

AmATH 567

HW 7

Nate Whybra

3.5.1)

b) The function  $f(z) = \frac{e^{2z} - 1}{z^2}$  fails to be analytic at  $z=0$ , so we find the Laurent series for  $f(z)$  centered at  $z_0=0$ ,

$$\begin{aligned} f(z) &= \frac{1}{z^2} (e^{2z} - 1) = \frac{1}{z^2} \left( \sum_{n=0}^{\infty} \frac{2^n}{n!} z^n - 1 \right) \\ &= \frac{1}{z^2} \sum_{n=1}^{\infty} \frac{2^n}{n!} z^n \\ &= \sum_{n=1}^{\infty} \frac{2^n}{n!} z^{n-2} \end{aligned}$$

Put  $n-2 = k$ , then the above

$$= \sum_{k=-1}^{\infty} \frac{2^{k+2}}{(k+2)!} z^k$$

So  $f(z)$  has a pole of order 1 with strength  $= \frac{2^{-1+2}}{(-1+2)!} = \frac{2^1}{1!} = 2$  at  $z=0$ .

(c) From example 3.5.3 in the textbook, the function  $g(z) = \tan(z)$  has <sup>( $m \in \mathbb{Z}$ )</sup> simple poles of strength  $-1$  at  $z = \frac{\pi}{2} + m\pi$ , and the Laurent series centered at  $z_0 = \frac{\pi}{2} + m\pi$  starts like

$$g(z) = \tan(z) = -\frac{1}{z-z_0} + \frac{1}{3}(z-z_0) + \text{rest of terms}$$

So

$$\begin{aligned} f(z) = e^{\tan(z)} &= \sum_{n=0}^{\infty} \frac{(\tan(z))^n}{n!} = \sum_{n=0}^{\infty} \frac{1}{n!} \left( -\frac{1}{z-z_0} + \frac{1}{3}(z-z_0) + \dots \right)^n \\ &= \sum_{n=0}^{\infty} \frac{1}{n!} \left( -\frac{1}{z-z_0} \right)^n + \dots \quad \swarrow \text{the } \frac{1}{n!} \text{ is here too} \\ &= \sum_{n=0}^{\infty} \frac{1}{n!} (-1)^n (z-z_0)^{-n} + \dots \\ &= \sum_{n=-\infty}^0 \frac{1}{(-n)!} (-1)^n (z-z_0)^n + \dots \end{aligned}$$

The series has infinitely many terms  $(z-z_0)^n$  with  $n$  being negative, hence  $f(z) = e^{\tan(z)}$  has an essential singularity at  $z = \frac{\pi}{2} + m\pi$  for every  $m \in \mathbb{Z}$ .

(d) we have the function

$$f(z) = \frac{z^3}{z^2 + z + 1} = \frac{z^3}{(z - \omega_+)(z - \omega_-)}$$

where  $\omega_{\pm} = -\frac{1}{2} \pm \frac{\sqrt{3}}{2}i = \pm e^{i\frac{2\pi}{3}}$ . The function has a simple pole for  $z = \omega_+$  and  $z = \omega_-$ . By similar argument to (Example 3.5.1) in the textbook, the leading term of the Laurent series near  $z = \omega_+$  is (the strength)

$$\frac{\omega_+^3}{\omega_+ - \omega_-} = \frac{e^{i2\pi}}{2e^{i2\pi/3}} = \frac{1}{2}e^{i4\pi/3}$$

and the strength of the pole at  $z = \omega_-$  is

$$\frac{\omega_-^3}{\omega_- - \omega_+} = \frac{(-e^{i2\pi/3})^3}{-2e^{i2\pi/3}} = \frac{1}{2}e^{i4\pi/3}$$

3.5.3)

$$a) f(z) = \frac{z}{z^4 + 2} = \frac{z}{(z - w_0)(z - w_1)(z - w_2)(z - w_3)}$$

where  $w_j = \sqrt[4]{2} e^{i(\frac{\pi}{4} + \frac{\pi}{2}j)}$  for  $j = 0, 1, 2, 3$ .

