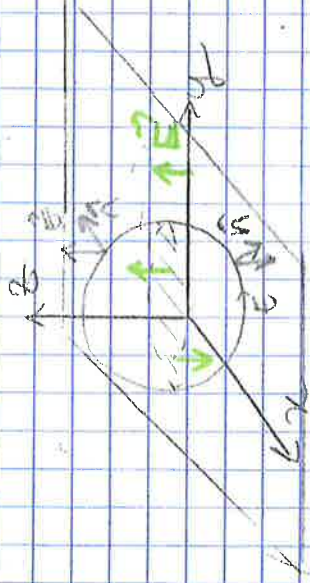


## 2.6 Distribution planaire de charge



① Gauss : Forme integrale

$$\oint \vec{E} \cdot \vec{n} \, dS = \frac{Q_{\text{int}}}{\epsilon_0}$$

avec  $Q_{\text{int}} = \rho_s \pi R^2$

Donc  $\oint \vec{E} \cdot \vec{n} \, dS = \frac{\rho_s \pi R^2}{\epsilon_0}$

② Calcul de l'integrale

On choisit le système de coordonnées sphérique

$S_1$  est la surface de la demi-sphère sup.

$$\oint_S \vec{E} \cdot \vec{n} \, dS = \int_{S_1} \vec{E} \cdot \vec{n}_1 \, dS + \int_{S_2} \vec{E} \cdot \vec{n}_2 \, dS$$

$$= \int_{S_1} \frac{\rho_s \hat{e}_z \cdot \hat{e}_r \, dS}{2\epsilon_0} - \int_{S_2} \frac{\rho_s \hat{e}_r \cdot \hat{e}_r \, dS}{2\epsilon_0}$$

$$= \int_{S_1} \frac{\rho_s}{2\epsilon_0} \cos \theta \, dS - \int_{S_2} \frac{\rho_s}{2\epsilon_0} \cos \theta \, dS$$

$$= \frac{\rho_s}{2\epsilon_0} \left( \int_{S_1} \cos \theta \, dS \right) - \left( \int_{S_2} \cos \theta \, dS \right)$$

$$dS = R^2 \sin \theta \, d\theta \, d\varphi$$

$$\oint_S \vec{E} \cdot \vec{n} \, dS = \frac{\rho_s}{2\epsilon_0} \left( \int_0^{2\pi} \int_0^{\frac{\pi}{2}} R^2 \sin \theta \cos \theta \, d\theta \, d\varphi \right)$$

$$= \frac{\rho_s}{2\epsilon_0} \int_0^{2\pi} \int_0^{\frac{\pi}{2}} R^2 \sin \theta \cos \theta \, d\theta \, d\varphi$$

$$= \frac{\rho_s}{2\epsilon_0} \int_0^{2\pi} \left( \int_0^{\frac{\pi}{2}} \sin \theta \cos \theta \, d\theta \right) d\varphi$$

$$= \frac{\rho_s \pi R^2}{\epsilon_0} \left( +\frac{1}{2} + \frac{1}{2} \right)$$

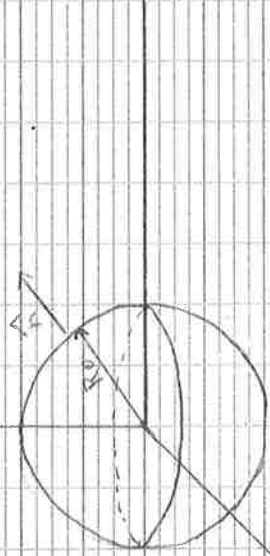
$$= \frac{\rho_s \pi R^2}{\epsilon_0}$$





# Electromagnetisme TD4

## 4.1 Une sphere chargée



$$\vec{E}(\vec{r}) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{Q_0}{R_0^3} \vec{r} & r \leq R_0 \\ \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r^2} \vec{r} & r > R_0 \end{cases}$$

