



Sandia
National
Laboratories

Exceptional service in the national interest

ME469: Geometric Fidelity

Stefan P. Domino^{1,2}

¹ Computational Thermal and Fluid Mechanics, Sandia National Laboratories

² Institute for Computational and Mathematical Engineering, Stanford

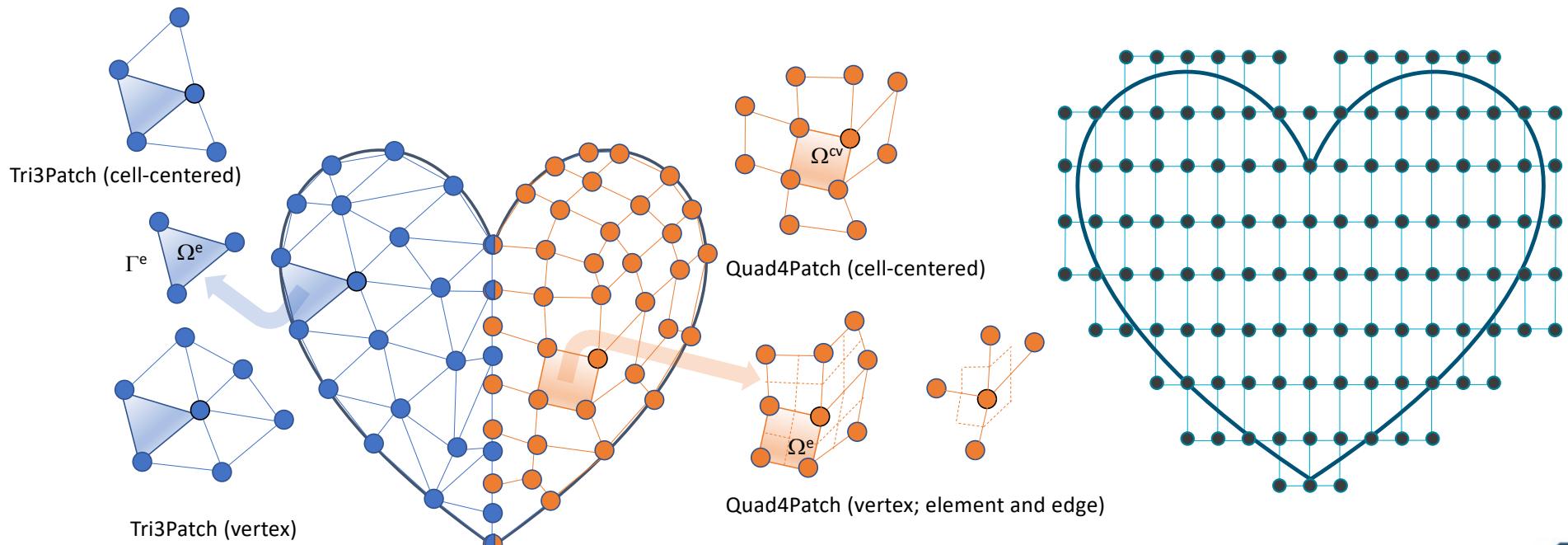
This presentation has been authored by an employee of National Technology & Engineering Solutions of Sandia, LLC under Contract No. DE-NA0003525 with the U.S. Department of Energy (DOE). The employee owns all right, title and interest in and to the presentation and is solely responsible for its contents. The United States Government retains and the publisher, by accepting the article for publication, acknowledges that the United States Government retains a non-exclusive, paid-up, irrevocable, world-wide license to publish or reproduce the published form of this article or allow others to do so, for United States Government purposes. The DOE will provide public access to these results of federally sponsored research in accordance with the DOE Public Access Plan <http://www.energy.gov/downloads/doe-public-access-plan>.

SAND2018-4536 PE



Review of Discretization Options

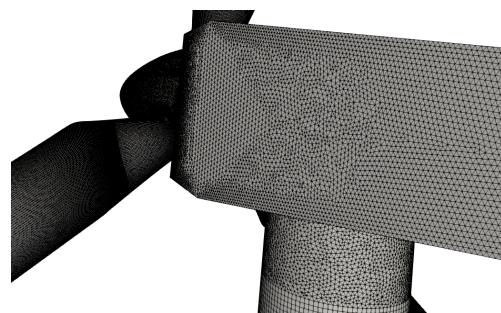
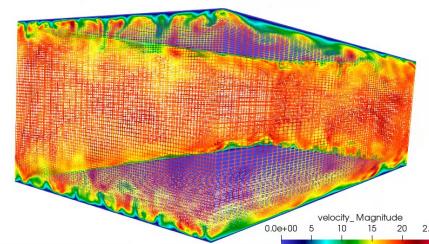
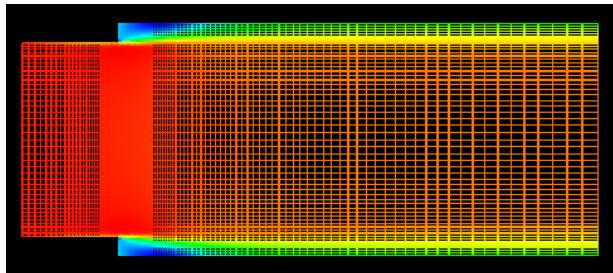
- Degree-of-freedom (DOF) for:
 - Cell-centered: Stencil is based on an element:face:element
 - DOFs at vertices of elements, or “nodes”, element:node (CVFEM, FEM), edge:node (EBVC)





Simple vs Complex Domains

- Some Geometries are simple, while others are complex





Simple One-Dimensional Diffusion Example

Let us introduce (briefly) the Method of Manufactured Solutions (MMS)

- Consider a simple diffusion equation:

$$\frac{d^2\phi}{dx^2} = 0$$

with a presumed, or manufactured solution:

$$\phi^{MMS}(x) = x^2$$

$$\frac{d^2\phi^{MMS}}{dx^2} = S^{MMS} = 2$$

Gauss-Divergence

$$\int \frac{d^2\phi^h}{dx^2} dV = \int S^{MMS} dV \quad \int \frac{d\phi^h}{dx} dS = \int S^{MMS} dV$$

