

Navdeep Singh Dhindsa **02/08/2024**

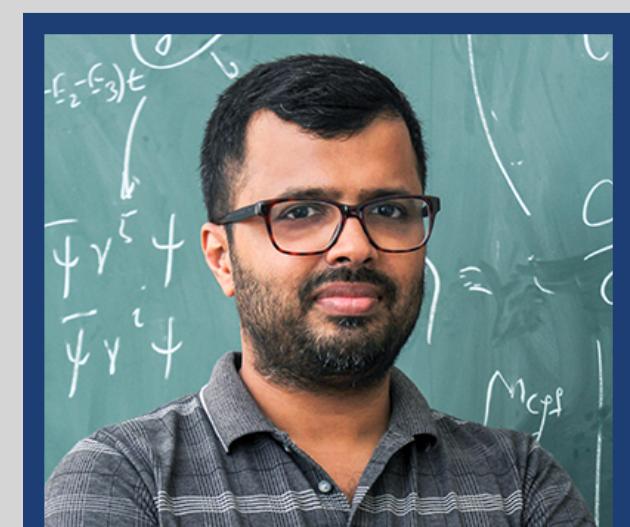
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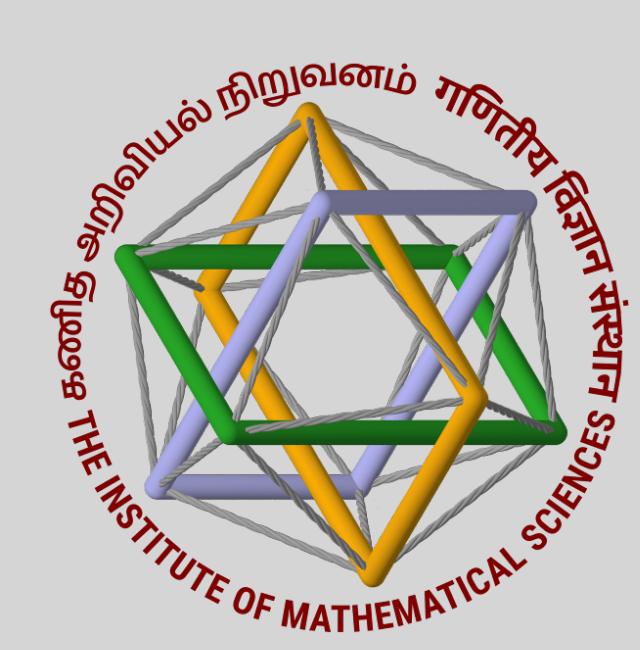
Exploring Single-Flavor Dibaryons: A lattice perspective

Funding resources



Work in collaboration with
M. Padmanath (IMSc Chennai) and Nilmani Mathur (TIFR Mumbai)





Dibaryons - Status

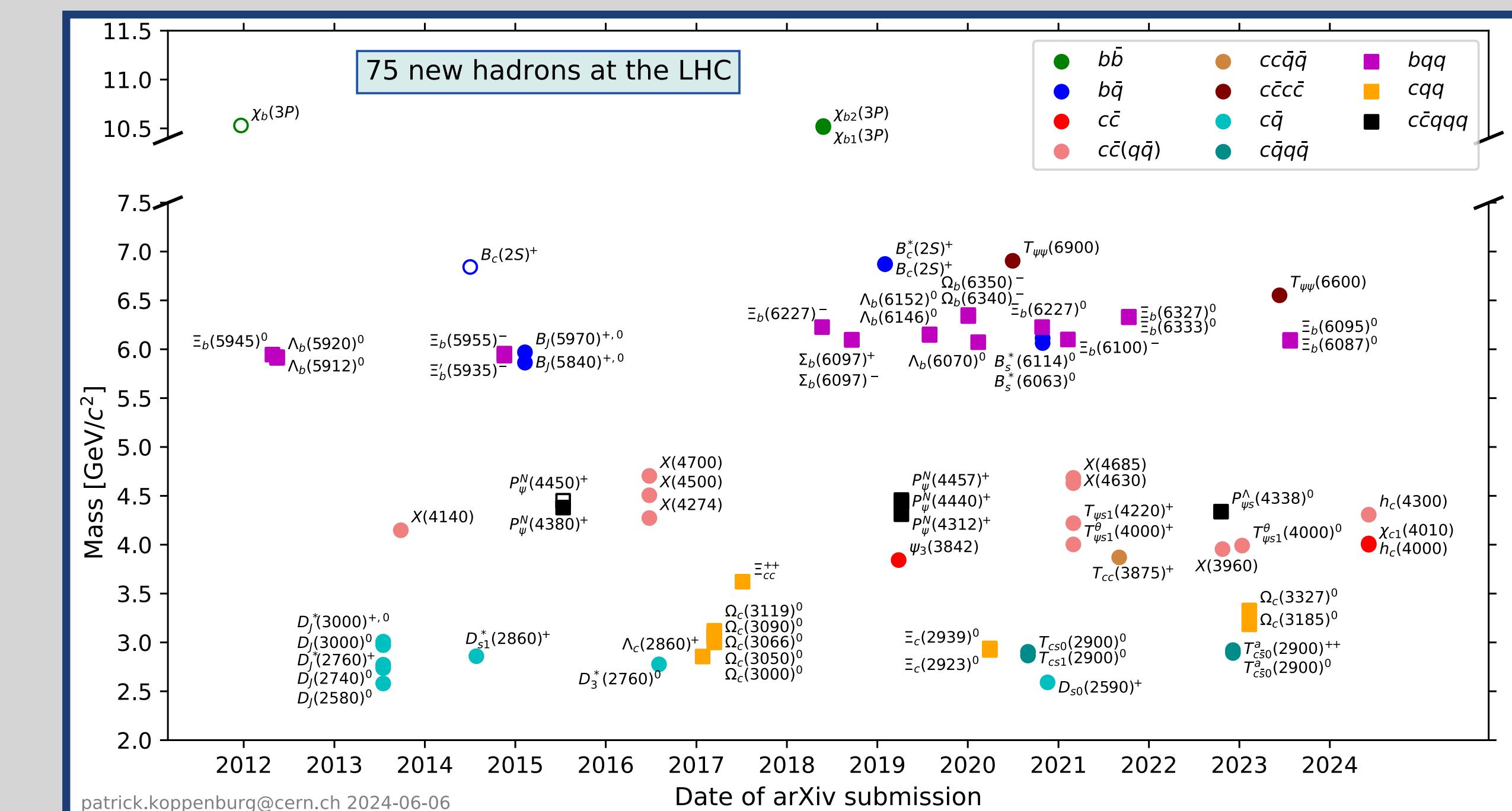
- Deuteron only stable dibaryon
- 1950's - Many predictions of various dibaryon states but failed experimental checks
- Experimental evidence of existence of d^*
- Recent renewal in interest due to discoveries of complex quark systems

Based on theory of strong interactions, we cannot rule out more dibaryons in nature.

Hexaquark - Any object with Baryon number 2

Composed of six valence quarks

Can be molecular or compact



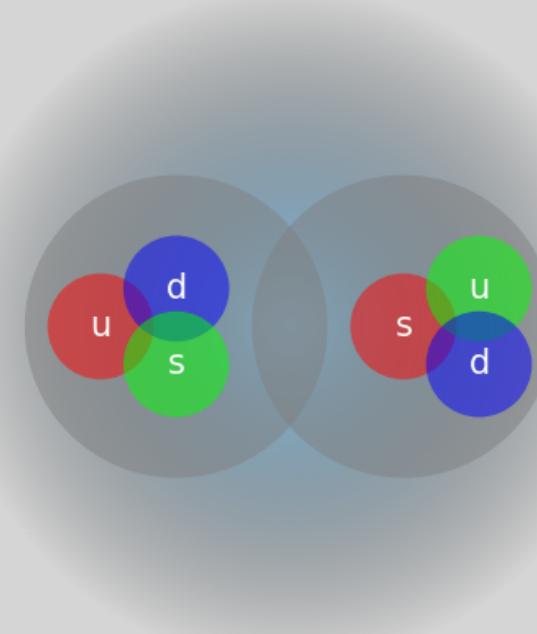
Review by Clement
PPNP (2016)



Dibaryons - Status

Hyperon-Nucleon experimental results indicates:

- ΛN attractive, though less than NN , no strange deuteron.
- ΣN interaction even weaker than ΛN .



Hyperon-Hyperon:

- $\Lambda\Lambda$ experiments does not rule out bound system.
- Jaffe prediction of dihyperons

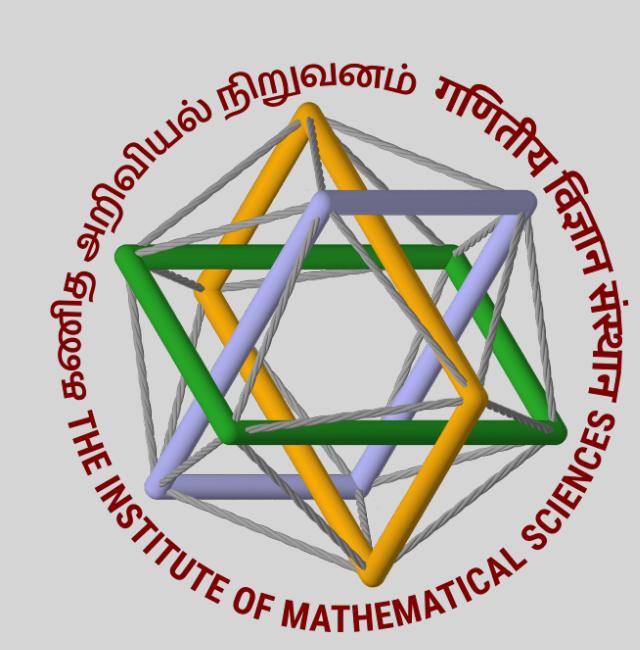
Jaffe
PRL 138 (1977) 195

*Dihyperon other calculations vary from very deep bound (even more than Jaffe's prediction) to unbound.

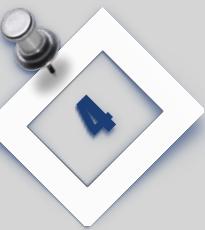
* Lattice QCD - gives bound result (8 MeV) - Large pion mass used.