Pà l'infini

Soit M un point de l'espace

$$V(\vec{r}_P) = 0$$

$$V(\vec{r}_M) - V(\vec{r}_P) = - \int_{\vec{r}_P \rightarrow \vec{r}_M} \vec{E} \cdot \vec{r} \, dl$$

Cas  $r > R_0$ : (Extérieur)

$$V(\vec{r}_M) = - \int_{\vec{r}_P \rightarrow \vec{r}_M} \vec{E} \cdot \vec{r} \, dl$$

couleur irrégulière



13,25

U

14,55 - 14,85

17,33

14,25

16,08

20

12,50

8,52

11,50

17,04

7,52

U

8,82 - 10,16

$$V(\vec{r}_M) = - \int \vec{E} \cdot \vec{r} \, dl = \int_r^\infty E \cdot dr$$

$$\text{or } \vec{E} = \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r^2} \vec{e}_r \text{ done } V(\vec{r}_M) = \int_{r_M}^\infty \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r^2} dr$$

$$dl = dr$$

$$V(\vec{r}_M) = \frac{1}{4\pi\epsilon_0} \int_{r_M}^\infty \frac{1}{r^2} dr$$

$$\int_{r_M}^\infty \frac{1}{r^2} dr = \lim_{x \rightarrow \infty} \int_{r_M}^x \frac{1}{r^2} = \lim_{x \rightarrow \infty} \left[ -\frac{1}{r} \right]_{r_M}^x$$

$$= \lim_{x \rightarrow \infty} \left( -\frac{1}{x} + \frac{1}{r_M} \right) = \frac{1}{r_M}$$

$$\text{Done } V(\vec{r}_M) = \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r_M}$$

$$\text{cas } r \leq R_0 :$$

On appelle  $A$  le point de la droite  $(PM)$  qui coupe la sphère. A est à une distance  $R_0$  de l'origine.

$$V(\vec{r}_M) - V(\vec{r}_P) = - \int_{\vec{r}_P \rightarrow \vec{r}_M} \vec{E} \cdot \vec{r} d\ell$$

on  $V(\vec{r}_P) = 0$  et  $\vec{r}_P \rightarrow \vec{r}_M = \underbrace{\vec{r}_P \rightarrow \vec{r}_A}_{\text{hors de la sphère}} + \underbrace{\vec{r}_A \rightarrow \vec{r}_M}_{\text{dans la sphère}}$

$$V(\vec{r}_M) = - \left( \int_{\vec{r}_P \rightarrow \vec{r}_A} \vec{E} \cdot \vec{r} d\ell + \int_{\vec{r}_A \rightarrow \vec{r}_M} \vec{E} \cdot \vec{r} d\ell \right)$$

$$= - \left( \int_{\vec{r}_P \rightarrow \vec{r}_A} \underbrace{\vec{E} \cdot \vec{r} d\ell}_{V(\vec{r}_A)} - \int_{\vec{r}_A \rightarrow \vec{r}_M} \vec{E} \cdot \vec{r} d\ell \right)$$

$$V_{r_M} = \frac{Q_0}{4\pi\epsilon_0 R_0} - \int_{\vec{r}_A \rightarrow \vec{r}_M} \vec{E} \cdot \vec{r} d\ell \quad (*)$$

Calculons :

$$- \int_{\vec{r}_A \rightarrow \vec{r}_M} \vec{E} \cdot \vec{r} d\ell : \text{ Dans la sphère}$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{Q_0 r}{R_0^3} \hat{r} \quad r = -a \text{ et } d\ell = dr$$

$$r \cdot \hat{r} = -1$$

$$\begin{aligned}
 - \int_{\vec{r}_M \rightarrow \vec{r}_A} \vec{E} \cdot \vec{r} d\vec{r} &= \int_{r_M}^{r_A} \frac{1}{4\pi\epsilon_0} \frac{Q_0}{R_0^3} R_0^2 dr \\
 &= \frac{Q_0}{4\pi\epsilon_0 R_0^3} \int_{r_M}^{r_A} r dr \\
 &= \frac{Q_0}{4\pi\epsilon_0 R_0^3} \left[ \frac{r^2}{2} \right]_{r_M}^{r_A} \\
 &= \frac{Q_0}{4\pi\epsilon_0 R_0^3} \frac{R_0^2 - r_M^2}{2} \quad (**)
 \end{aligned}$$

