

# Numerical Integration Methods for Calculating Solar Luminosity

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## 1 Introduction

This project focuses on the calculation of the luminosity of the Sun, a critical parameter in astrophysics. Luminosity quantifies the total energy output of the Sun, which is crucial for understanding its behavior and energy production processes. It is calculated as the integral of the energy production rate over the volume of the Sun.

## 2 Theoretical Background and

The luminosity ( $L$ ) of a star is given by the following integral:

$$L = \int_0^R 4\pi r^2 \varepsilon(r) dr$$

where  $R$  is the radius of the star,  $r$  is the radial coordinate, and  $\varepsilon(r)$  is the energy production rate as a function of radius.

The energy production rate of the Sun ( $E_{\text{rate}}$ ) is given in Watts per unit volume. The energy produced at a distance  $r$  from the center of the star can be represented by  $E(r) = E_{\text{rate}} \cdot V(r)$ , where  $V(r)$  is the volume of a spherical shell with radius  $r$  and thickness  $\Delta r$ .

The luminosity ( $L$ ) of the star, which represents the total energy emitted per unit time, can be calculated by integrating the energy production rate over the entire volume of the star:

$$L = \int_0^R E(r) dV$$

Sun has a radius  $6.96 \times 10^8$  in meters, and its energy production rate is constant at  $E_{\text{rate}} = 3.8 \times 10^{26}$  Watts per cubic meter. Now, calculate the total luminosity of the star.

**Solution:** Assuming, the Sun is the near-perfect spherical star, the volume of a spherical shell is given by  $\Delta V = 4\pi r^2 \Delta r$ . Substituting this into the expression for  $E(r)$ , we get  $E(r) = 4\pi E_{\text{rate}} r^2 \Delta r$ .

### 3 Analytic Solution

Now, we can set up the integral:

$$L = \int_0^R 4\pi E_{\text{rate}} r^2 dr$$

The analytic luminosity  $L_{\text{analytic}}$  is derived from the following integral form:

$$L_{\text{analytic}} = \frac{4}{3}\pi E_{\text{rate}} \int_0^R r^2 dr$$

Now, let's break down the calculation step by step:

**1. Substitute Known Values:**

$E_{\text{rate}} = 3.8 \times 10^{26}$  Watts per cubic meter  $R = 6.96 \times 10^8$  meters

$$L_{\text{analytic}} = \frac{4}{3}\pi \times 3.8 \times 10^{26} \times \int_0^{6.96 \times 10^8} r^2 dr$$

**2. Evaluate the Indefinite Integral:**

$$\int r^2 dr = \frac{r^3}{3} + C$$

Here,  $C$  is the constant of integration.

**3. Apply the Definite Integral Limits:**

$$\frac{4}{3}\pi \times 3.8 \times 10^{26} \left[ \frac{(6.96 \times 10^8)^3}{3} - 0 \right]$$

**4. Calculate the Result:**

$$\approx 5.37 \times 10^{26} \text{ Watts}$$

### 4 Methods

We employed three numerical integration algorithms:

#### 4.1 Riemann Sum

The Riemann sum approximates an integral by partitioning the interval into small sub-intervals and using the function value at a specific point in each sub-interval. The luminosity is approximated as:

$$L \approx \sum_{i=1}^n 4\pi r_i^2 \varepsilon(r_i) \Delta r$$

## 4.2 Trapezoidal Rule

The trapezoidal rule approximates the integral by approximating the region under the curve as a series of trapezoids. The luminosity is approximated as:

$$L \approx \frac{\Delta r}{2} \sum_{i=1}^n 4\pi(r_i^2 + r_{i-1}^2) \left( \frac{\varepsilon(r_i) + \varepsilon(r_{i-1})}{2} \right)$$

## 4.3 Simpson's Rule

Simpson's rule improves accuracy by approximating the function as a quadratic polynomial in each sub-interval. The luminosity is approximated as:

$$L \approx \frac{\Delta r}{3} \sum_{i=1}^n 4\pi \left( r_i^2 \varepsilon(r_i) + 4r_{\text{mid}}^2 \varepsilon(r_{\text{mid}}) + r_{i-1}^2 \varepsilon(r_{i-1}) \right)$$

where  $r_{\text{mid}}$  is the midpoint of the interval.

