

# Lecture-2: Variational Formulation

## MA 573: Finite Element Methods for PDEs

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In the previous lecture, we have introduced notion for weak derivatives. Then, we have used weak derivative to introduce Hilbertian Sobolev spaces  $H^m$  satisfying following inclusion relation

$$L^2(\Omega) \supseteq H^1(\Omega) \supseteq H^2(\Omega) \supseteq \dots \supseteq H^m(\Omega).$$

In this lecture, we will introduce variational formulation for boundary value problem, which is the first step towards the finite element approximation.

### 1 Introduction

We first introduce variational formulation for a simple ODE

$$-y'' + y = f, \quad x \in \Omega = (a, b) \quad (1.1)$$

with homogeneous Dirichlet boundary condition

$$y(a) = 0 \quad \& \quad y(b) = 0. \quad (1.2)$$

For  $f \in L^2(\Omega)$ , we assume that  $y \in H$  is the solution for the BVP (1.1)-(1.2). Here,  $H$  is a suitable function space, which need to be introduced.

For the variational formulation, we multiply the equation (1.1) by  $v \in V$  and then integrate from  $a$  to  $b$  to have

$$\int_a^b \{-y'' + y\}v dx = \int_a^b f v dx.$$

Integration by parts, we obtain

$$-[y'v]_a^b + \int_a^b \{y'v' + yv\} dx = \int_a^b f v dx. \quad (1.3)$$

Note that we don't have the information of  $y'$  at  $a$  and  $b$ , therefore, we set  $v(a) = 0 = v(b)$  so that (1.4) yields

$$\int_a^b \{y'v' + yv\} dx = \int_a^b f v dx \quad \text{OR} \quad A(y, v) = (f, v), \quad (1.4)$$

where  $A(y, v) = \int_a^b \{y'v' + yv\} dx$  &  $(f, v) = \int_a^b f v dx$ .

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**Remark 1.1 (Multiplier Space)** Observe, due to Cauchy-Schwarz inequality, that

$$\begin{aligned}
|A(y, v)| &\leq \left[ \int_a^b |y'|^2 dx \right]^{\frac{1}{2}} \left[ \int_a^b |v'|^2 dx \right]^{\frac{1}{2}} + \left[ \int_a^b |y|^2 dx \right]^{\frac{1}{2}} \left[ \int_a^b |v|^2 dx \right]^{\frac{1}{2}} \\
&= \|y'\|_{L^2(\Omega)} \|v'\|_{L^2(\Omega)} + \|y\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \\
&\leq (\|y'\|_{L^2(\Omega)} + \|y\|_{L^2(\Omega)}) (\|v'\|_{L^2(\Omega)} + \|v\|_{L^2(\Omega)}) \\
&= \|y\|_{H^1(\Omega)} \|v\|_{H^1(\Omega)}.
\end{aligned} \tag{1.5}$$

Clearly,  $|A(y, v)| < \infty$  for  $y$  and  $v$  belongs to  $H^1(\Omega)$ . Again  $y(a) = 0 = y(b)$  and  $v(a) = 0 = v(b)$  that is  $y$  and  $v$  vanishes on the boundary of domain  $\Omega = (a, b)$ . Therefore, we assume  $v \in H_0^1(\Omega)$  and hence, multiplier space  $V = H_0^1(\Omega)$ . Next question arises whether  $y \in H_0^1(\Omega)$  and satisfies (1.4). More precisely, we need to discuss the existence and uniqueness of the problem: *Find  $y \in H_0^1(\Omega)$  such that*

$$A(y, v) = L(v) \quad \forall v \in H_0^1(\Omega), \tag{1.6}$$

where  $A : H_0^1(\Omega) \times H_0^1(\Omega) \rightarrow \mathbb{R}$  and  $L : H_0^1(\Omega) \rightarrow \mathbb{R}$  are functions defined by

$$A(y, v) = \int_a^b \{y'v' + yv\} dx \quad \& \quad L(v) = (f, v) = \int_a^b f v dx.$$

**Remark 1.2** Problem defined by (1.6) is known as variational formulation (problem) or weak formulation or integral formulation of the original BVP (1.1)-(1.2). Original equation (1.1) demands solution  $y$  to be twice differentiable, but variational formulation (1.6) requires first derivative of solution  $y$ . Due to this weaker requirement of the derivative, solution  $y$  satisfying (1.6) is called **weak solution** and the formulation is called weak formulation.