(\*) at (\*\*) )

$$\begin{aligned}
 V(\vec{r}_M) &= \frac{Q_0}{4\pi\epsilon_0} \frac{1}{R_0} + \frac{Q_0}{4\pi\epsilon_0 R_0^3} \left( \frac{R_0^2 - r_M^2}{2} \right) \\
 &= \frac{Q_0}{4\pi\epsilon_0} \left( \frac{1}{R_0} + \frac{R_0^2 - r_M^2}{2R_0^3} \right) \\
 &= \frac{Q_0}{4\pi\epsilon_0} \left( \frac{2R_0^2 + R_0^2 - r_M^2}{2R_0^3} \right) \\
 &= \frac{Q_0}{4\pi\epsilon_0} \left( \frac{3R_0^2 - r_M^2}{2R_0^3} \right)
 \end{aligned}$$



# Electromagnetisme TD 4

4.1 (Suite)

$$[r > R_0]$$

Raccourci?

Condition s. :  $\vec{E} = E_r(r) \vec{e}_r$

On sait que

$$\begin{aligned} \vec{E}(x, y, z) &= -\text{grad}(V(r, \varphi, \theta)) \\ &= -\left( \frac{\partial V}{\partial x} \vec{e}_x + \frac{\partial V}{\partial y} \vec{e}_y + \frac{\partial V}{\partial z} \vec{e}_z \right) \end{aligned}$$

$$\begin{aligned} \vec{E}(r, \varphi, \theta) &= -\text{grad}(V(r, \varphi, \theta)) \\ &= -\left( \frac{\partial V}{\partial r} \vec{e}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \vec{e}_\theta + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \varphi} \vec{e}_\varphi \right) \end{aligned}$$

Comme  $\vec{E} = E_r \vec{e}_r$  alors :  $\frac{1}{r} \frac{\partial V}{\partial \theta} = 0$  et  $\frac{1}{r \sin \theta} \frac{\partial V}{\partial \varphi} = 0$

$$\vec{E} = E_r \vec{e}_r = -\frac{\partial V}{\partial r} \vec{e}_r$$

$$V(r) = - \int E_r(r) dr + \text{cst}$$

$$= - \int \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r^2} dr + \text{cst}$$

$$= -\frac{1}{4\pi\epsilon_0} Q_0 \int \frac{1}{r^2} dr = \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r} + \text{cst}$$

$$\begin{array}{c} V(r) \\ 11 \\ 0 \end{array}$$



## Electromag:

### Fin exercice 4.1:

Raccourci pour  $r \leq R_0$ :

$$\vec{E} = E_r(r) \hat{e}_r$$

∴ m chose que <sup>cas</sup> précédent

$$V(r) = - \int E_r(r) dr + \text{conste}$$

$$= - \int \frac{1}{4\pi\epsilon_0} \times \frac{Q_0}{R_0^3} r dr + \text{Cste}$$

$$= - \frac{Q_0}{4\pi\epsilon_0} \times \frac{1}{R_0^3} \int r dr + \text{Cste}$$

$$= - \frac{1}{4\pi\epsilon_0} \times \frac{Q_0 r^2}{2 R_0^3} + \text{Cste}$$

Or  $V$  est une fonction continue, et on se place en  $r = R_0$ . Ainsi  $V_{\text{int}}(R_0) = V_{\text{ext}}(R_0)$ .

