

Keldysh vertices and mfRG flow equations

Santiago Aguirre & Elias Walter

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General properties of the four-point vertex

$$\Gamma = \Gamma_{1'2'|12} = \Gamma_{\sigma'_1\sigma'_2|\sigma_1\sigma_2}^{\alpha'_1\alpha'_2|\alpha_1\alpha_2}(q'_1q'_2|q_1q_2, \nu'_1\nu'_2|\nu_1\nu_2) \quad (1)$$

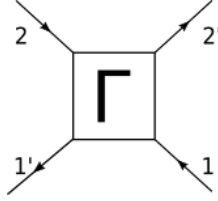


Figure 1: Convention for the vertex

The Keldysh indices take values in the set $\{1, 2\}$, with the convention $1 = q$ and $2 = c$, where in some texts c stands for 'classical component' and q for 'quantum component'.

Spin conservation:

Spin conservation: $\sigma_{1'} + \sigma_{2'} = \sigma_1 + \sigma_2$

Overall spin conservation of the vertex dictates that there are two different cases:

All spins equal

Sum of the incoming (or outgoing) spins equals either -1 or 1 : $\Gamma = \Gamma_{\sigma\sigma|\sigma\sigma}$

Incoming spins reversed

Sum of incoming spins equals 0 : $\Gamma = \Gamma_{\sigma\bar{\sigma}|\sigma\bar{\sigma}}$ or $\Gamma = \Gamma_{\sigma\bar{\sigma}|\bar{\sigma}\sigma}$ These last two cases can be linked through a spin-inversion transformation and, hence, aren't considered to be (computationally) different, although they represent distinct physical processes.

Frequency conservation:

Conservation of momentum implies that one of the frequencies, conventionally chosen to be ν_2 , will be dependent on the other three.

$$\nu_{1'} + \nu_{2'} = \nu_1 + \nu_2 \quad (2)$$

$$\implies \nu_2 = \nu_{1'} + \nu_{2'} - \nu_1 \quad (3)$$

Hence, we will omit ν_2 altogether from all formulas it would have appeared in.

Causality:

Causality, in the Keldysh formalism, translates to $\Gamma^{22|22} = 0$ always.

Particle exchange and complex conjugation:

Here is how the vertex components are related under particle exchange:

$$\Gamma_{\sigma'_2 \sigma'_1 | \sigma_1 \sigma_2}^{\alpha'_2 \alpha'_1 | \alpha_1 \alpha_2} (q'_2 q'_1 | q_1 q_2, \nu'_2 \nu'_1 | \nu_1 \nu_2) \quad (4)$$

$$= \Gamma_{\sigma'_1 \sigma'_2 | \sigma_2 \sigma_1}^{\alpha'_1 \alpha'_2 | \alpha_2 \alpha_1} (q'_1 q'_2 | q_2 q_1, \nu'_1 \nu'_2 | \nu_2 \nu_1) \quad (5)$$

$$= -\Gamma_{\sigma'_2 \sigma'_1 | \sigma_2 \sigma_1}^{\alpha'_2 \alpha'_1 | \alpha_2 \alpha_1} (q'_2 q'_1 | q_2 q_1, \nu'_2 \nu'_1 | \nu_2 \nu_1) \quad (6)$$

$$= -\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \quad (7)$$

And under complex conjugation:

$$\begin{aligned} & \left[\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \right]^* \\ &= (-1)^{(1+\alpha'_1+\alpha'_2+\alpha_1+\alpha_2)} \Gamma_{\sigma_1 \sigma_2 | \sigma'_1 \sigma'_2}^{\alpha_1 \alpha_2 | \alpha'_1 \alpha'_2} (q_1 q_2 | q'_1 q'_2, \nu_1 \nu_2 | \nu'_1 \nu'_2) \end{aligned} \quad (8)$$

Channel decomposition & reducibility

In the mfRG formalism we have that the full vertex can be written as the sum of the irreducible part $R = U + \mathcal{O}(u^4)$ and reducible contributions γ_r in three channels: $r = a$ for *anti-parallel*, $r = p$ for *parallel* and $r = t$ for *transverse*. These names allude to the relation of the cut propagator lines have, were one to make a diagram disconnected. Since any reducible diagram can only belong to one of these classes[], one can then write the following equation for the full vertex:

$$\Gamma = R + \sum_r \gamma_r = R + \gamma_a + \gamma_p + \gamma_t \quad (9)$$

This reducibility of the diagrams then gives rise to the “bubble” objects, Π_r , which are the pair of propagators that, should they be cut, would disconnect the diagram. These pairs of propagators effectively carry a (bosonic) transfer frequency between the two sides of the bubble. Knowing this, we can tailor an explicit frequency dependence that highlights and exploits this fact, in order to then simplify calculations. Hence, we introduce the following structures:

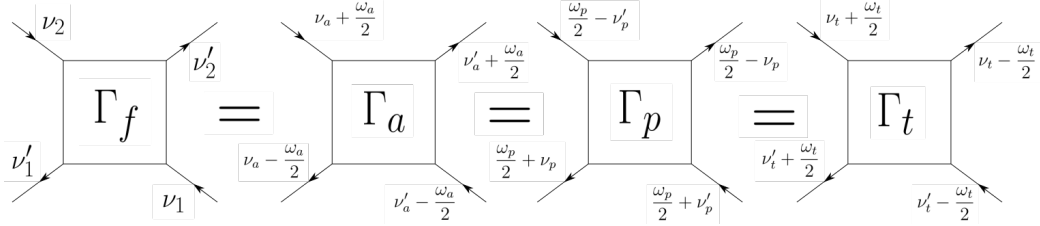


Figure 2: Convention for the introduction of channel-specific frequencies

In this notation we explicitly describe the full vertex function Γ , which is naturally a function of three arguments, as slight variations that depend on the convention for the respective channel. This means that all Γ_r represent the same object, namely the full vertex, but each instantiation receives the arguments ω_r , ν_r and ν_r' , the bosonic and the two fermionic frequencies, respectively, already expressed in the channel's natural parametrization. This clarification is important since these objects are fundamentally different from the notationally-similar γ_r , which are the reducible vertices in the r -channel.

The frequencies in Fig. 2 lead then to the following conversion rules between the channels, where the conventional vertex depicted in Fig. 1 is referred to as the 'fermionic' vertex, Γ_f :

a-channel

From and to the fermionic channel:

$$\left. \begin{aligned} \nu_{1'} &= \nu_a - \frac{\omega_a}{2} \\ \nu_{2'} &= \nu_a' + \frac{\omega_a}{2} \\ \nu_1 &= \nu_a' - \frac{\omega_a}{2} \end{aligned} \right\} \implies \left\{ \begin{aligned} \omega_a &= \nu_2 - \nu_{1'} = \nu_{2'} - \nu_1 \\ \nu_a &= \frac{1}{2}(2\nu_{1'} + \nu_{2'} - \nu_1) \\ \nu_a' &= \frac{1}{2}(\nu_{2'} + \nu_1) \end{aligned} \right. \quad (10)$$

