Personal Notes: Piece-wise linear shape/basis functions

August, 2018

Jan Vermaak Rev 1.00

1 Fundamental strategy

In the finite element method the basic poisson equation,

$$-\nabla^2 \phi + \sigma_a \phi = q,$$

is multiplied by a trial function v_i to obtain

$$-\int_{V} v_{i} \nabla^{2} \phi. dV + \int_{V} v_{i}. \sigma_{a} \phi. dV = \int_{V} v_{i}. q. dV.$$

We then integrate by parts

$$-\int_{V} \nabla(v_{i} \nabla \phi) . dV + \int_{V} \nabla v_{i} \cdot \nabla \phi . dV + \int_{V} v_{i} . \sigma_{a} \phi . dV = \int_{V} v_{i} . q . dV$$

and apply Gauss' divergence theorem to get

$$-\int_{S} \hat{\mathbf{n}} \cdot (v_{i} \nabla \phi) . dA + \int_{V} \nabla v_{i} \cdot \nabla \phi . dV + \int_{V} v_{i} . \sigma_{a} \phi . dV = \int_{V} v_{i} . q . dV$$

We then expand ϕ into basis functions and coefficients $\phi = \sum \phi_j b_j$ after which we get

$$\sum_{j} - \int_{S} \hat{\mathbf{n}} \cdot (v_{i} \phi_{j} \nabla b_{j}) . dA + \int_{V} \phi_{j} \nabla v_{i} \cdot \nabla b_{j} . dV + \int_{V} \phi_{j} \sigma_{a} v_{i} . b_{j} . dV = \int_{V} v_{i} . q . dV$$

The goal of this exercise is to define the trial functions v_i and basis functions b_j

2 Two Dimensional Triangle

We start with the two dimensional reference triangle as shown in Figure 1 below. We wish to map any triangular cell C to this reference triangle in order to develop a consistent method for cells of any orientation.

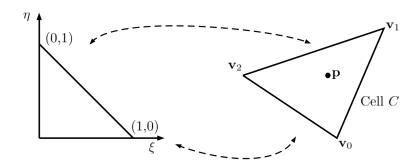


Figure 1: Natural coordinates used as a reference 2D element.

We first see that we can project any point $\mathbf{p}=(x,y)$, within C, using a set of dot products

$$\xi = (\mathbf{p} - \mathbf{v}_0) \cdot (\mathbf{v}_1 - \mathbf{v}_0) \tag{2.1}$$

$$\eta = (\mathbf{p} - \mathbf{v}_0) \cdot (\mathbf{v}_2 - \mathbf{v}_0). \tag{2.2}$$

Conversely we can map any point, in the natural coordinates triangle, to a point in C as

$$\mathbf{p} = \mathbf{v}_0 + \xi(\mathbf{v}_1 - \mathbf{v}_0) + \eta(\mathbf{v}_2 - \mathbf{v}_0). \tag{2.3}$$

In expanded form this is

$$x = x_0 + \xi(x_1 - x_0) + \eta(x_2 - x_0)$$
$$y = y_0 + \xi(y_1 - y_0) + \eta(y_2 - y_0)$$

from which we determine

$$\begin{split} \frac{\partial x}{\partial \xi} &= \mathbf{x}_1 - \mathbf{x}_0 & \frac{\partial x}{\partial \eta} &= \mathbf{x}_2 - \mathbf{x}_0 \\ \frac{\partial y}{\partial \xi} &= \mathbf{y}_1 - \mathbf{y}_0 & \frac{\partial y}{\partial \eta} &= \mathbf{y}_2 - \mathbf{y}_0. \end{split}$$

We can then define the Jacobian matrix of the transformation

$$J = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \xi} & \frac{\partial y}{\partial \eta} \end{bmatrix} = \begin{bmatrix} x_1 - x_0 & x_2 - x_0 \\ y_1 - y_0 & y_2 - y_0 \end{bmatrix}$$
(2.4)

which gives us the simplified form

and

$$\begin{bmatrix} \xi \\ \eta \end{bmatrix} = J^{-1} \begin{bmatrix} x - x_0 \\ y - y_0 \end{bmatrix}. \tag{2.6}$$

