

L6 Hamiltonian II

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$$(q^1, \dots, q^n, \dot{q}^1, \dots, \dot{q}^n) \quad p_i = \frac{\partial L}{\partial \dot{q}^i} \quad \text{Legendre Transformation}$$
$$(q^1, \dots, q^n, p_1, \dots, p_n)$$

对于 $f(x) \quad df = \frac{\partial f}{\partial x} dx$ (下面看为什么勒让得变换)

$$\therefore df = pdx. \quad d(xp - f) = dx(p + Xdp) - \frac{\partial f}{\partial x} dx = Xdp.$$

$$\text{定义 } f^* = xp - f. \quad \therefore \frac{\partial f^*}{\partial p} = X.$$

Variational Non-holonomic & Nonholonomic

$$L_A = L + \lambda_1(\) + \lambda_2(\) + \dots$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} + \sum_{j=1}^{n-m} \lambda_j \dot{q}_j$$

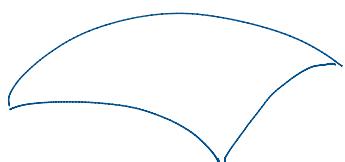
如果约束是 Holonomic 的. $\dot{q}_j = \frac{d\phi^j}{dt}$ 那么两者一样.

例: 圆上的粒子.

$$L = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) \quad \text{限制在 } x^2 + y^2 = 1 \text{ 上.}$$

$$L_A = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) + \lambda(\dot{x}x + \dot{y}y) \Rightarrow \begin{cases} m\ddot{x} = \lambda x \\ m\ddot{y} = \lambda y. \end{cases}$$

$$T = \sum_{ij} \bar{g}_{ij}(q) \dot{q}_i \dot{q}_j \quad \mathbb{R}^n: Riemannian \text{ Mechanism.}$$



- dynamics due to take.

- geometric flow

"Least" Action

- Modified Hamilton Principal

$$\delta \int_{t_1}^{t_2} L dt \quad \text{Hamiltonian: } \delta \int_{t_1}^{t_2} \left(\sum_i p_i \dot{q}_i - H(p, q) \right) dt$$

$$\int_{t_1}^{t_2} \left(\sum_i \delta p_i \left(\dot{q}_i - \frac{\partial H}{\partial p_i} \right) - \delta q_i \left(\dot{p}_i + \frac{\partial H}{\partial q_i} \right) \right) dt = 0.$$

$$\text{所以 } \dot{q}_i = \frac{\partial H}{\partial p_i} \quad \dot{p}_i = -\frac{\partial H}{\partial q_i} \quad (\text{这里 } p \text{ 和 } q \text{ 是独立的}).$$

历史上叫 Least Action.

$$A = \int_{t_1}^{t_2} \sum_i p_i \dot{q}_i dt \quad L_A(\text{critical}) = \Delta \int_{t_1}^{t_2} \sum_i p_i \dot{q}_i dt$$

- 允许改变在终点的时间, 但约束能量

$$\text{如果 } L = T - V \quad \text{而 } V=0. \quad T = \frac{1}{2} \sum_{ij} g_{ij} \dot{q}_i \dot{q}_j \quad \sum_i p_i \dot{q}_i = 2T$$

$$\text{Least Action} \quad \Delta \int T dt \quad \Delta \int dt \Rightarrow \text{用最短时间} \quad (\text{从费马原理出发的})$$

$$\text{Hamilton: } \dot{q}_i^i = \frac{\partial H}{\partial p_i} \quad \dot{p}_i = -\frac{\partial H}{\partial q_i} \quad H(q_i, p_i)$$

定义: 对于两个独立的平滑函数 $F(q^i, p_i), G(q^i, p_i)$

$$\text{我们定义相括号 } \sum_{i=1}^n \left(\frac{\partial F}{\partial q_i} \frac{\partial G}{\partial p_i} - \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} \right) = \{F, G\}$$

Note: ① 相括号是双线性的 $\{F, G+H\} = \{F, G\} + \{F, H\}$

$$\textcircled{2} \quad (\text{Skew}) \quad \{F, G\} = -\{G, F\}$$

$$\textcircled{3} \quad (\text{Jacobi Identity}) \quad \{F, \{G, H\}\} + \{G, \{H, F\}\} + \{H, \{F, G\}\} = 0.$$

$$\text{我们发现} \begin{cases} \dot{q}_i^i = \{q^i, H\} = \frac{\partial H}{\partial p_i} \\ \dot{p}_i = \{p_i, H\} = -\frac{\partial H}{\partial q_i} \end{cases}$$

对于任何的守恒量, 我们可以用 $\{G, H\}$ 来检测 G 是否守恒

Rigid Body

Euler Equations : 转动惯量 I_1, I_2, I_3 , 角速度 $\Omega_1, \Omega_2, \Omega_3$. 令 $I_1 > I_2 > I_3$

$$\begin{cases} I_1 \dot{\Omega}_1 = (I_2 - I_3) \Omega_2 \Omega_3 \\ I_2 \dot{\Omega}_2 = (I_3 - I_1) \Omega_1 \Omega_3 \\ I_3 \dot{\Omega}_3 = (I_1 - I_2) \Omega_1 \Omega_2 \end{cases} \quad \text{Quadratic.}$$

Energy 和 Momentum 都守恒.

$$H = \frac{1}{2} \sum I_i \Omega_i^2 \quad M = \sum I_i^2 \Omega_i^2 \quad \dot{H} = \dot{M} = 0 \quad \text{along flow}$$

$$\pi_i = I_i \Omega_i, \quad i=1,2,3. \quad \text{Rewrite} \quad \dot{\pi}_i = \frac{(I_2 - I_3) \pi_2 \pi_3}{I_2 I_3} \quad \Rightarrow \quad \dot{\pi} = \pi \times \Omega$$

Hamilton with respect to bracket

$$\{F, G\}(\pi) = -\pi (\nabla F \times \nabla G).$$

$$\text{Euler equations} \quad \dot{F} = \underbrace{\{F, H\}}_{\text{Lie-Poisson}}$$

$$\text{对于 Momentum } \sum_i \pi_i^2 = C(\pi) \quad \text{Casimir} \quad \Rightarrow \text{相容括号的性质,}$$

$$\dot{C} = \{C, H\} = 0 \quad \text{因为 } C = F, \quad \nabla C = \nabla F = \pi \quad \text{所以 } \dot{C} = 0.$$