And from any other channel to the a-channel:

$$\begin{aligned} \omega_a &= \omega_a & \omega_a &= -\nu_p - \nu_p' & \omega_a &= \nu_t - \nu_t' \\ \nu_a &= \nu_a & \nu_a &= \frac{\omega_p + \nu_p - \nu_p'}{2} & \nu_a &= \frac{\omega_t + \nu_t + \nu_t'}{2} \\ \nu_a' &= \nu_a' & \nu_a' &= \frac{\omega_p - \nu_p + \nu_p'}{2} & \nu_a' &= \frac{-\omega_t + \nu_t + \nu_t'}{2} \end{aligned} \quad (11)$$

p-channel

From and to the fermionic channel:

$$\left. \begin{array}{l} \nu_{1'} = \frac{\omega_p}{2} + \nu_p \\ \nu_{2'} = \frac{\omega_p}{2} - \nu_p \\ \nu_1 = \frac{\omega_p}{2} + \nu'_p \end{array} \right\} \implies \left\{ \begin{array}{l} \omega_p = \nu_{1'} + \nu_{2'} = \nu_1 + \nu_2 \\ \nu_p = \frac{1}{2}(\nu_{1'} - \nu_{2'}) \\ \nu'_p = \frac{1}{2}(2\nu_1 - \nu_{1'} - \nu_{2'}) \end{array} \right. \quad (12)$$

And from any other channel to the p-channel:

$$\begin{array}{lll} \omega_p = \nu_a + \nu'_a & \omega_p = \omega_p & \omega_p = \nu_t + \nu'_t \\ \nu_p = \frac{-\omega_a + \nu_a - \nu'_a}{2} & \nu_p = \nu_p & \nu_p = \frac{\omega_t - \nu_t + \nu'_t}{2} \\ \nu'_p = \frac{-\omega_a - \nu_a + \nu'_a}{2} & \nu'_p = \nu'_p & \nu'_p = \frac{-\omega_t - \nu_t + \nu'_t}{2} \end{array} \quad (13)$$

t-channel

From and to the fermionic channel

$$\left. \begin{array}{l} \nu_{1'} = \nu'_t + \frac{\omega_t}{2} \\ \nu_{2'} = \nu_t - \frac{\omega_t}{2} \\ \nu_1 = \nu'_t - \frac{\omega_t}{2} \end{array} \right\} \implies \left\{ \begin{array}{l} \omega_t = \nu_{1'} - \nu_1 = \nu_2 - \nu_{2'} \\ \nu_t = \frac{1}{2}(2\nu_{2'} + \nu_{1'} - \nu_1) \\ \nu'_t = \frac{1}{2}(\nu_{1'} + \nu_1) \end{array} \right. \quad (14)$$

And from any other channel to the t-channel:

$$\begin{array}{lll} \omega_t = \nu_a - \nu'_a & \omega_t = \nu_p - \nu'_p & \omega_t = \omega_t \\ \nu_t = \frac{\omega_a + \nu_a + \nu'_a}{2} & \nu_t = \frac{\omega_p - \nu_p - \nu'_p}{2} & \nu_t = \nu_t \\ \nu'_t = \frac{-\omega_a + \nu_a + \nu'_a}{2} & \nu'_t = \frac{\omega_p + \nu_p + \nu'_p}{2} & \nu'_t = \nu'_t \end{array} \quad (15)$$

Transformations

Define the following transformations, based on the previous properties:

Spin flip

$$T_S \Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) = \Gamma_{\bar{\sigma}'_1 \bar{\sigma}'_2 | \bar{\sigma}_1 \bar{\sigma}_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \quad (16)$$

Exchange symmetries

Exchange of the Keldysh and spin indices of the incoming (T_1), outgoing (T_2) and incoming and outgoing (T_3) legs lead to the following vertex-internal dependencies:

Incoming legs exchange - T_1

$$T_1 \left(\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \right) = \Gamma_{\sigma'_1 \sigma'_2 | \sigma_2 \sigma_1}^{\alpha'_1 \alpha'_2 | \alpha_2 \alpha_1} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \quad (17)$$

$$= -\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_2 q_1, \nu'_1 \nu'_2 | \nu_2 \nu_1) \quad (18)$$

Outgoing legs exchange - T_2

$$T_2 \left(\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \right) = \Gamma_{\sigma'_2 \sigma'_1 | \sigma_1 \sigma_2}^{\alpha'_2 \alpha'_1 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \quad (19)$$

$$= -\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_2 q'_1 | q_1 q_2, \nu'_2 \nu'_1 | \nu_1 \nu_2) \quad (20)$$

Incoming and outgoing legs exchange - T_3

$$T_3 \left(\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \right) = \Gamma_{\sigma'_2 \sigma'_1 | \sigma_2 \sigma_1}^{\alpha'_2 \alpha'_1 | \alpha_2 \alpha_1} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \quad (21)$$

$$= \Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_2 q'_1 | q_2 q_1, \nu'_2 \nu'_1 | \nu_2 \nu_1) \quad (22)$$

Complex conjugation (exchange of incoming and outgoing legs) - T_C

$$T_C(\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2)) = \Gamma_{\sigma_1 \sigma_2 | \sigma'_1 \sigma'_2}^{\alpha_1 \alpha_2 | \alpha'_1 \alpha'_2} (q'_1 q'_2 | q_1 q_2, \nu'_1 \nu'_2 | \nu_1 \nu_2) \quad (23)$$

$$= (-1)^{1+\alpha_1+\alpha_2+\alpha'_1+\alpha'_2} [\Gamma_{\sigma'_1 \sigma'_2 | \sigma_1 \sigma_2}^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} (q_1 q_2 | q'_1 q'_2, \nu_1 \nu_2 | \nu'_1 \nu'_2)]^* \quad (24)$$

Notice the T_i transformations ($i \in \{1, 2, 3\}$) form a group which is isomorphic to the Klein group (every element has order 2 and $(T_i)^{-1} = T_i$). This fact, in combination with the action of T_C , the equations that define the multiplication map of the “expanded” group are:

$$\begin{aligned} T_1 T_2 &= T_3 & T_1 T_C &= T_C T_2 \\ T_2 T_1 &= T_3 & T_C &= T_1 T_C T_2 = T_2 T_C T_1 \end{aligned} \quad (25)$$

The inclusion of T_S expands the group, but does not yield any interesting structure.

Independent spin and Keldysh components

Keldysh components which are equal due to diagrammatic structure of \mathcal{K}_1 , \mathcal{K}_2 , \mathcal{K}_3

Class \mathcal{K}_1 :

$$\begin{aligned} (\mathcal{K}_1^a)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\mathcal{K}_1^a)^{\bar{\alpha}'_1 \alpha'_2 | \alpha_1 \bar{\alpha}_2} = (\mathcal{K}_1^a)^{\alpha'_1 \bar{\alpha}'_2 | \bar{\alpha}_1 \alpha_2} = (\mathcal{K}_1^a)^{\bar{\alpha}'_1 \bar{\alpha}'_2 | \bar{\alpha}_1 \bar{\alpha}_2} \\ (\mathcal{K}_1^p)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\mathcal{K}_1^p)^{\bar{\alpha}'_1 \bar{\alpha}'_2 | \alpha_1 \alpha_2} = (\mathcal{K}_1^p)^{\alpha'_1 \alpha'_2 | \bar{\alpha}_1 \bar{\alpha}_2} = (\mathcal{K}_1^p)^{\bar{\alpha}'_1 \bar{\alpha}'_2 | \bar{\alpha}_1 \bar{\alpha}_2} \\ (\mathcal{K}_1^t)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\mathcal{K}_1^t)^{\bar{\alpha}'_1 \alpha'_2 | \bar{\alpha}_1 \alpha_2} = (\mathcal{K}_1^t)^{\alpha'_1 \bar{\alpha}'_2 | \alpha_1 \bar{\alpha}_2} = (\mathcal{K}_1^t)^{\bar{\alpha}'_1 \bar{\alpha}'_2 | \bar{\alpha}_1 \bar{\alpha}_2} \end{aligned}$$