Beane et al.
PRL 106(2011) 162001

Inoue et al.
PRL 106(2011) 162002



Dyson and Xuong
PRL 13 (1964) 815



Dibaryons - Status

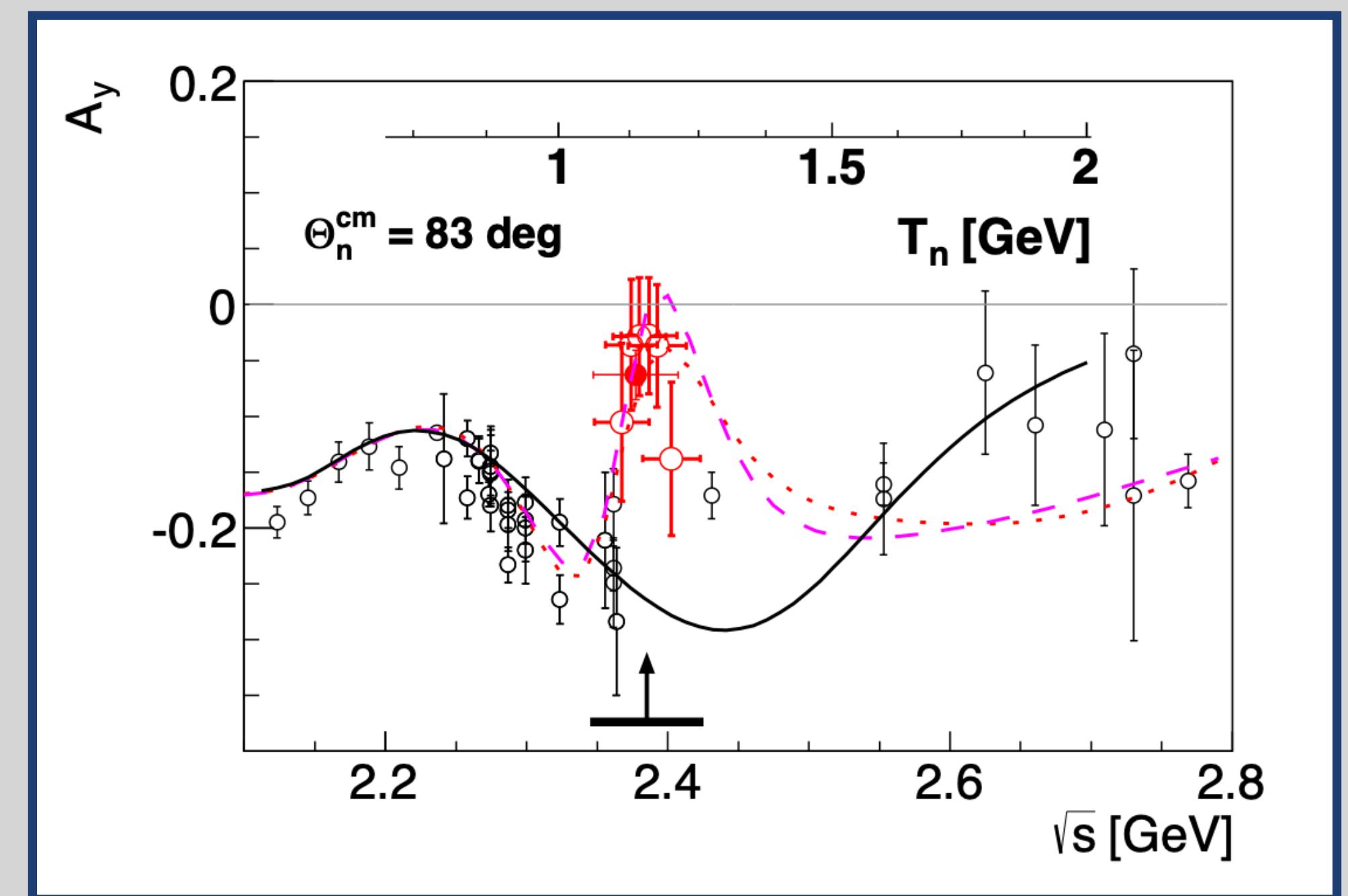
H dibaryon

- * NAGARA event - constraint on binding energy
- * Dedicated experiments for H dibaryon indicates existence unlikely but its existence not ruled out yet.

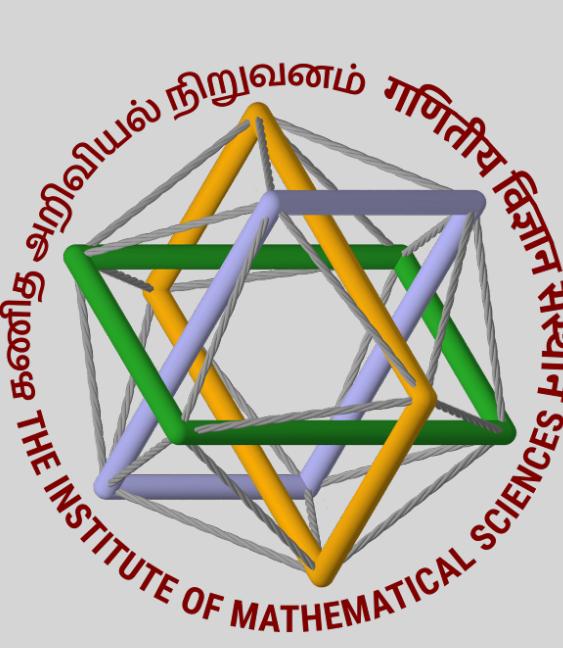
Next talk by Jeremy Green

1964 - Prediction of possible bound states

They predicted mass of D_{03} (close to d^* which was found later)



WASA @ COSY collaboration
SAID Data Analysis Center
PRL 112 (2014) 202301



Dibaryons - Lattice

Larger Pion Mass

Signal (m_N hadron)
 $\propto e^{-m_N t}$

Deuteron, Dineutron - Bound or Unbound ??

Tension between results

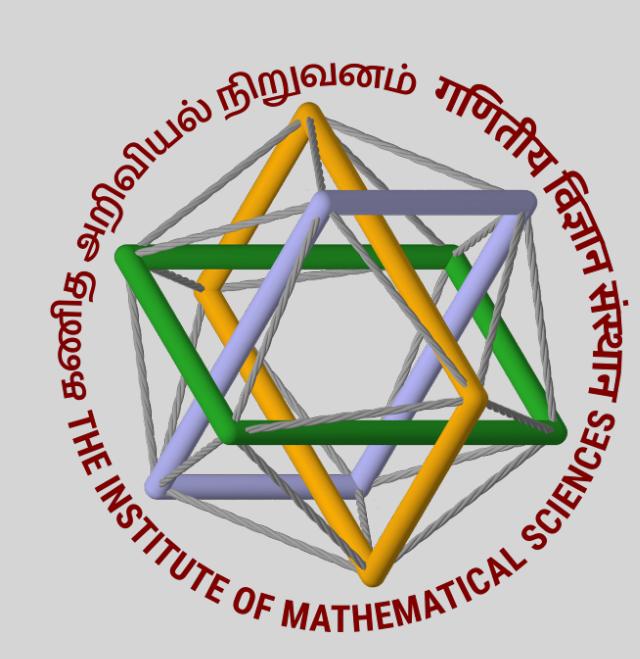
Noise
 $\propto e^{-\frac{3}{2}m_\pi t}$

Error in propagator correlation function
dominated by pions because of virtue of
lower energy states

Signal to Noise ratio exponentially degrades for $m_q \rightarrow 0$

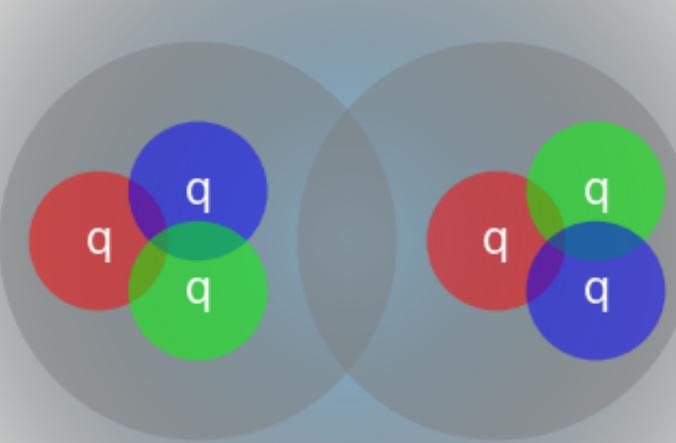
Lepage, TASI (1989)

- * Multiple ensembles for continuum limit.
- * Even more computing power in contractions for exotic hadrons.
- * $m_\pi L \geq 4$, to constraint finite volume effects.
- * Finer lattices for lesser discretisation errors.

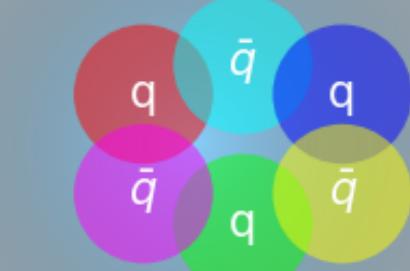


Hexaquark - Dibaryon

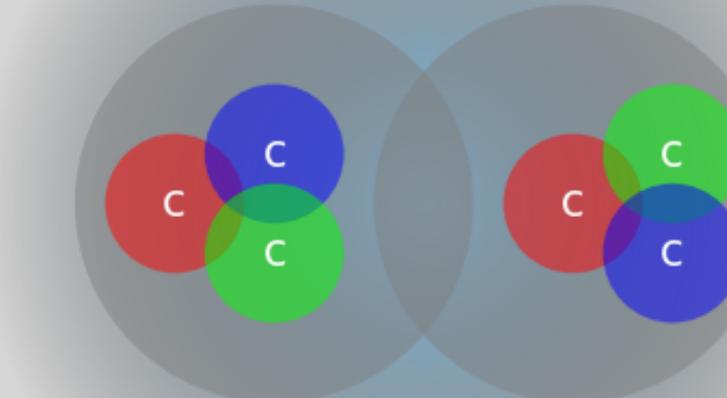
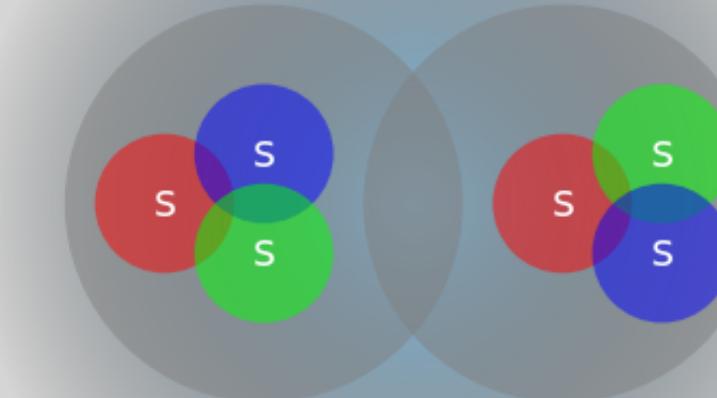
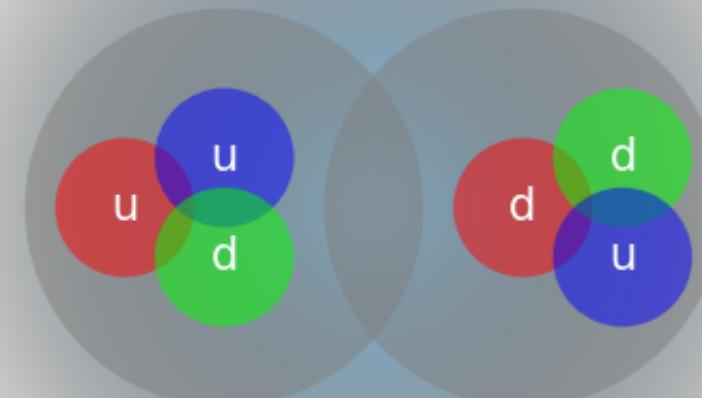
Extensive studies of deuteron like heavy dibaryons using
Lattice QCD