Discrete flux-based form at the sub-control volume surface (scs) and sub-control volume IP:

With error: $\epsilon = \phi^{MMS} - \phi^h$

$$\sum_{scsIP} \frac{d\phi^h}{dx} A_{scsIP} = \sum_{scvIP} S^{MMS} V_{scvIP}$$

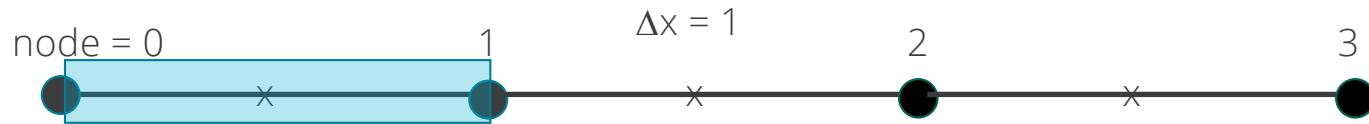


First....

- Let's review matrix assembly from a previous lecture, with a focus on a verification "patch test"
- Recall, a linear basis was able to reconstruct a quadratic solution perfectly when the norm was computed at the degree of freedom locations

Simple One-D MMS Example; Edge-based Assemble of Diffusion

We will use a linear basis and solve this system over a small patch of linear bar elements



Recall, simplified equation with RHS from source term already provided

$$\sum_{scsIP} \frac{d\phi^h}{dx}_{scsIP} = \sum_{scvIP} S^{MMS} \frac{\Delta x}{2}$$

Iterate elements with a simple left-hand side rule: $+= L$ and $-= R$
 (note implicit conservation statement)

$$\frac{d\phi^h}{dx}_{scsIP} = \frac{\phi_1^h - \phi_0^h}{\Delta x} + C_0(\Delta x) + \dots$$

$$LHS_0+ = \frac{\phi_1^h - \phi_0^h}{\Delta x}$$

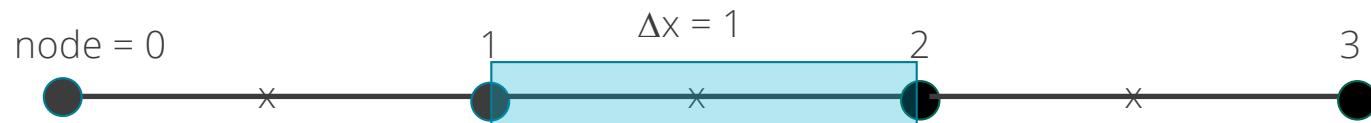
$$LHS_1- = \frac{\phi_1^h - \phi_0^h}{\Delta x}$$

Edge 1



Simple One-D MMS Example; Edge-based Assemble of Diffusion

We will use a linear basis and solve this system over a small patch of linear bar elements



Recall, simplified equation with RHS from source term already provided

$$\sum_{scsIP} \frac{d\phi^h}{dx}_{scsIP} = \sum_{scvIP} S^{MMS} \frac{\Delta x}{2}$$

Iterate elements with a simple left-hand side rule: $+= L$ and $-= R$

(note implicit conservation statement)

$$\frac{d\phi^h}{dx}_{scsIP} = \frac{\phi_2^h - \phi_1^h}{\Delta x} + C_0(\Delta x) + \dots$$

$$LHS_1+ = \frac{\phi_2^h - \phi_1^h}{\Delta x}$$

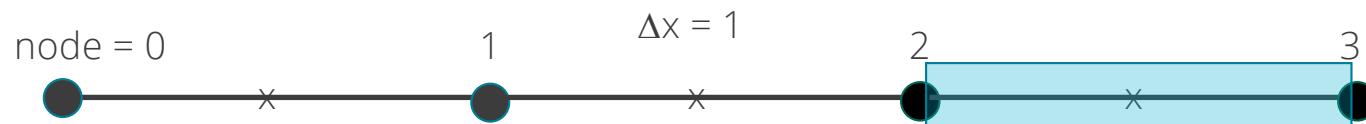
$$LHS_2- = \frac{\phi_2^h - \phi_1^h}{\Delta x}$$

Edge 2



Simple One-D MMS Example; Edge-based Assemble of Diffusion

We will use a linear basis and solve this system over a small patch of linear bar elements



Recall, simplified equation with RHS from source term already provided

$$\sum_{scsIP} \frac{d\phi^h}{dx}_{scsIP} = \sum_{scvIP} S^{MMS} \frac{\Delta x}{2}$$

Iterate elements with a simple left-hand side rule: += L and -=R

(note implicit conservation statement)

$$\frac{d\phi^h}{dx}_{scsIP} = \frac{\phi_3^h - \phi_2^h}{\Delta x} + C_0(\Delta x) + \dots$$

$$LHS_2+ = \frac{\phi_3^h - \phi_2^h}{\Delta x}$$

$$LHS_3- = \frac{\phi_3^h - \phi_2^h}{\Delta x}$$

Edge 3



Simple One-D MMS Example; Collect all of the terms

$$LHS_0+ = \frac{\phi_1^h - \phi_0^h}{\Delta x}$$

$$LHS_1- = \frac{\phi_1^h - \phi_0^h}{\Delta x}$$

$$LHS_1+ = \frac{\phi_2^h - \phi_1^h}{\Delta x}$$

$$LHS_2- = \frac{\phi_2^h - \phi_1^h}{\Delta x}$$

$$LHS_2+ = \frac{\phi_3^h - \phi_2^h}{\Delta x}$$

$$LHS_3- = \frac{\phi_3^h - \phi_2^h}{\Delta x}$$

$$A = \frac{1}{\Delta x} \begin{bmatrix} -1 & +1 & 0 & 0 \\ +1 & -2 & +1 & 0 \\ 0 & +1 & -2 & 1 \\ 0 & 0 & +1 & -1 \end{bmatrix}$$