**Remark 1.3** Clearly, a solution to the BVP (1.1)-(1.2) is also a solution to weak formulation (1.6). What about the converse? Converse is not always true that is a weak solution need not be a solution to the BVP (1.1)-(1.2) because original equation requires higher derivative of the solution. Therefore, a solution to the equation (1.1) is called **strong solution**.

**Remark 1.4** For a continuous function  $f$ , we know that there exists a unique  $y \in C^2[a, b]$  satisfying BVP (1.1)-(1.2). For  $f \in L^2(\Omega)$ , we do not know the existence theory for the BVP (1.1)-(1.2). So, we first try to ensure existence of a weak solution and then through that weak solution we try to prove existence of a strong solution.

Existence of a weak solution requires a celebrated result known as **Lax-Milgram Lemma**, which is an extension of Riesz representation lemma. Before, we proceed further, let us introduce few basic definitions.

**Definition 1.1** Let  $H$  be a normed linear space with norm  $\|\cdot\|_H$ . A map  $A : H \times H \rightarrow \mathbb{R}$  is said to be:

(a) **Bi-linear** if  $A$  is linear in both the arguments, that is, for any scalars  $c_1$  and  $c_2$ , we have

$$\begin{aligned}
T(c_1 w_1 + c_2 w_2, v) &= c_1 T(w_1, v) + c_2 T(w_2, v) \quad \forall w_1, w_2, v \in H \text{ and} \\
T(w, c_1 v_1 + c_2 v_2) &= c_1 T(w, v_1) + c_2 T(w, v_2) \quad \forall w, v_1, v_2 \in H.
\end{aligned}$$

(b) **Continuous/Bounded** if there exists a constant  $C > 0$  such that

$$|A(u, v)| \leq C \|u\|_H \|v\|_H \quad \forall u, v \in H.$$

(c) **Positive/Coercive** if there exists a constant  $M > 0$  such that

$$A(u, u) \geq M \|u\|_H^2 \quad \forall u \in H.$$

**Definition 1.2** Let  $H$  be a normed linear space with norm  $\|\cdot\|_H$ . A map  $L : H \rightarrow \mathbb{R}$  is said to be continuous/bounded if there exists a positive constant  $m$  such that

$$|L(v)| \leq m \|v\|_H \quad \forall v \in H.$$

Now, it is a good time to recall the following Lax-Milgram theorem which tells about the existence of a unique weak solution.

**Theorem 1.1** For a given Hilbert space  $H$ , a continuous, coercive bilinear map  $A(\cdot, \cdot) : H \times H \rightarrow \mathbb{R}$  and a continuous linear functional  $L : H \rightarrow \mathbb{R}$ , there exists an unique  $y \in H$  such that

$$A(y, v) = L(v) \quad \forall v \in H.$$

It is easy to verify that map  $A(\cdot, \cdot)$  defined by (1.6) is bi-linear. Again, from (1.5), it is clear that  $A(\cdot, \cdot)$  is continuous in  $H^1(\Omega)$  and so it is continuous in  $H_0^1(\Omega) \subset H^1(\Omega)$ . Further, we observe that

$$\begin{aligned} A(v, v) &= \int_a^b |v'|^2 dx + \int_a^b |v|^2 dx \\ &= \|v'\|_{L^2(\Omega)}^2 + \|v\|_{L^2(\Omega)}^2 \\ &= \|v\|_{H^1(\Omega)}^2 \geq m \|v\|_{H^1(\Omega)}^2, \quad m = \frac{1}{2}. \end{aligned} \quad (1.7)$$

