

# RC 7: Steady Magnetic Fields

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## Magnetic Field Intensity, Relative Permeability and Magnetic circuit

Considering the effect of both the internal dipole moment and the induced magnetic moment in a magnetic material, the magnetic field intensity  $H$  is defined as:

$$H = \frac{B}{\mu_0} - M \quad (1)$$

$$\nabla \times H = J \quad (2)$$

directly relates the magnetic field intensity with the density of free charge. The integral form of which is then,

$$\oint_C H \cdot dl = I \quad (3)$$

It is another form of Ampere's circuital law: the circulation of the magnetic field intensity around any closed path is equal to the free current flowing through the surface bounded by the path.

If the closed path  $C$  is chosen to enclose  $N$  turns of a winding carrying a current  $I$  that excites a magnetic circuit, we have

$$\oint_C H \cdot dl = NI = V_m \quad (4)$$

$V_m$  is analogous to electromotive force (emf) and is called magnetomotive force (mmf).

If we define:

$$M = \chi_m H \quad (5)$$

where  $\chi_m$  is magnetic susceptibility, Then

$$B = \mu_0(1 + \chi_m)H = \mu_0\mu_r H = \mu H \quad (6)$$

$$H = \frac{1}{\mu} B \quad (7)$$

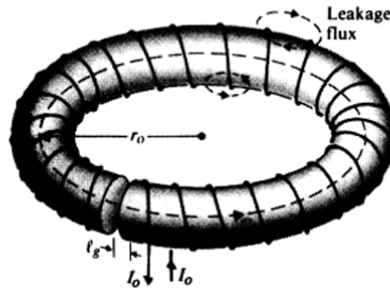
$$\mu_r = 1 + \chi_m = \frac{\mu}{\mu_0} \quad (8)$$

where  $\mu_r$  is the relative permeability of the medium.  $\mu = \mu_r \mu_0$  is the absolute permeability/permeability.

## Exercise

Assume that  $N$  turns of wire are wound around a toroidal core of a ferromagnetic material with permeability  $\mu$ . The core has a mean radius  $r_o$ , a circular cross section of radius  $a$  ( $a \ll r_o$ ), and a narrow air gap of length  $l_g$ , as shown in the following figure. A steady current  $I_o$  flows in the wire. Determine:

- (a) the magnetic flux density,  $B_f$ , in the ferromagnetic core;
- (b) the magnetic field intensity,  $H_f$ , in the core;
- (c) the magnetic field intensity,  $H_g$ , in the air gap.



## Behavior of Magnetic Materials

1. **Diamagnetic**, if  $\mu_r \lesssim 1$  ( $\chi_m$  is a very small negative number)
2. **Paramagnetic**, if  $\mu_r \gtrsim 1$  ( $\chi_m$  is a very small positive number)
3. **Ferromagnetic**, if  $\mu_r \gg 1$  ( $\chi_m$  is a large positive number)

## Boundary Conditions for Magnetostatic Fields

1. The normal component of  $B$  is continuous across an interface,

$$B_{1n} = B_{2n} \quad (9)$$

which is

$$\mu_1 H_{1n} = \mu_2 H_{2n} \quad (10)$$

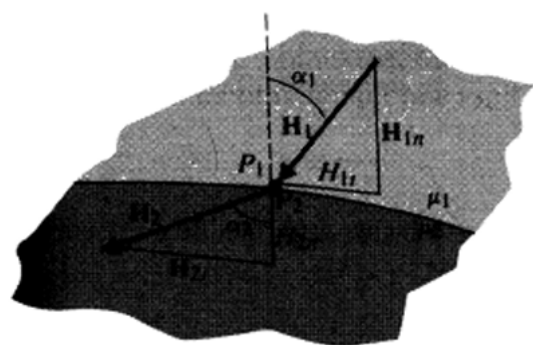
2. The tangential component of the  $H$  field is discontinuous across an interface where a free surface current exists.

$$\mathbf{a}_n \times (H_1 - H_2) = J_s \quad (11)$$

When the conductivities of both media are finite, currents are defined by volume current densities and free surface currents do not exist on the interface. Hence  $J_s = 0$ , the tangential component of  $H$  is continuous across the boundary of almost all physical media; it is discontinuous only when an interface with an ideal perfect conductor or a superconductor is assumed.

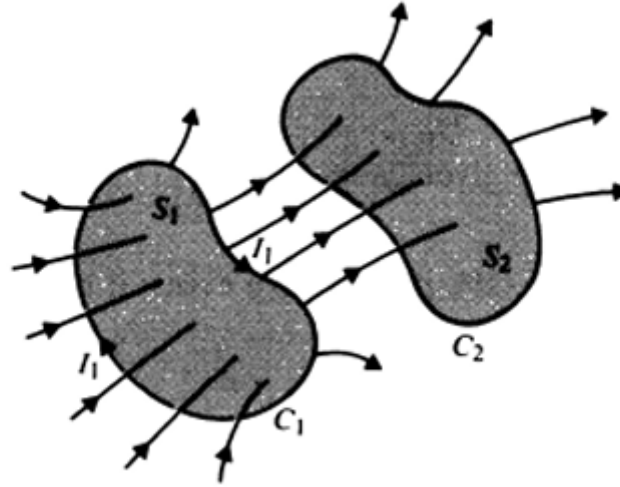
## Exercise

Two magnetic media with permeabilities  $\mu_1$  and  $\mu_2$  have a common boundary, as shown in the following figure. The magnetic field intensity in medium 1 at the point  $P_1$  has a magnitude  $H_1$  and makes an angle  $\alpha_1$  with the normal. Determine the magnitude and the direction of the magnetic field intensity at point  $P_2$  in medium 2.