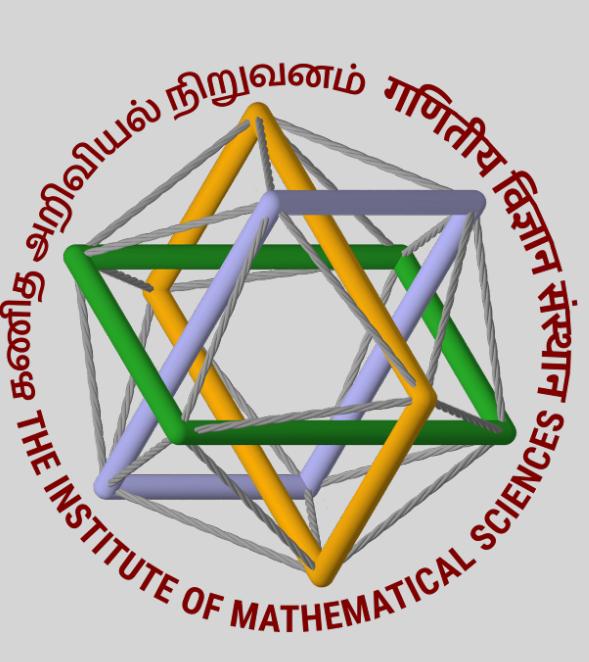
Hadron with 6 quarks



Deuteron



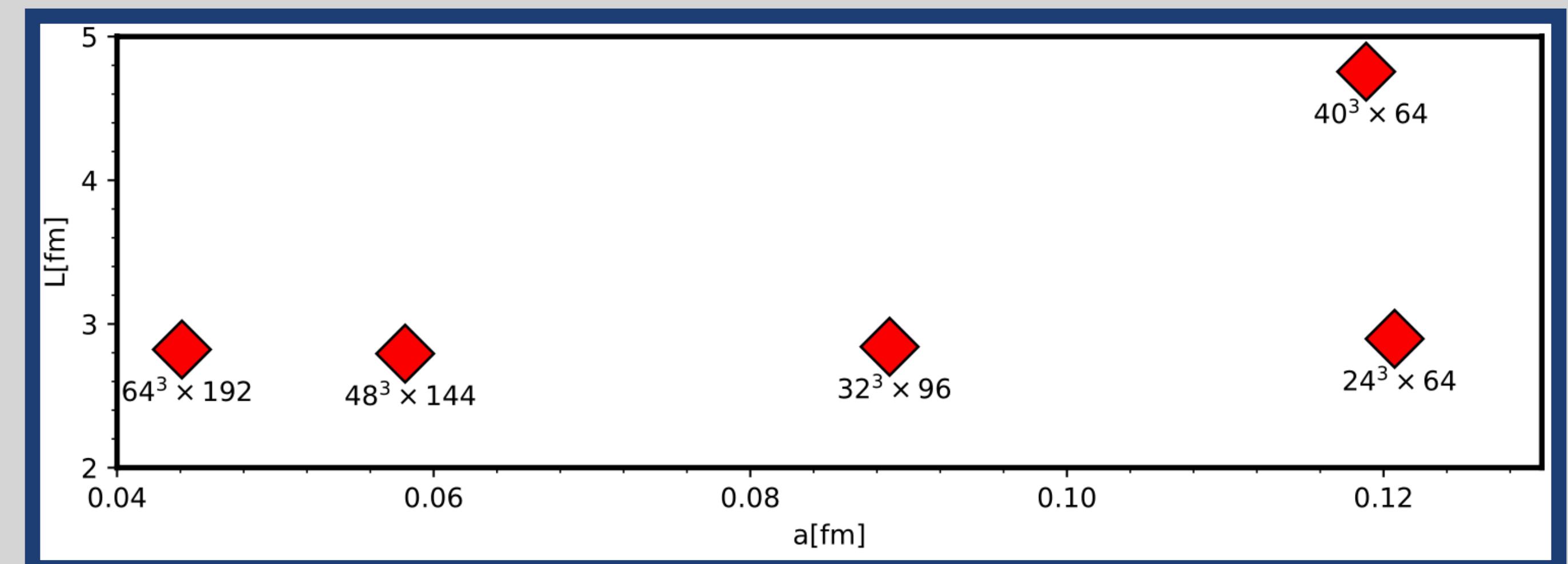
We work with single flavored dibaryons composed of strange
and charm quarks named as \mathcal{D}_{6s} and \mathcal{D}_{6c} respectively



Lattice Setup

Bazavov et al., PRD 87 (2013) 5, 054505

- Overlap action on background of Highly Improved Staggered Quark (HISQ) gauge configurations.
 - Strange and charm masses set at physical values.
 - Up and down set as degenerate masses heavier than physical values.
 - Finest lattice used $a \approx 0.044 \text{ fm}$ with Volume as $64^3 \times 192$



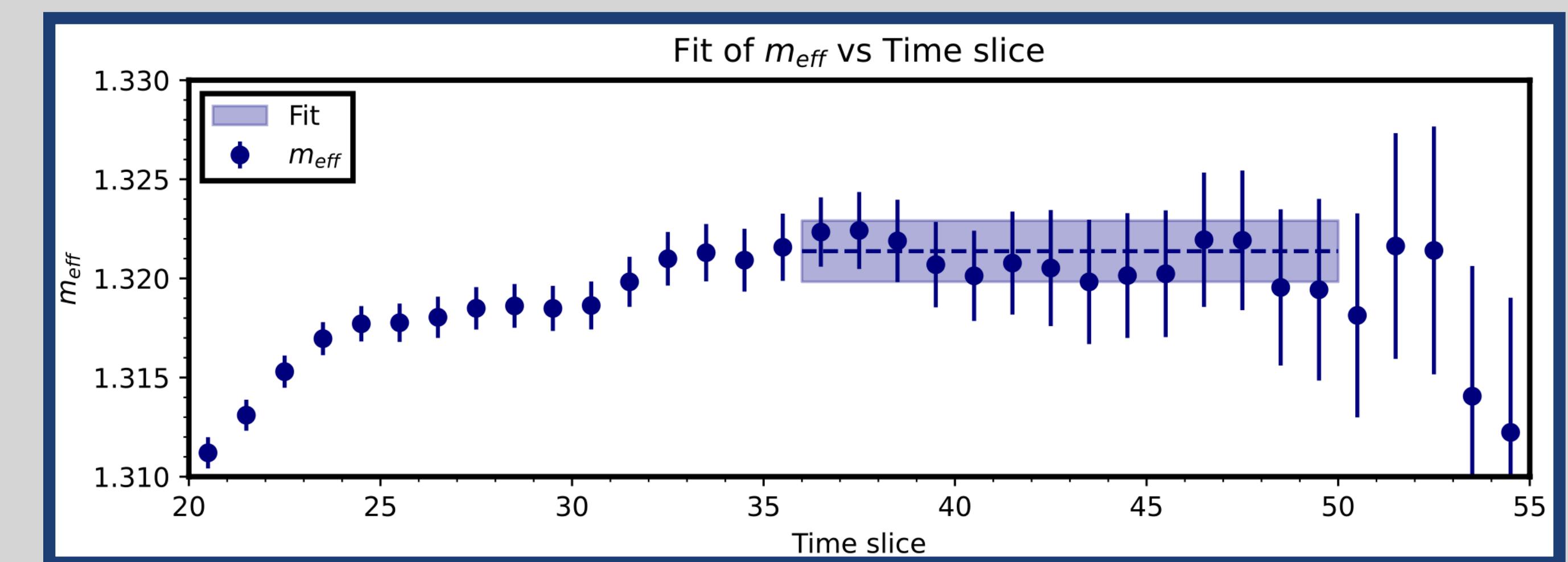
Masses from Lattice

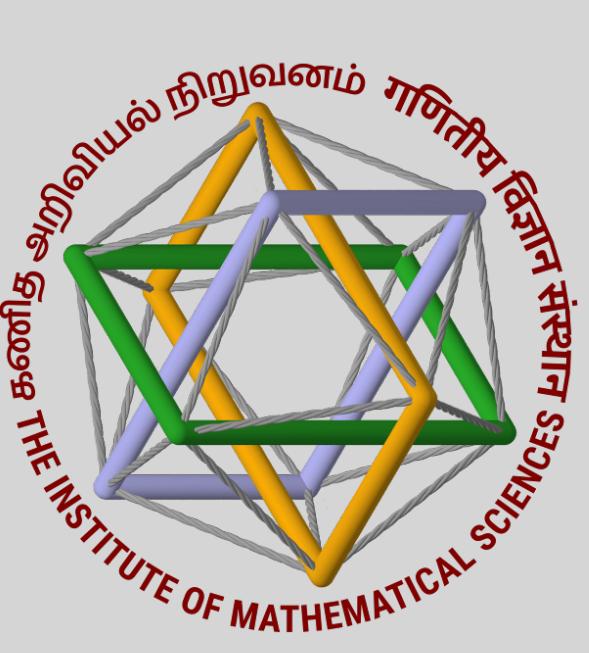
- * Euclidean two point correlator as: $C_{ji}(t_f - t_i) = \langle 0 | O_j(t_f) \bar{O}_i(t_i) | 0 \rangle = \sum_n \frac{Z_i^{n*} Z_j^n}{2m_n} e^{-m_n(t_f - t_i)}$
- * $O_j(t_f)$ and $\bar{O}_i(t_i)$ are the desired interpolating operators and $Z_j^n = \langle 0 | O_j | n \rangle$

- * Effective mass = $\log \left[\frac{C(t)}{C(t+1)} \right]$

Wall source to point sink.

Cross checked results with boxesink.