Class \mathcal{K}_2 , $\bar{\mathcal{K}}_2$:

$$\begin{aligned} (\mathcal{K}_2^a)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\mathcal{K}_2^a)^{\alpha'_1 \bar{\alpha}'_2 | \bar{\alpha}_1 \alpha_2}, & (\bar{\mathcal{K}}_2^a)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\bar{\mathcal{K}}_2^a)^{\bar{\alpha}'_1 \alpha'_2 | \alpha_1 \bar{\alpha}_2} \\ (\mathcal{K}_1^p)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\mathcal{K}_1^p)^{\alpha'_1 \alpha'_2 | \bar{\alpha}_1 \bar{\alpha}_2}, & (\bar{\mathcal{K}}_1^p)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\bar{\mathcal{K}}_1^p)^{\bar{\alpha}'_1 \bar{\alpha}'_2 | \alpha_1 \alpha_2} \\ (\mathcal{K}_2^t)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\mathcal{K}_2^t)^{\bar{\alpha}'_1 \alpha'_2 | \bar{\alpha}_1 \alpha_2}, & (\bar{\mathcal{K}}_2^t)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} &= (\bar{\mathcal{K}}_2^t)^{\alpha'_1 \bar{\alpha}'_2 | \alpha_1 \bar{\alpha}_2} \end{aligned}$$

Equal spins $\sigma\sigma|\sigma\sigma$

Symmetries under particle exchange, complex conjugation

Class \mathcal{K}_1 :

- particle exchange:

– channels a , t :

$$\begin{aligned} \mathcal{K}_1^a(1'2'|12) &= -\mathcal{K}_1^t(1'2'|21) \Rightarrow (\mathcal{K}_1^a)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} = T_1(\mathcal{K}_1^t)^{\alpha'_1 \alpha'_2 | \alpha_2 \alpha_1} \\ &= -\mathcal{K}_1^t(2'1'|12) \Rightarrow (\mathcal{K}_1^a)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} = T_2(\mathcal{K}_1^t)^{\alpha'_2 \alpha'_1 | \alpha_1 \alpha_2} \\ \mathcal{K}_1^a(1'2'|12) &= \mathcal{K}_1^a(2'1'|21) \Rightarrow (\mathcal{K}_1^a)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} = T_3(\mathcal{K}_1^a)^{\alpha'_2 \alpha'_1 | \alpha_2 \alpha_1} \\ \mathcal{K}_1^t(1'2'|12) &= \mathcal{K}_1^t(2'1'|21) \Rightarrow (\mathcal{K}_1^t)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} = T_3(\mathcal{K}_1^t)^{\alpha'_2 \alpha'_1 | \alpha_2 \alpha_1} \end{aligned}$$

– channel p :

$$\begin{aligned} \mathcal{K}_1^p(1'2'|12) &= -\mathcal{K}_1^p(1'2'|21) \Rightarrow (\mathcal{K}_1^p)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} = T_1(\mathcal{K}_1^p)^{\alpha'_1 \alpha'_2 | \alpha_2 \alpha_1} \\ &= -\mathcal{K}_1^p(2'1'|12) \Rightarrow (\mathcal{K}_1^p)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} = T_2(\mathcal{K}_1^p)^{\alpha'_2 \alpha'_1 | \alpha_1 \alpha_2} \\ &= \mathcal{K}_1^p(2'1'|21) \Rightarrow (\mathcal{K}_1^p)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} = T_3(\mathcal{K}_1^p)^{\alpha'_2 \alpha'_1 | \alpha_2 \alpha_1} \end{aligned}$$

- complex conjugation:

$$\mathcal{K}_1^r(1'2'|12) = -(-1)^{\sum_j (\alpha_j + \bar{\alpha}_j)} (\mathcal{K}_1^r(12|1'2'))^* \Rightarrow (\mathcal{K}_1^r)^{\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2} = T_C(\mathcal{K}_1^r)^{\alpha_1 \alpha_2 | \alpha'_1 \alpha'_2}$$

Class $\mathcal{K}_2, \bar{\mathcal{K}}_2$:

- particle exchange:

– channels a, t :

$$\begin{aligned}
\mathcal{K}_2^a(1'2'|12) &= -\bar{\mathcal{K}}_2^t(1'2'|21) \Rightarrow (\mathcal{K}_2^a)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_1(\bar{\mathcal{K}}_2^t)^{\alpha'_1\alpha'_2|\alpha_2\alpha_1} \\
&= -\mathcal{K}_2^t(2'1'|12) \Rightarrow (\mathcal{K}_2^a)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_2(\mathcal{K}_2^t)^{\alpha'_2\alpha'_1|\alpha_1\alpha_2} \\
\mathcal{K}_2^a(1'2'|12) &= \bar{\mathcal{K}}_2^a(2'1'|21) \Rightarrow (\mathcal{K}_2^a)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_3(\bar{\mathcal{K}}_2^a)^{\alpha'_2\alpha'_1|\alpha_2\alpha_1} \\
\mathcal{K}_2^t(1'2'|12) &= \bar{\mathcal{K}}_2^t(2'1'|21) \Rightarrow (\mathcal{K}_2^t)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_3(\bar{\mathcal{K}}_2^t)^{\alpha'_2\alpha'_1|\alpha_2\alpha_1}
\end{aligned}$$

– channel p :

$$\begin{aligned}
\mathcal{K}_2^p(1'2'|12) &= -\mathcal{K}_2^p(1'2'|21) \Rightarrow (\mathcal{K}_2^p)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_1(\mathcal{K}_2^p)^{\alpha'_1\alpha'_2|\alpha_2\alpha_1} \\
&= -\mathcal{K}_2^p(2'1'|12) \Rightarrow (\mathcal{K}_2^p)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_2(\mathcal{K}_2^p)^{\alpha'_2\alpha'_1|\alpha_1\alpha_2} \\
&= \mathcal{K}_2^p(2'1'|21) \Rightarrow (\mathcal{K}_2^p)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_3(\mathcal{K}_2^p)^{\alpha'_2\alpha'_1|\alpha_2\alpha_1}
\end{aligned}$$

(similarly for $\bar{\mathcal{K}}_2^p$)

- complex conjugation:

$$\begin{aligned}
\mathcal{K}_2^{a,p}(1'2'|12) &= -(-1)^{\sum_j(\alpha_j+\bar{\alpha}_j)}(\bar{\mathcal{K}}_2^{a,p}(12|1'2'))^* \Rightarrow (\mathcal{K}_2^{a,p})^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_C(\bar{\mathcal{K}}_2^{a,p})^{\alpha_1\alpha_2|\alpha'_1\alpha'_2} \\
\mathcal{K}_2^t(1'2'|12) &= -(-1)^{\sum_j(\alpha_j+\bar{\alpha}_j)}(\mathcal{K}_2^t(12|1'2'))^* \Rightarrow (\mathcal{K}_2^t)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_C(\mathcal{K}_2^t)^{\alpha_1\alpha_2|\alpha'_1\alpha'_2} \\
\bar{\mathcal{K}}_2^t(1'2'|12) &= -(-1)^{\sum_j(\alpha_j+\bar{\alpha}_j)}(\bar{\mathcal{K}}_2^t(12|1'2'))^* \Rightarrow (\bar{\mathcal{K}}_2^t)^{\alpha'_1\alpha'_2|\alpha_1\alpha_2} = T_C(\bar{\mathcal{K}}_2^t)^{\alpha_1\alpha_2|\alpha'_1\alpha'_2}
\end{aligned}$$

