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so (1): $E = \frac{S}{S + P_m}$

$$\tau = T - t$$

$$V(S, t) = (S + P_m) \bar{V}(\epsilon, \tau)$$

$$S = \frac{P_m \epsilon}{1 - \epsilon}, \quad S + P_m = \frac{P_m}{1 - \epsilon}$$

$$\Rightarrow \frac{d\epsilon}{dS} = \frac{P_m}{(S + P_m)^2} = \frac{(1 - \epsilon)^2}{P_m}$$

because

$$\frac{\partial V}{\partial t} = \frac{\partial}{\partial \tau} [(S + P_m) \bar{V}(\epsilon, \tau)] = (S + P_m) \frac{\partial \bar{V}}{\partial \tau} = -\frac{P_m}{1 - \epsilon} \frac{\partial \bar{V}}{\partial \tau}$$

$$\frac{\partial V}{\partial S} = \frac{\partial}{\partial S} [(S + P_m) \bar{V}(\epsilon, \tau)] = (S + P_m) \frac{\partial \bar{V}}{\partial \epsilon} \frac{d\epsilon}{dS} + \bar{V}$$

$$= (1 - \epsilon) \frac{\partial \bar{V}}{\partial \epsilon} + \bar{V}$$

$$\text{and } \frac{\partial^2 V}{\partial S^2} = \frac{\partial}{\partial \epsilon} \left((1 - \epsilon) \frac{\partial \bar{V}}{\partial \epsilon} + \bar{V} \right) \frac{d\epsilon}{dS} = \frac{(1 - \epsilon)^3}{P_m} \frac{\partial^2 \bar{V}}{\partial \epsilon^2}$$

$$\text{and assume } \bar{\sigma}(\epsilon) = \sigma(S(\epsilon)) = \sigma\left(\frac{P_m \epsilon}{1 - \epsilon}\right)$$

Hence, the given PDE reduces to

$$\frac{P_m}{1 - \epsilon} \frac{\partial \bar{V}}{\partial \tau} = \frac{\bar{\sigma}^2(\epsilon) P_m \epsilon^2 (1 - \epsilon)}{2} \frac{\partial^2 \bar{V}}{\partial \epsilon^2} + (r - D_0) P_m \epsilon \frac{\partial \bar{V}}{\partial \epsilon} + \frac{(1 - D_0)(\epsilon - r)}{1 - \epsilon} P_m \bar{V}$$

$$= \frac{\partial \bar{V}}{\partial \tau} = \frac{\sigma^2(\epsilon) \epsilon^2 (1-\epsilon)^2}{2} \frac{\partial^2 \bar{V}}{\partial \epsilon^2} + (r - D_0) \epsilon (1-\epsilon) \frac{\partial \bar{V}}{\partial \epsilon} - [r(1-\epsilon) + D_0 \epsilon] \bar{V} \quad \text{where } \epsilon \in [0, 1] \cdot (0 \leq \epsilon < 1)$$

Assuming \bar{V} is a smooth fn of ϵ , then the ^{above} eqⁿ holds for $\epsilon=1$ also.

$$\therefore V(S, T) = (S + P_m) \bar{V}(\epsilon, 0) = \bar{V}(\epsilon, 0) \frac{P_m}{1-\epsilon}$$

$V(S, T) = V_T(S) \Rightarrow$ can be rewritten as ~~$\bar{V}(\epsilon, 0)$~~

$$\bar{V}(\epsilon, 0) = V_T \left(\frac{P_m \epsilon}{1-\epsilon} \right) \left(\frac{1-\epsilon}{P_m} \right)$$

Hence the original eqⁿ could be written as:-

$$\left\{ \begin{aligned} \frac{\partial \bar{V}}{\partial \tau} &= \frac{1}{2} \sigma^2(\epsilon) \epsilon^2 (1-\epsilon)^2 \frac{\partial^2 \bar{V}}{\partial \epsilon^2} + (r - D_0) \epsilon (1-\epsilon) \frac{\partial \bar{V}}{\partial \epsilon} - [r(1-\epsilon) + D_0 \epsilon] \bar{V} \quad 0 \leq \epsilon \leq 1, 0 \leq \tau \\ \bar{V}(\epsilon, 0) &= \frac{(1-\epsilon)}{P_m} V_T \left(\frac{P_m \epsilon}{1-\epsilon} \right) \end{aligned} \right.$$

Hence we can obtain finite domain eqⁿ from infinite domain. ~~Parabolic eqⁿ defined on finite domain to have a unique solⁿ besides an initial condⁿ.~~ For parabolic eqⁿs boundary condⁿ also needed with the initial condⁿs.

