = \mathbf{U}_c^n + (\mathbf{x}_f - \mathbf{x}_c) \cdot \{\nabla \mathbf{U}\}_c^n \quad (2.12)$$

where \mathbf{x}_f is the face-centroid.

2.4 Step 4 - Execute a series of Riemann solvers

A Riemann solver computes the interface fluxes, $\mathbf{F}_f^{*\mathcal{R},n+\frac{1}{2}}$, using $\mathbf{U}_c^{n+\frac{1}{2}}$, where \mathcal{R} denotes the specific Riemann solver. These interface fluxes are then used to advance the conserved variables by a single timestep as

$$\mathbf{U}_c^{n+1} = \mathbf{U}_c^n - \frac{\Delta t^n}{V_c} \sum_{f=0}^{N_{f,c}-1} \left(A_f \mathbf{F}_f^{*\mathcal{R},n+\frac{1}{2}} \right) + \Delta t^n \mathbf{Q}. \quad (2.13)$$

The discontinuity across a face is treated as a one dimensional problem with a left and right side having different values for \mathbf{U} , i.e., $\mathbf{U}_L = \mathbf{U}_{f,c}^{n+\frac{1}{2}}$ and $\mathbf{U}_R = \mathbf{U}_{f,cn}^{n+\frac{1}{2}}$ for the left and right side respectively. The face values of \mathbf{U} are extrapolated using the gradient $\{\nabla \mathbf{U}\}^n$ such that

$$\mathbf{U}_{f,c}^n = \mathbf{U}_c^{n+\frac{1}{2}} + (\mathbf{x}_f - \mathbf{x}_c) \cdot \{\nabla \mathbf{U}\}_c^n \quad (2.14a)$$

$$\mathbf{U}_{f,cn}^n = \mathbf{U}_{cn}^{n+\frac{1}{2}} + (\mathbf{x}_f - \mathbf{x}_c) \cdot \{\nabla \mathbf{U}\}_{cn}^n. \quad (2.14b)$$

When using the HLLC solver, as described in [1], the Riemann solver will compute $\mathbf{F}_f^{*hllc,n+\frac{1}{2}}$ after which we compute the conserved variables at $n+1$ from

$$\mathbf{U}_c^{n+1} = \mathbf{U}_c^n - \frac{\Delta t^n}{V_c} \sum_{f=0}^{N_{f,c}-1} \left(A_f \mathbf{F}_f^{*hllc,n+\frac{1}{2}} \right) + \Delta t^n \mathbf{Q}. \quad (2.15)$$

The HLLC Riemann solver is detailed in section 3.

3 The HLLC Approximate Riemann Solver

The Harten, Lax and van Leer (HLL) solver scheme was developed in 1983 [1] and requires estimates for the fastest wave/signal/shock velocities emerging from a discontinuity. Later Toro, Spruce and Speares proposed the Harten, Lax, van Leer, *Contact* (HLLC) scheme [1] which adds another wave to the problem.

The first input, required by the HLLC Riemann solver, is the face normal, \mathbf{n}_f , which allows us to compute the transformation matrix, $T_{\mathbf{i}}$, as per section 1.3. The other input parameters are then

$$\begin{aligned} \mathbf{U}_L &= T_{\mathbf{i}} \mathbf{U}_{f,c}^{n+\frac{1}{2}}, \\ \mathbf{U}_R &= T_{\mathbf{i}} \mathbf{U}_{f,cn}^{n+\frac{1}{2}}, \\ \mathbf{F}_L &= \mathbf{F}(T_{\mathbf{i}} \mathbf{U}_{f,c}^{n+\frac{1}{2}}), \\ \mathbf{F}_R &= \mathbf{F}(T_{\mathbf{i}} \mathbf{U}_{f,cn}^{n+\frac{1}{2}}), \\ p_L &= p_c, \quad \gamma_L = \gamma_c, \\ p_R &= p_{cn}, \quad \gamma_R = \gamma_{cn}, \end{aligned} \tag{3.1}$$

where the quantities denoted with c denotes those belonging to the cell which maintains a negative sense with respect to face f , and conversely cn denotes the quantities associated with the cell maintaining a positive sense with respect to the face.