Enumeration of the components

To make the whole thing a little easier and be able to refer to the components of the vertex in a concrete and “linear” manner, we convert the indices $\alpha'_1\alpha'_2|\alpha_1\alpha_2$ to numbers in the set $\mathcal{N} := \{0, 1, \dots, 15\}$. The following formula does the trick:

$$i_{\mathcal{N}} = 2^3(\alpha'_1 - 1) + 2^2(\alpha'_2 - 1) + 2^1(\alpha_1 - 1) + 2^0(\alpha_2 - 1) \quad (26)$$

$$= 8(\alpha'_1 - 1) + 4(\alpha'_2 - 1) + 2\alpha_1 - 1 + (\alpha_2 - 1) \quad (27)$$

There’s also an inverse function for this, which involves using both the floor and the mod functions:

$$\alpha_2 = \left\lfloor \frac{(i_{\mathcal{N}} \bmod 2)}{1} \right\rfloor + 1 = (i_{\mathcal{N}} \bmod 2) + 1 \quad (28)$$

$$\alpha_1 = \left\lfloor \frac{(i_{\mathcal{N}} \bmod 4)}{2} \right\rfloor + 1 \quad (29)$$

$$\alpha'_2 = \left\lfloor \frac{(i_{\mathcal{N}} \bmod 8)}{4} \right\rfloor + 1 \quad (30)$$

$$\alpha'_1 = \left\lfloor \frac{(i_{\mathcal{N}} \bmod 16)}{8} \right\rfloor + 1 \quad (31)$$

We summarize the results of these functions for future reference:

$$\begin{array}{ll} 11|11 \longleftrightarrow 0 & 21|11 \longleftrightarrow 8 \\ 11|12 \longleftrightarrow 1 & 21|12 \longleftrightarrow 9 \\ 11|21 \longleftrightarrow 2 & 21|21 \longleftrightarrow 10 \\ 11|22 \longleftrightarrow 3 & 21|22 \longleftrightarrow 11 \\ \\ 12|11 \longleftrightarrow 4 & 22|11 \longleftrightarrow 12 \\ 12|12 \longleftrightarrow 5 & 22|12 \longleftrightarrow 13 \\ 12|21 \longleftrightarrow 6 & 22|21 \longleftrightarrow 14 \\ 12|22 \longleftrightarrow 7 & 22|22 \longleftrightarrow 15 \end{array} \quad (32)$$

This allows us to refer to the components in the matrix as i.e. “the 7th component”, where now it’s unambiguously clear that the element referred to is $\Gamma^{12|22}$.

The above numeration and transformation properties of the respective diagrammatic classes yield the following tables:

		$\sigma\sigma \sigma\sigma$			$\sigma\bar{\sigma} \sigma\bar{\sigma}$			$\sigma\bar{\sigma} \bar{\sigma}\sigma$		
		\mathcal{K}_1^a	\mathcal{K}_1^p	\mathcal{K}_1^t	\mathcal{K}_1^a	\mathcal{K}_1^p	\mathcal{K}_1^t	\mathcal{K}_1^a	\mathcal{K}_1^p	\mathcal{K}_1^t
1111	0	0	0	0	0	0	0	0	0	0
1112	1	B_1^a	B_1^p	B_1^t	\bar{B}_1^a	\bar{B}_1^p	\bar{B}_1^t	$T_S T_2 \bar{B}_1^t$	$T_1 \bar{B}_1^p$	$T_S T_2 \bar{B}_1^a$
1121	2	$T_3 B_1^a$	B_1^p	$T_3 B_1^t$	$T_S T_3 \bar{B}_1^a$	\bar{B}_1^p	$T_S T_3 \bar{B}_1^t$	$T_1 \bar{B}_1^t$	$T_1 \bar{B}_1^p$	$T_1 \bar{B}_1^a$
1122	3	C_1^a	0	C_1^t	\bar{C}_1^a	0	\bar{C}_1^t	$T_1 \bar{C}_1^t$	0	$T_1 \bar{C}_1^a$
1211	4	$T_3 B_1^a$	$T_C B_1^p$	B_1^t	$T_S T_3 \bar{B}_1^a$	$T_C \bar{B}_1^p$	\bar{B}_1^t	$T_1 \bar{B}_1^t$	$T_1 T_C \bar{B}_1^p$	$T_S T_2 \bar{B}_1^a$
1212	5	C_1^a	D_1^p	0	\bar{C}_1^a	\bar{D}_1^p	0	$T_1 \bar{C}_1^t$	$T_1 \bar{D}_1^p$	0
1221	6	0	D_1^p	C_1^t	0	\bar{D}_1^p	\bar{C}_1^t	0	$T_1 \bar{D}_1^p$	$T_1 \bar{C}_1^a$
1222	7	B_1^a	$T_C B_1^p$	$T_3 B_1^t$	\bar{B}_1^a	$T_C \bar{B}_1^p$	$T_S T_3 \bar{B}_1^t$	$T_S T_2 \bar{B}_1^t$	$T_1 T_C \bar{B}_1^p$	$T_1 \bar{B}_1^a$
2111	8	B_1^a	$T_C B_1^p$	$T_3 B_1^t$	\bar{B}_1^a	$T_C \bar{B}_1^p$	$T_S T_3 \bar{B}_1^t$	$T_S T_2 \bar{B}_1^t$	$T_1 T_C \bar{B}_1^p$	$T_1 \bar{B}_1^a$
2112	9	0	D_1^p	C_1^t	0	\bar{D}_1^p	\bar{C}_1^t	0	$T_1 \bar{D}_1^p$	$T_1 \bar{C}_1^a$
2121	10	C_1^a	D_1^p	0	\bar{C}_1^a	\bar{D}_1^p	0	$T_1 \bar{C}_1^t$	$T_1 \bar{D}_1^p$	0
2122	11	$T_3 B_1^a$	$T_C B_1^p$	B_1^t	$T_S T_3 \bar{B}_1^a$	$T_C \bar{B}_1^p$	\bar{B}_1^t	$T_1 \bar{B}_1^t$	$T_1 T_C \bar{B}_1^p$	$T_S T_2 \bar{B}_1^a$
2211	12	C_1^a	0	C_1^t	\bar{C}_1^a	0	\bar{C}_1^t	$T_1 \bar{C}_1^t$	0	$T_1 \bar{C}_1^a$
2212	13	$T_3 B_1^a$	B_1^p	$T_3 B_1^t$	$T_S T_3 \bar{B}_1^a$	\bar{B}_1^p	$T_S T_3 \bar{B}_1^t$	$T_1 \bar{B}_1^t$	$T_1 \bar{B}_1^p$	$T_1 \bar{B}_1^a$
2221	14	B_1^a	B_1^p	B_1^t	\bar{B}_1^a	\bar{B}_1^p	\bar{B}_1^t	$T_S T_2 \bar{B}_1^t$	$T_1 \bar{B}_1^p$	$T_S T_2 \bar{B}_1^a$
2222	15	0	0	0	0	0	0	0	0	0