Thus,  $A(\cdot, \cdot)$  is positive in  $H^1(\Omega)$  and so it is positive in  $H_0^1(\Omega) \subset H^1(\Omega)$ . For the operator  $L$  defined in (1.6), we find that

$$\begin{aligned} |L(v)| &= \left| \int_a^b f v dx \right| \leq \left( \int_a^b f^2 dx \right)^{\frac{1}{2}} \left( \int_a^b v^2 dx \right)^{\frac{1}{2}} \\ &= \|f\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \\ &\leq C \|v\|_{H^1(\Omega)}, \quad C = \|f\|_{L^2(\Omega)} < \infty. \end{aligned} \quad (1.8)$$

Here, we have used the facts that  $f \in L^2(\Omega)$  and  $\|v\|_{L^2(\Omega)} \leq \|v\|_{H^1(\Omega)}$ .

**Remark 1.5 (Weak Solution Space)** All assumptions in Lax-Milgram Lemma are satisfied in  $H_0^1(\Omega)$  by the bilinear map  $A(\cdot, \cdot)$  and operator  $L$  defined as in (1.6). Therefore, there exists unique  $y \in H_0^1(\Omega)$  such that

$$A(y, v) = L(v) \quad \forall v \in H_0^1(\Omega).$$

Hence, weak solution corresponding to the original BVP (1.1)-(1.2) belongs to  $H_0^1(\Omega)$ . **If we somehow know that strong solution to the BVP (1.1)-(1.2) exists.** Let  $y$  be the strong solution. Then the strong solution will be a weak solution and it will belong to function space  $H_0^1(\Omega)$ . For the strong solution  $y$ , from (1.1), we have

$$y'' = f - y' \in L^2(\Omega) \quad \text{since } f \in L^2(\Omega) \text{ \& } y \in H_0^1(\Omega).$$

Thus strong solution, if it exists, belongs to the space  $H^2(\Omega) \cap H_0^1(\Omega)$ . **Have we proved that strong solution to the BVP (1.1)-(1.2) exists? Not yet!** We proceed as discussed in Remark 1.4.

For  $f \in L^2(\Omega)$ , any function  $y \in H_0^1(\Omega)$  satisfying variational problem (1.6) belongs to function space  $H^2(\Omega)$ . This result is known as **Elliptic Regularity** theorem. That is a weak solution becomes twice differentiable. Then from (1.6) and integration by parts, we obtain

$$\left[ y'v \right]_a^b - \int_a^b \{ y'' - y \} v dx = \int_a^b f v dx \quad \forall v \in H_0^1(\Omega).$$

Using the fact that  $v(a) = 0 = v(b)$ , we have

$$\int_a^b \{ -y'' + y \} v dx = \int_a^b f v dx \quad \forall v \in H_0^1(\Omega).$$

Equivalently

$$\int_a^b \{ -y'' + y - f \} v dx = 0 \quad \forall v \in H_0^1(\Omega).$$

Due to arbitrariness of  $v$ , we find that

$$-y'' + y - f = 0.$$

Hence, a weak solution is also a strong solution.

**Remark 1.6** *In general, a weak solution need not be a strong solution. But, in the present example we have seen that a weak solution is also a strong solution. Now onwards, we will assume that original BVP and its weak formulation are equivalent.*

**Example 1.1** For  $f \in L^2(\Omega)$ , consider following BVP

$$-y''(x) + P(x)y'(x) + Q(x)y(x) = f(x), \quad x \in \Omega$$

with boundary condition

$$y(a) = \alpha \quad \& \quad y(b) = \beta.$$

Here, coefficient functions  $P$  and  $Q$  are sufficiently smooth function. For suitable choice of  $P$  and  $Q$ , justify that above BVP has an unique weak solution.