Fully assembled matrix

$$\frac{1}{\Delta x} \begin{bmatrix} -1 & +1 & 0 & 0 \\ +1 & -2 & +1 & 0 \\ 0 & +1 & -2 & 1 \\ 0 & 0 & +1 & -1 \end{bmatrix} \begin{bmatrix} \phi_0 \\ \phi_1 \\ \phi_2 \\ \phi_3 \end{bmatrix} = \begin{bmatrix} 2 \\ 2 \\ 2 \\ 2 \end{bmatrix}$$

Fully assembled system

$$\frac{1}{\Delta x} \begin{bmatrix} +1 & 0 & 0 & 0 \\ +1 & -2 & +1 & 0 \\ 0 & +1 & -2 & 1 \\ 0 & 0 & 0 & +1 \end{bmatrix} \begin{bmatrix} \phi_0 \\ \phi_1 \\ \phi_2 \\ \phi_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 2 \\ 2 \\ 9 \end{bmatrix}$$

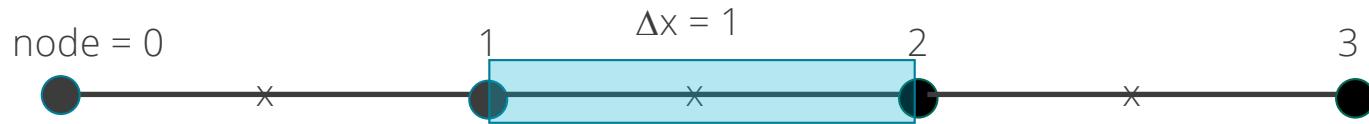
Correct for BCs

$$\phi^h = \begin{bmatrix} 0 \\ 1 \\ 4 \\ 9 \end{bmatrix}$$

Solution is exact.... Linear basis exactly captures a quadratic solution

Simple One-D Advection; Edge-based Assembly

Similar approach: Iteration of edges, definition of Left and Right node, etc.



$$\int w \frac{\partial \rho u_j \phi}{\partial x_j} dV \approx \sum_{ip} (\rho u_j)_{ip} \phi_{ip} n_j dS \approx \sum_{ip} \dot{m}_{ip} \phi_{ip}$$

$$\dot{m}_{ip} = \rho u_j n_j dS|_{ip} = \rho u_j A_j|_{ip}$$

We may also create a local matrix data structure, A[4]: With:

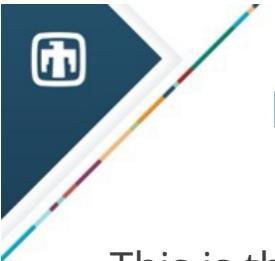
- A[0] += lhsFac (row node 1, "central-coeff")
- A[1] += lhsFac (row node 1, "east-coeff")
- A[2] -= lhsFac (row node 2, "west-coeff")
- A[3] -= lhsFac (row node 2, "central-coeff")

$$\text{With: } lhsFac = \frac{\dot{m}}{2}$$

$$LHS_1+ = \frac{\dot{m}}{2} (\phi_2^h + \phi_2^h)$$

$$LHS_2- = \frac{\dot{m}}{2} (\phi_1^h + \phi_2^h)$$

Edge 2

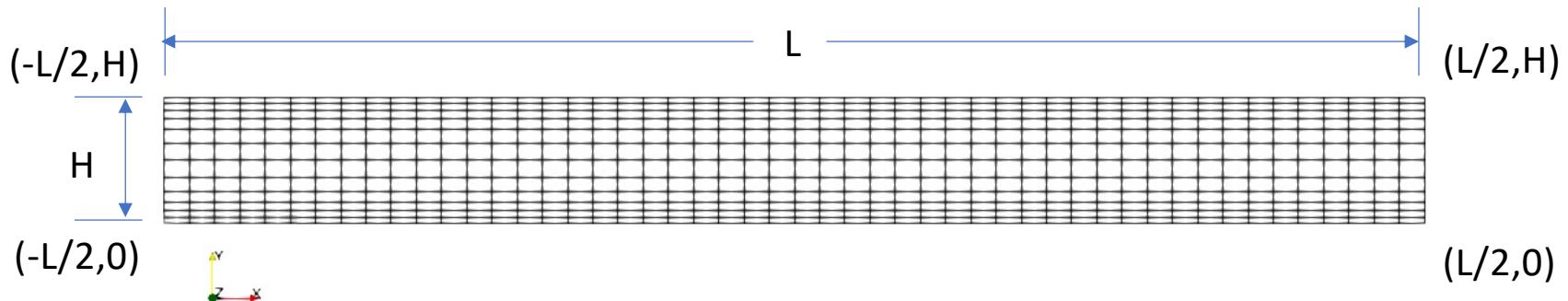


Let's Revisit a Classic 2-dimensional Channel Flow (laminar)

This is the 2d_quad4_channel laboratory:

https://github.com/NaluCFD/Nalu/blob/master/reg_tests/test_files/laboratory/2d_quad4_channel/write_up/2d_quad4_channel_laboratory.pdf

- No-slip (walls) top and bottom
- Periodic (left and right) with a constant body force ...or... Open (left and right pressure drop)