For a 2×2 matrix the inverse of the Jacobian, J^{-1} , can easily be determined as

$$J^{-1} = \frac{1}{|J|} \begin{bmatrix} \frac{\partial y}{\partial \eta} & -\frac{\partial x}{\partial \eta} \\ -\frac{\partial y}{\partial \xi} & \frac{\partial x}{\partial \xi} \end{bmatrix}$$
 (2.7)

We can now define a basis function \bar{b} for each vertex of the reference triangle

$$\bar{b}_0(\xi, \eta) = 1 - \xi - \eta$$

$$\bar{b}_1(\xi, \eta) = \xi$$

$$\bar{b}_2(\xi, \eta) = \eta$$
(2.8)

and map it to cartesian coordinates to find

$$b_{0}(x,y) = 1 - J_{11}^{-1}(x - x_{0}) - J_{12}^{-1}(y - y_{0}) - J_{21}^{-1}(x - x_{0}) - J_{22}^{-1}(y - y_{0})$$

$$b_{1}(x,y) = J_{11}^{-1}(x - x_{0}) + J_{12}^{-1}(y - y_{0})$$

$$b_{2}(x,y) = J_{21}^{-1}(x - x_{0}) + J_{22}^{-1}(y - y_{0}).$$
(2.9)

2.1 Integrating the basis functions

The finite element method will in some form require the integration of the basis functions over either a volume or an area. Integrating the basis functions over an area can be accomplished as

$$\int \int b_j(x,y).dx.dy = \int \int \bar{b}_j.|J|.d\xi.d\eta.$$

This integration involves linear functions for which we can use a quadrature rule.

2.2 Integrating the gradient of the basis functions

In any finite element method there would be the need for the integration of the shape functions as well as the gradients of the shape functions. The shape functions are now already defined in this context but we still need the gradients of the shape functions

$$\vec{\nabla} f = \begin{bmatrix} \frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y} \end{bmatrix}$$

Now from the chain rule we have

$$\frac{\partial f}{\partial x} = \frac{\partial f}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial f}{\partial \eta} \frac{\partial \eta}{\partial x}$$
$$\frac{\partial f}{\partial y} = \frac{\partial f}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial f}{\partial \eta} \frac{\partial \eta}{\partial y}$$

Which we can write as

$$\begin{bmatrix} \frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y} \end{bmatrix} = \begin{bmatrix} \frac{\partial \xi}{\partial x} & \frac{\partial \eta}{\partial x} \\ \frac{\partial \xi}{\partial y} & \frac{\partial \eta}{\partial y} \end{bmatrix} \begin{bmatrix} \frac{\partial f}{\partial \xi} \\ \frac{\partial f}{\partial \eta} \end{bmatrix}$$

But recall the mapping in equation 2.6 for which we have

$$\xi = J_{11}^{-1}(x-x_0) + J_{12}^{-1}(y-y_0)$$

$$\eta = J_{21}^{-1}(x-x_0) + J_{22}^{-1}(y-y_0)$$

Therefore

$$\begin{bmatrix} \frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y} \end{bmatrix} = J^{-1} \begin{bmatrix} \frac{\partial f}{\partial \xi} \\ \frac{\partial f}{\partial \eta} \end{bmatrix}$$

3 Polygon

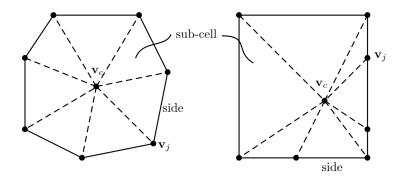


Figure 2: Application of sub-cells for polygons

References

- [1] Blender a 3D modelling and rendering package, Blender Online Community, Blender Foundation, Blender Institute, Amsterdam, 2018
- [2] Cheng et al, Delaunay Mesh Generation, Chapman & Hall/CRC Computer & Information Science Series, 2013