So again, by similar argument to Example 3.5.1, there is a simple pole for every root

$z = w_j$ . Denote the strengths of each pole as  $s_j$ , then

$$s_0 = \frac{w_0}{(w_0 - w_1)(w_0 - w_2)(w_0 - w_3)}$$

Note that with  $j \neq k$

$$\begin{aligned} w_j - w_k &= 2^{1/4} e^{i\pi/4} (e^{\frac{\pi}{2}ij} - e^{\frac{\pi}{2}ik}) \\ &= 2^{1/4} e^{i\pi/4} (i^j - i^k) \end{aligned}$$

$$\text{So, } s_0 = \frac{2^{1/4} e^{i\pi/4}}{2^{3/4} e^{i3\pi/4} (1 - i)(1 + i)(1 + i)}$$

$$\begin{aligned} &= \frac{2^{1/4} e^{i\pi/4}}{2 \cdot 2^{3/4} e^{i3\pi/4}} = \frac{2^{-5/2} e^{-i\pi/2}}{2} = -i 2^{-5/2} \\ &= \frac{-i}{2^{5/2}} \end{aligned}$$

$$\begin{aligned}
 S_1 &= \frac{\omega_1}{(\omega_1 - \omega_0)(\omega_1 - \omega_2)(\omega_1 - \omega_3)} \\
 &= \frac{2^{1/4} e^{i\frac{3\pi}{4}}}{2^{3/4} e^{i\frac{3\pi}{4}} (i-1)(i+1)(i+i)} \\
 &= \frac{1}{2^{1/2} (-2)(2i)} = \frac{i}{2^{1/2} \cdot 2^2} = \frac{i}{2^{5/2}}
 \end{aligned}$$

$$\begin{aligned}
 S_2 &= \frac{\omega_2}{(\omega_2 - \omega_0)(\omega_2 - \omega_1)(\omega_2 - \omega_3)} \\
 &= \frac{2^{1/4} e^{i5\pi/4}}{2^{3/4} e^{i\frac{3\pi}{4}} (-1-1)(-1-i)(-1+i)} \\
 &= \frac{e^{i\pi/2}}{2^{1/2} \cdot 2 \cdot (i-1)(i+1)} \\
 &= \frac{i}{2^{3/2} \cdot (-2)} = \frac{-i}{2^{5/2}}
 \end{aligned}$$

$$\begin{aligned}
 S_3 &= \frac{\omega_3}{(\omega_3 - \omega_0)(\omega_3 - \omega_1)(\omega_3 - \omega_2)} \\
 &= \frac{2^{1/4} e^{i7\pi/4}}{2^{3/4} e^{i3\pi/4} (-i-1)(-i-i)(-i+1)} \\
 &= \frac{-1}{2^{1/2} (1-i)(1+i) 2i} = \frac{i}{2^{1/2} \cdot 2^1 \cdot 2^1} = \frac{i}{2^{5/2}}
 \end{aligned}$$



c) For the function  $f(z) = \frac{z}{\sin^2(z)}$ , there are singularities whenever  $z = m\pi$  for  $m \in \mathbb{Z}$ . We can write  $\sin^2(z) = \frac{1 - \cos(2z)}{2}$  so that

$$f(z) = \frac{2z}{1 - \cos(2z)} = \frac{2z}{\sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{(2n)!} \cdot 2^{2n} z^{2n}} \quad (1)$$

$$= \frac{\cancel{z} \cdot 2}{\cancel{z} \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{(2n)!} 2^{2n} z^{2n-1}}$$

$$= \frac{2}{z(2 - 4z^2 + \dots)}$$

So  $f(z)$  has a pole of order 1 at  $z=0$  with strength  $\frac{2}{2} = 1$ . Now if we plug  $z = z_0 + m\pi$  into (1), ( $m \neq 0$ )