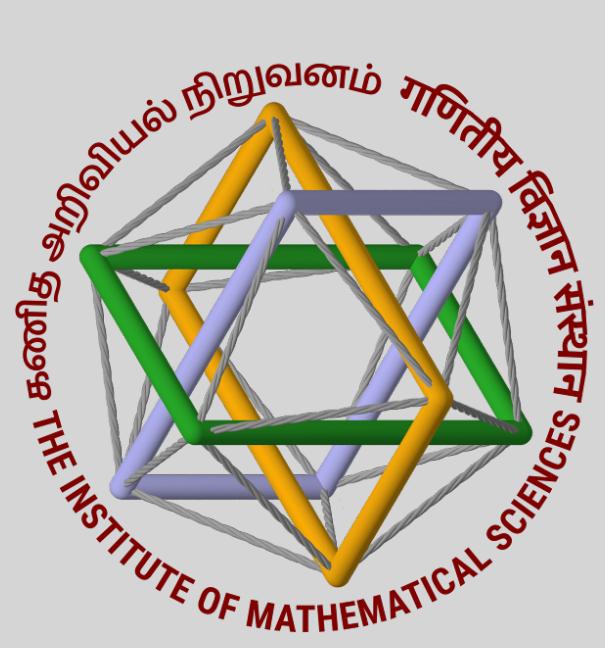
Dibaryon Operators

$$\mathcal{O} = \epsilon_{abc} q_{\mu_1}^a q_{\mu_2}^b q_{\mu_3}^c$$

$$\mathcal{O}_d = \mathcal{O}_1 \cdot CG \cdot \mathcal{O}_2$$

- Total wave function anti-symmetric under exchange of baryons.
- Single flavor baryons (symmetric).
- Assume only s wave interactions (symmetric) in dibaryon system.
- Color singlet baryons (symmetric)
- Hence Spin must be anti-symmetric which is in case of even spin (Spin 0 and 2)

- Use reduction coefficients to project continuum based operators to suitable octahedral group.
- $S = 0$ continuum spin subduces to one dimensional A_1^+ irrep.
- $S = 2$ continuum spin subduces to two dimensional E^+ and three dimensional T_2^+ irrep.



Dibaryon Operators

a,b - different embeddings

$$\mathcal{O}_{d,A_1,1}^{[0]} = \frac{1}{2} \left({}^a H_{3/2} {}^b H_{-3/2} - {}^a H_{1/2} {}^b H_{-1/2} + {}^a H_{-1/2} {}^b H_{1/2} - {}^a H_{-3/2} {}^b H_{3/2} \right)$$

For Spin 0, dibaryon operator corresponding to one dimensional A_1^+ irrep.

For Spin 2, 5 such operators corresponding to E^+ and T_2^+ irrep.

S_z	Operator	State
3/2	${}^1 H_{3/2}$	111
1/2	${}^1 H_{1/2}$	112+121+211
-1/2	${}^1 H_{-1/2}$	122+212+221
-3/2	${}^1 H_{-3/2}$	222

Non Relativistic Embedding

S_z	Operator	State
3/2	${}^2 H_{3/2}$	133+313+331
1/2	${}^2 H_{1/2}$	233+323+332+134+341+413+143+431+314
-1/2	${}^2 H_{-1/2}$	144+414+441+234+342+423+243+432+324
-3/2	${}^2 H_{-3/2}$	244+424+442

Relativistic Embedding

Operator Contraction

Two baryons at source and two at sink

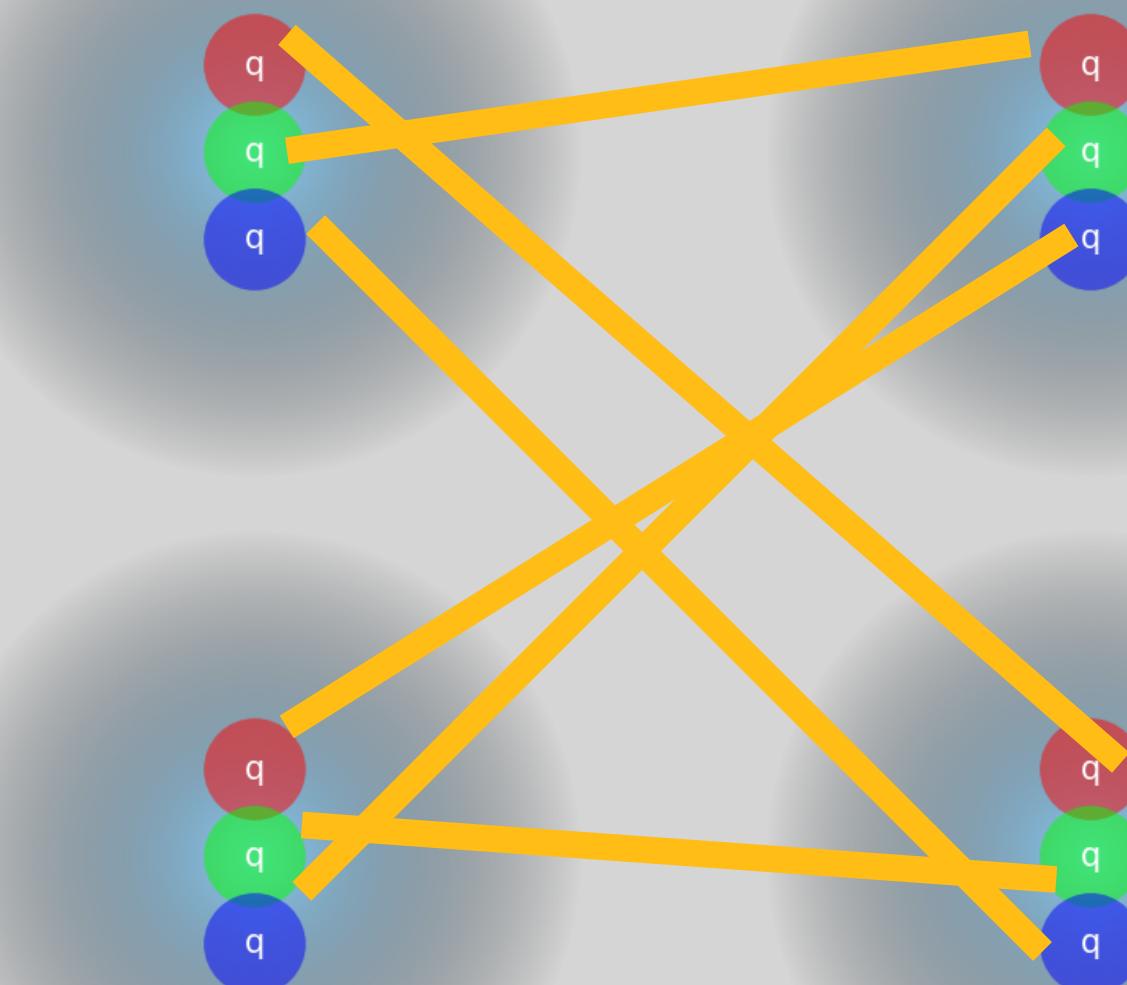


$6! = 720$ contraction possibilities

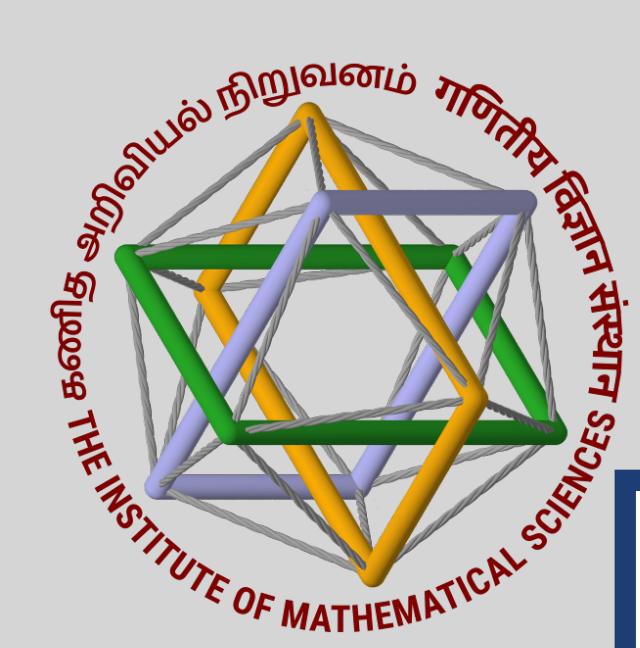


720 contractions can happen in 16 different ways depending on different embeddings

N-N-N-N	N-N-N-R	N-N-R-N	N-N-R-R
N-R-N-N	N-R-N-R	N-R-R-N	N-R-R-R
R-N-N-N	R-N-N-R	R-N-R-N	R-N-R-R
R-R-N-N	R-R-N-R	R-R-R-N	R-R-R-R



Some of these are degenerate and some of these do not contribute at all for different Spin cases



Dibaryons results from Lattice

$\mathcal{D}_{6b}, S = 0$

Mathur, Padmanath, Chakraborty
PRL, 130, 111901 (2023)

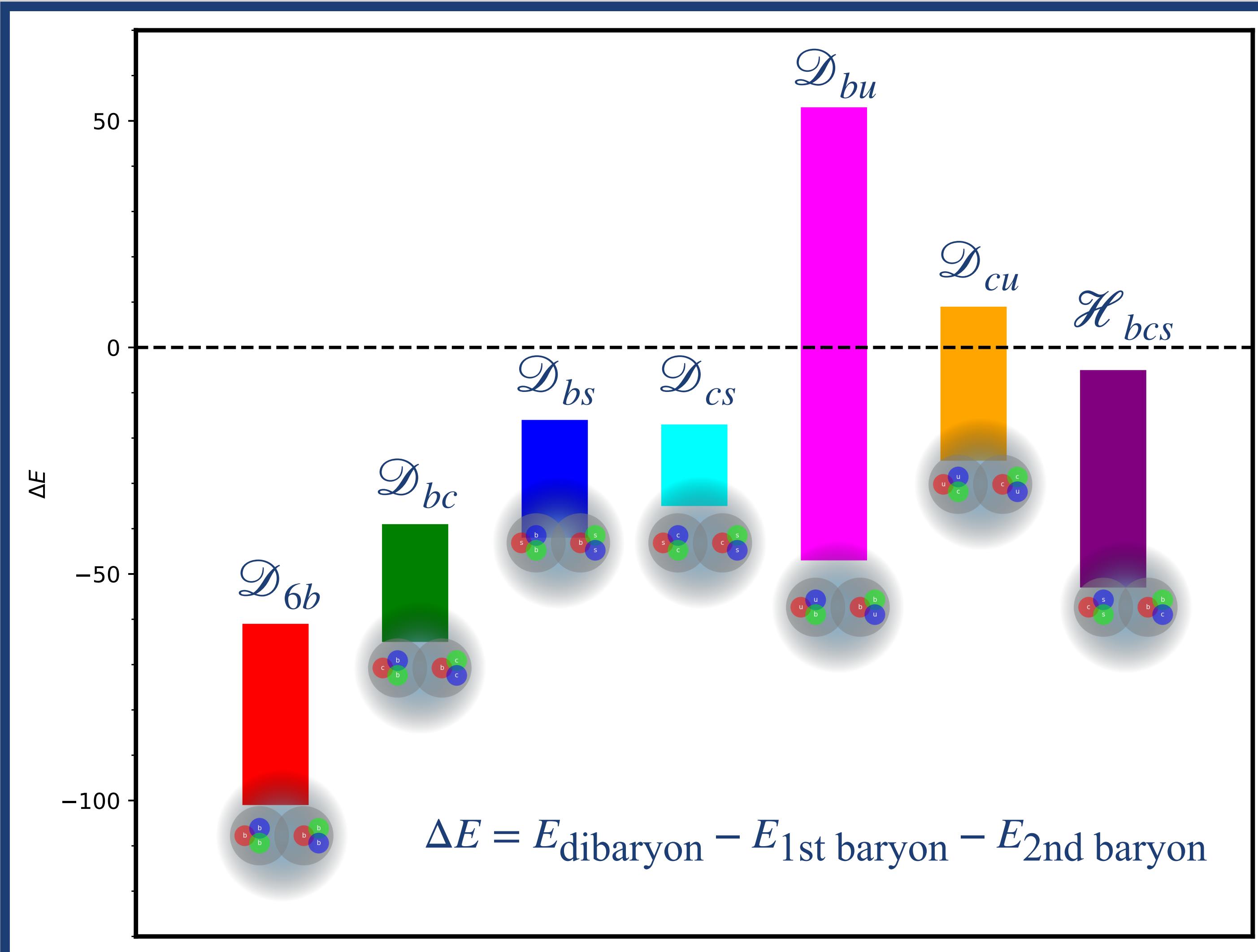
$\mathcal{D}_{bc}, \mathcal{D}_{bs}, \mathcal{D}_{cs}, \mathcal{D}_{bu}, \mathcal{D}_{cu}$
 $S = 1$

