

# Counter-diabatic driving using Floquet engineering

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## 1 Introduction

$$H = H_0 + H_1$$

$$H = J \sum_j (c_{j+1}^\dagger c_j + h.c.) + \cos(\omega t) \sum_j A_j c_j^\dagger c_j \quad (1)$$

$$V = \exp(-i \sin(\omega t)/\omega \sum_j A_j c_j^\dagger c_j)$$

$$\begin{aligned} H_{rot} &= V^\dagger H V - i V^\dagger \dot{V} \\ &= V^\dagger H_0 V + \cos(\omega t) \sum_j A_j c_j^\dagger c_j + i^2 \cos(\omega t) \sum_j A_j c_j^\dagger c_j \\ &= V^\dagger H_0 V \\ &= V^\dagger c_{j+1}^\dagger V V^\dagger c_j V \end{aligned}$$

Using  $[n_j, c_j] = -c_j$  and  $[n_j, c_j^\dagger] = c_j^\dagger$

$$H_{rot} = J \sum_j (g^{j,j+1} c_{j+1}^\dagger c_j + h.c.)$$

where  $g^{j,j+1} = \exp\left(i \sin(\omega t) \frac{A_{j+1} - A_j}{\omega}\right)$

$$\begin{aligned} H_F^{(0)} &= \frac{1}{T} \int_{t_0}^{T+t_0} (c_{j+1}^\dagger c_j \exp\left(i \sin(\omega t) \frac{A_{j+1} - A_j}{\omega}\right) dt + h.c.) \\ &= J_{eff} (c_{j+1}^\dagger c_j + h.c.) \end{aligned}$$

where  $J_{eff} = J \mathcal{J}_0\left(\frac{A_{j+1} - A_j}{\omega}\right)$

Few points: Bessel function is always less than 1 while CD hopping  $J$  is always greater than 1. Furthermore, driving frequency would be high so that for all values of  $A_{j+1} - A_j$ , we would have  $\frac{A_{j+1} - A_j}{\omega} \ll 1$  so that we are close to origin and away from zeros of Bessel's functions.

## 2 Magnus expansion

For a Hamiltonian which is periodic in time, it's unitary operator over a full driving cycle is given by:

$$U(T + t_0, t_0) = \mathcal{T}_t \exp\left(-\frac{i}{\hbar} \int_{t_0}^T dt H(t)\right) = \exp\left(-\frac{i}{\hbar} H_F[t_0]T\right) \quad (2)$$

$H_F[t_0] = \sum_n H_F^{(n)}[t_0]$  where

$$H_F^{(0)} = \frac{1}{T} \int_{t_0}^{T+t_0} H(t) dt$$

$$H_F^{(1)} = \frac{1}{2!T i \hbar} \int_{t_0}^{T+t_0} dt_1 \int_{t_0}^{t_1} dt_2 [H(t_1), H(t_2)]$$

### 3 Bessel's function of first kind

Integral representation of Bessel's function of first kind  $\mathcal{J}_n(x)$  is given by:

$$\mathcal{J}_n(x) = \frac{1}{2\pi} \int_{-\pi}^{\pi} d\tau e^{i(n\tau - x \sin \tau)} = \frac{1}{T} \int_{-T}^T d\tau e^{i(n\omega\tau - x \sin \omega\tau)} \quad (3)$$

For  $x \ll 1$ ,  $\mathcal{J}_0(x) = 1 - \frac{x^2}{2}$