However coefficients of $\frac{\partial^2 \bar{V}}{\partial \varepsilon^2}$, $\frac{\partial \bar{V}}{\partial \varepsilon}$ at $\varepsilon=0, \varepsilon=1$ are 0.

Actually at $\varepsilon=0 \Rightarrow$

$$\frac{\partial \bar{V}(0, z)}{\partial z} = -\gamma \bar{V}(0, z)$$

$$\Rightarrow \bar{V}(0, z) = \bar{V}(0, 0) e^{-\gamma z}$$

Similarly at $\varepsilon=1$ we have

$$\frac{\partial \bar{V}(1, z)}{\partial z} = -D_0 \bar{V}(1, z)$$

$$\Rightarrow \bar{V}(1, z) = \bar{V}(1, 0) e^{-D_0 z}$$

Hence, we obtain two solutions of the ordinary differential equation and providing value at boundaries.

Hence eqⁿ is reduced to finite domain and boundary condⁿ are obtained at $\varepsilon=0, \varepsilon=1$.

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Solⁿ (2): $K \rightarrow$ set of all family of admissible functions

for $v \in K$, let y be the exact solⁿ of linear complementarity problem for american put options. Since y is the solⁿ for this partial differential inequality $\Rightarrow y \rightarrow C^2$ smooth on the continuation region and $y \in K$.

Hence we can write: $K = \{v \in C^0 : \frac{\partial v}{\partial x} \sim C^0\}$ previous

$$\Rightarrow \text{given the problem } v \geq y, \quad \frac{\partial y}{\partial z} - \frac{\partial^2 y}{\partial x^2} \geq 0 \quad \left| \begin{array}{l} v(x,t) \geq y(x,t) \\ \text{for } x,t \\ v(x,0) = y(x,0) \end{array} \right.$$

$$\Rightarrow \int_{x_{\min}}^{x_{\max}} \left(\frac{\partial y}{\partial z} - \frac{\partial^2 y}{\partial x^2} \right) (v-y) dx \geq 0 \rightarrow (1) \quad \left| \begin{array}{l} v(x_{\max}, t) \\ = y(x_{\max}, t) \end{array} \right.$$

Invoking the complementarity, $\left| \begin{array}{l} \text{and} \\ v(x_{\min}, t) = y(x_{\min}, t) \end{array} \right.$

$$\int_{x_{\min}}^{x_{\max}} \left(\frac{\partial y}{\partial z} - \frac{\partial^2 y}{\partial x^2} \right) (y-y) dx = 0 \rightarrow (2)$$

and subtraction given $\Rightarrow (1) - (2) \Rightarrow$

$$\int_{x_{\min}}^{x_{\max}} \left(\frac{\partial y}{\partial z} - \frac{\partial^2 y}{\partial x^2} \right) (v-y) dx \geq 0$$

Applying integration by parts: we have

$$\int_{x_{\min}}^{x_{\max}} \left(\frac{\partial y}{\partial z} (v-y) + \frac{\partial y}{\partial x} \left(\frac{\partial v}{\partial x} - \frac{\partial y}{\partial x} \right) \right) dx - \frac{\partial y}{\partial x} (v-y) \Big|_{x_{\min}}^{x_{\max}} \geq 0$$

* $\frac{\partial \psi}{\partial n}(U-y)$ vanishes because at x_{\min}, x_{\max}

In view of $v=g, y=g$ then $\Rightarrow v=y$.

Hence

\mathcal{Q}_g^h reduces to

$$\int_{x_{\min}}^{x_{\max}} \left(\frac{\partial \psi}{\partial z}(v-y) + \frac{\partial \psi}{\partial x} \left(\frac{\partial v}{\partial x} - \frac{\partial y}{\partial x} \right) \right) dx \geq 0 \quad \forall v \in K$$

holds which is $I(y; v)$.

