

Experimental Methods Lecture 16

October 26th, 2020

A final note: multiple stochastic processes

- Assume you have two different stochastic processes x(t) and y(t).
- Their statistics in known if for any N and any two sets of times $\{t_1, t_N\}$ and $\{t_1', t_M'\}$ we know the joint probability density $f_{x(t_1), y(t_1'), x(t_2), y(t_2') \dots x(t_N), y(t_M')} \left(\chi_1, \psi_1, \chi_2, \psi_2 \dots \chi_N, \psi_M\right)$

• One can define in addition to moments for each separate process, the joint moments: cross-correlation

$$R_{x,y}(t,t') = \langle x(t)y(t')\rangle = R_{y,x}(t',t)$$

Cross covariance

$$C_{x,y}(t,t') = R_{x,y}(t,t') - \eta_x(t)\eta_y(t)$$

• The autocorrelation is just a special case when y(t)=x(t)

$$R_{x,x}(t,t') = \langle x(t)x(t')\rangle = R_{x,x}(t',t)$$

• Two processes are joint stationary if all their density functions are invariant under a time origin shift. They are normal if joint densities are normal etc.

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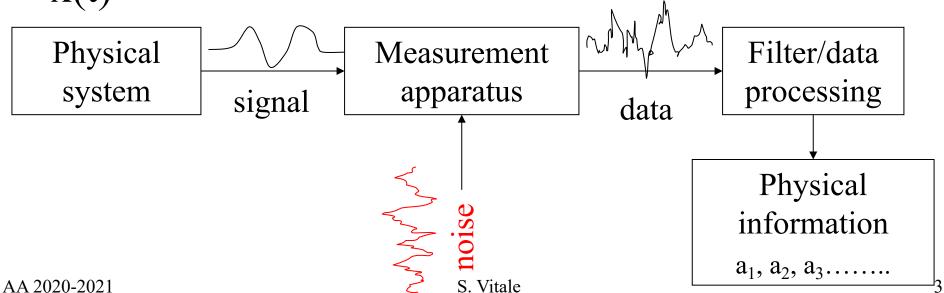
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Noise

Noise in physical experiments is described as a random signal x(t):

Independent (ensemble) repetitions of the same experiment produce different functions of time x(t)



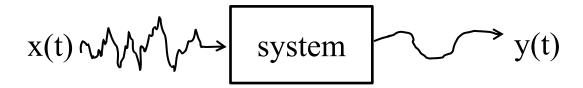


Noise and systems

- We now develop a model for the noise in physical apparatuses where noise is treated as consisting of stochastic signals.
- The signals are fed to systems and produce stochastic signals at output
- For linear systems, linear response theory allows to calculate the statistical properties of output given those at input.



• For a single execution of an experiment, a stochastic process is just an ordinary signal



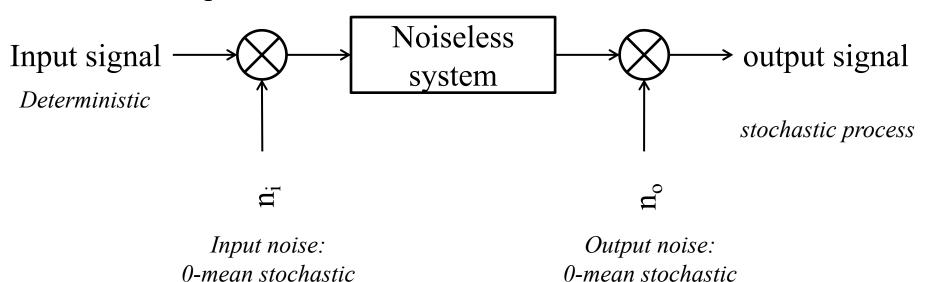
- The output is a signal
- The input is a random signal, then the output is also random
- The output is a stochastic process



The conceptual scheme for an apparatus in the presence of noise

- Disturbances in physical systems are successfully described as stochastic processes acting at input and at output of an intrinsically noiseless system:
- Notice: this implies that noise is independent of signal and signal levels.
- Not true for parametric noise

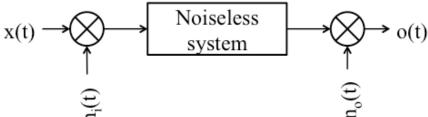
process



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Stochastic processes and linear systems the interior of the control of the contro