$$f(z) = \frac{2z_0}{(z_0 + m\pi)^2 [2 - 4(z_0 + m\pi)^2 + \dots]} + \frac{2m\pi}{(z_0 + m\pi)^2 [2 - 4(z_0 + m\pi)^2 + \dots]}$$

if  $z = m\pi$ , i.e.  $z_0 = 0$ , we see the above becomes

$$= \frac{2\pi m}{z^2 [2 - 4z^2 + \dots]}$$

So there is a pole of order 2 at  $z = m\pi$  with strength  $= \frac{2\pi m}{2} = \pi m \quad \forall m \neq 0$ .

d) we have

$$f(z) = \frac{e^z - 1 - z}{z^4} = \frac{\sum_{n=2}^{\infty} \frac{z^n}{n!}}{z^4}$$

$$= \sum_{n=2}^{\infty} \frac{z^{n-4}}{n!}$$

$$= \sum_{n=-2}^{\infty} \frac{z^n}{(n+4)!}$$

since  $f$  has singularities at  $z=0$ , we  
see  $f$  has a pole of order 2 at  
 $z=0$  with strength  $= \frac{1}{(1+2)!} = \frac{1}{2}$



3.6.6)

a) Taking the principle branch of the logarithm, recall that  $\frac{d}{dz} \ln(f(z)) = \frac{f'(z)}{f(z)}$ , so

$$\frac{d}{dz} \ln\left(\frac{1}{f(z)}\right) = \frac{d}{dz} \ln(f(z)^{-1}) = \frac{d}{dz} -\ln(f(z)) = -\frac{f'(z)}{f(z)}$$

So if,

$$\frac{1}{\Gamma(z)} = z e^{\gamma z} \prod_{n=1}^{\infty} \left(1 + \frac{z}{n}\right) e^{-z/n}$$

$$-\log(\Gamma(z)) = \log(z) + \gamma z + \sum_{n=1}^{\infty} \log\left(1 + \frac{z}{n}\right) - \frac{z}{n}$$

$$\xrightarrow{\frac{d}{dz}} -\frac{\Gamma'(z)}{\Gamma(z)} = \frac{1}{z} + \gamma + \sum_{n=1}^{\infty} \frac{\frac{1}{n}}{1 + z/n} - \frac{1}{n}$$

$$= \frac{1}{z} + \gamma + \sum_{n=1}^{\infty} \frac{1}{z+n} - \frac{1}{n}$$

$$\rightarrow \frac{\Gamma'(z)}{\Gamma(z)} = -\frac{1}{z} - \gamma - \sum_{n=1}^{\infty} \frac{1}{z+n} - \frac{1}{n}$$

as desired.

b) we have

$$\frac{\Gamma'(z+1)}{\Gamma(z+1)} - \frac{\Gamma'(z)}{\Gamma(z)} = -\frac{1}{z}$$

which we'd like to show  $= 0$ . So to that effect, we can show  $\frac{\Gamma'(z+1)}{\Gamma(z+1)} = \frac{1}{z} + \frac{\Gamma'(z)}{\Gamma(z)}$ .

From (a), we see

$$\frac{\Gamma'(z+1)}{\Gamma(z+1)} = -\frac{1}{z+1} - \gamma - \sum_{n=1}^{\infty} \frac{1}{z+1+n} - \frac{1}{z}$$

and

$$\frac{1}{z} + \frac{\Gamma'(z)}{\Gamma(z)} = \frac{1}{z} - \frac{1}{z} - \gamma - \sum_{n=1}^{\infty} \frac{1}{z+n} - \frac{1}{z}$$

$$= -\gamma - \sum_{n=1}^{\infty} \frac{1}{z+n} - \frac{1}{z}$$

$$= -\frac{1}{z+1} - \gamma - \left( \sum_{n=2}^{\infty} \frac{1}{z+n} - \frac{1}{z} \right) + 1$$