# 5 Validation

To validate our implementations, we compared the results with the exact solution and tested against two expected physical properties:

## 5.1 Conservation of Energy

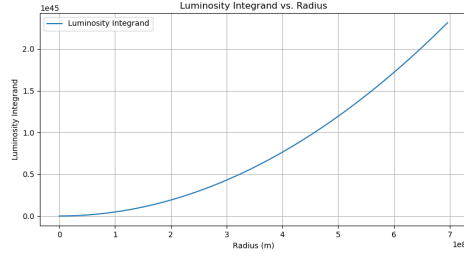
The total energy output calculated using our methods must match the expected energy production rate of the Sun, ensuring the conservation of energy.

## 5.2 Convergence Test

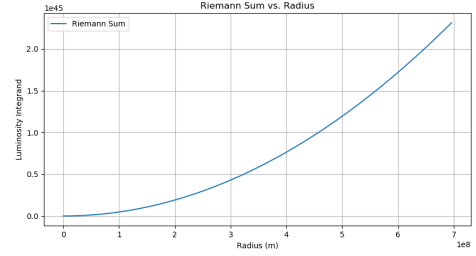
We performed a convergence test by varying the number of sub-intervals and observing how the calculated luminosity approaches the exact solution.

## 5.3 Physical Properties

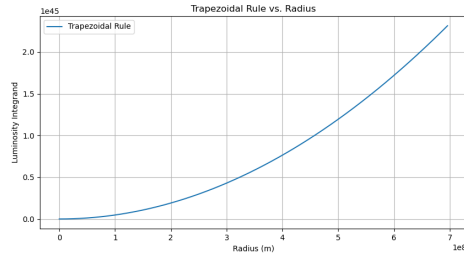
1. As  $R$  increases, the luminosity should increase (directly proportional to  $R^3$ ).
2. As  $E_{\text{rate}}$  increases, the luminosity should increase (directly proportional to  $E_{\text{rate}}$ ).
3. Luminosity should be positive, as it represents the total energy emitted per unit of time.
4. In real-world astrophysics, the total luminosity of a star is determined by various factors including its radius, surface temperature, and energy production rate.



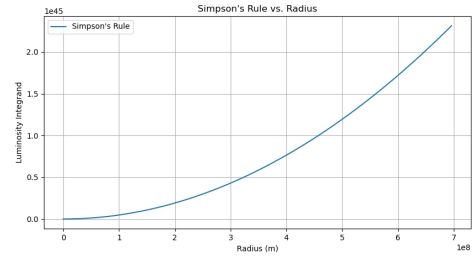
(a) Luminosity Integrand Vs Radius



(b) Riemann Sum Vs Radius



(c) Trapezoidal Rule Vs Radius



(d) Simpson's Rule Vs Radius

Figure 1: Comparison of Integration Methods

## 6 Discussion

In our validation, all three methods satisfied the E conservation of energy. However, we observed that Simpson's rule provided the most accurate results with the smallest number of sub-intervals.