		$\sigma\sigma \sigma\sigma$					
		\mathcal{K}_2^a	$\bar{\mathcal{K}}_2^a$	\mathcal{K}_2^p	$\bar{\mathcal{K}}_2^p$	\mathcal{K}_2^t	$\bar{\mathcal{K}}_2^t$
1111	0	A_2^a	$T_3 A_2^a$	A_2^p	$T_C A_2^p$	$T_2 A_2^a$	$T_1 A_2^a$
1112	1	B_2^a	$T_3 C_2^a$	B_2^p	$T_C C_2^p$	$T_2 B_2^a$	$T_1 C_2^a$
1121	2	C_2^a	$T_3 B_2^a$	B_2^p	$T_C T_3 C_2^p$	$T_2 C_2^a$	$T_1 B_2^a$
1122	3	D_2^a	$T_3 D_2^a$	A_2^p	0	$T_2 D_2^a$	$T_1 D_2^a$
1211	4	C_2^a	$T_C B_2^a$	C_2^p	$T_C B_2^p$	$T_C T_2 B_2^a$	$T_1 C_2^a$
1212	5	D_2^a	$T_C D_2^a$	D_2^p	$T_C D_2^p$	0	$T_1 A_2^a$
1221	6	A_2^a	0	D_2^p	$T_C T_3 D_2^p$	$T_C T_2 D_2^a$	$T_1 D_2^a$
1222	7	B_2^a	$T_3 F_2^a$	C_2^p	$T_C F_2^p$	$T_2 F_2^a$	$T_1 B_2^a$
2111	8	$T_C T_3 B_2^a$	$T_3 C_2^a$	$T_3 C_2^p$	$T_C B_2^p$	$T_2 C_2^a$	$T_C T_1 B_2^a$
2112	9	0	$T_3 A_2^a$	$T_3 D_2^p$	$T_C D_2^p$	$T_2 D_2^a$	$T_C T_1 D_2^a$
2121	10	$T_C T_3 D_2^a$	$T_3 D_2^a$	$T_3 D_2^p$	$T_C T_3 D_2^p$	$T_2 A_2^a$	0
2122	11	F_2^a	$T_3 B_2^a$	$T_3 C_2^p$	$T_C F_2^p$	$T_2 B_2^a$	$T_1 F_2^a$
2211	12	$T_C T_3 D_2^a$	$T_C D_2^a$	0	$T_C A_2^p$	$T_C T_2 D_2^a$	$T_C T_1 D_2^a$
2212	13	F_2^a	$T_C B_2^a$	F_2^p	$T_C C_2^p$	$T_2 F_2^a$	$T_C T_1 B_2^a$
2221	14	$T_C T_3 B_2^a$	$T_3 F_2^a$	F_2^p	$T_C T_3 C_2^p$	$T_C T_2 B_2^a$	$T_1 F_2^a$
2222	15	0	0	0	0	0	0

		$\sigma\bar{\sigma} \sigma\bar{\sigma}$					
		\mathcal{K}_2^a	$\bar{\mathcal{K}}_2^a$	\mathcal{K}_2^p	$\bar{\mathcal{K}}_2^p$	\mathcal{K}_2^t	$\bar{\mathcal{K}}_2^t$
1111	0	\bar{A}_2^a	$T_S T_3 \bar{A}_2^a$	\bar{A}_2^p	$T_C \bar{A}_2^p$	\bar{A}_2^t	$T_S T_3 \bar{A}_2^t$
1112	1	\bar{B}_2^a	$T_S T_3 \bar{C}_2^a$	\bar{B}_2^p	$T_C \bar{C}_2^p$	\bar{B}_2^t	$T_S T_3 \bar{C}_2^t$
1121	2	\bar{C}_2^a	$T_S T_3 \bar{B}_2^a$	\bar{B}_2^p	$T_S T_C T_3 \bar{C}_2^p$	\bar{C}_2^t	$T_S T_3 \bar{B}_2^t$
1122	3	\bar{D}_2^a	$T_S T_3 \bar{D}_2^a$	\bar{A}_2^p	0	\bar{D}_2^t	$T_S T_3 \bar{D}_2^t$
1211	4	\bar{C}_2^a	$T_C \bar{B}_2^a$	\bar{C}_2^p	$T_C \bar{B}_2^p$	$T_C \bar{B}_2^t$	$T_S T_3 \bar{C}_2^t$
1212	5	\bar{D}_2^a	$T_C \bar{D}_2^a$	\bar{D}_2^p	$T_C \bar{D}_2^p$	0	$T_S T_3 \bar{A}_2^t$
1221	6	\bar{A}_2^a	0	\bar{D}_2^p	$T_S T_C T_3 \bar{D}_2^p$	$T_C \bar{D}_2^t$	$T_S T_3 \bar{D}_2^t$
1222	7	\bar{B}_2^a	$T_S T_3 \bar{F}_2^a$	\bar{C}_2^p	$T_C \bar{F}_2^p$	\bar{F}_2^t	$T_S T_3 \bar{B}_2^t$
2111	8	$T_S T_C T_3 \bar{B}_2^a$	$T_S T_3 \bar{C}_2^a$	$T_S T_3 \bar{C}_2^p$	$T_C \bar{B}_2^p$	\bar{C}_2^t	$T_S T_C T_3 \bar{B}_2^t$
2112	9	0	$T_S T_3 \bar{A}_2^a$	$T_S T_3 \bar{D}_2^p$	$T_C \bar{D}_2^p$	\bar{D}_2^t	$T_S T_C T_3 \bar{D}_2^t$
2121	10	$T_S T_C T_3 \bar{D}_2^a$	$T_S T_3 \bar{D}_2^a$	$T_S T_3 \bar{D}_2^p$	$T_S T_C T_3 \bar{D}_2^p$	\bar{A}_2^t	0
2122	11	\bar{F}_2^a	$T_S T_3 \bar{B}_2^a$	$T_S T_3 \bar{C}_2^p$	$T_C \bar{F}_2^p$	\bar{B}_2^t	$T_S T_3 \bar{F}_2^t$
2211	12	$T_S T_C T_3 \bar{D}_2^a$	$T_C \bar{D}_2^a$	0	$T_C \bar{A}_2^p$	$T_C \bar{D}_2^t$	$T_S T_C T_3 \bar{D}_2^t$
2212	13	\bar{F}_2^a	$T_C \bar{B}_2^a$	\bar{F}_2^p	$T_C \bar{C}_2^p$	\bar{F}_2^t	$T_S T_C T_3 \bar{B}_2^t$
2221	14	$T_S T_C T_3 \bar{B}_2^a$	$T_S T_3 \bar{F}_2^a$	\bar{F}_2^p	$T_S T_C T_3 \bar{C}_2^p$	$T_C \bar{B}_2^t$	$T_S T_3 \bar{F}_2^t$
2222	15	0	0	0	0	0	0