*Solution.* Consider the given BVP

$$-y''(x) + P(x)y'(x) + Q(x)y(x) = f(x), \quad x \in \Omega \tag{1.9}$$

with boundary condition

$$y(a) = \alpha \quad \& \quad y(b) = \beta. \tag{1.10}$$

For weak formulation, we multiply equation (1.9) by suitable  $v$  and then integrate in the interval  $[a, b]$  to obtain

$$\int_a^b \{ -y''(x) + P(x)y'(x) + Q(x)y(x) \} v dx = \int_a^b f v dx.$$

Integration by parts yields

$$-\left[y'v\right]_a^b + \int_a^b y'v'dx + \int_a^b Py'vdx + \int_a^b Qyv'dx = \int_a^b fvd x. \quad (1.11)$$

Here,  $y'(a)$  and  $y'(b)$  are not given, so set  $v(a) = 0 = v(b)$  in (1.11) to have

$$A(y, v) = L(v), \quad (1.12)$$

where

$$A(y, v) = \int_a^b \left\{ y'v' + Py'v + Qyv \right\} dx \quad \& \quad L(v) = \int_a^b fvd x. \quad (1.13)$$

Further, it is easy exercise to verify that  $A(\cdot, \cdot)$  is bilinear and  $L(\cdot)$  is linear.

Next, we try to fix the multiplier and weak solution space. Using Cauchy-Schwarz inequality, we have

$$\begin{aligned} |A(w, v)| &\leq \|y'\|_{L^2(\Omega)} \|v'\|_{L^2(\Omega)} + \|p\|_{\infty} \|y'\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \\ &\quad + \|Q\|_{\infty} \|y\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \\ &\leq \|y'\|_{L^2(\Omega)} \|v'\|_{L^2(\Omega)} + m_1 \|y'\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} + m_2 \|y\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \\ &\leq \max\{1, m_1, m_2\} \left( \|y'\|_{L^2(\Omega)} + \|y\|_{L^2(\Omega)} \right) \left( \|v'\|_{L^2(\Omega)} + \|v\|_{L^2(\Omega)} \right) \\ &\leq C \|y\|_{H^1(\Omega)} \|v\|_{H^1(\Omega)}. \end{aligned} \quad (1.14)$$

For the operator  $L(\cdot)$ , we have

$$|L(v)| \leq \|f\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \leq C \|v\|_{H^1(\Omega)}, \quad C = \|f\|_{L^2(\Omega)}. \quad (1.15)$$

From (1.14) and (1.15), we find that  $A(\cdot, \cdot)$  and  $L(\cdot)$  are well defined in  $H^1(\Omega) \times H^1(\Omega)$  and  $H^1(\Omega)$ , respectively. More precisely,  $A : H^1(\Omega) \times H^1(\Omega) \rightarrow \mathbb{R}$  and  $L : H^1(\Omega) \rightarrow \mathbb{R}$ , defined by (1.13), are continuous in  $H^1(\Omega)$ . Thus, multiplier  $v \in H_0^1(\Omega)$  as  $v(a) = 0 = v(b)$ . But,  $y \in H^1(\Omega)$  not  $y \in H_0^1(\Omega)$  as  $y(a) \neq 0$  and  $y(b) \neq 0$ . Hence, multiplier space is  $H_0^1(\Omega)$  and weak solution space is  $H^1(\Omega)$ .

In order to use Lax-Milgram theorem, we try to convert the weak solution space to  $H_0^1(\Omega)$  by a meaningful way. For given  $\alpha$  and  $\beta$ , we find a linear polynomial  $g$  such that  $g(a) = \alpha$  and  $g(b) = \beta$ . Then the function

$$w = y - g \in H_0^1(\Omega)$$

and satisfies following equation

$$A(w + g, v) = L(v), \quad (1.16)$$

which gives

$$A(w, v) = L(v) - A(g, v) = \tilde{L}(v) \quad \forall v \in H_0^1(\Omega), \quad (1.17)$$

where  $\tilde{L}(v) = L(v) - A(g, v)$ . It is worth to note that existence of  $w \in H_0^1(\Omega)$  satisfying (1.17) gives existence of weak solution  $y = w + g$  satisfying (1.12). **So, we concentrate on the existence of  $w$  satisfying (1.17).**

From (1.14), we observe that bilinear map  $A(\cdot, \cdot)$  is continuous in  $H_0^1(\Omega)$ . Now, it remains to verify positivity of bilinear map  $A(\cdot, \cdot)$ . Integration by parts, we obtain

$$\int_a^b P v' v dx = \int_a^b P \frac{1}{2} \frac{d}{dx} (v^2) dx = \left[ \frac{P}{2} v^2 \right]_a^b - \int_a^b \left( \frac{P}{2} \right)' v^2 dx.$$