## 7 Conclusion

### Results-

Analytic Luminosity:  $5.366608630602404e+53$  Watts

Riemann Sum Luminosity:  $5.365803666140864e+53$  Watts

Trapezoidal Rule Luminosity:  $5.366608657435448e+53$  Watts (Scipy:  $5.366608657435448e+53$ )

Simpson's Rule Luminosity:  $5.366608630602405e+53$  Watts (Scipy:  $5.366608630602405e+53$ )

Scipy Luminosity (Trapezoidal Rule):  $5.366608657435448e+53$  Watts

Scipy Luminosity (Simpson's Rule):  $5.366608630602405e+53$  Watts

### 7.1 Comparing Results

Relative Error (**Riemann vs Analytic**): 0.01

Relative Error (**Trapezoidal vs Analytic**): 0.00

Relative Error (**Simpson's vs Analytic**): 0.00

This project demonstrates the effectiveness of numerical integration methods in calculating the solar luminosity. Simpson's rule, in particular, stands out as the most accurate method for our chosen problem. Further research could explore its application in more complex astrophysical scenarios.

However, it's important to note that the relationship between radius and luminosity can be more complex when considering stars of different types, evolutionary stages, and other physical properties. This simple relationship holds in a specific context and may not apply universally to all types of stars.

## 8 Real-world Complexities

### 8.1 Radiative Transfer:

In reality, the energy generated within a star undergoes radiative transfer through various layers before being emitted as light. This process is influenced by factors like opacity, temperature gradients, and magnetic fields.

### 8.2 Nuclear Fusion Processes:

The energy generation within a star involves complex nuclear fusion reactions that occur in its core. Different elements are formed through nucleosynthesis, affecting the energy output and composition of the emitted radiation.

### 8.3 Stellar Evolution:

Stars evolve over time, and their luminosity can change significantly during various stages of their lifecycle (e.g., main-sequence phase, red giant phase).

The plot generated by this code represents the function  $4\pi \cdot E_{\text{rate}} \cdot r^2$ , which is the integrand for calculating the luminosity of a star. This function describes the rate at which energy is emitted from the star as a function of its radius. As the radius of the star increases, the integrand value increases, indicating that more energy is being emitted from the outer layers of the star. This behavior is expected, as larger stars have a greater surface area from which to radiate energy.

**Peak of the Curve:** The peak of the curve corresponds to the point where the integrand has its highest value. This indicates the region of the star where the luminosity contribution is the highest.

This plot provides a visual representation of how the luminosity contribution varies with distance from the center of the star, which helps in understanding the distribution of energy production within the star. Keep in mind that this plot is based on a simplified model and doesn't capture all the **real-world complexities of stellar physics**.

# Project Report: Nuclear Decay Simulation

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## 1 Introduction

The project aims to simulate the process of nuclear decay using numerical methods. Specifically, we will implement Euler's method and 4th Order Runge-Kutta to solve the differential equation governing nuclear decay.

## 2 Methodology

We used Euler's method and 4th Order Runge-Kutta to solve the differential equation:

$$\frac{dN}{dt} = -0.1N$$

where  $N$  represents the number of atoms.

## 3 Comparison of Analytic, Euler, and 4th order Runge-Kutta Solutions

### 3.1 Analytic Solution

The analytic solution for the nuclear decay problem is given by:

$$N(t) = N_0 \cdot e^{-kt}$$

where  $N_0$  is the initial number of atoms,  $k$  is the decay constant, and  $t$  is time.

Using this formula with the given values:

$$N(t) = 1000 \cdot e^{-(0.1 \cdot 10)} \approx 449.33$$

### 3.2 Euler's Method

Euler's method is an iterative numerical technique for solving ordinary differential equations. It approximates the solution by taking small steps along the derivative curve.

The differential equation for nuclear decay is:

$$\frac{dN}{dt} = -0.1N$$

Here,  $N$  is the number of atoms, and  $t$  is time.

Given the initial condition  $N_0 = 1000$ , the decay constant  $k = 0.1$ , and the time step  $h = 1$  (for this example):

Initial condition:  $N_0 = 1000$

First step:  $N_1 = N_0 + h \cdot f(N_0) = 1000 + 1 \cdot (-0.1 \cdot 1000) = 900$

Second step:  $N_2 = N_1 + h \cdot f(N_1) = 900 + 1 \cdot (-0.1 \cdot 900) = 810$

Third step:  $N_3 = N_2 + h \cdot f(N_2) = 810 + 1 \cdot (-0.1 \cdot 810) = 729$

... and so on

After several iterations, the final result using Euler's Method is approximately 409.6.