3.1 Left and right wave speed estimation

The HLLC Riemann solver is predicated on knowing an estimate for wave speeds S_L and S_R , which we estimate as

$$S_L = \min(u_L - a_L, u_R - a_R) \tag{3.2}$$

and

$$S_R = \max(u_L + a_L, u_R + a_R). \tag{3.3}$$

where a_L and a_R are the sound speeds associated with the left- and right conserved variables as

$$a = \sqrt{\frac{\gamma p}{\rho}}. \tag{3.4}$$

Next we require the contact wave speed.

3.2 Contact wave speed

The contact wave speed, S_* , is given by

$$S_* = \frac{p_R - p_L + \rho_L u_L (S_L - u_L) - \rho_R u_R (S_R - u_R)}{\rho_L (S_L - u_L) - \rho_R (S_R - u_R)}. \tag{3.5}$$

3.3 Intermediate fluxes

As per [1] the intermediate fluxes, \mathbf{F}_{*L} and \mathbf{F}_{*R} are given by

$$F_{*K} = \frac{S_*(S_K \mathbf{U}_K - \mathbf{F}_K) + S_K(p_K + \rho_L(S_K - u_K)(S_* - u_K))D_*}{S_K - S_*} \tag{3.6}$$

for $K = L$ and $K = R$. The vector \mathbf{D}_* is a vector such that

$$\mathbf{F}(\mathbf{U}) = u\mathbf{U} + p\mathbf{D}, \tag{3.7}$$

therefore

$$\mathbf{D}_* = [0, 1, 0, 0, S_*]^T \tag{3.8}$$

3.4 Interface flux

The interface flux \mathbf{F}_f^{*hllc} is now given by

$$\mathbf{F}_f^{*hllc} = T_{\mathbf{i}}^{-1} \mathbf{F}_f^{hllc}$$

$$\mathbf{F}_f^{hllc} = \begin{cases} \mathbf{F}_L & , \text{ if } S_L \geq 0, \\ \mathbf{F}_{*L} & , \text{ if } S_L \leq 0 \leq S_*, \\ \mathbf{F}_{*R} & , \text{ if } S_* \leq 0 \leq S_R, \\ \mathbf{F}_R & , \text{ if } S_R \leq 0 \end{cases} \quad (3.9)$$

4 Verification - Sod shock tube problem

The Sod shock tube problem is a simple problem with the following specifications:

- The one dimensional problem domain has a total size of 1.0 spanning $x \in [-\frac{1}{2}, \frac{1}{2}]$.
- The left-half of the problem has an initial state denoted with L and the right-half has an initial state denoted with R .
- At time $t = 0$:
 - Densities: $\rho_L = 1, \rho_R = 0.125$
 - Pressures: $p_L = 1, p_R = 0.1$
 - Velocity: $u = 0 \forall x$
 - Ratio of specific heats: $\gamma = 1.4 \forall x$
 - Boundary conditions: Transmissive

As time evolves the problem exhibits a shock-wave, a contact-wave and a rarefaction-wave. The analytical solution has been obtain from a Fortran code by Timmers [4]. At time $t = 0.2$ the analytical solution for ρ, p, u and e is tabulated over 500 points in appendix B.

The MHM-HLLC scheme employed above has been executed with the following inputs:

- $\Delta t_{max} = 1e - 2$
- $CFL = 0.3$
- Maximum number of timesteps, 2000
- Maximum total time, 0.2
- $\Delta x = 0.01$ or 100 cells.

The program reached the maximum total time after 73 iterations and the results at $t = 0.2$ is compared to the analytical solution in Figure 2.