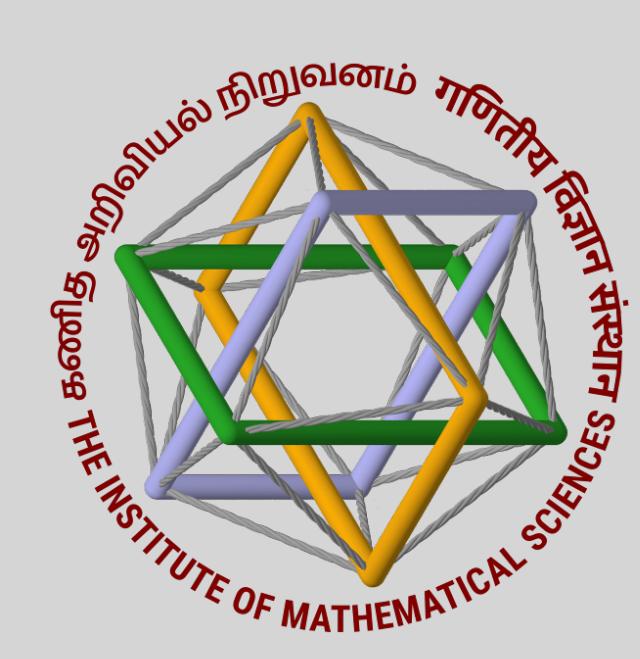
Junnarkar, Mathur PRL, 123, 162003 (2019)

$\mathcal{H}_{bcs}, S = 0$

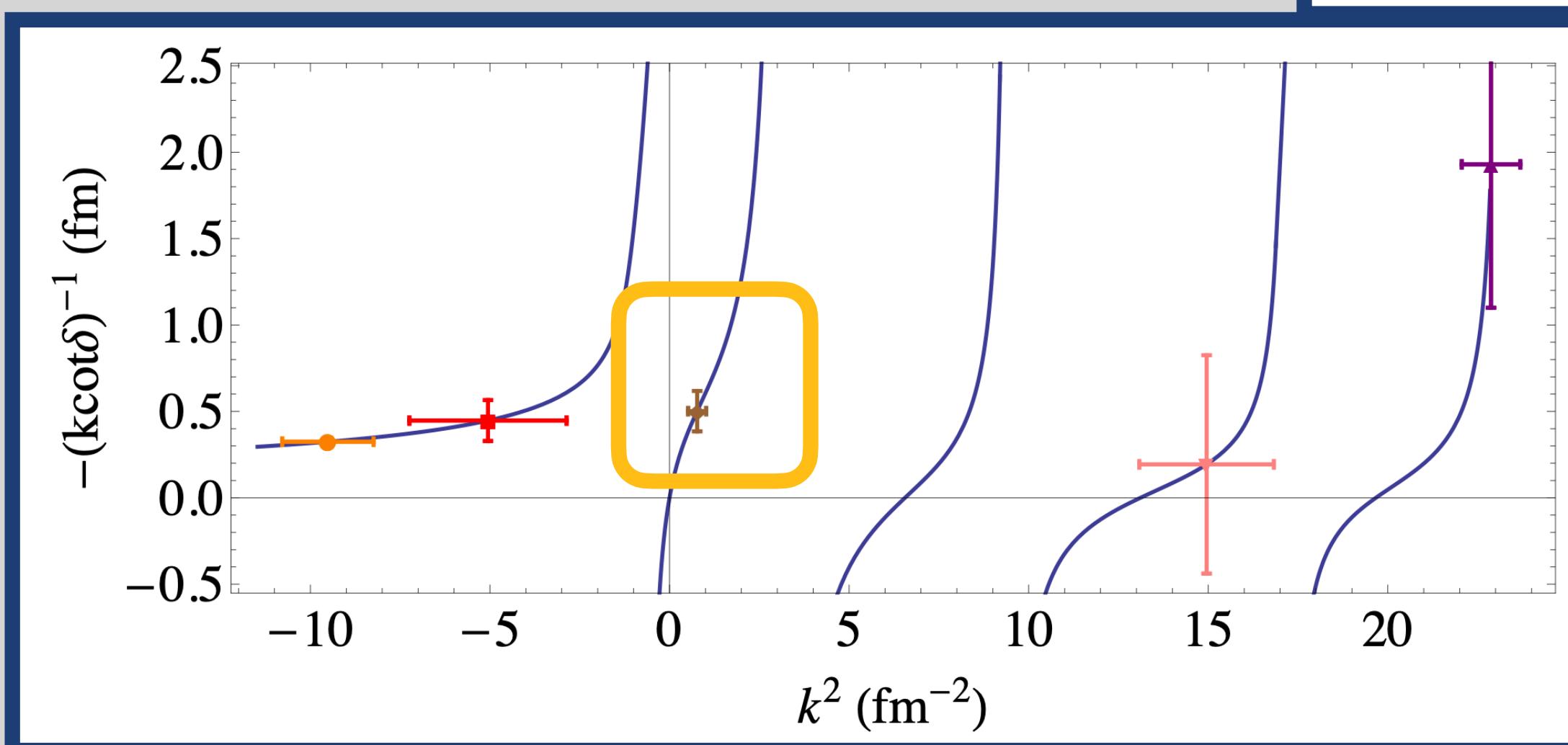
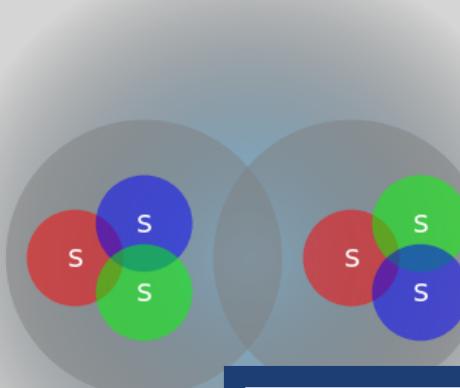
Junnarkar, Mathur PRD, 106, 054511 (2022)

Where does $\mathcal{D}_{6c}, \mathcal{D}_{6s}$ stand ??