$$\Leftrightarrow - \frac{1}{4\pi\epsilon_0} \times \frac{Q_0 R_0^2}{2 R_0^3} + \text{Cste} = \frac{1}{4\pi\epsilon_0} \times \frac{Q_0}{R_0}$$

$$\Leftrightarrow \text{Cste} = \frac{1}{4\pi\epsilon_0} \left( \frac{Q_0}{R_0} + \frac{Q_0}{2R_0} \right)$$

$$\Leftrightarrow \text{Cste} = \frac{Q_0}{4\pi\epsilon_0 R_0} \left( 1 + \frac{1}{2} \right)$$

$$\Leftrightarrow \text{Cste} = \frac{3}{2} \frac{Q_0}{4\pi\epsilon_0 R_0}$$

Conclusion: Pour  $r \leq R_0$ .

$$V(r) = - \frac{1}{4\pi\epsilon_0} \frac{Q_0 r^2}{2 R_0^3} + \frac{3}{2} \frac{Q_0}{4\pi\epsilon_0 R_0}$$



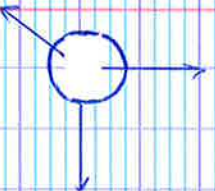
$$\begin{aligned}
 J &= \int_{V_{\text{ext}}} \epsilon_0 \cdot E^2(\vec{r}) dV \\
 &= \epsilon_0 \int_{V_{\text{ext}}} \left( \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r} \right)^2 dV \\
 &= \frac{\epsilon_0 Q_0^2}{16\pi^2 \epsilon_0^2} \int_{V_{\text{ext}}} \frac{1}{r^2} dV \\
 &= \frac{Q_0^2}{16\pi^2 \epsilon_0} \int_R^{+\infty} \frac{1}{r^2} dr \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\varphi \\
 &= \frac{Q_0^2}{16\pi^2 \epsilon_0} \frac{1}{R} 4\pi \\
 &= \frac{Q_0^2}{4\pi \epsilon_0 R}.
 \end{aligned}$$

Finalement :  $W_e = \frac{1}{2} (I + J)$

$$\begin{aligned}
 &= \frac{1}{2} \left( \frac{Q_0^2}{20\pi \epsilon_0 R} + \frac{Q_0^2}{4\pi \epsilon_0 R} \right) \\
 &= \frac{Q_0^2}{10\pi \epsilon_0 R} \left( \frac{1}{4} + \frac{1}{8} \right)
 \end{aligned}$$

$$W_e = \frac{3}{20\pi \epsilon_0 R} Q_0^2$$

Exercice 4.2 :



a) Symétrie : axiale  
 Coordonnée : sphérique  
 $\vec{E}(\vec{r}) = \begin{cases} E(r) \vec{e}_r & r > R \\ -E(r) \vec{e}_r & r \leq R \end{cases}$

Prendre une sphère fictive de rayon  $r \leq R$ .

Gauss :  $\int \vec{E}(\vec{r}) \cdot \vec{n} dS = \frac{Q_{\text{int}}}{\epsilon_0}$

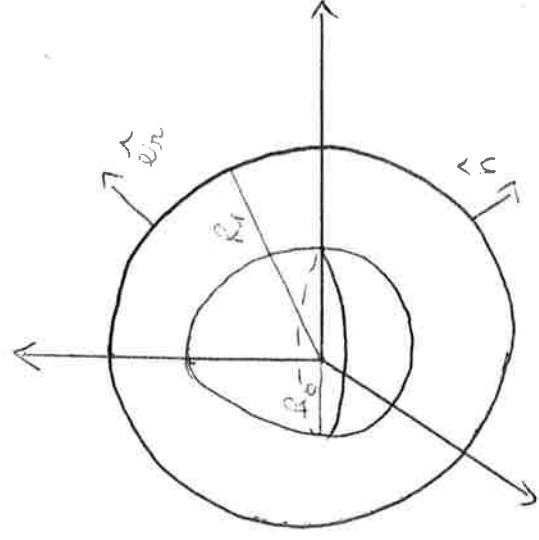
ou  $Q_{\text{int}} = 0$   
 donc  $\vec{E}(\vec{r}) = \vec{0}$  si  $r \leq R$