• Suppose the system is linear. The output is

$$o(t) = n_o(t) + \int_{-\infty}^{\infty} h(t,t') \left[x(t') + n_i(t') \right] dt'$$

• As the system obeys the principle of superposition:

$$o(t) = n_o(t) + \int_{-\infty}^{\infty} h(t,t')x(t')dt' + \int_{-\infty}^{\infty} h(t,t')n_i(t')dt'$$

• Then the output can be written as $o(t) = y(t) + \tilde{n}(t)$ where

$$\underbrace{y(t)}_{\text{Signal}} = \int_{-\infty}^{\infty} h(t,t')x(t')dt' \qquad \underbrace{\tilde{n}(t)}_{\text{Noise}} = n_o(t) + \int_{-\infty}^{\infty} h(t,t')n_i(t')dt'$$

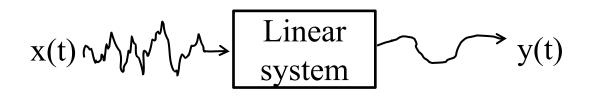
• The response to noise is independent of the presence and the nature of the signal. Noise properties can be treated independently of signals. A reasonably accurate model for most systems.(not for parametric noise)

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• Suppose a stochastic process x(t) is fed to the input of a linear system whose impulse response is h(t,t')



• The output is

$$y(t) = \int_{-\infty}^{\infty} h(t,t')x(t')dt'$$

- One could in principle derive the joint probability densities of any order for y(t)
- An easier task is to derive moments

A parenthesis on mean value and linear operators

• Assume you have N random variables $x_1, x_2, ... x_N$ with a joint distribution $f_{x_1,x_2,...x_N}(\chi_1,\chi_2...\chi_N)$

• Take now M functions of these random variables $y_1=g_1(x_1, x_2, ...x_N)$),... $y_M = g_M(x_1, x_2, ...x_N)$, and calculate the mean value of some linear

combination of them
$$\langle c_1 y_1 + c_2 y_2 + ... + c_N y_N \rangle =$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left[c_1 g_1 \left(\chi_1, \chi_2...\chi_N \right) + c_2 g_1 \left(\chi_1, \chi_2...\chi_N \right) + ... + c_N g_1 \left(\chi_1, \chi_2...\chi_N \right) \right] \times$$

$$\times f_{x_1, x_2, ...x_N} \left(\chi_1, \chi_2...\chi_N \right) d\chi_1 d\chi_2...d\chi_N$$

$$= c_1 \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} g_1 \left(\chi_1, \chi_2...\chi_N \right) f_{x_1, x_2, ...x_N} \left(\chi_1, \chi_2...\chi_N \right) d\chi_1 d\chi_2...d\chi_N + ...$$

 $+c_{N}\int_{-\infty}^{\infty}\int_{-\infty}^{\infty}....\int_{-\infty}^{\infty}g_{N}(\chi_{1},\chi_{2}..\chi_{N})f_{\chi_{1},\chi_{2}...\chi_{N}}(\chi_{1},\chi_{2}...\chi_{N})d\chi_{1}d\chi_{2}..d\chi_{N}+...$ • Then

Then
$$\langle c_1 y_1 + c_2 y_2 + ... + c_N y_N \rangle == c_1 \langle y_1 \rangle + c_2 \langle y_2 \rangle + ... + c_N \langle y_N \rangle$$

In particular

 $\left\langle c_{_{1}}x_{_{1}} + c_{_{2}}x_{_{2}} + ... + c_{_{N}}x_{_{N}} \right\rangle == c_{_{1}}\left\langle x_{_{1}}\right\rangle + c_{_{2}}\left\langle x_{_{2}}\right\rangle + ... + c_{_{N}}\left\langle x_{_{N}}\right\rangle$



Let's go back to our linear system

$$y(t) = \int_{-\infty}^{\infty} h(t,t')x(t')dt' \qquad x(t) \text{ (t) inear system} \qquad y(t)$$

• And calculate the mean value of the output. Let's rewrite

$$y(t) = \lim_{\substack{\Delta t \to 0 \\ N \to \infty}} \sum_{k=-N}^{N} h(t, k\Delta t) x(k\Delta t) \Delta t$$

• Then y(t) is a linear combination of the $x(k\Delta t)$ and

$$\langle y(t) \rangle = \lim_{\substack{\Delta t \to 0 \\ N}} \sum_{k=-N}^{N} h(t, k\Delta t) \langle x(k\Delta t) \rangle \Delta t$$