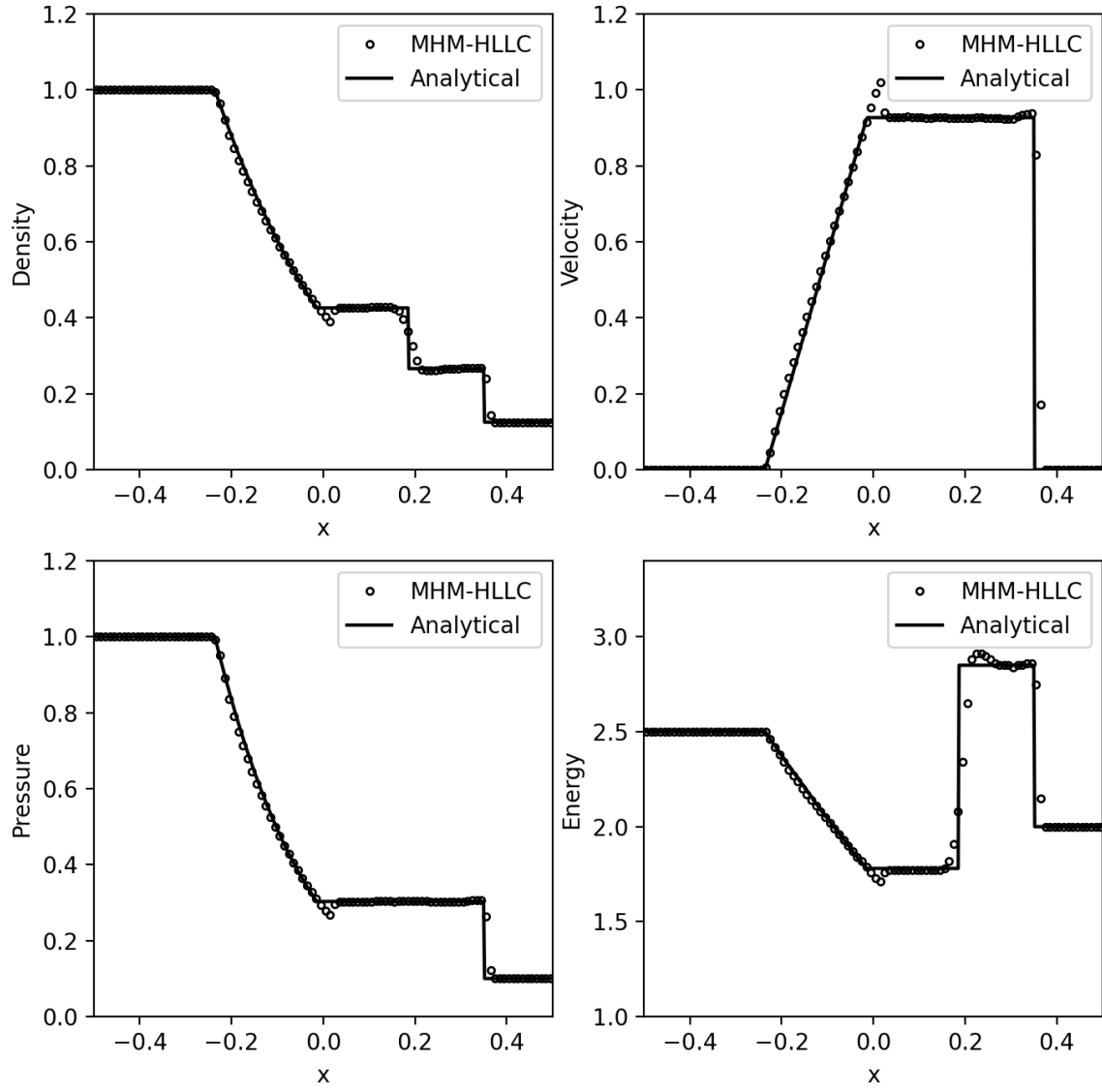


Figure 2: Numerical results compared to the analytical results at $t = 0.2$.