Derivation

3.1 Analytical Velocity Profile

Given the assumptions provided in the introduction, the streamwise velocity equation reduces to,

$$\frac{dP}{dx} = \mu \frac{d^2 u_x}{dy^2}. \quad (4)$$

This equation can be integrated twice to obtain,

$$u_x(y) = \frac{1}{\mu} \frac{dP}{dx} \frac{y^2}{2} + k_1 y + k_2, \quad (5)$$

where k_1 and k_2 are constants of integration that are obtained through the application of boundary conditions, $u_x(y = 0) = 0$ and $u_x(y = H) = 0$. Therefore, the final expression for the streamwise velocity is,

$$u_x(y) = \frac{1}{2\mu} \frac{dP}{dx} [y^2 - Hy]. \quad (6)$$

The above equation can be used to determine the vertical location at which the maximum velocity is found via solving $\frac{du_x}{dy} = 0$, or u_x^{max} occurs at $H/2$. The functional form for the maximum velocity is,

$$u_x^{max} = \frac{1}{8\mu} \frac{dP}{dx} H^2. \quad (7)$$



Derivation

3.1 Analytical Velocity Profile

Given the assumptions provided in the introduction, the streamwise velocity equation reduces to,

$$\frac{dP}{dx} = \mu \frac{d^2 u_x}{dy^2}. \quad (4)$$

This equation can be integrated twice to obtain,

$$u_x(y) = \frac{1}{\mu} \frac{dP}{dx} \frac{y^2}{2} + k_1 y + k_2, \quad (5)$$

where k_1 and k_2 are constants of integration that are obtained through the application of boundary conditions, $u_x(y = 0) = 0$ and $u_x(y = H) = 0$. Therefore, the final expression for the streamwise velocity is,

$$u_x(y) = \frac{1}{2\mu} \frac{dP}{dx} [y^2 - Hy]. \quad (6)$$

The above equation can be used to determine the vertical location at which the maximum velocity is found via solving $\frac{du_x}{dy} = 0$, or u_x^{max} occurs at $H/2$. The functional form for the maximum velocity is,

$$u_x^{max} = \frac{1}{8\mu} \frac{dP}{dx} H^2. \quad (7)$$



Sizing

- Let's run at $Re_\tau = 10$

Reynolds number:

- Ratio of inertial to viscous forces

$$Re_D = \frac{\rho u D}{\mu}$$

- In this channel, define Reynolds number based on the wall friction velocity, u_τ (at $H/2$ length scale)

$$Re_\tau = \frac{\rho u_\tau H/2}{\mu}$$

4 Results

Let us test a simulation in which the Reynolds number based on wall friction velocity, u^τ , and half-channel height $H/2$, is ten: $Re^\tau = 10$. By constraining the Reynolds number, wall friction velocity, and density, the consistent viscosity is obtained via,

$$\mu = \frac{\rho u^\tau H}{2Re^\tau}. \quad (8)$$

To obtain the required pressure gradient, we exercise the relationship,

$$u^\tau = \sqrt{\frac{\tau_w}{\rho}}, \quad (9)$$

along with global momentum balance,

$$\int \frac{dP}{dx} dV = \int \tau_w dA, \quad (10)$$

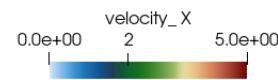
with $dV = LH$ and $dA = 2L$, to obtain the relationship between required pressure gradient and wall shear stress, $\frac{dP}{dx} = 2\tau_w$.



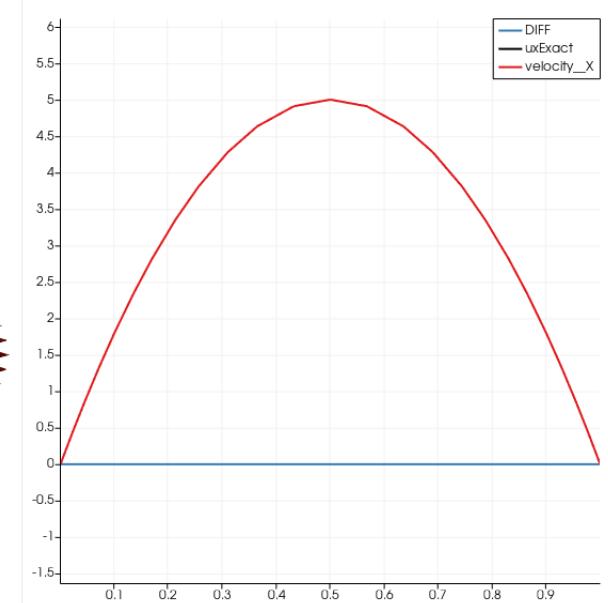
Simulation Results

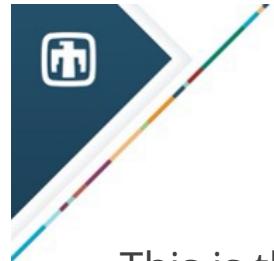
Analytical Expression:

$$u_x(y) = \frac{1}{2\mu} \frac{dP}{dx} [y^2 - Hy].$$



Zero Error when evaluating the norm at the nodes (in general, we may integrate over a domain and choose integration points that are not *nodally-lumped*)



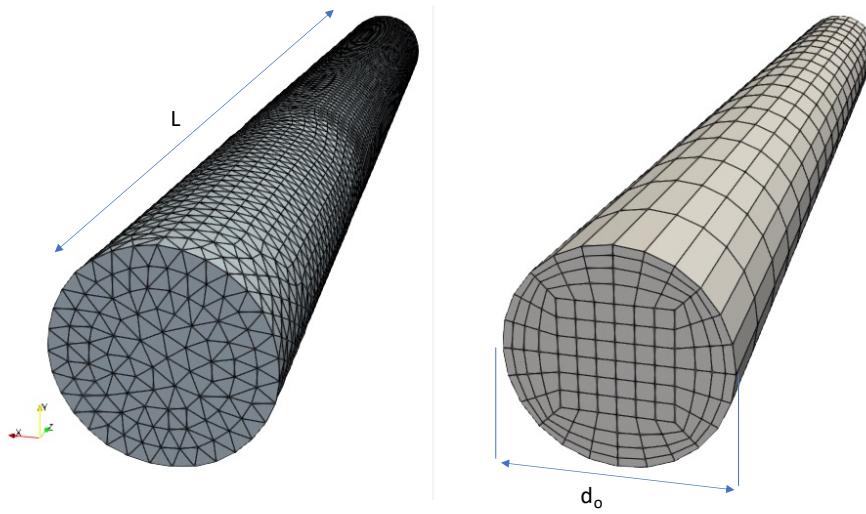


Let's Revisit a Classic 3-dimensional Pipe Flow (laminar)