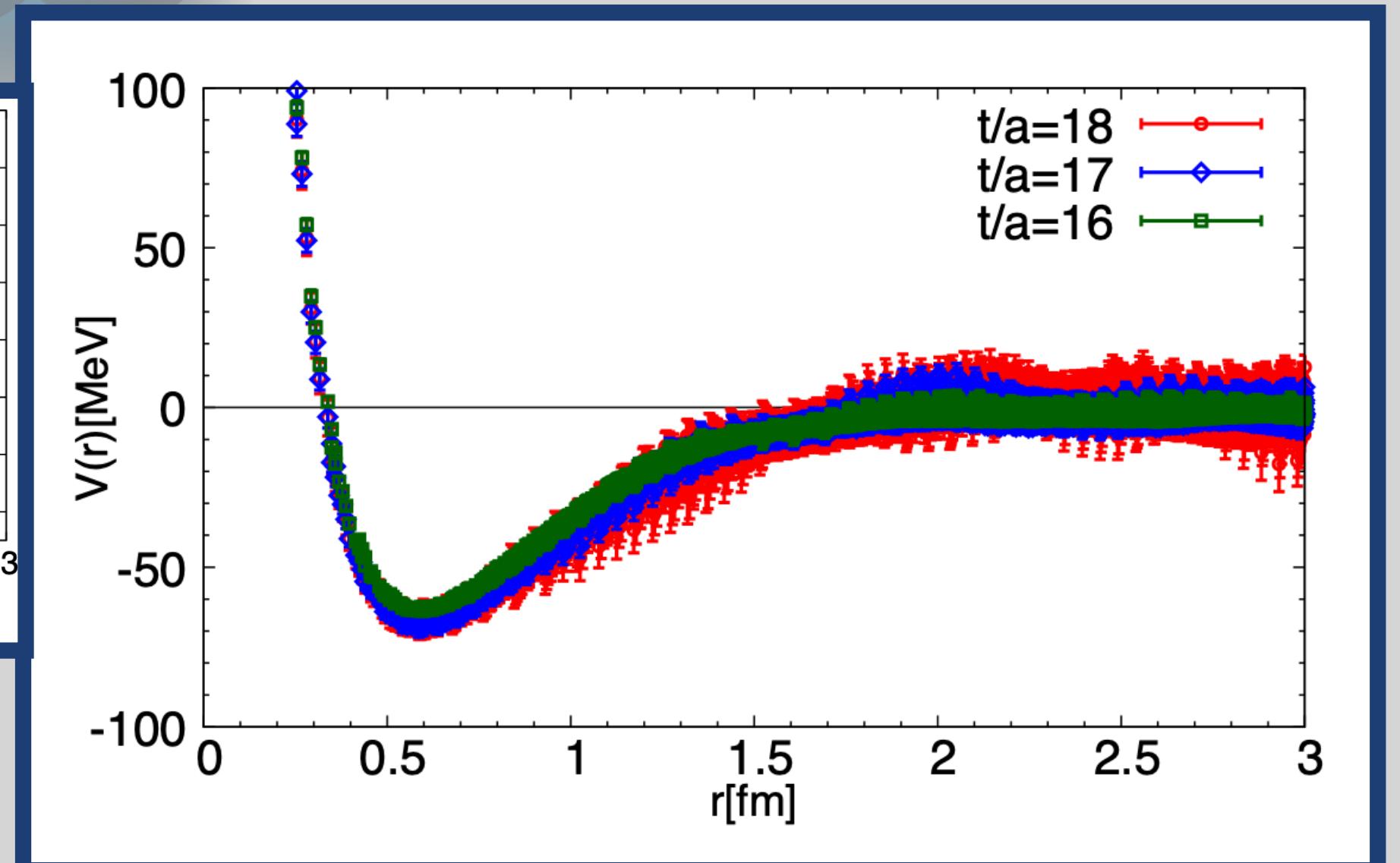
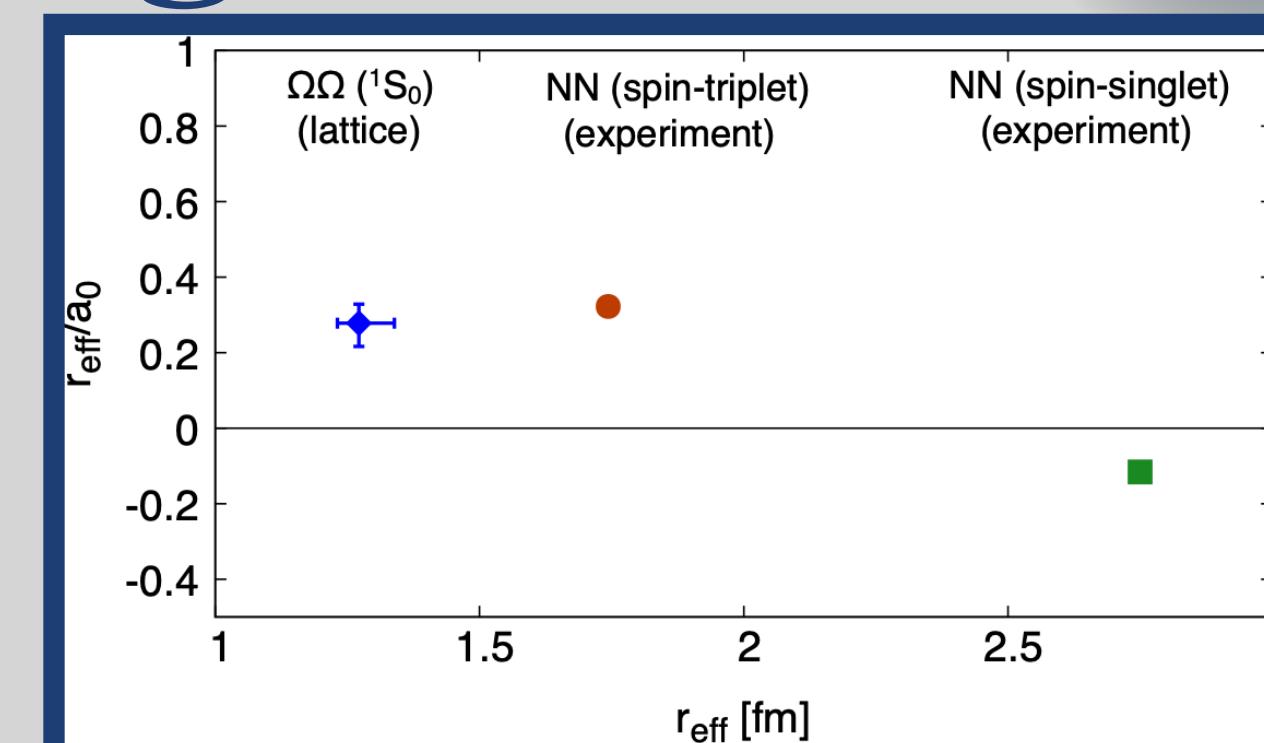




D_{6s} Existing Results



- Weakly repulsive in Spin 0 H^+H^+ irrep, No bound state
- Attractive in Spin 1,2 $G_1^+H^+$ but only single volume used.

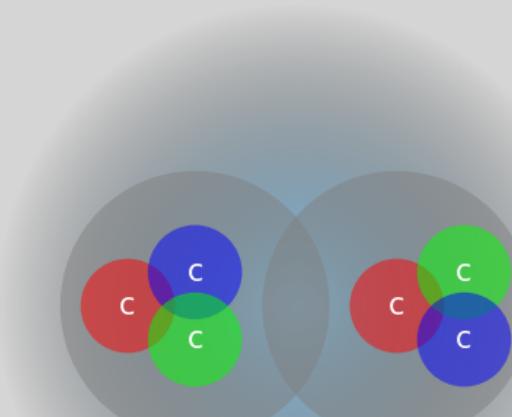


- Weakly attractive in Spin 0, hence bound state

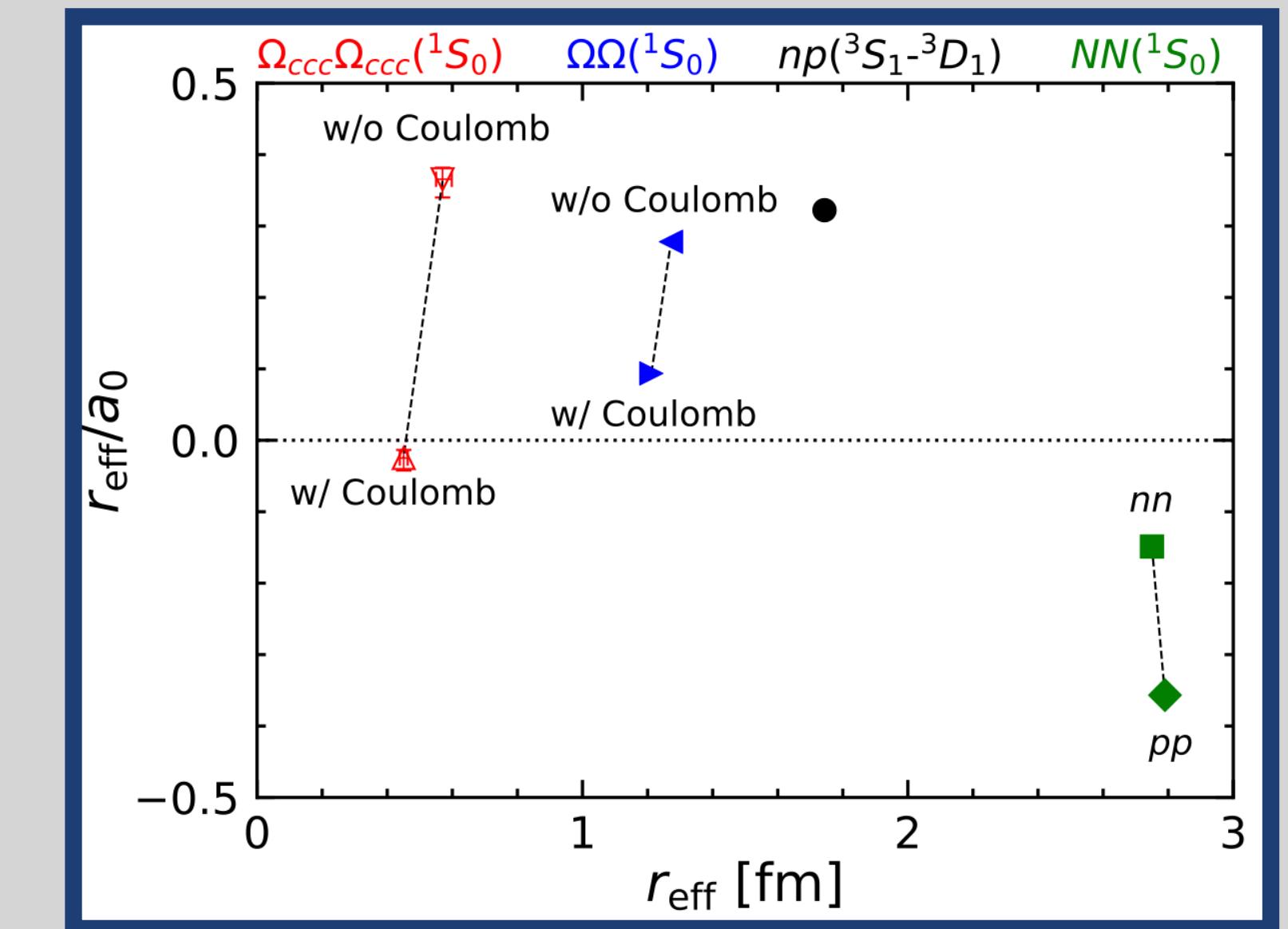
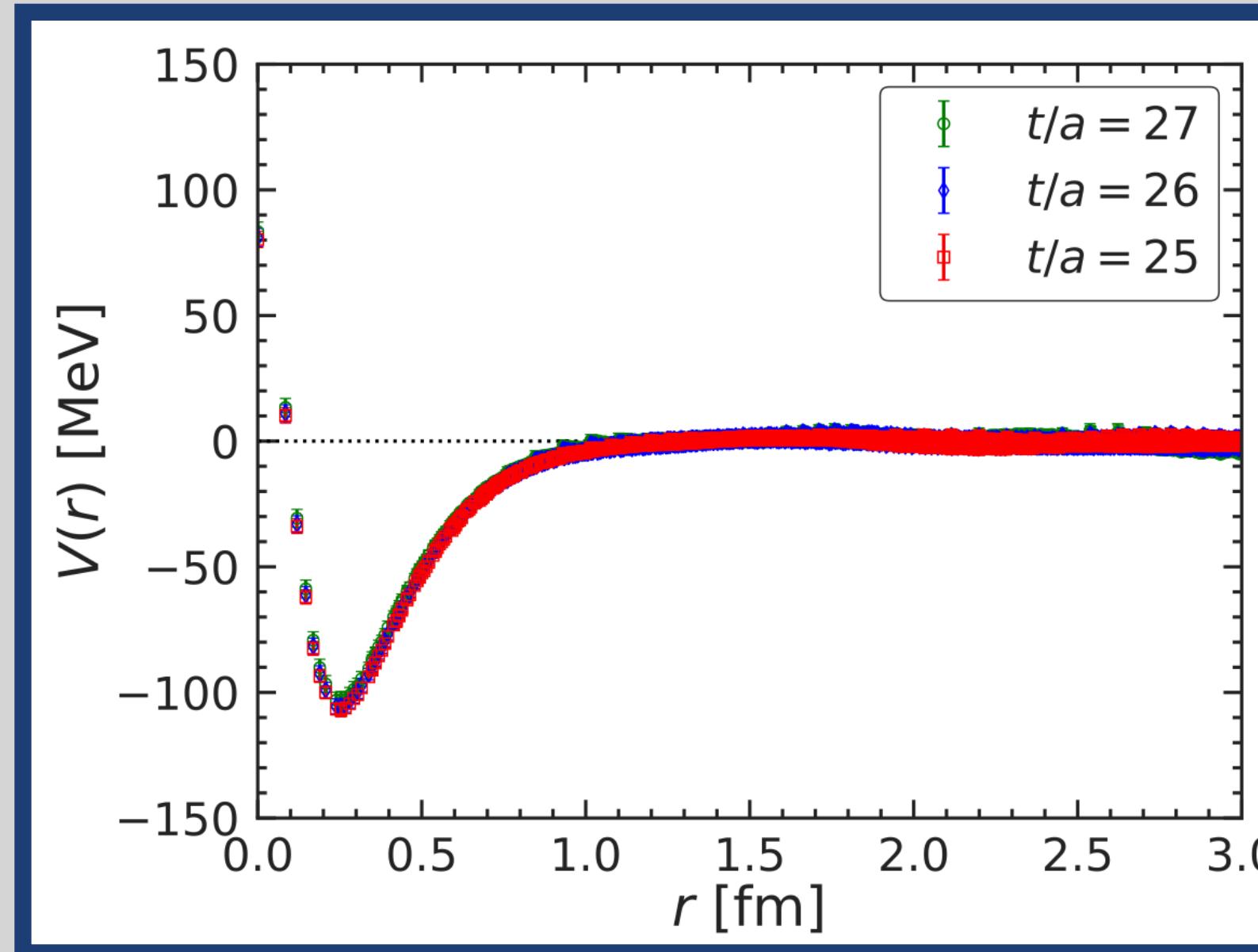
HALQCD
PRL 120, 212001 (2018)