For  $v \in H_0^1(\Omega)$ , we obtain

$$\int_a^b P v' v dx = - \int_a^b \left( \frac{P}{2} \right)' v^2 dx,$$

which together with definition of  $A(\cdot, \cdot)$ , for all  $v \in H_0^1(\Omega)$ , leads to

$$\begin{aligned} A(v, v) &= \int_a^b \left\{ (v')^2 + \left( Q - \frac{P'}{2} \right) v^2 \right\} dx \\ &= \int_a^b \left\{ (v')^2 + H v^2 \right\} dx, \quad H = Q - \frac{P'}{2} \\ &\geq \int_a^b (v')^2 dx, \quad \text{assuming } H = Q - \frac{P'}{2} \geq 0 \\ &\geq C \|v\|_{H^1(\Omega)}^2, \quad \text{due to Poincaré inequality.} \end{aligned} \quad (1.18)$$

Hence,  $A(\cdot, \cdot)$  is positive in  $H_0^1(\Omega)$ . For **Poincaré inequality**, see equation (2.7).

For the continuity of  $\tilde{L}$ , we use (1.14)-(1.15) to have

$$\begin{aligned} |\tilde{L}(v)| \leq |L(v)| + |A(g, v)| &\leq C \|v\|_{H^1(\Omega)} + C \|g\|_{H^1(\Omega)} \|v\|_{H^1(\Omega)} \\ &\leq \tilde{C} \|v\|_{H^1(\Omega)}. \end{aligned} \quad (1.19)$$

Thus,  $\tilde{L}$  is linear and continuous in  $H_0^1(\Omega)$ , which together with the fact that bilinear map  $A(\cdot, \cdot)$  is continuous and positive in  $H_0^1(\Omega)$  yields that there exists unique  $w \in H_0^1(\Omega)$  such that

$$A(w, v) = \tilde{L}(v) \quad \forall v \in H_0^1(\Omega).$$

Finally, existence of  $w$  gives existence of  $y = w + g \in H^1(\Omega)$  satisfying (1.12) and boundary condition  $y(a) = \alpha$  and  $y(b) = \beta$ .

## 2 Variation Formulation for Poisson's Equation

In this section, we consider a linear elliptic problem of the form

$$-\nabla \cdot (\nabla u(x, y)) = f(x, y) \quad \text{in } \Omega \quad (2.1)$$

subject to the boundary condition

$$u = 0 \quad \text{on } \partial\Omega. \quad (2.2)$$

The above problem is also called Dirichlet boundary value problem. The equations of the form (2.1)-(2.2) are often encountered in stationary heat conduction problems, material sciences and fluid dynamics. As a model problem, we consider the stationary heat conduction problem in  $\Omega$ . The smoothness of the solution  $u$ , depends on the smoothness of the given data  $f$  which is known as **elliptic regularity** of the solution. Concerning the problem (2.1)-(2.2), we have the following regularity result.

**Theorem 2.1 Regularity Result :** *For any  $f \in L^2(\Omega)$ , the Dirichlet problem has an unique solution  $u \in H^2(\Omega) \cap H_0^1(\Omega)$  and satisfies the following stability estimate*

$$\|u\|_{H^2(\Omega)} \leq C \|f\|_{L^2(\Omega)}.$$

As a first step towards the finite element approximation to (2.1)-(2.2), we first introduce the weak formulation. The following Green's formula is borrowed from Lecture-1 for our convenience.

$$\int_{\Omega} \Delta u v dx = - \int_{\Omega} \nabla u \cdot \nabla v dx + \int_{\partial\Omega} \frac{\partial u}{\partial \mathbf{n}} v ds, \quad (2.3)$$

where  $\frac{\partial u}{\partial \mathbf{n}} = \nabla u \cdot \mathbf{n}$  is the outward normal derivative of  $u$  on  $\partial\Omega$ . Let  $v \in H_0^1(\Omega)$ , then multiplying the given elliptic problem by  $v$  and further using Green's theorem, we have

$$\int_{\Omega} \nabla u \cdot \nabla v dx - \int_{\partial\Omega} \frac{\partial u}{\partial \mathbf{n}} v ds = \int_{\Omega} f v dx. \quad (2.4)$$