This is the 3d_tet4_pipe laboratory:

https://github.com/NaluCFD/Nalu/tree/master/reg_tests/test_files/laboratory/3d_tet4_pipe/write_up

- No-slip (walls) outer pipe
- Periodic (left and right) with a constant body force ...or... Open (left and right pressure drop)
- We have Tet4 (left) and a Hex8 (right) mesh





Derivation

3.1 Analytical Velocity Profile

Given the assumptions provided in the introduction, the axial momentum equation reduces to,

$$\frac{dP}{dz} = \frac{1}{r} \frac{\partial r \tau_{rz}}{\partial r}, \quad (11)$$

where $\tau_{rz} = \mu \frac{\partial u_z}{\partial r}$. Integration once results in the following equation,

$$\frac{r^2}{2} \frac{dP}{dz} = r \mu \frac{\partial u_z}{\partial r} + k_1, \quad (12)$$

which, based on $\frac{\partial u_z}{\partial r} = 0$ at the centerline $r = 0$, yields $k_1 = 0$. A second integration results in,

$$u_z(r) = \frac{r^2}{4\mu} \frac{dP}{dz} + k_2. \quad (13)$$

Applying $u_z(r = R) = 0$, yields the following expression for the axial velocity as a function of radial position,

$$u_z(r) = \frac{1}{4\mu} \frac{dP}{dz} [r^2 - R^2]. \quad (14)$$

The above equation can be used to determine the location at which the maximum velocity is found via solving $\frac{du_z}{dr} = 0$, or u_z^{max} occurs at $r = 0$. The functional form for the maximum velocity is,

$$u_x^{max} = -\frac{R^2}{4\mu} \frac{dP}{dZ}. \quad (15)$$



Sizing

- Let's run at $Re_\tau = 20$

Reynolds number:

- Ratio of inertial to viscous forces

$$Re_D = \frac{\rho u D}{\mu}$$

- In this channel, define Reynolds number based on the wall friction velocity, u_τ (at d_o length scale)

$$Re_\tau = \frac{\rho u_\tau d_o}{\mu}$$

4 Results

Let us test a simulation in which the Reynolds number based on wall friction velocity, u^τ , and diameter of the pipe (0.02 m) is $Re^\tau = 20$. In this mesh configuration, the length L is 0.2 m . By constraining the Reynolds number, wall friction velocity, and density, the consistent viscosity is obtained via,

$$\mu = \frac{\rho u^\tau d_o}{Re^\tau}. \quad (16)$$

To obtain the required pressure gradient, we exercise the relationship,

$$u^\tau = \sqrt{\frac{\tau_w}{\rho}}, \quad (17)$$

along with global momentum balance,

$$\int \frac{dP}{dz} dV = \int \tau_w dA, \quad (18)$$

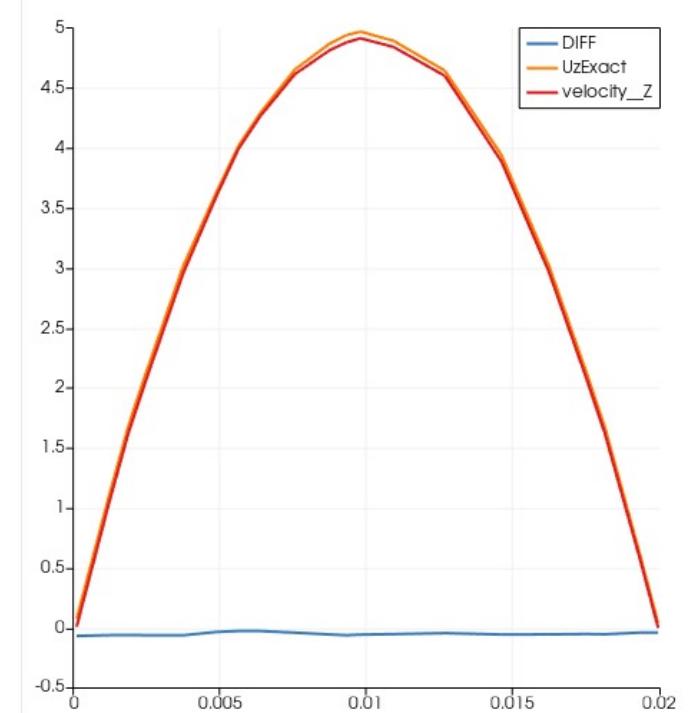
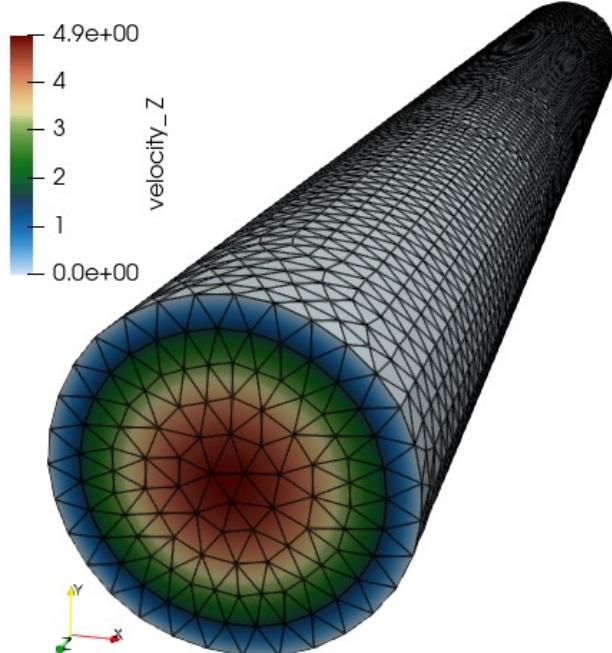
with $dV = \pi r^2 L$ and $dA = 2\pi r L$, to obtain the relationship between required pressure gradient and wall shear stress, $\frac{dP}{dz} = 2 \frac{\tau_w}{r}$.