## Inductances and Inductors

For two loops  $C_1, C_2$ :



For linear media where permeability does not change with  $I_1$ ,

$$L_{12} = \frac{\Lambda_{12}}{I_1} \quad (12)$$

where  $\Lambda_{12}$  is the flux linkage,  $L_{12}$  is the mutual inductance between loops  $C_1$  and  $C_2$ . A more general definition for  $L_{12}$  is then:

$$L_{12} = \frac{d\Lambda_{12}}{dI_1} \quad (13)$$

The total flux linkage with  $C_1$  caused by  $I_1$  is calculated as:

$$\Lambda_{11} = N_1 \Phi_{11} > N_1 \Phi_{12} \quad (14)$$

The self-inductance of loop  $C_1$  could be calculated as:

$$L_{11} = \frac{\Lambda_{11}}{I_1} \quad (15)$$

For linear medium, in general:

$$L_{11} = \frac{d\Lambda_{11}}{dI_1} \quad (16)$$

A conductor arranged in an appropriate shape (e.g., a conducting wire wound as a coil) to supply a certain amount of self-inductance is called an inductor. It can store magnetic energy.

Method to determine the self-inductance of an inductor:

1. Choose an appropriate coordinate system for the given geometry.
2. Assume a current  $I$  in the conducting wire.
3. Find  $B$  from  $I$  by Ampere's circuital law,

$$\oint_C B \cdot dl = \mu_0 I \quad (17)$$

if symmetry exists; if not, Biot-Savart law

$$B = \frac{\mu_0 I}{4\pi} \oint_{C'} \frac{dl' \times \mathbf{a}_R}{R^2} \quad (18)$$

must be used.

4. Find the flux linking with each turn,  $\Phi$ , from  $B$  by integration:

$$\Phi = \int_S B \cdot ds \quad (19)$$

where  $S$  is the area over which  $B$  exists and links with the assumed current.

5. Find the flux linkage  $\Lambda$  by multiplying  $\Phi$  by the number of turns.
6. Find  $L$  by taking the ratio  $L = \frac{\Lambda}{I}$ .

To determine the mutual inductance  $L_{12}$  between two circuits, after choosing an appropriate coordinate system, assume  $I_1 \rightarrow$  Find  $B_1 \rightarrow$  Find  $\Phi_{12}$  by integrating  $B_1$  over surface  $S_2 \rightarrow$  Find flux linkage  $\Lambda_{12} = N_2 \Phi_{12} \rightarrow$  Find  $L_{12} = \frac{\Lambda_{12}}{I_1}$ .

In high-frequency applications, current tends to concentrate in the "skin" of the inner conductor as a surface current, internal self-inductance tends to zero. Neumann formula for mutual inductance:

$$L_{12} = L_{21} = \frac{\mu_0}{4\pi} \oint_{C_1} \oint_{C_2} \frac{dl_1 \cdot dl_2}{R} \quad (20)$$

## Magnetic Energy

For a system of  $N$  loops carrying currents,

$$W_m = \frac{1}{2} \sum_{j=1}^N \sum_{k=1}^N L_{jk} I_j I_k \quad (21)$$

For a single inductor,

$$W_m = \frac{1}{2}LI^2 \quad (22)$$

$$W_m = \frac{1}{2} \sum_{k=1}^N I_k \Phi_k \quad (23)$$

## Magnetic Energy In Terms of Field Quantities

Generally,

$$W_m = \frac{1}{2} \int_{V'} \mathbf{A} \cdot \mathbf{J} dv' = \frac{1}{2} \int_{V'} \mathbf{H} \cdot \mathbf{B} dv' = \frac{1}{2} \int_{V'} \frac{B^2}{\mu} dv' = \frac{1}{2} \int_{V'} \mu H^2 dv' \quad (24)$$

If we define a magnetic energy density  $w_m$  such that

$$W_m = \int_{V'} w_m dv' \quad (25)$$

then

$$w_m = \frac{1}{2} \mathbf{H} \cdot \mathbf{B} = \frac{B^2}{2\mu} = \frac{1}{2} \mu H^2 \quad (26)$$

And we have

$$L = \frac{2W_m}{I^2} \quad (27)$$

## Reference

- Fan Hu, RC Slides, VE230
- Nana Liu, Slides, VE230

**Thanks!**