Since the information of  $\frac{\partial u}{\partial \mathbf{n}} = \nabla u \cdot \mathbf{n}$  is not given on the boundary  $\partial\Omega$ , we select  $v = 0$  on  $\partial\Omega$ , so that equation (2.4) reduces to

$$A(u, v) = L(v), \quad (2.5)$$

where

$$A(u, v) = \int_{\Omega} \nabla u \cdot \nabla v dx \text{ and } L(v) = \int_{\Omega} f v dx.$$

**Remark 2.1** The equation (2.5) involve only the first order partial derivative of  $u$  where as the original equation demands second order partial derivative of  $u$ . Therefore, (2.5) is called weak form of the original equation. Since the BVP (2.1)-(2.2) has a strong solution  $u$ , so it is also a weak solution. Now, we confirm that such weak solution satisfying (2.5) is unique.

Applying Cauchy-Schwartz inequality, we have

$$\begin{aligned} |A(u, v)| &\leq \int_{\Omega} |\nabla u| |\nabla v| dx \\ &\leq \left( \int_{\Omega} |\nabla u|^2 dx \right)^{\frac{1}{2}} \left( \int_{\Omega} |\nabla v|^2 dx \right)^{\frac{1}{2}} \\ &\leq \|\nabla u\|_{L^2(\Omega)} \|\nabla v\|_{L^2(\Omega)} \\ &\leq (\|u\|_{L^2(\Omega)} + \|\nabla u\|_{L^2(\Omega)}) (\|\nabla v\|_{L^2(\Omega)} + \|v\|_{L^2(\Omega)}) \\ &\leq \|u\|_{H^1(\Omega)} \|v\|_{H^1(\Omega)}. \end{aligned} \quad (2.6)$$

Thus,  $|A(u, v)| < \infty$  provided  $u, v \in H^1(\Omega)$ . Therefore  $A(.,.) : H^1(\Omega) \times H^1(\Omega) \rightarrow \mathbb{R}$  defined by

$$A(u, v) = \int_{\Omega} \nabla u \cdot \nabla v dx dy$$

is well defined. Further, it is easy to verify that  $L(\cdot)$  is linear in  $H^1(\Omega)$  and satisfies the following estimate

$$\begin{aligned} |L(v)| &\leq \left( \int_{\Omega} |f|^2 dx \right)^{\frac{1}{2}} \left( \int_{\Omega} |v|^2 dx \right)^{\frac{1}{2}} \leq \|f\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \\ &\leq \|f\|_{L^2(\Omega)} (\|\nabla v\|_{L^2(\Omega)} + \|v\|_{L^2(\Omega)}) \\ &\leq C \|v\|_{H^1(\Omega)}, \quad C = \|f\|_{L^2(\Omega)} < \infty. \end{aligned}$$

Clearly,  $L(\cdot)$  is well defined linear map in  $H^1(\Omega)$ . In fact,  $L$  is continuous linear map in  $H^1(\Omega)$ .

Regarding  $A(\cdot, \cdot)$ , we have the following observations:

**Observation 1:** For any scalars  $\alpha$  and  $\beta$ , we have

$$A(\alpha w_1 + \beta w_2, v) = \alpha A(w_1, v) + \beta A(w_2, v)$$

that is  $A(\cdot, \cdot)$  is bilinear map.