$$= -\frac{1}{z+1} - \gamma - \left( \sum_{n=1}^{\infty} \frac{1}{z+n+1} - \frac{1}{n+1} \right) + 1$$

$$= -\frac{1}{z+1} - \gamma - \left[ \left( \frac{1}{z+2} - \frac{1}{2} \right) + \left( \frac{1}{z+3} - \frac{1}{3} \right) + \dots + -1 \right]$$

$$= -\frac{1}{z+1} - \gamma - \left[ \left( \frac{1}{z+2} - 1 \right) + \frac{1}{2} + \left( \frac{1}{z+3} - \frac{1}{2} \right) + \frac{1}{6} + \dots - 1 \right]$$

$$= -\frac{1}{z+1} - \gamma + 1 - \sum_{n=1}^{\infty} \frac{1}{z+1+n} - \frac{1}{n} + \left( \frac{1}{n} - \frac{1}{n+1} \right)$$

$$\stackrel{(*)}{=} -\frac{1}{z+1} - \gamma + 1 - \left( \sum_{n=1}^{\infty} \frac{1}{z+1+n} - \frac{1}{n} \right) - \sum_{n=1}^{\infty} \left( \frac{1}{n} - \frac{1}{n+1} \right)$$

where the splitting of the above sum is allowed  
 since  $\sum_{n=1}^{\infty} \frac{1}{n} - \frac{1}{n+1} = \sum_{n=1}^{\infty} \frac{1}{n(n+1)} < \sum_{n=1}^{\infty} \frac{1}{n^2} = \frac{\pi^2}{6} < \infty$

But notice

$$\sum_{n=1}^{\infty} \frac{1}{n} - \frac{1}{n+1} = 1 - \frac{1}{2} + \frac{1}{2} - \frac{1}{3} + \frac{1}{3} + \dots$$

$$= \lim_{N \rightarrow \infty} \sum_{n=1}^N \frac{1}{n} - \frac{1}{n+1}$$

$$= \lim_{N \rightarrow \infty} 1 - \frac{1}{N+1}$$

$$= 1$$

telescopes

So  $(*)$

$$= -\frac{1}{z+1} - \gamma - \sum_{n=1}^{\infty} \frac{1}{z+1+n} - \frac{1}{n} = \frac{\Gamma'(z+1)}{\Gamma(z)} \quad \text{as}$$

desired, establishing

$$\frac{\Gamma'(z+1)}{\Gamma(z+1)} - \frac{\Gamma'(z)}{\Gamma(z)} - \frac{1}{z} = 0$$

By integrating both sides of the equation  
we see  $\int w/$  respect to  $z$ )  $\downarrow$  (some constant)

$$\ln(\Gamma(z+1)) - \ln(\Gamma(z)) - \ln(z) = C$$

$$\ln\left(\frac{\Gamma(z+1)}{z\Gamma(z)}\right) = C$$

$$\frac{\Gamma(z+1)}{z\Gamma(z)} = e^C = e^C$$

$$\Gamma(z+1) = e^C \cdot z \cdot \Gamma(z) = C' z \cdot \Gamma(z)$$

also as desired,

$$c) \lim_{z \rightarrow 0} z \cdot \Gamma(z) = \lim_{z \rightarrow 0} \frac{z}{\frac{1}{\Gamma(z)}}$$

$$= \lim_{z \rightarrow 0} \frac{z}{e^{\gamma z} \prod_{n=1}^{\infty} \left(1 + \frac{z}{n}\right) e^{-z/n}}$$

$$= \frac{1}{e^{\gamma(0)} \prod_{n=1}^{\infty} 1} = \frac{1}{1} = 1$$

$$\text{So } \lim_{z \rightarrow 0} \Gamma(z+1) = \lim_{z \rightarrow 0} C' z \cdot \Gamma(z)$$

$$\rightarrow \Gamma(1) = C' \cdot 1 = C'$$

as desired.