Simulation Results

Analytical Expression:

$$u_z(r) = \frac{1}{4\mu} \frac{dP}{dz} [r^2 - R^2].$$

Non-zero error when evaluating the norm at the nodes...



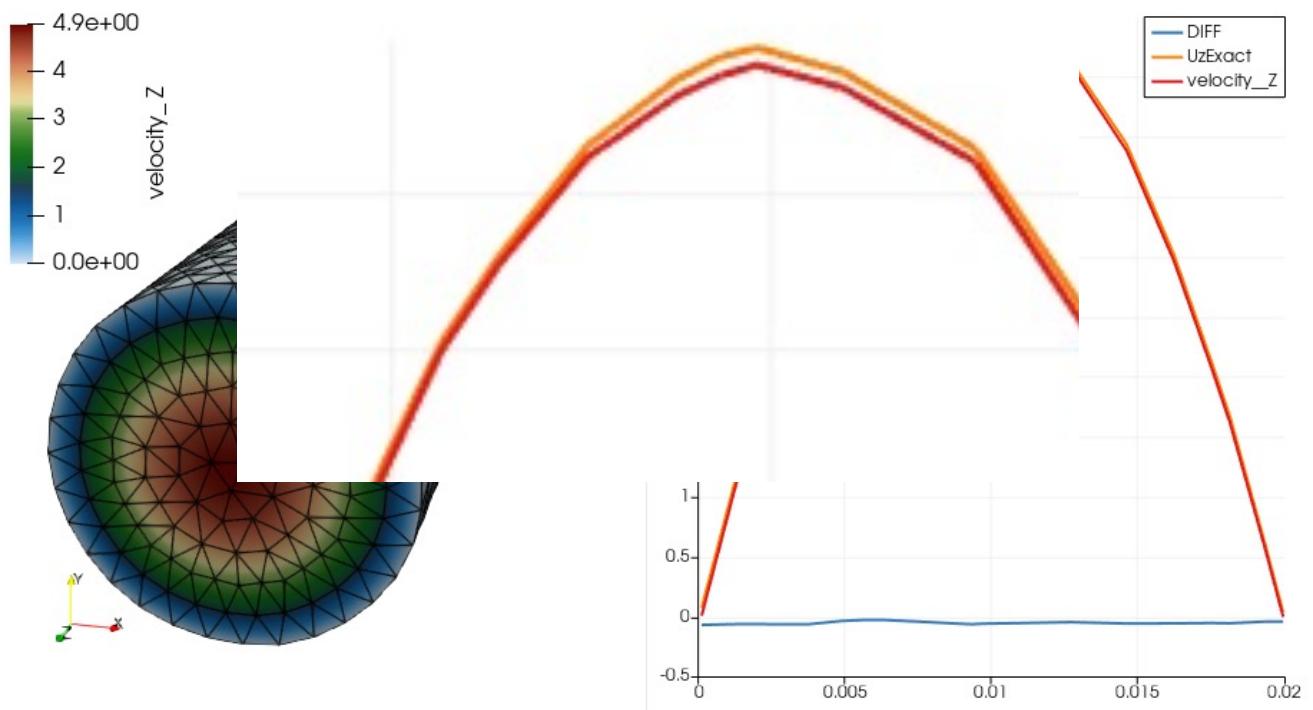
Simulation Results

Analytical Expression:

$$u_z(r) = \frac{1}{4\mu} \frac{dP}{dz} [r^2 - R^2].$$

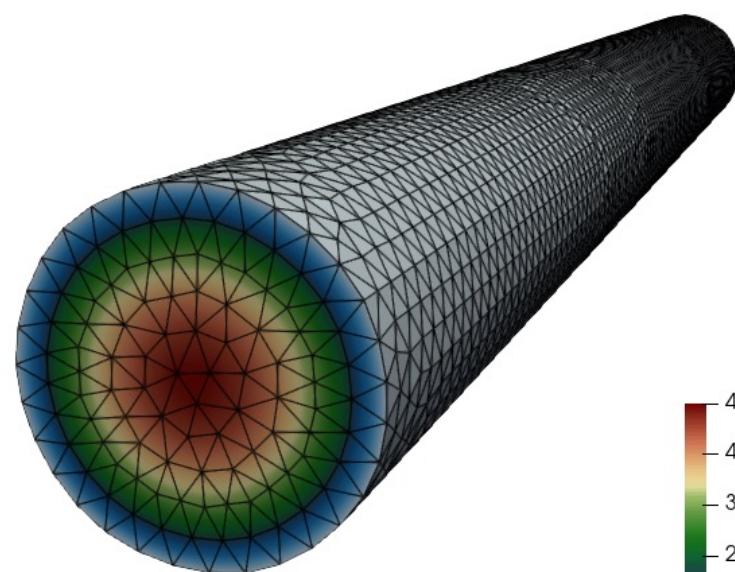
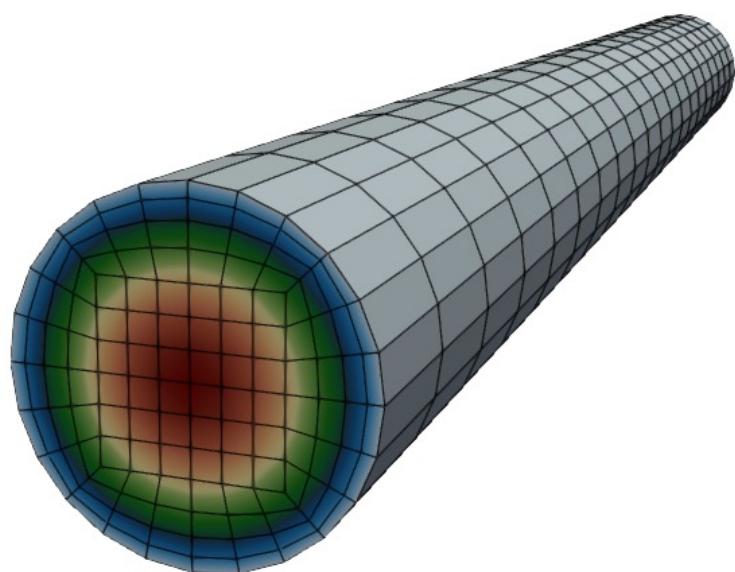
Non-zero error when evaluating the norm at the nodes...