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A Roderigues's formula

Roderigues' formula for the rotation of a vector \mathbf{v} about a unit vector \mathbf{a} with right-hand rule

$$\mathbf{v}_{rotated} = \cos \theta \mathbf{v} + (\mathbf{a} \cdot \mathbf{v})(1 - \cos \theta) \mathbf{a} + \sin \theta (\mathbf{a} \times \mathbf{v}) \quad (\text{A.1})$$

In matrix form

$$\mathbf{v}_{rotated} = A \mathbf{v} \quad (\text{A.2})$$

where

$$A = \begin{bmatrix} 0 & -a_z & a_y \\ a_z & 0 & -a_x \\ -a_y & a_x & 0 \end{bmatrix} \quad (\text{A.3})$$

and

$$R = I + \sin \theta A + (1 - \cos \theta) A^2 \quad (\text{A.4})$$

B Sod shock tube problem - analytical solution at $t = 0.2$

i	x	density	pressure	velocity	energy	i	x	density	pressure	velocity	energy
1	-5.00E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	251	1.00E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00
2	-4.98E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	252	3.01E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00
3	-4.96E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	253	5.01E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00
4	-4.94E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	254	7.01E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00
5	-4.92E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	255	9.02E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00
6	-4.90E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	256	1.10E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
7	-4.88E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	257	1.30E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
8	-4.86E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	258	1.50E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
9	-4.84E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	259	1.70E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
10	-4.82E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	260	1.90E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
11	-4.80E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	261	2.10E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
12	-4.78E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	262	2.30E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
13	-4.76E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	263	2.51E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
14	-4.74E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	264	2.71E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
15	-4.72E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	265	2.91E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
16	-4.70E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	266	3.11E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
17	-4.68E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	267	3.31E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
18	-4.66E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	268	3.51E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
19	-4.64E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	269	3.71E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
20	-4.62E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	270	3.91E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
21	-4.60E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	271	4.11E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
22	-4.58E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	272	4.31E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
23	-4.56E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	273	4.51E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
24	-4.54E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	274	4.71E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
25	-4.52E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	275	4.91E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
26	-4.50E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	276	5.11E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
27	-4.48E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	277	5.31E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
28	-4.46E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	278	5.51E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
29	-4.44E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	279	5.71E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
30	-4.42E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	280	5.91E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
31	-4.40E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	281	6.11E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
32	-4.38E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	282	6.31E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
33	-4.36E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	283	6.51E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
34	-4.34E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	284	6.71E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
35	-4.32E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	285	6.91E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
36	-4.30E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	286	7.11E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
37	-4.28E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	287	7.31E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
38	-4.26E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	288	7.52E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
39	-4.24E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	289	7.72E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
40	-4.22E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	290	7.92E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
41	-4.20E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	291	8.12E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
42	-4.18E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	292	8.32E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
43	-4.16E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	293	8.52E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
44	-4.14E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	294	8.72E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
45	-4.12E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	295	8.92E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
46	-4.10E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	296	9.12E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
47	-4.08E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	297	9.32E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
48	-4.06E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	298	9.52E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
49	-4.04E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	299	9.72E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
50	-4.02E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	300	9.92E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00
51	-4.00E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	301	1.01E-01	4.26E-01	3.03E-01	9.27E-01	1.78E+00
52	-3.98E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	302	1.03E-01	4.26E-01	3.03E-01	9.27E-01	1.78E+00
53	-3.96E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	303	1.05E-01	4.26E-01	3.03E-01	9.27E-01	1.78E+00
54	-3.94E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	304	1.07E-01	4.26E-01	3.03E-01	9.27E-01	1.78E+00
55	-3.92E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	305	1.09E-01	4.26E-01	3.03E-01	9.27E-01	1.78E+00
56	-3.90E-01	1.00E+00	1.00E+00	0.00E+00	2.50E+00	306	1.11E-01	4.26E-01	3.03E-01	9.27E-01	1.78E+00