		$\sigma\bar{\sigma} \bar{\sigma}\sigma$					
		\mathcal{K}_2^a	$\bar{\mathcal{K}}_2^a$	\mathcal{K}_2^p	$\bar{\mathcal{K}}_2^p$	\mathcal{K}_2^t	$\bar{\mathcal{K}}_2^t$
1111	0	$T_S T_2 \bar{A}_2^t$	$T_1 \bar{A}_2^t$	$T_1 \bar{A}_2^p$	$T_S T_C T_1 \bar{A}_2^p$	$T_S T_2 \bar{A}_2^a$	$T_1 \bar{A}_2^a$
1112	1	$T_S T_2 \bar{B}_2^t$	$T_1 \bar{C}_2^t$	$T_1 \bar{B}_2^p$	$T_S T_C T_1 \bar{C}_2^p$	$T_S T_2 \bar{B}_2^a$	$T_1 \bar{C}_2^a$
1121	2	$T_S T_2 \bar{C}_2^t$	$T_1 \bar{B}_2^t$	$T_1 \bar{B}_2^p$	$T_1 T_C \bar{C}_2^p$	$T_S T_2 \bar{C}_2^a$	$T_1 \bar{B}_2^a$
1122	3	$T_S T_2 \bar{D}_2^t$	$T_1 \bar{D}_2^t$	$T_1 \bar{A}_2^p$	0	$T_S T_2 \bar{D}_2^a$	$T_1 \bar{D}_2^a$
1211	4	$T_S T_2 \bar{C}_2^t$	$T_1 T_C \bar{B}_2^t$	$T_1 \bar{C}_2^p$	$T_1 T_C \bar{B}_2^p$	$T_1 T_C \bar{B}_2^a$	$T_1 \bar{C}_2^a$
1212	5	$T_S T_2 \bar{D}_2^t$	$T_1 T_C \bar{D}_2^t$	$T_1 \bar{D}_2^p$	$T_S T_C T_1 \bar{D}_2^p$	0	$T_1 \bar{A}_2^a$
1221	6	$T_S T_2 \bar{A}_2^t$	0	$T_1 \bar{D}_2^p$	$T_1 T_C \bar{D}_2^p$	$T_1 T_C \bar{D}_2^a$	$T_1 \bar{D}_2^a$
1222	7	$T_S T_2 \bar{B}_2^t$	$T_1 \bar{F}_2^t$	$T_1 \bar{C}_2^p$	$T_S T_C T_1 \bar{F}_2^p$	$T_S T_2 \bar{F}_2^a$	$T_1 \bar{B}_2^a$
2111	8	$T_S T_C T_1 \bar{B}_2^t$	$T_1 \bar{C}_2^t$	$T_S T_2 \bar{C}_2^p$	$T_1 T_C \bar{B}_2^p$	$T_S T_2 \bar{C}_2^a$	$T_S T_C T_1 \bar{B}_2^a$
2112	9	0	$T_1 \bar{A}_2^t$	$T_S T_2 \bar{D}_2^p$	$T_S T_C T_1 \bar{D}_2^p$	$T_S T_2 \bar{D}_2^a$	$T_S T_C T_1 \bar{D}_2^a$
2121	10	$T_S T_C T_1 \bar{D}_2^t$	$T_1 \bar{D}_2^t$	$T_S T_2 \bar{D}_2^p$	$T_1 T_C \bar{D}_2^p$	$T_S T_2 \bar{A}_2^a$	0
2122	11	$T_S T_2 \bar{F}_2^t$	$T_1 \bar{B}_2^t$	$T_S T_2 \bar{C}_2^p$	$T_S T_C T_1 \bar{F}_2^p$	$T_S T_2 \bar{B}_2^a$	$T_1 \bar{F}_2^a$
2211	12	$T_S T_C T_1 \bar{D}_2^t$	$T_1 T_C \bar{D}_2^t$	0	$T_S T_C T_1 \bar{A}_2^p$	$T_1 T_C \bar{D}_2^a$	$T_S T_C T_1 \bar{D}_2^a$
2212	13	$T_S T_2 \bar{F}_2^t$	$T_1 T_C \bar{B}_2^t$	$T_1 \bar{F}_2^p$	$T_S T_C T_1 \bar{C}_2^p$	$T_S T_2 \bar{F}_2^a$	$T_S T_C T_1 \bar{B}_2^a$
2221	14	$T_S T_C T_1 \bar{B}_2^t$	$T_1 \bar{F}_2^t$	$T_1 \bar{F}_2^p$	$T_1 T_C \bar{C}_2^p$	$T_1 T_C \bar{B}_2^a$	$T_1 \bar{F}_2^a$
2222	15	0	0	0	0	0	0

		$\sigma\sigma \sigma\sigma$			$\sigma\bar{\sigma} \sigma\bar{\sigma}$			$\sigma\bar{\sigma} \bar{\sigma}\sigma$		
		\mathcal{K}_3^a	\mathcal{K}_3^p	\mathcal{K}_3^t	\mathcal{K}_3^a	\mathcal{K}_3^p	\mathcal{K}_3^t	\mathcal{K}_3^a	\mathcal{K}_3^p	\mathcal{K}_3^t
1111	0	A_3^a	A_3^p	$T_2 A_3^a$	\bar{A}_3^a	\bar{A}_3^p	\bar{A}_3^t	$T_1 \bar{A}_3^t$	$T_1 \bar{A}_3^p$	$T_1 \bar{A}_3^a$
1112	1	B_3^a	B_3^p	$T_2 B_3^a$	\bar{B}_3^a	\bar{B}_3^p	\bar{B}_3^t	$T_S T_2 \bar{B}_3^t$	$T_S T_2 \bar{B}_3^p$	$T_S T_2 \bar{B}_3^a$
1121	2	$T_3 B_3^a$	$T_3 B_3^p$	$T_1 B_3^a$	$T_S T_3 \bar{B}_3^a$	$T_S T_3 \bar{B}_3^p$	$T_S T_3 \bar{B}_3^t$	$T_1 \bar{B}_3^t$	$T_1 \bar{B}_3^p$	$T_1 \bar{B}_3^a$
1122	3	C_3^a	C_3^p	$T_2 C_3^a$	\bar{C}_3^a	\bar{C}_3^p	\bar{C}_3^t	$T_1 \bar{C}_3^t$	$T_1 \bar{C}_3^p$	$T_1 \bar{C}_3^a$
1211	4	$T_C B_3^a$	$T_C B_3^p$	$T_1 T_C B_3^a$	$T_C \bar{B}_3^a$	$T_C \bar{B}_3^p$	$T_C \bar{B}_3^t$	$T_1 T_C \bar{B}_3^t$	$T_1 T_C \bar{B}_3^p$	$T_1 T_C \bar{B}_3^a$
1212	5	D_3^a	D_3^p	D_3^t	\bar{D}_3^a	\bar{D}_3^p	\bar{D}_3^t	$T_1 \bar{E}_3^t$	$T_1 \bar{E}_3^p$	$T_1 \bar{E}_3^a$
1221	6	$T_1 D_3^t$	$T_1 D_3^p$	$T_1 D_3^a$	\bar{E}_3^a	\bar{E}_3^p	\bar{E}_3^t	$T_1 \bar{D}_3^t$	$T_1 \bar{D}_3^p$	$T_1 \bar{D}_3^a$
1222	7	F_3^a	F_3^p	$T_1 F_3^a$	\bar{F}_3^a	\bar{F}_3^p	\bar{F}_3^t	$T_1 \bar{F}_3^t$	$T_1 \bar{F}_3^p$	$T_1 \bar{F}_3^a$
2111	8	$T_C T_3 B_3^a$	$T_C T_3 B_3^p$	$T_C T_1 B_3^a$	$T_S T_C T_3 \bar{B}_3^a$	$T_S T_C T_3 \bar{B}_3^p$	$T_S T_C T_3 \bar{B}_3^t$	$T_S T_C T_1 \bar{B}_3^t$	$T_S T_C T_1 \bar{B}_3^p$	$T_S T_C T_1 \bar{B}_3^a$
2112	9	$T_2 D_3^t$	$T_2 D_3^p$	$T_2 D_3^a$	$T_S T_3 \bar{E}_3^a$	$T_S T_3 \bar{E}_3^p$	$T_S T_3 \bar{E}_3^t$	$T_S T_2 \bar{D}_3^t$	$T_S T_2 \bar{D}_3^p$	$T_S T_2 \bar{D}_3^a$
2121	10	$T_3 D_3^a$	$T_3 D_3^p$	$T_3 D_3^t$	$T_S T_3 \bar{D}_3^a$	$T_S T_3 \bar{D}_3^p$	$T_S T_3 \bar{D}_3^t$	$T_S T_2 \bar{E}_3^t$	$T_S T_2 \bar{E}_3^p$	$T_S T_2 \bar{E}_3^a$
2122	11	$T_3 F_3^a$	$T_3 F_3^p$	$T_2 F_3^a$	$T_S T_3 \bar{F}_3^a$	$T_S T_3 \bar{F}_3^p$	$T_S T_3 \bar{F}_3^t$	$T_S T_2 \bar{F}_3^t$	$T_S T_2 \bar{F}_3^p$	$T_S T_2 \bar{F}_3^a$
2211	12	$T_C C_3^a$	$T_C C_3^p$	$T_1 T_C C_3^a$	$T_C \bar{C}_3^a$	$T_C \bar{C}_3^p$	$T_C \bar{C}_3^t$	$T_1 T_C \bar{C}_3^t$	$T_1 T_C \bar{C}_3^p$	$T_1 T_C \bar{C}_3^a$
2212	13	$T_C F_3^a$	$T_C F_3^p$	$T_2 T_C F_3^a$	$T_C \bar{F}_3^a$	$T_C \bar{F}_3^p$	$T_C \bar{F}_3^t$	$T_S T_C T_1 \bar{F}_3^t$	$T_S T_C T_1 \bar{F}_3^p$	$T_S T_C T_1 \bar{F}_3^a$
2221	14	$T_C T_3 F_3^a$	$T_C T_3 F_3^p$	$T_C T_2 F_3^a$	$T_S T_C T_3 \bar{F}_3^a$	$T_S T_C T_3 \bar{F}_3^p$	$T_S T_C T_3 \bar{F}_3^t$	$T_1 T_C \bar{F}_3^t$	$T_1 T_C \bar{F}_3^p$	$T_1 T_C \bar{F}_3^a$
2222	15	0	0	0	0	0	0	0	0	0