What happened?





What About Tet4 vs Hex8?

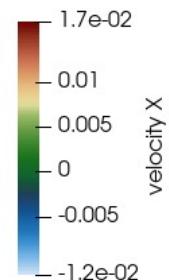
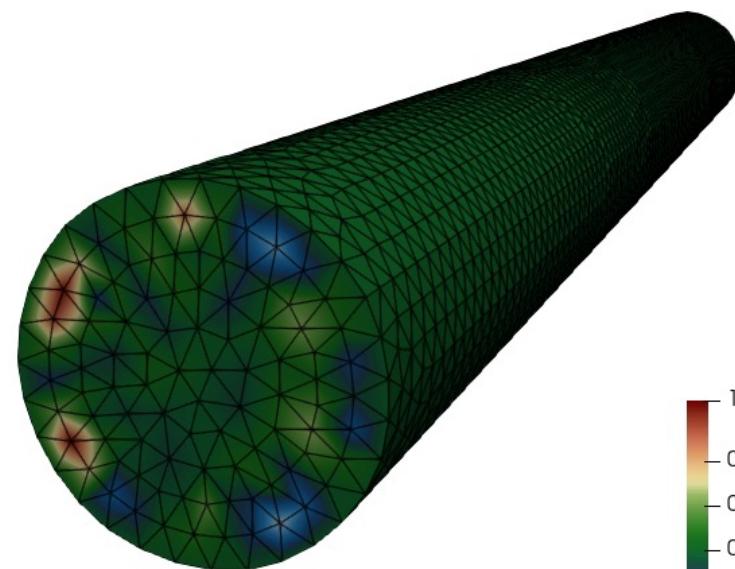
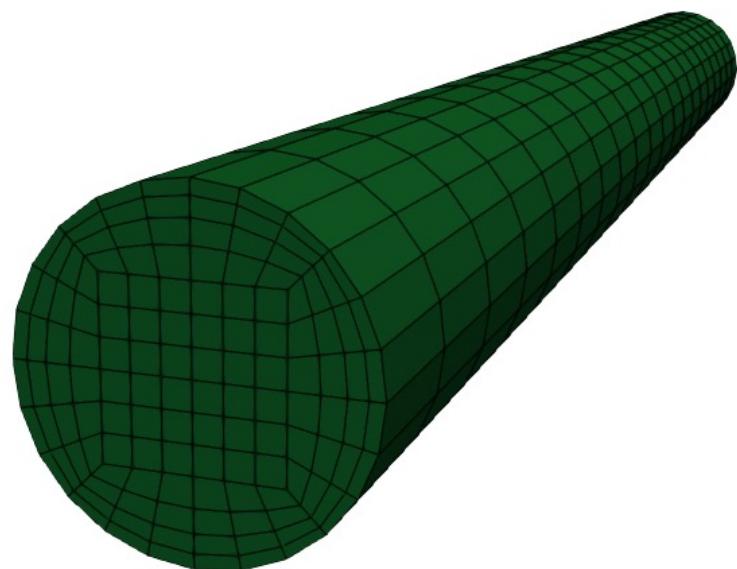


velocity z

4.9e+00
4
3
2
1
0.0e+00



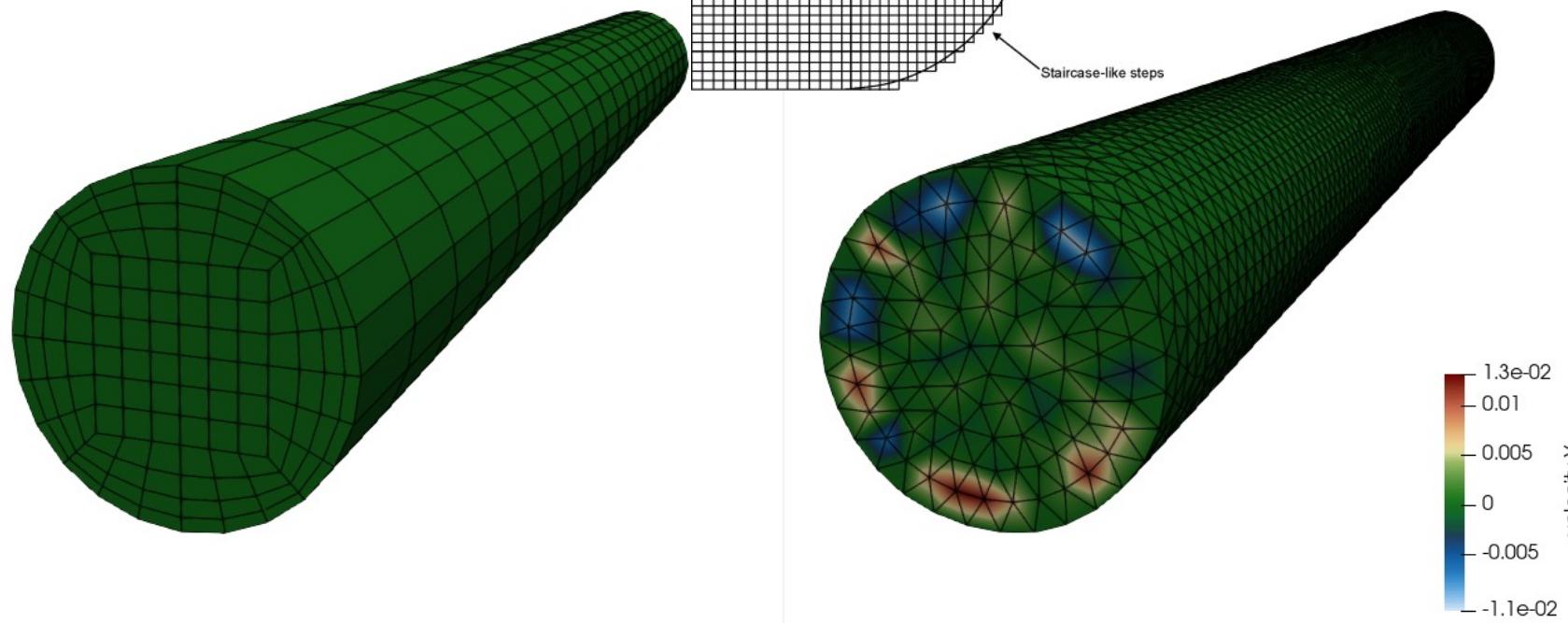
What About Tet4 vs Hex8?