140	-2.21E-01	9.48E-01	9.27E-01	6.33E-02	2.45E+00	390	2.80E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
141	-2.19E-01	9.41E-01	9.18E-01	7.17E-02	2.44E+00	391	2.82E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
142	-2.17E-01	9.34E-01	9.09E-01	8.00E-02	2.43E+00	392	2.84E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
143	-2.15E-01	9.28E-01	9.00E-01	8.84E-02	2.43E+00	393	2.86E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
144	-2.13E-01	9.21E-01	8.91E-01	9.67E-02	2.42E+00	394	2.88E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
145	-2.11E-01	9.14E-01	8.82E-01	1.05E-01	2.41E+00	395	2.90E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
146	-2.09E-01	9.08E-01	8.73E-01	1.13E-01	2.41E+00	396	2.92E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
147	-2.07E-01	9.01E-01	8.65E-01	1.22E-01	2.40E+00	397	2.94E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
148	-2.05E-01	8.95E-01	8.56E-01	1.30E-01	2.39E+00	398	2.96E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
149	-2.03E-01	8.88E-01	8.47E-01	1.38E-01	2.38E+00	399	2.98E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
150	-2.01E-01	8.82E-01	8.39E-01	1.47E-01	2.38E+00	400	3.00E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
151	-1.99E-01	8.76E-01	8.30E-01	1.55E-01	2.37E+00	401	3.02E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
152	-1.97E-01	8.69E-01	8.22E-01	1.64E-01	2.36E+00	402	3.04E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
153	-1.95E-01	8.63E-01	8.14E-01	1.72E-01	2.36E+00	403	3.06E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
154	-1.93E-01	8.57E-01	8.05E-01	1.80E-01	2.35E+00	404	3.08E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
155	-1.91E-01	8.50E-01	7.97E-01	1.89E-01	2.34E+00	405	3.10E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
156	-1.89E-01	8.44E-01	7.89E-01	1.97E-01	2.34E+00	406	3.12E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
157	-1.87E-01	8.38E-01	7.81E-01	2.05E-01	2.33E+00	407	3.14E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
158	-1.85E-01	8.32E-01	7.73E-01	2.14E-01	2.32E+00	408	3.16E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
159	-1.83E-01	8.26E-01	7.65E-01	2.22E-01	2.32E+00	409	3.18E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
160	-1.81E-01	8.20E-01	7.57E-01	2.30E-01	2.31E+00	410	3.20E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
161	-1.79E-01	8.14E-01	7.50E-01	2.39E-01	2.30E+00	411	3.22E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
162	-1.77E-01	8.08E-01	7.42E-01	2.47E-01	2.30E+00	412	3.24E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
163	-1.75E-01	8.02E-01	7.34E-01	2.55E-01	2.29E+00	413	3.26E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
164	-1.73E-01	7.96E-01	7.27E-01	2.64E-01	2.28E+00	414	3.28E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
165	-1.71E-01	7.90E-01	7.19E-01	2.72E-01	2.28E+00	415	3.30E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
166	-1.69E-01	7.84E-01	7.12E-01	2.80E-01	2.27E+00	416	3.32E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
167	-1.67E-01	7.79E-01	7.04E-01	2.89E-01	2.26E+00	417	3.34E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
168	-1.65E-01	7.73E-01	6.97E-01	2.97E-01	2.26E+00	418	3.36E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
169	-1.63E-01	7.67E-01	6.90E-01	3.05E-01	2.25E+00	419	3.38E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
170	-1.61E-01	7.61E-01	6.83E-01	3.14E-01	2.24E+00	420	3.40E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
171	-1.59E-01	7.56E-01	6.76E-01	3.22E-01	2.24E+00	421	3.42E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
172	-1.57E-01	7.50E-01	6.69E-01	3.31E-01	2.23E+00	422	3.44E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
173	-1.55E-01	7.45E-01	6.62E-01	3.39E-01	2.22E+00	423	3.46E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
174	-1.53E-01	7.39E-01	6.55E-01	3.47E-01	2.22E+00	424	3.48E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
175	-1.51E-01	7.33E-01	6.48E-01	3.56E-01	2.21E+00	425	3.50E-01	2.66E-01	3.03E-01	9.27E-01	2.85E+00
176	-1.49E-01	7.28E-01	6.41E-01	3.64E-01	2.20E+00	426	3.52E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
177	-1.47E-01	7.23E-01	6.34E-01	3.72E-01	2.20E+00	427	3.54E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
178	-1.45E-01	7.17E-01	6.28E-01	3.81E-01	2.19E+00	428	3.56E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
179	-1.43E-01	7.12E-01	6.21E-01	3.89E-01	2.18E+00	429	3.58E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
180	-1.41E-01	7.06E-01	6.15E-01	3.97E-01	2.18E+00	430	3.60E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