## Exercice 4.2:

### a) Méthode de Gauss:

Symétrie axiale : coord. sphériques

$$\vec{E}(\vec{r}) = \begin{cases} E(r) \hat{e}_r & r > R_0 \\ -E(r) \hat{e}_r & r \leq R_0 \end{cases}$$



Sphère rayon  $r \leq R_0$ :

$$\oint_S \vec{E}(\vec{r}) \cdot \hat{n} dS = \frac{Q_{int}}{\epsilon_0}$$

$$\text{Car } Q_{int} = 0$$

$$\text{Donc } \vec{E}(\vec{r}) = \vec{0} \quad \text{si } r \leq R_0$$

Sphère rayon  $R_1 > R_0$ :

$$\oint_S \vec{E}(\vec{r}) \cdot \hat{n} dS = \frac{Q_{int}}{\epsilon_0} = \frac{Q_0}{\epsilon_0} \quad (Q_{int} = Q_0)$$

$$\Rightarrow \oint_S E(r) \hat{e}_r \cdot \hat{n} dS = \frac{Q_0}{\epsilon_0}$$

$$\Rightarrow E(R_1) \oint_S dS = \frac{Q_0}{\epsilon_0}$$

$$\Rightarrow E(R_1) S = \frac{Q_0}{\epsilon_0}$$

$$\Rightarrow E(R_1) S = \frac{Q_0}{\epsilon_0}$$

$$\text{avec } S = 4 \pi R_1^2$$

$$\Rightarrow E(R_1) = \frac{Q_0}{4 \pi \epsilon_0 R_1^2}$$

valable pour tout  $R_1 > R_0$  (on l'appelle  $r$ )

$$\text{Donc } \vec{E}(\vec{r}) = \frac{1}{4 \pi \epsilon_0} \frac{Q_0}{r^2} \hat{e}_r$$

Donc  $\vec{E}(\vec{r}) = \int \vec{0}$  si  $r \leq R_0$

$$\begin{cases} \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r^2} e_r & \text{si } r > R_0 \end{cases}$$

b)  $V(\vec{r}_B) - V(\vec{r}_A) = - \int_{\vec{r}_A \rightarrow \vec{r}_B} \vec{E}(\vec{r}) \cdot \hat{r} d\ell = \frac{W_{A \rightarrow B}}{q}$

$V(\vec{r}_A) = 0$  car A est à l'infini

Pour  $r > R_0$ :

Donc  $V(\vec{r}_B) = - \int_{\vec{r}_A \rightarrow \vec{r}_B} \vec{E}(\vec{r}) \cdot \hat{r} d\ell = - \int_{\infty}^r \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r^2} \underbrace{e_r \cdot \hat{r}}_{=1} d\ell$

$$= \int_r^\infty \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r^2} d\ell$$

$$= \frac{Q_0}{4\pi\epsilon_0} \int_r^\infty \frac{1}{r^2} d\ell \quad \text{Or } d\ell = dr$$

Donc  $V(\vec{r}_B) = \frac{Q_0}{4\pi\epsilon_0} \left[ \frac{-1}{r} \right]_r^\infty = \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r}$  valable pour tout  $r > R_0$

Pour  $r \leq R_0$ :

$$V(\vec{r}_B) = - \int_\infty^{R_0} \vec{E}(\vec{r}) \cdot \hat{r} d\ell - \underbrace{\int_{R_0}^r \vec{E}(\vec{r}) \cdot \hat{r} d\ell}_0$$

$$= - \int_\infty^{R_0} \vec{E}(\vec{r}) \cdot \hat{r} d\ell = V(R_0) = \frac{1}{4\pi\epsilon_0} \frac{Q_0}{R_0} \quad \text{valable pour tout } r \leq R_0$$