*Proof.:* For any scalars  $\alpha$  and  $\beta$ , we have

$$\begin{aligned} A(\alpha w_1 + \beta w_2, v) &= \int_{\Omega} \nabla(\alpha w_1 + \beta w_2) \cdot \nabla v dx \\ &= \int_{\Omega} (\alpha \nabla w_1 + \beta \nabla w_2) \cdot \nabla v dx \\ &= \int_{\Omega} \alpha \nabla w_1 \cdot \nabla v dx + \int_{\Omega} \beta \nabla w_2 \cdot \nabla v dx \\ &= \alpha \int_{\Omega} \nabla w_1 \cdot \nabla v dx + \beta \int_{\Omega} \nabla w_2 \cdot \nabla v dx = \alpha A(w_1, v) + \beta A(w_2, v). \end{aligned}$$

Therefore  $A(\cdot, \cdot)$  is a linear map in first argument  $w$ . Similarly, it can be shown that  $A(\cdot, \cdot)$  is also linear in second argument  $v$ . This completes the rest of the proof.  $\square$

**Observation 2:** The bilinear map  $A(\cdot, \cdot)$  is continuous in  $H^1(\Omega)$ .

*Proof.:* The bilinear map  $A(\cdot, \cdot)$  is said to be continuous in  $H^1(\Omega)$  if there exists a constant  $C$  such that

$$|A(u, v)| \leq C \|u\|_{H^1(\Omega)} \|v\|_{H^1(\Omega)},$$

which is true due to (2.6). Therefore,  $A(\cdot, \cdot)$  is continuous in  $H^1(\Omega)$ .  $\square$

**Remark 2.2** We know that  $\int_{\Omega} \left(\frac{dw}{dx}\right)^2 dx$  can not leads to norm, due to the fact that  $\int_{\Omega} \left(\frac{dw}{dx}\right)^2 dx = 0$  does not imply that  $w = 0$ . For instance, set  $w = 1 \neq 0$ , but,  $\int_{\Omega} \left(\frac{dw}{dx}\right)^2 dx = 0$ . In this regards, the Sobolev space  $H_0^1(\Omega) = \left\{v \in H^1(\Omega) : v|_{\partial\Omega} = 0\right\}$  is special due one of the most important inequality, know as **Poincaré inequality**.

**Lemma 2.1 (Poincaré Inequality)** For any  $w \in H_0^1(\Omega)$ ,  $\Omega \subset \mathbb{R}^d$ , following inequality

$$\|w\|_{H^1(\Omega)} \leq C \|\nabla w\|_{L^2(\Omega)}, \quad (2.7)$$

holds for some constant  $C > 0$ .

**Observation 3:** Bilinear map  $A(\cdot, \cdot)$  is positive in  $H_0^1(\Omega)$ . More precisely, for any  $w \in H_0^1(\Omega)$ , we have

$$A(w, w) \geq C \|w\|_{H^1(\Omega)}^2,$$

for some constant  $C > 0$ .

*Proof.:* Using Poincare's inequality, we have

$$\begin{aligned} A(w, w) &= \int_{\Omega} |(\nabla w)(\nabla w)| dx = \int_{\Omega} |\nabla w|^2 dx = \|\nabla w\|_{L^2(\Omega)}^2 \\ &\geq C \|w\|_{H^1(\Omega)}^2. \quad \square \end{aligned}$$



From the above discussion, we find that  $A : H_0^1(\Omega) \times H_0^1(\Omega) \rightarrow \mathbb{R}$  and  $L : H_0^1(\Omega) \rightarrow \mathbb{R}$  satisfies all assumptions of Lax-Milgram lemma. Therefore, **there exists unique  $u \in H_0^1(\Omega)$  such that**

$$A(u, v) = L(v) \quad \forall v \in H_0^1(\Omega).$$

**Remark 2.3** Suppose  $u_1$  and  $u_2$  are strong solutions for the BVP (2.1)-(2.2). Hence, they are also weak solutions and satisfies equation (2.5). But, weak solution to the BVP (2.1)-(2.2) is unique. So, we must have  $u_1 = u_2$ .

**Example 2.1 (Homework)** Suppose  $u$  is a strong solution for the BVP

$$-\Delta u + u = f \quad \text{in } \Omega \quad \& \quad u = 0 \quad \text{on } \partial\Omega,$$

where  $f \in L^2(\Omega)$  and  $\Omega \subset \mathbb{R}^2$ . Determine the weak formulation and justify that weak solution is unique. Further, show that the strong solution is unique and belongs to the function space  $H^2(\Omega) \cap H_0^1(\Omega)$ .