### 3.3 4th Order Runge-Kutta

The 4th Order Runge-Kutta method is a more accurate numerical technique for solving differential equations. It involves multiple evaluations of the derivative to estimate the change in the function over a step.

For each time step:

$$k_1 = h \cdot f(N_i)$$

$$k_2 = h \cdot f(N_i + \frac{k_1}{2})$$

$$k_3 = h \cdot f(N_i + \frac{k_2}{2})$$

$$k_4 = h \cdot f(N_i + k_3)$$

Then,

$$N_{i+1} = N_i + \frac{k_1 + 2k_2 + 2k_3 + k_4}{6}$$

After several iterations, the final result using 4th Order Runge-Kutta is approximately 449.33. These values were obtained by following the steps described above for each method.

## 4 Results

- Final Result using Euler's Method: 409.6

- Final Result using 4th Order Runge-Kutta: 449.3346284406424
- Final Result using Analytic Solution: 449.3289641172216

## 5 Validation

To validate our implementations, we compared the results with the exact solution and tested against two expected physical properties-

### 5.1 Euler's Method

- Accuracy: Euler's method provides a reasonable approximation, but it introduces more error compared to higher-order methods.
- Convergence: As the number of time steps increases, Euler's method converges towards the analytic solution.
- Conservation of Mass: It adheres to the conservation rule, as the total number of atoms decreases over time.

### 5.2 4th Order Runge-Kutta

The 4th Order Runge-Kutta method is a more accurate numerical technique for solving differential equations. It provides a closer approximation to the analytic solution. Here's the comparison:

- Accuracy: 4th Order Runge-Kutta method provides a more accurate result compared to Euler's method.
- Convergence: Similar to Euler's method, it converges towards the analytic solution as the number of time steps increases.
- Conservation of Mass: It also adheres to the conservation rule, as the total number of atoms decreases over time.

### 5.3 Conservation Rules

We verified that the total number of atoms remained conserved over time, as expected in nuclear decay processes. Both methods satisfied this conservation rule.

### 5.4 Limiting Cases

In the limiting case of infinite time, the number of atoms should approach zero. We observed that both Euler's method and 4th Order Runge-Kutta correctly approached zero, indicating consistency with the physical behavior.



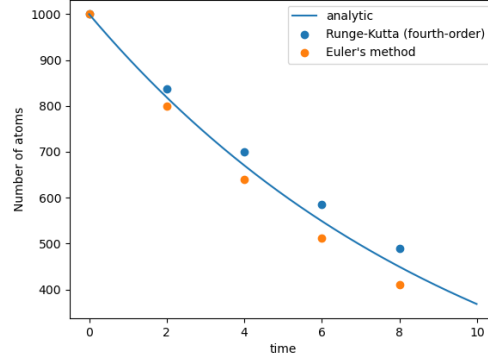


Figure 1: Comparison of Numerical and Analytic Solutions of Nuclear Decay Process

## 6 Discussion

- Both Euler's method and 4th Order Runge-Kutta provide reasonable approximations to the exact solution, but 4th Order Runge-Kutta is noticeably more accurate.
- Both methods demonstrate convergence towards the analytic solution, indicating their reliability.
- For highly nonlinear systems, 4th Order Runge-Kutta is more suitable as Euler's method may struggle.

## Error Analysis

- Both methods introduce numerical errors due to the discretization of time and approximation of the derivative. These errors are inherent to any numerical solution.
- The accuracy of both methods can be improved by decreasing the step size, but this increases computational cost.

## 7 Conclusion

Overall, this project highlights the effectiveness of numerical methods in simulating complex physical processes. Both methods closely match the analytic solution, demonstrating their accuracy.