$I(y; v)$ holds for all $v \in K$. Considering the case when $v=y \Rightarrow I(y; v)$ takes its minimum value.

$$\min_{v \in K, v=y} I(y; v) = I(y; y) = 0.$$

(a) To obtain the fully discrete version of above minimization problem:

We assume approximations for \hat{g} and v in similar forms:

$$\sum_1^N w_i(\tau) \phi_i(x) \quad \text{for } \hat{g}$$

$$\sum_1^N u_i(\tau) \phi_i(x) \quad \text{for } v$$

The reduced smoothness match the requirements for K .

We assumed that variables τ, x ~~are independent~~ and ~~which~~ which are independent are to be separated.

Hence same x -grid is applied for all $\tau \Rightarrow$ rectangular (x, τ) grid. The time ~~gradient~~ is introduced

In for w_i, v_i . Since the basis function ϕ_i represent x_i grid. Putting values in the eqⁿ formed in previous part:

$$\int \left\{ \left(\sum_i \frac{dw_i}{dz} \phi_i \right) \left(\sum_j (v_j - w_j) \phi_j \right) + \left(\sum_i w_i \phi_i' \right) \left(\sum_j (v_j - w_j) \phi_j' \right) \right\} dx$$

$$= \sum_i \sum_j \frac{dw_i}{dz} (v_j - w_j) \int \phi_i \phi_j dx + \sum_i \sum_j w_i (v_j - w_j) \int \phi_i' \phi_j' dx$$

$$\int \phi_i' \phi_j' dx \geq 0$$

We can write above in vector notation like

$$\left(\frac{dw}{dz} \right)^T B (v - w) + w^T A (v - w) \geq 0$$

$$\Rightarrow (v - w)^T \left(B \frac{dw}{dz} + A w \right) \geq 0 \quad \text{where } A, B \text{ are corresponding matrices}$$

As time is discretized, define

$$w^{(n)} := w(\tau_n)$$

$$w^{(n)} := w(\tau_n), \quad v^{(n)} := v(\tau_n)$$

Using substitution and θ -averaging $A w$ term in (3):

$$(v^{(n+1)} - w^{(n+1)})^T \left(B \frac{1}{\Delta z} (w^{(n+1)} - w^{(n)}) + \theta (A w^{(n+1)}) + (1 - \theta) A w^{(n)} \right) \geq 0$$

$\forall n$

$$= (v^{(n+1)} - w^{(n+1)})^T \left((B + \Delta z \theta A) w^{(n+1)} + (\Delta z (1 - \theta) A - B) w^{(n)} \right) \geq 0$$

Using the variables given in ques:

$$r := (B - \Delta z (1-\theta) A) w^{(n)}$$

$$C := B + \Delta z \theta A$$

we get

$$(u^{(n+1)} - w^{(n+1)})^T (C w^{(n+1)} - r) \geq 0$$

Here is fully discrete version of above minimised problem.

Solⁿ (3) - To show that following problems are equivalent \Rightarrow

Consider the 1st case:

a) (FDM) solⁿ implies (FEM) solⁿ.

Let y ~~solⁿ~~ be the solⁿ of FDM. $\Rightarrow y \geq g$.

$$I = \cancel{y} \quad (u-y)^T (cy-r) = (u-g)^T (cy-r) - (y-g)^T (cy-r)$$

$$\therefore (u-g)^T (cy-r) \geq 0; \quad (y-g)^T (cy-r) = 0$$

Hence $I \geq 0 \quad \forall y \in u \geq g$.

Hence solⁿ of FDM also satisfies that of FEM.

(b) (FEM) solⁿ \Rightarrow (FDM) solⁿ

Let y ~~solⁿ~~ be solⁿ of FEM eqⁿ. $\Rightarrow y \geq g$.

$$v^T (cy-r) \geq y^T (cy-r) \quad \forall v \in R.$$

Suppose

let's say ~~any~~ k^{th} component of $cy-r < 0$ and make v_k very large.

Hence $v^T (cy-r)$ becomes very small as compared which is not possible as LHS \geq RHS.

Hence $\Rightarrow cy-r \geq 0$

Since we have $y \geq g \Rightarrow (y-g)^T (cy-r) \geq 0$.

Put in FEM eqⁿ: $u=g \Rightarrow (y-g)^T (cy-r) \leq 0$

$\Rightarrow (y-g)^T (cy-r) = 0$ hence solⁿ of FEM

also satisfies FDM. Hence we could say both problems are equivalent.