181	-1.39E-01	7.01E-01	6.08E-01	4.06E-01	2.17E+00	431	3.62E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
182	-1.37E-01	6.96E-01	6.02E-01	4.14E-01	2.16E+00	432	3.64E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
183	-1.35E-01	6.90E-01	5.95E-01	4.22E-01	2.16E+00	433	3.66E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
184	-1.33E-01	6.85E-01	5.89E-01	4.31E-01	2.15E+00	434	3.68E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
185	-1.31E-01	6.80E-01	5.83E-01	4.39E-01	2.14E+00	435	3.70E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
186	-1.29E-01	6.75E-01	5.77E-01	4.47E-01	2.14E+00	436	3.72E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
187	-1.27E-01	6.70E-01	5.71E-01	4.56E-01	2.13E+00	437	3.74E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
188	-1.25E-01	6.65E-01	5.64E-01	4.64E-01	2.12E+00	438	3.76E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
189	-1.23E-01	6.60E-01	5.58E-01	4.72E-01	2.12E+00	439	3.78E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
190	-1.21E-01	6.55E-01	5.52E-01	4.81E-01	2.11E+00	440	3.80E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
191	-1.19E-01	6.50E-01	5.47E-01	4.89E-01	2.10E+00	441	3.82E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
192	-1.17E-01	6.45E-01	5.41E-01	4.98E-01	2.10E+00	442	3.84E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
193	-1.15E-01	6.40E-01	5.35E-01	5.06E-01	2.09E+00	443	3.86E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
194	-1.13E-01	6.35E-01	5.29E-01	5.14E-01	2.08E+00	444	3.88E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
195	-1.11E-01	6.30E-01	5.23E-01	5.23E-01	2.08E+00	445	3.90E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
196	-1.09E-01	6.25E-01	5.18E-01	5.31E-01	2.07E+00	446	3.92E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
197	-1.07E-01	6.20E-01	5.12E-01	5.39E-01	2.06E+00	447	3.94E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
198	-1.05E-01	6.15E-01	5.07E-01	5.48E-01	2.06E+00	448	3.96E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
199	-1.03E-01	6.11E-01	5.01E-01	5.56E-01	2.05E+00	449	3.98E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
200	-1.01E-01	6.06E-01	4.96E-01	5.64E-01	2.05E+00	450	4.00E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
201	-9.92E-02	6.01E-01	4.90E-01	5.73E-01	2.04E+00	451	4.02E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
202	-9.72E-02	5.96E-01	4.85E-01	5.81E-01	2.03E+00	452	4.04E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
203	-9.52E-02	5.92E-01	4.80E-01	5.89E-01	2.03E+00	453	4.06E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
204	-9.32E-02	5.87E-01	4.74E-01	5.98E-01	2.02E+00	454	4.08E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
205	-9.12E-02	5.83E-01	4.69E-01	6.06E-01	2.01E+00	455	4.10E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
206	-8.92E-02	5.78E-01	4.64E-01	6.14E-01	2.01E+00	456	4.12E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
207	-8.72E-02	5.73E-01	4.59E-01	6.23E-01	2.00E+00	457	4.14E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
208	-8.52E-02	5.69E-01	4.54E-01	6.31E-01	2.00E+00	458	4.16E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
209	-8.32E-02	5.64E-01	4.49E-01	6.39E-01	1.99E+00	459	4.18E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
210	-8.12E-02	5.60E-01	4.44E-01	6.48E-01	1.98E+00	460	4.20E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
211	-7.92E-02	5.56E-01	4.39E-01	6.56E-01	1.98E+00	461	4.22E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
212	-7.72E-02	5.51E-01	4.34E-01	6.65E-01	1.97E+00	462	4.24E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
213	-7.52E-02	5.47E-01	4.29E-01	6.73E-01	1.96E+00	463	4.26E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
214	-7.31E-02	5.42E-01	4.25E-01	6.81E-01	1.96E+00	464	4.28E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
215	-7.11E-02	5.38E-01	4.20E-01	6.90E-01	1.95E+00	465	4.30E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
216	-6.91E-02	5.34E-01	4.15E-01	6.98E-01	1.94E+00	466	4.32E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
217	-6.71E-02	5.30E-01	4.11E-01	7.06E-01	1.94E+00	467	4.34E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
218	-6.51E-02	5.25E-01	4.06E-01	7.15E-01	1.93E+00	468	4.36E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
219	-6.31E-02	5.21E-01	4.02E-01	7.23E-01	1.93E+00	469	4.38E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
220	-6.11E-02	5.17E-01	3.97E-01	7.31E-01	1.92E+00	470	4.40E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
221	-5.91E-02	5.13E-01	3.93E-01	7.40E-01	1.91E+00	471	4.42E-01	1.25E-01	1.00E-01		