Donc  $V(\vec{r}) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{Q_0}{R_0} & \text{si } r \leq R_0 \\ \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r} & \text{si } r > R_0 \end{cases}$



$$c) U_e = \frac{1}{2} \int_V \rho(\vec{r}) V(\vec{r}) dV \quad (4)$$

$$U_e = \frac{1}{2} \int_V \epsilon_0 E(\vec{r})^2 dV \quad (5)$$

$$\begin{aligned} (4) \quad U_e &= \frac{1}{2} \int_V \rho(\vec{r}) V(\vec{r}) dV \\ &= \frac{1}{2} \oint_S \rho(\vec{r}) V(\vec{r}) dS \\ &= \frac{\rho}{2} \oint_S V(\vec{r}) dS \\ &= \frac{1}{2} \rho V(R_0) \oint_S dS \\ &= \frac{1}{2} \rho V(R_0) S \\ &= \frac{1}{2} \frac{Q_0}{S} V(R_0) S \\ &= \frac{Q_0}{2} \frac{1}{4\pi\epsilon_0} \frac{Q_0}{R_0} \\ &= \frac{1}{8\pi\epsilon_0} \frac{Q_0^2}{R_0} \end{aligned}$$

$$d) R_0 = 30 \text{ cm}$$

$$E_{\text{max}} = 3 \text{ MV} \cdot \text{m}^{-1}$$

$$\vec{E}(\vec{r}) = \frac{Q_0}{4\pi\epsilon_0} \frac{1}{r^2} \quad r > R_0$$

On fait tendre  $r$  vers  $R_0$

$$\text{Donc } \vec{E}(\vec{r}) = \frac{Q_0}{4\pi\epsilon_0} \frac{1}{R_0^2}$$

$$\begin{aligned} \text{Donc } Q_0 &= 4\pi\epsilon_0 R_0^2 \vec{E}(\vec{r}) = 4\pi \times 8,85 \times 10^{-12} \times (30 \times 10^{-2})^2 \times 3 \times 10^6 \\ &= 30 \mu\text{C} \end{aligned}$$

$$(5) U_e = \frac{1}{2} \int_V \epsilon_0 E(\vec{r})^2 dV$$

$$\begin{aligned} &= \frac{1}{2} \epsilon_0 \int_V \left( \frac{1}{4\pi\epsilon_0} \frac{Q_0}{r^2} \right)^2 dV \\ &= \frac{Q_0^2}{32\pi^2\epsilon_0} \int_V \frac{1}{r^4} dV \end{aligned}$$

$$\text{avec } dV = r^2 \sin\theta \, dr \, d\theta \, d\varphi$$

$$\begin{aligned} \text{Donc } U_e &= \frac{Q_0^2}{32\pi^2\epsilon_0} \int_{R_0}^{\infty} \frac{1}{r^2} dr \int_0^\pi \sin\theta \, d\theta \int_0^{2\pi} d\varphi \\ &= \frac{Q_0^2}{32\pi^2\epsilon_0} \frac{1}{R_0} 2 \times 2\pi \end{aligned}$$

$$= \frac{1}{8\pi^2\epsilon_0} \frac{Q_0^2}{R_0}$$

$$V(R_0) = \frac{1}{4\pi\epsilon_0 R_0} \frac{Q_0}{R_0} = \frac{1}{4\pi\epsilon_0 \times 8,85 \times 10^{-12}} \times \frac{30 \times 10^{-6}}{30 \times 10^{-2}} = 900 \text{ kV}$$

$$\frac{Q_0}{V(R_0)} = \frac{30 \times 10^{-6}}{900 \times 10^3} = 33,33 \text{ pF}$$

$$V(1m) = \frac{1}{4\pi\epsilon_0} \frac{Q_0}{1} = \frac{30 \times 10^{-6}}{4\pi \times 8,85 \times 10^{-12}} = 270 \text{ kV}$$

$$V(2m) = \frac{V(1m)}{2} = 135 \text{ kV}$$

$$V(3m) = \frac{V(1m)}{3} = 90 \text{ kV}$$