~0.3% error in non-streamwise velocity components in the tet-based result, ~0% for Hex8!!!



What About Tet4 vs Hex8?

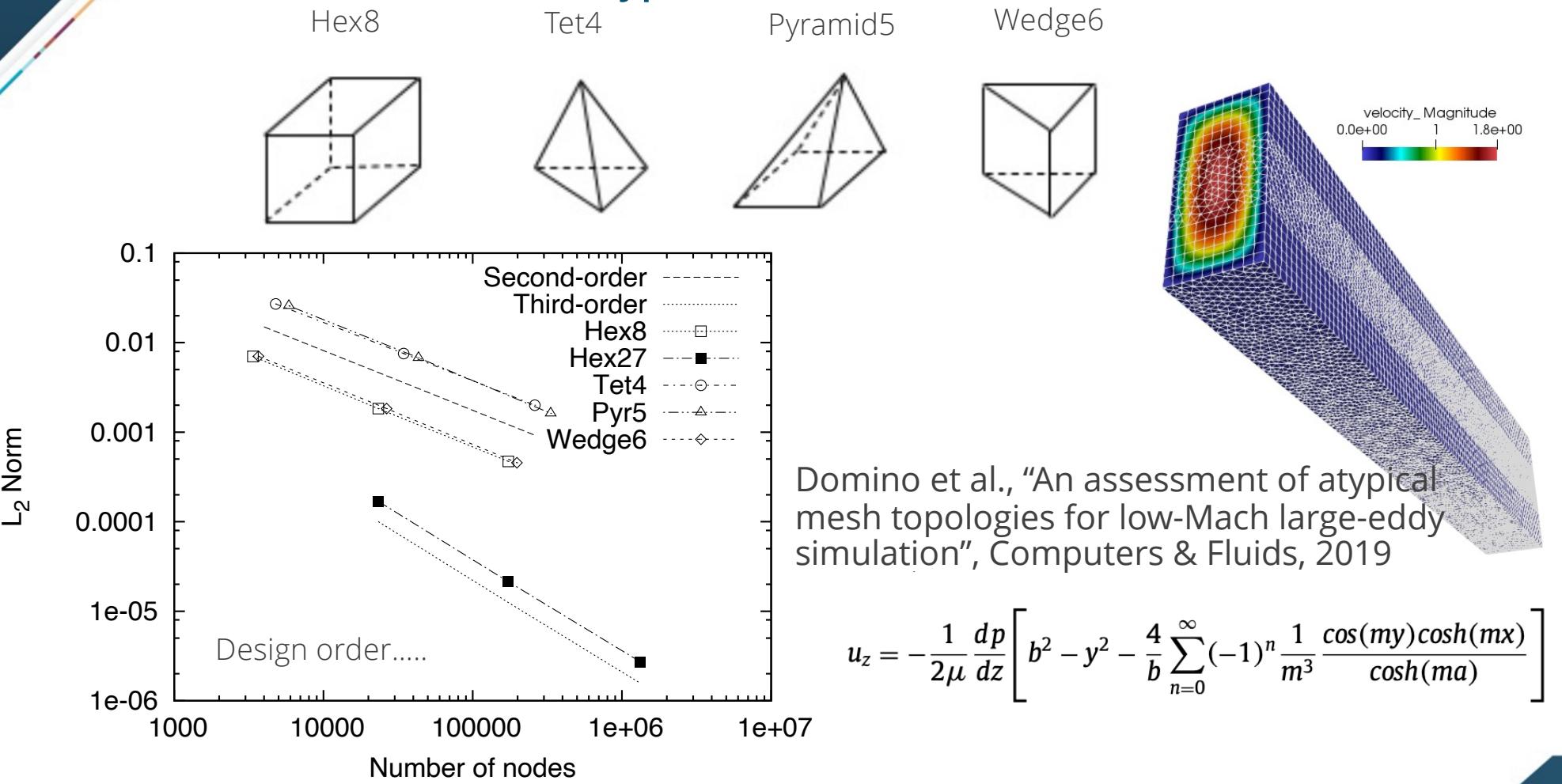


Jiyuan Tu, Guan-Heng Yeoh, Chaoqun Liu,
Chapter 4 - CFD Mesh Generation: A Practical Guideline,
Editor(s): Jiyuan Tu, Guan-Heng Yeoh, Chaoqun Liu,
Computational Fluid Dynamics (Third Edition),
Butterworth-Heinemann, 2018, Pages 125-154,

~0.3% error in non-streamwise velocity components in the tet-based result, ~0% for Hex8!!!



Unstructured Mesh Types (Review) and Results





Conclusions

- Geometric fidelity for curved surfaces can be very important
- Unstructured meshing generally affords easier meshing, with increase throughput (especially when simulating complex geometries)
- Hex-based errors may be lower than Tet-based simulation results, however, balanced with ease of meshing
- The key is that all methods are “design-order”, i.e., demonstrate expected reduction of errors as a function of mesh size