d) If we want  $\Gamma(1)=1$ , that is equivalent to  $\frac{1}{\Gamma(1)}=1$ , so when  $z=1$

$$1 := \frac{1}{\Gamma(1)} = e^{\gamma} \prod_{n=1}^{\infty} \left(1 + \frac{1}{n}\right) e^{-1/n}$$

$$\rightarrow e^{-\gamma} = \prod_{n=1}^{\infty} \left(1 + \frac{1}{n}\right) e^{-1/n} \quad (\text{as desired}).$$

e) we have,

$$\prod_{n=1}^{\infty} \left(1 + \frac{1}{n}\right) e^{-1/n} = \prod_{n=1}^{\infty} \frac{n+1}{n} e^{-1/n}$$

$$= \lim_{N \rightarrow \infty} \prod_{n=1}^N \frac{n+1}{n} e^{-1/n} \quad \downarrow \text{grouping terms}$$

$$= \lim_{N \rightarrow \infty} \left( \frac{2}{1} \frac{3}{2} \frac{4}{3} \dots \frac{N+1}{N} \right) \left( e^{-1/1} e^{-1/2} \dots e^{-1/N} \right)$$

$$= \lim_{N \rightarrow \infty} \left( \frac{2}{1} \frac{3}{2} \frac{4}{3} \dots \frac{N+1}{N} \right) e^{-S(N)} \quad (\text{as desired})$$

$$= \lim_{N \rightarrow \infty} (N+1) e^{-S(N)} \quad \begin{array}{l} (\text{as also desired}) \\ (\text{by cancelling fractions}) \end{array}$$



So,

$$e^{-\gamma} = \prod_{n=1}^{\infty} \left(1 + \frac{1}{n}\right) e^{-1/n} = \lim_{n \rightarrow \infty} (n+1) e^{-S(n)}$$

$$-\gamma = \log \left( \lim_{n \rightarrow \infty} (n+1) e^{-S(n)} \right)$$

$$= \lim_{n \rightarrow \infty} \log \left( (n+1) e^{-S(n)} \right) \quad \left[ \begin{array}{l} \text{as the input} \\ \text{is } > 0 \text{ and} \\ \log \text{ is continuous} \end{array} \right]$$

$$= \lim_{n \rightarrow \infty} \log(n+1) - S(n)$$

$$\rightarrow \gamma = \lim_{n \rightarrow \infty} S(n) - \log(n+1)$$

$$= \lim_{n \rightarrow \infty} \left[ \sum_{k=1}^n \frac{1}{k} - \log(n+1) \right]$$

as desired.



4) without loss of generality, assume  $a_n \neq 0$  for any  $n$ , then (wlog since finitely many  $z$ 's with  $|z| < 1$  won't bother the convergence of the product).

$$H(z) = \prod_{n=0}^{\infty} B(a_n, z)$$

$$= \prod_{n=0}^{\infty} (1 + (-1 + B(a_n, z)))$$

$$= \prod_{n=0}^{\infty} (1 + \underbrace{\left( \frac{|a_n|}{a_n} \frac{a_n - z}{1 - \bar{a}_n z} - 1 \right)}_{(*)})$$

we'd like to use Weierstrass M-Test, so we'd like to bound  $(*)$  by some constant  $M_n$  such that  $\sum_{n=0}^{\infty} M_n < \infty$ . So

$$|(*)| = \left| \frac{|a_n| a_n - |a_n| z - a_n + |a_n|^2 z}{a_n (1 - \bar{a}_n z)} \right|$$

find common denominator

$$= \left| \frac{(|a_n| - 1)(a_n - |a_n| z)}{a_n (1 - \bar{a}_n z)} \right|$$

$$= \left| \frac{(|a_n| - 1)(1 - \bar{a}_n z)}{1 - \bar{a}_n z} \right|$$

divide the right term by  $a_n$

$$\leq \frac{|1 - |a_n|| |1 - \bar{a}_n z|}{|1 - \bar{a}_n z|}$$

$$= 1 - |a_n|$$

Since we are given that  $\sum_{n=0}^{\infty} 1 - |a_n| < \infty$ ,

we choose our  $M_n$ 's to be  $1 - |a_n|$ . Each

$B(a_n, z)$  is analytic everywhere except when

$z = \frac{1}{a_n}$ , so by the M-test, we get that  $H(z)$  is analytic  $\forall z$  in the unit disk other than at the points just mentioned.