223	-5.51E-02	5.05E-01	3.84E-01	7.56E-01	1.90E+00	473	4.46E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
224	-5.31E-02	5.01E-01	3.79E-01	7.65E-01	1.90E+00	474	4.48E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
225	-5.11E-02	4.96E-01	3.75E-01	7.73E-01	1.89E+00	475	4.50E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
226	-4.91E-02	4.92E-01	3.71E-01	7.81E-01	1.88E+00	476	4.52E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
227	-4.71E-02	4.88E-01	3.67E-01	7.90E-01	1.88E+00	477	4.54E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
228	-4.51E-02	4.85E-01	3.63E-01	7.98E-01	1.87E+00	478	4.56E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
229	-4.31E-02	4.81E-01	3.58E-01	8.06E-01	1.86E+00	479	4.58E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
230	-4.11E-02	4.77E-01	3.54E-01	8.15E-01	1.86E+00	480	4.60E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
231	-3.91E-02	4.73E-01	3.50E-01	8.23E-01	1.85E+00	481	4.62E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
232	-3.71E-02	4.69E-01	3.46E-01	8.32E-01	1.85E+00	482	4.64E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
233	-3.51E-02	4.65E-01	3.42E-01	8.40E-01	1.84E+00	483	4.66E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
234	-3.31E-02	4.61E-01	3.38E-01	8.48E-01	1.83E+00	484	4.68E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
235	-3.11E-02	4.57E-01	3.35E-01	8.57E-01	1.83E+00	485	4.70E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
236	-2.91E-02	4.54E-01	3.31E-01	8.65E-01	1.82E+00	486	4.72E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
237	-2.71E-02	4.50E-01	3.27E-01	8.73E-01	1.82E+00	487	4.74E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
238	-2.51E-02	4.46E-01	3.23E-01	8.82E-01	1.81E+00	488	4.76E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
239	-2.30E-02	4.43E-01	3.19E-01	8.90E-01	1.80E+00	489	4.78E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
240	-2.10E-02	4.39E-01	3.16E-01	8.98E-01	1.80E+00	490	4.80E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
241	-1.90E-02	4.35E-01	3.12E-01	9.07E-01	1.79E+00	491	4.82E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
242	-1.70E-02	4.32E-01	3.08E-01	9.15E-01	1.79E+00	492	4.84E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
243	-1.50E-02	4.28E-01	3.05E-01	9.23E-01	1.78E+00	493	4.86E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
244	-1.30E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00	494	4.88E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
245	-1.10E-02	4.26E-01	3.03E-01	9.27E-01	1.78E+00	495	4.90E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
246	-9.02E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00	496	4.92E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
247	-7.01E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00	497	4.94E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
248	-5.01E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00	498	4.96E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
249	-3.01E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00	499	4.98E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00
250	-1.00E-03	4.26E-01	3.03E-01	9.27E-01	1.78E+00	500	5.00E-01	1.25E-01	1.00E-01	0.00E+00	2.00E+00