Buchoff, Luu, Wasem
PRD 85, 094511 (2012)

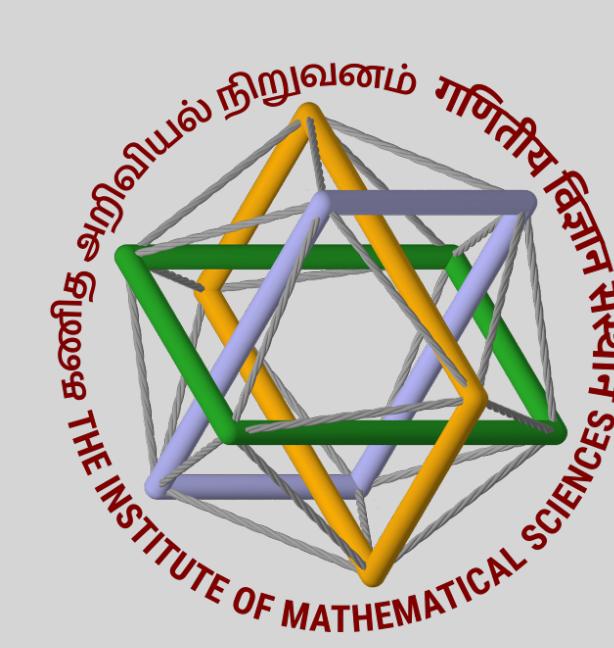
\mathcal{D}_{6c} Existing Results



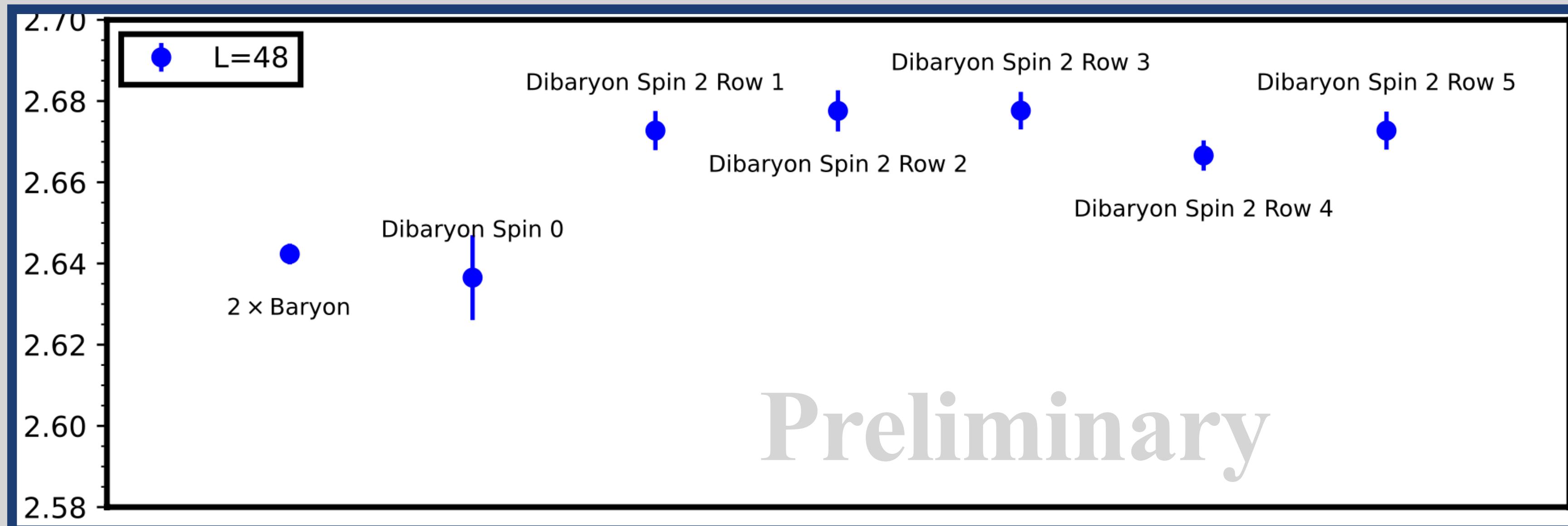
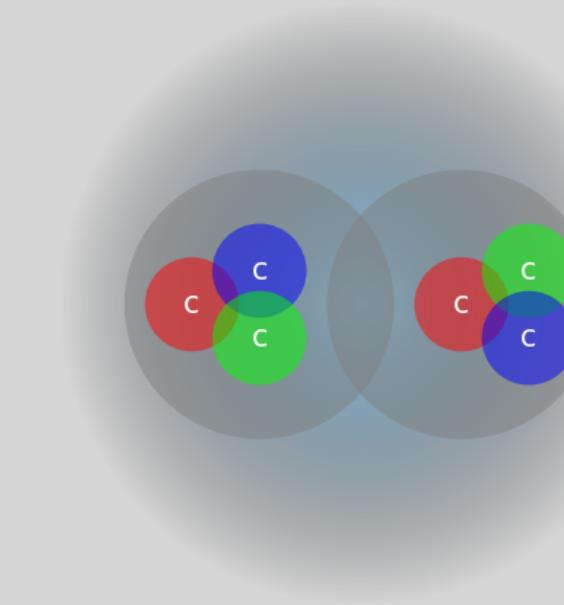
HALQCD
PRL 127, 072003 (2021)



- Simulation with physical charm mass and near physical light quark mass.
- Dibaryon existence without Coulomb interaction.
- Near unitary region with Coulomb interaction (scattering length less than corresponding strange dibaryon calculation).