To see the zeroes of  $H$ , this occurs when

$$B(a_n, z) = 0$$

$$\frac{|a_n|}{a_n} \frac{a_n - z}{1 - \bar{a}_n z} = 0$$

$$\Leftrightarrow z = a_n.$$

And if  $a = 0$ , then  $z = 0$  also is a zero.

So  $H$  has countably many zeroes inside the unit disk.

$H(z)$  need not be 0  $\forall |z| \leq 1$  because our function at  $a=1$  is 1 for all  $z$ .

Since  $\sum_{n=0}^{\infty} 1 - |a_n| < \infty$ , we get for free that

$\lim_{n \rightarrow \infty} 1 - |a_n| = 0 \rightarrow \lim_{n \rightarrow \infty} |a_n| = 1$ . Since our function

has 0's at every  $a_n$  but not at the limit point  $a=1$ , we have that the zeros of  $H$  are not dense inside the unit disk.

5) a) we can write with  $\omega = (\omega_1, \omega_2)$

$$p(z, \omega) = \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(z + \omega \cdot i)^2} - \frac{1}{(\omega \cdot i)^2}$$

Put  $z' = z + M\omega_1 + N\omega_2 = z + T \cdot \omega$  ( $T = (M, N)$ ),

then

$$p(z', \omega) = \frac{1}{(z + T \cdot \omega)^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(z + T \cdot \omega + i \cdot \omega)^2} - \frac{1}{(\omega \cdot i)^2}$$

$$= \frac{1}{(z + T \cdot \omega)^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(z + (T+i) \cdot \omega)^2} - \frac{1}{(\omega \cdot i)^2}$$

Put  $l = (T+i)$ , then the above becomes

$$= \frac{1}{(z + T \cdot \omega)^2} + \sum_{\substack{l=(j,k) \\ l \neq T}} \frac{1}{(z + l \cdot \omega)^2} - \frac{1}{(\omega \cdot (T-l))^2}$$

$$= \frac{1}{z^2} + \sum_{\substack{l=(j,k) \\ l \neq 0}} \frac{1}{(z + l \cdot \omega)^2} - \frac{1}{(\omega \cdot (T-l))^2}$$

swapping  
1 term  
from  
inside  
and  
the  
one  
outside

$$= \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq 0}} \frac{1}{(z + i \cdot \omega)^2} - \frac{1}{(\omega \cdot i)^2}$$

replacing  
 $T-l$  with  $i$   
by the slab  
definitions and  
realizing  $l$  and  $i$   
now have the  
same index range

$$= p(z, \omega) \text{ (as desired).}$$

So  $p$  is doubly periodic.



b)

$$p(-z, w) = \frac{1}{(-z)^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(-z - w \cdot i)^2} - \frac{1}{(w \cdot i)^2}$$

Replacing  $i$  with  $-i$  won't change the sum, so

$$= \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(-z + w \cdot i)^2} - \frac{1}{(w \cdot i)^2}$$

↙ this won't change b/c of square

$$= \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \left[ \frac{1}{(-1)(z - w \cdot i)} \right]^2 - \frac{1}{(w \cdot i)^2}$$

$$= \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(z - w \cdot i)^2} - \frac{1}{(w \cdot i)^2}$$

$$= p(z, w)$$

so  $p$  is an even function in  $z$ .