Bubbles

Keldysh structure of the bubbles

When considering the derivative of the full vertex, there are contributions coming only from each one of the channels, i.e. $\Gamma^\Lambda = \sum_r \dot{\gamma}_r$ and

$$\dot{\gamma}_r = \Gamma \circ \dot{\Pi}_r \circ \Gamma + \text{multi-loop terms.} \quad (33)$$

What now becomes important is determining the Keldysh indices of the whole equation 33, since automating this step plays a crucial role in the computation of the derivative. First, one has to look at the Keldysh structure of the bubbles $\Pi_r^{\alpha'_3 \alpha'_4 | \alpha_3 \alpha_4}$. To do so, we recur first to the following bubble structures:

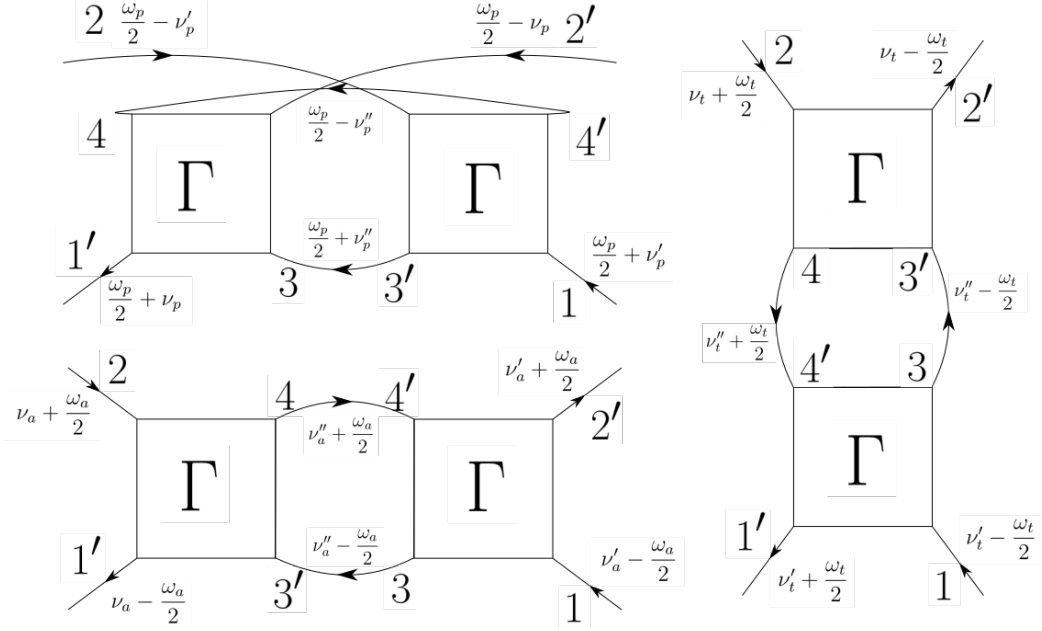


Figure 3: Bubbles in the respective channels

Thus, first of all, we have the frequency arguments that go into both propagators of the bubble. Now, as for the sole bubbles, their Keldysh structure gets simplified thanks to the causality constraint imposed by $G^{1|1} = 0$. Thus, from the possible total of 16 Keldysh components, only 9 are not always zero and, if one chooses the following convention to define the indexation, these are the same for all channels, though the Keldysh structure of these elements not necessarily have to be the same. More precisely, the bubbles in the p-channel have a slightly different structure than the a- and t-channels that, due to their "conjugacy", take exactly the same structure.

For the a-channel, we have

$$\Pi_a^{\alpha'_3 \alpha'_4 | \alpha_3 \alpha_4} = G^{\alpha'_3 | \alpha_3} \left(\nu_a'' - \frac{\omega_a}{2} \right) G^{\alpha'_4 | \alpha_4} \left(\nu_a'' + \frac{\omega_a}{2} \right) \quad (34)$$

$$\begin{pmatrix} 0 & 1 & 2 & 3 \\ 4 & 5 & 6 & 7 \\ 8 & 9 & 10 & 11 \\ 12 & 13 & 14 & 15 \end{pmatrix} = \begin{pmatrix} 1111 & 1112 & 1121 & 1122 \\ 1211 & 1212 & 1221 & 1222 \\ 2111 & 2112 & 2121 & 2122 \\ 2211 & 2212 & 2221 & 2222 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & AA \\ 0 & 0 & AR & AK \\ 0 & RA & 0 & KA \\ RR & RK & KR & KK \end{pmatrix}, \quad (35)$$