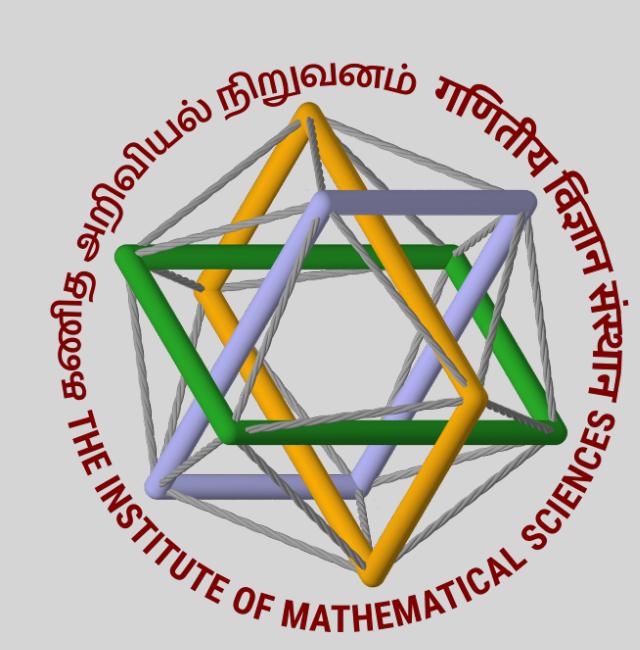


\mathcal{D}_{6c} Results

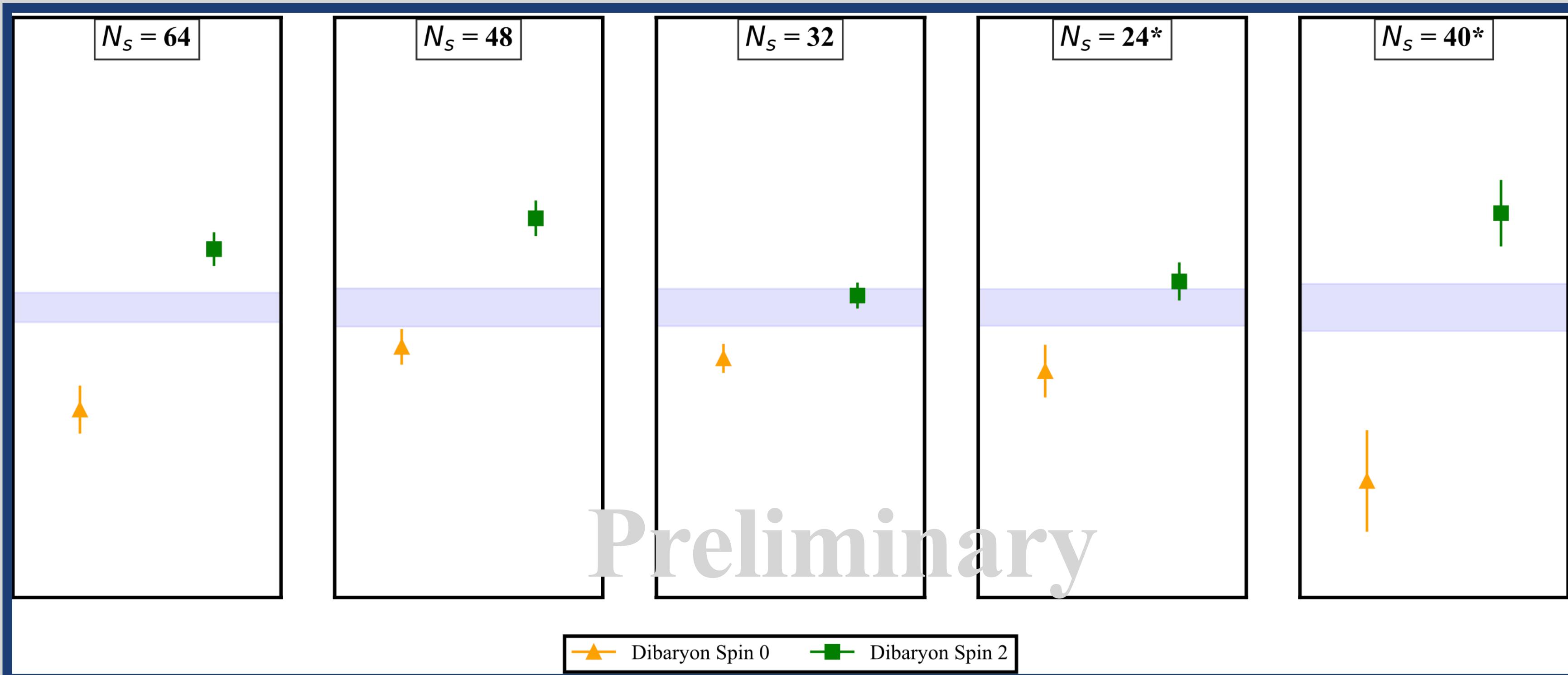
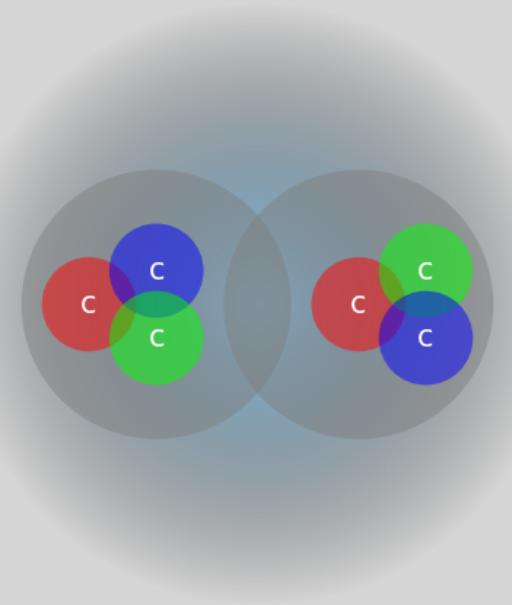


- Bound state, if exist, is shallow
- More probe using amplitude analysis

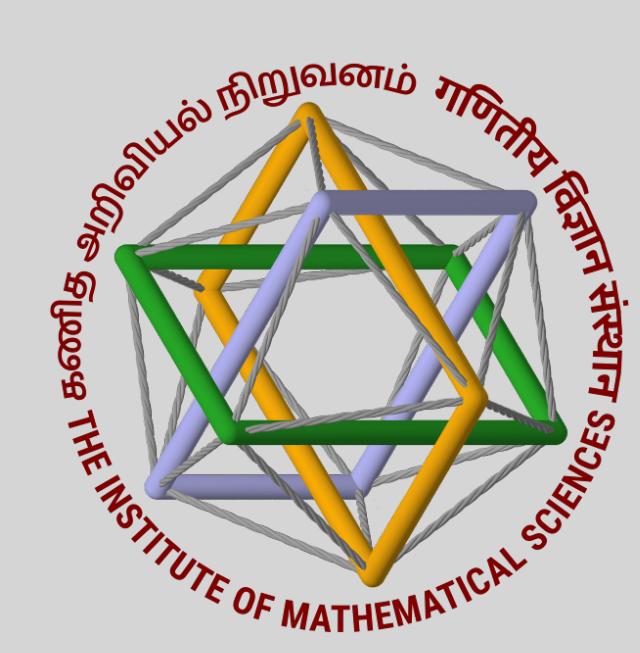
- Two lattice volumes, 4 lattice spacings, this plot with $L = 48$
- Spin 2 - repulsive interactions, Spin 0 dibaryon energy same as twice of baryon (within error)



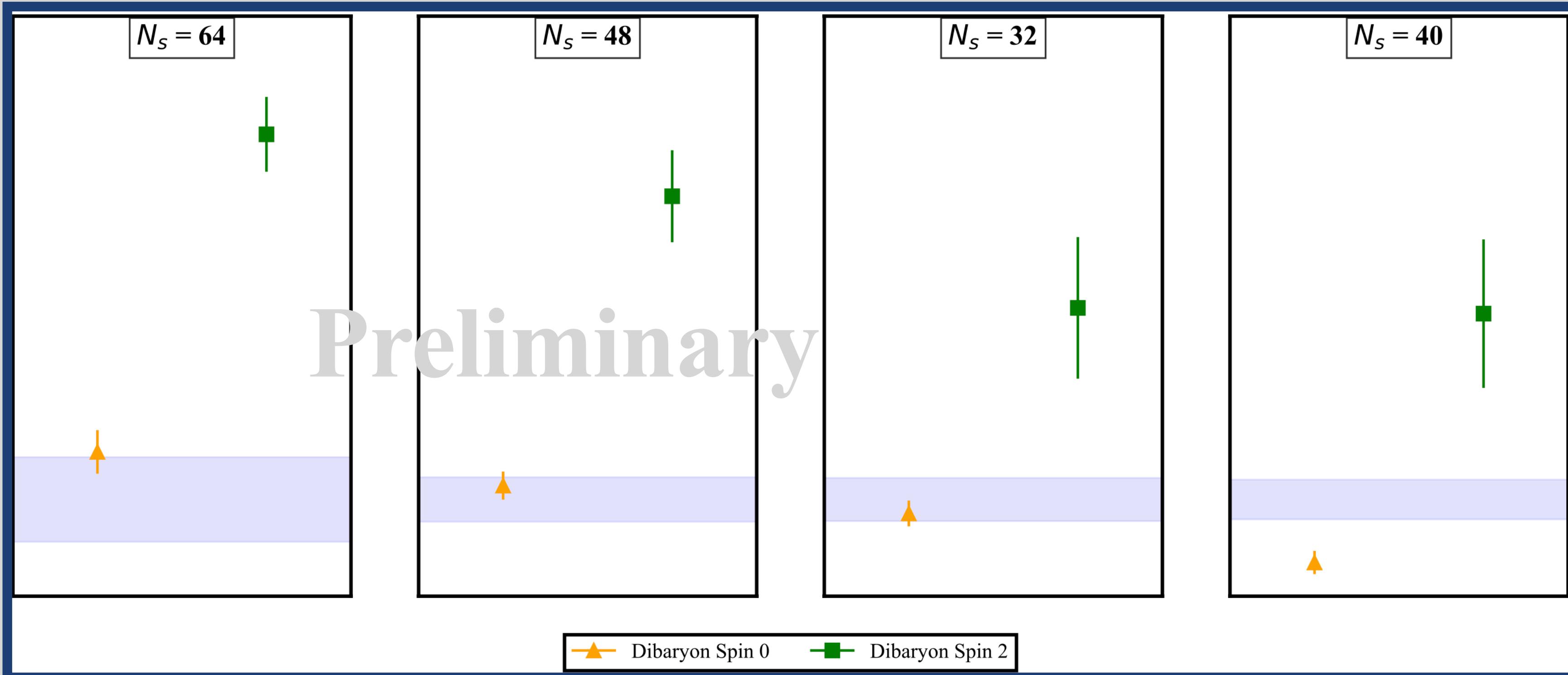
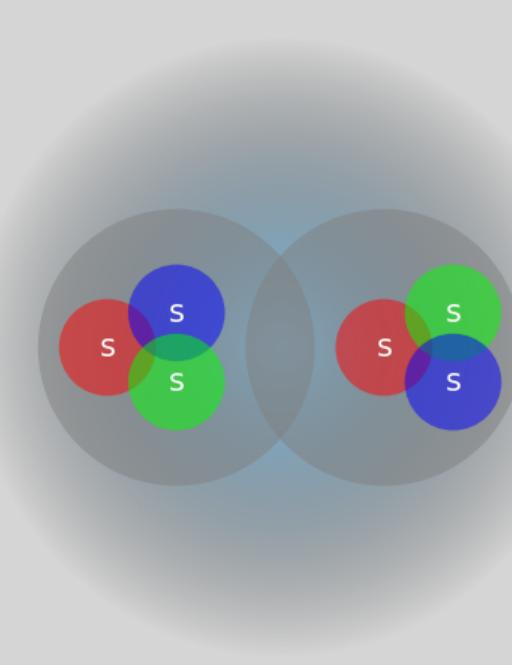
\mathcal{D}_{6c} Results



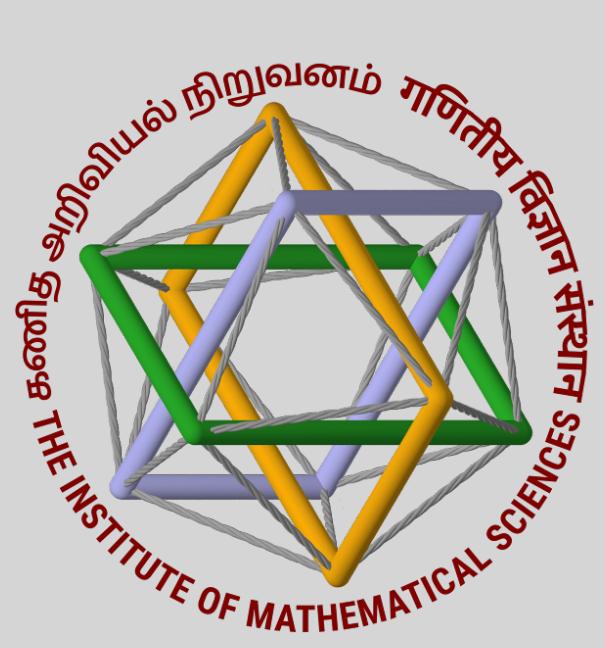
- Bound state, if exist, is shallow
- More probe using amplitude analysis



\mathcal{D}_{6s} Results



- No bound state
- More probe using amplitude analysis



Summary

- ★ We observe a positive shift in the S=2 channel, indicating a repulsive interaction and inability to host any bound state for both strange and charm systems.
- ★ In the charm sector, for spin zero, there is a slight tendency towards negative shifts, although these shifts have smaller magnitudes.
- ★ In the strange sector, for spin zero, the results generally suggest a non-interacting scenario, with weak interactions and potentially no bound states.
- ★ A more precise conclusion can only be drawn with larger statistics and a comprehensive finite-volume amplitude study.
- ★ Lattice estimation of $d^*(2380)\dots$

THANK YOU