$$c) \quad p(z, w) = \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(z - w \cdot i)^2} - \frac{1}{(w \cdot i)^2}$$

$$= \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(w \cdot i)^2} \left[ \left( \frac{z}{w \cdot i} - 1 \right)^2 - 1 \right]$$

Taylor series for  $\frac{1}{(z-1)^2}$

$$= \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{1}{(w \cdot i)^2} \left[ \sum_{n=1}^{\infty} n \frac{z^{n-1}}{(w \cdot i)^{n-1}} - 1 \right]$$

$$= \frac{1}{z^2} + \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \sum_{n=2}^{\infty} n \frac{z^{n-1}}{(w \cdot i)^{n+1}}$$

$$= \frac{1}{z^2} + \sum_{n=2}^{\infty} \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{n z^{n-1}}{(w \cdot i)^{n+1}}$$

$$= \frac{1}{z^2} + \sum_{n=1}^{\infty} \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{(2n+1) z^{2n}}{(w \cdot i)^{2n+2}}$$

since  $p$  is even

So our Laurent series has coefficients

$$a_{2n} = \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{(2n+1)}{(w \cdot i)^{2n+2}}$$

Let's find  $a_0$ ,  $a_2$ , and  $a_4$ .  $a_0 = 0$ .

$$a_2 = \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{3}{(w \cdot i)^4}, \quad a_4 = \sum_{\substack{i=(j,k) \\ i \neq (0,0)}} \frac{5}{(w \cdot i)^6}.$$

So that

$$p(z) = \frac{1}{z^2} + a_2 z^2 + a_4 z^4 + \dots$$

$$p'(z) = -\frac{2}{z^3} + \overbrace{2a_2}^{b_1} z + \overbrace{4a_4}^{b_2} z^3 + \dots$$

d) we have,

$$(p')^2 = \left( -\frac{2}{z^3} + a_2 z^2 + a_4 z^4 + \dots \right)^2$$

$$= \frac{4}{z^6} - \frac{8a_2}{z^2} - 16a_4 + f_1(z)$$

where  $f_1(z)$  has only positive powers of  $z$ .

Also,

$$p^3 = \left( \frac{1}{z^2} + a_2 z^2 + a_4 z^4 + \dots \right)^3$$

$$= \frac{1}{z^6} + \frac{3a_2}{z^2} + 3a_4 + f_2(z)$$

where  $f_2(z)$  also has only positive powers of  $z$ ,  
we can then see

$$(p')^2 - 4p^3 = \frac{-20a_2}{z^2} - 28a_4 + f_1(z) - 4f_2(z)$$

Since  $p(z) = \frac{1}{z^2} + a_2 z^2 + \dots$ ,  $-20a_2 \cdot p(z)$  only has  $\frac{1}{z^2}$  as a <sup>(non positive)</sup> negative power term, and we can write



$$(p')^2 - 4p^3 + 20a_2 p + 28a_4 = f_1(z) - 4f_2(z) + f_3(z)$$

where again,  $f_3(z)$  has only positive values of  $z$ . By construction, the above function has no poles, it is a polynomial with only positive powers of  $z$ . Since  $p$  is elliptic, so is the above formula, meaning it is doubly periodic. It is <sup>(analytic and hence)</sup> continuous and bounded on each

lattice parallelogram  $jw_1 + kw_2$  (since polynomials are entire), and since the function <sup>(periodic)</sup> repeats itself on each parallelogram, it must be analytic and bounded on all of  $\mathbb{C}$ . Liouville's theorem then implies the above is constant, i.e.,

$$(p')^2 - 4p^3 + 20a_2 p + 28a_4 = C \quad \text{for some } C \in \mathbb{C}$$

But plugging 0 into the RHS at the top of this page trivially returns 0 (only positive powers of  $z$ ).

Hence,

$$(p')^2 = 4p^3 - 20a_2 p - 28a_4$$

$$\text{So } a = 4, \quad b = 0, \quad c = -20a_2, \quad d = -28a_4.$$