for the p-channel

$$\begin{aligned} \Pi_p^{\alpha'_3 \alpha'_4 | \alpha_3 \alpha_4} &= G^{\alpha_3 | \alpha'_3} \left(\frac{\omega_p}{2} + \nu_p'' \right) G^{\alpha'_4 | \alpha_4} \left(\frac{\omega_p}{2} - \nu_p'' \right) \\ \begin{pmatrix} 0 & 1 & 2 & 3 \\ 4 & 5 & 6 & 7 \\ 8 & 9 & 10 & 11 \\ 12 & 13 & 14 & 15 \end{pmatrix} &= \begin{pmatrix} 1111 & 1112 & 1121 & 1122 \\ 1211 & 1212 & 1221 & 1222 \\ 2111 & 2112 & 2121 & 2122 \\ 2211 & 2212 & 2221 & 2222 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & AR \\ 0 & 0 & AA & AK \\ 0 & RR & 0 & KR \\ RA & RK & KA & KK \end{pmatrix}, \end{aligned} \quad (36)$$

and, lastly, for the t-channel, the same as for the a-channel:

$$\begin{aligned} \Pi_t^{\alpha'_3 \alpha'_4 | \alpha_3 \alpha_4} &= G^{\alpha'_3 | \alpha_3} \left(\nu_t'' - \frac{\omega_t}{2} \right) G^{\alpha'_4 | \alpha_4} \left(\nu_t'' + \frac{\omega_t}{2} \right) \\ \begin{pmatrix} 0 & 1 & 2 & 3 \\ 4 & 5 & 6 & 7 \\ 8 & 9 & 10 & 11 \\ 12 & 13 & 14 & 15 \end{pmatrix} &= \begin{pmatrix} 1111 & 1112 & 1121 & 1122 \\ 1211 & 1212 & 1221 & 1222 \\ 2111 & 2112 & 2121 & 2122 \\ 2211 & 2212 & 2221 & 2222 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & AA \\ 0 & 0 & AR & AK \\ 0 & RA & 0 & KA \\ RR & RK & KR & KK \end{pmatrix}. \end{aligned} \quad (37)$$

This is a good point to point out that it's only the internal structure of the bubble the one that is the same for the a- and t-channels. The way the legs are connected and, consequently, the following functions, do differ between both (and, actually, all three) channels. These aforementioned functions are a devise to exploit the zeros in the bubbles *and*, simultaneously, the fact that one only needs to store n_{K_i} Keldysh components per diagrammatic class. These are functions that then, based on the Keldysh index, i_0 , of the left hand side of Eq. (33) (i.e. $i_0 \in \{0, \dots, n_{K_i}\}$) and the Keldysh index, i_2 , of the list of non zero entries of the bubbles (i.e. $i_2 \in \{3, 6, 7, 9, 11, 12, 13, 14, 15\}$) determine precisely the corresponding components i_1 and i_3 for the vertices left and right of the bubble in Eq. (33). Condensed into an equation, the statement is as follows:

$$\dot{\gamma}_r^{i_0}(\omega_r, \nu_r, \nu_r') = \sum_{i_2 \in BK} \int d\nu_r'' \Gamma_1^{i_1^r(i_0, i_2)}(\omega_r, \nu_r, \nu_r'') \dot{\Pi}_r^{i_2}(\nu_1^r, \nu_2^r) \Gamma_3^{i_3^r(i_0, i_2)}(\omega_r, \nu_r'', \nu_r'), \quad (40)$$

where we have emphasized that the frequencies that go into the bubbles depend on the channel, as depicted above, that the indices i_1 and i_3 are functions of the other two indices and the way these are determined are also channel-dependent and $BK = \{3, 6, 7, 9, 11, 12, 13, 14, 15\}$. Writing down the exact form of these functions is a straight-forward task, which requires the conversion of i_0 and i_2 into indices in the $\{1111, \dots, 2222\}$ set. To visualize

the process, just look at the figures drawn above. There, it becomes apparent, that, for $i_0 \rightarrow \alpha'_1 \alpha'_2 | \alpha_1 \alpha_2$ and $i_2 \rightarrow \alpha'_3 \alpha'_4 | \alpha_3 \alpha_4$, the functions are as follows:

$$i_1^a(\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2, \alpha'_3 \alpha'_4 | \alpha_3 \alpha_3) = \alpha'_1 \alpha_4 | \alpha'_3 \alpha_2 \quad (41)$$

$$i_3^a(\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2, \alpha'_3 \alpha'_4 | \alpha_3 \alpha_3) = \alpha_3 \alpha'_2 | \alpha_1 \alpha'_4 \quad (42)$$

$$i_1^p(\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2, \alpha'_3 \alpha'_4 | \alpha_3 \alpha_3) = \alpha'_1 \alpha'_2 | \alpha'_3 \alpha_4 \quad (43)$$

$$i_3^p(\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2, \alpha'_3 \alpha'_4 | \alpha_3 \alpha_3) = \alpha_3 \alpha'_4 | \alpha_1 \alpha_2 \quad (44)$$

$$i_1^t(\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2, \alpha'_3 \alpha'_4 | \alpha_3 \alpha_3) = \alpha_4 \alpha'_2 | \alpha'_3 \alpha_2 \quad (45)$$

$$i_3^t(\alpha'_1 \alpha'_2 | \alpha_1 \alpha_2, \alpha'_3 \alpha'_4 | \alpha_3 \alpha_3) = \alpha'_1 \alpha_3 | \alpha_1 \alpha'_4 \quad (46)$$

$$(47)$$

Structure of the internal bubble - diagrammatic classes' contributions

Now, at this stage, the natural question to ask is how the *diagrammatic* decomposition of Eq. (40) looks like i.e., given that $\dot{\gamma}_r = \sum_i \dot{K}_i^r$, what are the diagrammatic classes that contribute to each of these derivatives? It should by now be clear that not all classes can contribute to all derivatives, since, for instance, diagrams in the K_3 class, by definition, do not fulfill the requirement of having, say, two lines connected to the same vertex on one side. Thus, no terms with K_3 can appear in \dot{K}_1 . The natural thing to do at this stage is, basically, classify by sides and lines connected to one vertex, what classes fulfill the requirements to contribute, at least on one side, to which ones. Thus, one can establish the following categorization:

	L	R
SV	$\Gamma_0, K_1^r, \bar{K}_2^r$	Γ_0, K_1^r, K_2^r
DV	$K_2^r, K_3^r, \gamma_{\bar{r}}$	$\bar{K}_2^r, K_3^r, \gamma_{\bar{r}}$

Here, SV and DV stand for "same vertex" and "different vertices", alluding to whether or not both lines to the right (R) of the left (L) of the vertex connect to the same or to different vertices. Γ_0 is the bare vertex. The inclusion already of $\gamma_{\bar{r}}$ is due to the foreshadowing of the multi-loop iterative loop i.e., we will be interested in feeding back diagrams of channels \bar{r} to calculate higher loop order contributions to channel r . Through this, one then gets the following Contribution Table to \dot{K}_i^r :

	L	R
\dot{K}_1^r	$\Gamma_0, K_1^r, \bar{K}_2^r$	Γ_0, K_1^r, K_2^r
\dot{K}_2^r	$K_2^r, K_3^r, \gamma_{\bar{r}}$	Γ_0, K_1^r, K_2^r
$\dot{\bar{K}}_2^r$	$\Gamma_0, K_1^r, \bar{K}_2^r$	$\bar{K}_2^r, K_3^r, \gamma_{\bar{r}}$
\dot{K}_3^r	$K_2^r, K_3^r, \gamma_{\bar{r}}$	$\bar{K}_2^r, K_3^r, \gamma_{\bar{r}}$

What this table summarizes is the possible terms that can appear on the RHS of Eq. (33), if one decomposes it into the diagrammatic classes.