Taking the limits

$$\langle y(t)\rangle = \int_{-\infty}^{\infty} h(t,t') \langle x(t')\rangle dt'$$

Mean value and integration commute!

$$y(t) = \int_{-\infty}^{\infty} h(t,t')x(t')dt' \qquad x(t) \text{ (t) Linear system} \qquad y(t)$$

Autocorrelation. Use always

$$y(t) = \lim_{\substack{\Delta t \to 0 \\ N \to \infty}} \sum_{k=-N}^{N} h(t, k\Delta t) x(k\Delta t) \Delta t$$

- Then $\left\langle y(t)y(t')\right\rangle = \lim_{\substack{\Delta t \to 0 \\ N \to \infty}} \left\langle \sum_{k,j=-N}^{N} h(t,k\Delta t)h(t',j\Delta t)x(k\Delta t)x(j\Delta t)\Delta t^2 \right\rangle$ Using superposition

$$\left\langle y\!\left(t\right)\!y\!\left(t'\right)\right\rangle\!=\!\lim_{\Delta t\to 0 \atop N\to\infty}\sum_{k,j=-N}^{N}h\!\left(t,k\Delta t\right)\!h\!\left(t',j\Delta t\right)\!\!\Delta t^{2}\left\langle x\!\left(k\Delta t\right)\!x\!\left(j\Delta t\right)\right\rangle$$

Taking the limits

$$\langle y(t)y(t')\rangle = \int_{-\infty}^{\infty} h(t,t'')h(t',t''')\langle x(t'')x(t''')\rangle dt''dt'''$$

Or

$$R_{yy}(t,t') = \int_{-\infty}^{\infty} h(t,t'')h(t',t''')R_{xx}(t'',t''')dt''dt'''$$

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$$y(t) = \int_{-\infty}^{\infty} h(t,t')x(t')dt' \qquad x(t) \text{ (t) Linear system} \qquad y(t)$$

Holds for higher moments

$$\langle y(t)y(t')y(t'')\rangle =$$

$$= \int_{-\infty}^{\infty} h(t,t''')h(t',t'''')h(t'',t'''')\langle x(t''')x(t'''')x(t''''')\rangle dt'''dt'''''$$

- Etc.
- It's a general rule for all linear operators (rather obvious if you allow for impulse response that are distributions)

$$\begin{split} & \left\langle L_{t} \left[i(t) \right] \right\rangle = L_{t} \left\langle i(t) \right\rangle \\ & \left\langle L_{1,t} \left[i(t) \right] L_{2,t'} \left[i(t') \right] \right\rangle = L_{1t} L_{2t'} \left\langle i(t) i(t') \right\rangle \\ & \left\langle L_{1,t} \left[i(t) \right] L_{2,t'} \left[i(t') \right] L_{3,t''} \left[i(t'') \right] \right\rangle = L_{1t} L_{2t'} L_{3t''} \left\langle i(t) i(t') i(t'') \right\rangle \end{split}$$

Additional examples: Derivatives

$$y(t) = \frac{dx(t)}{dt} = \lim_{dt \to 0} \frac{x(t+dt) - x(t)}{dt}$$

• The mean value $\lim_{dt\to 0} \left\langle \frac{x(t+dt)-x(t)}{dt} \right\rangle = \lim_{dt\to 0} \frac{\left\langle x(t+dt)\right\rangle - \left\langle x(t)\right\rangle}{dt}$

- Then $\left\langle \frac{\mathrm{dx}}{\mathrm{dt}} \right\rangle = \frac{\mathrm{d} \left\langle x(t) \right\rangle}{\mathrm{dt}}$
- Higher order $\left\langle \frac{dx}{dt} \middle| x(t') \right\rangle = \lim_{dt \to 0} \left\langle \frac{x(t+dt)-x(t)}{dt} x(t') \right\rangle$

$$= \lim_{dt \to 0} \frac{\left\langle x(t+dt)x(t')\right\rangle - \left\langle x(t)x(t')\right\rangle}{dt} = \frac{\partial R_{x,x}(t,t')}{\partial t}$$

• Same way $\left\langle \frac{dx}{dt} \middle|_{t} \frac{dx}{dt} \middle|_{t'} \right\rangle = \frac{\partial^{2} R_{x,x}(t,t')}{\partial t \partial t'}$



Example: shot noise

Shot noise is the derivative of the Poisson process

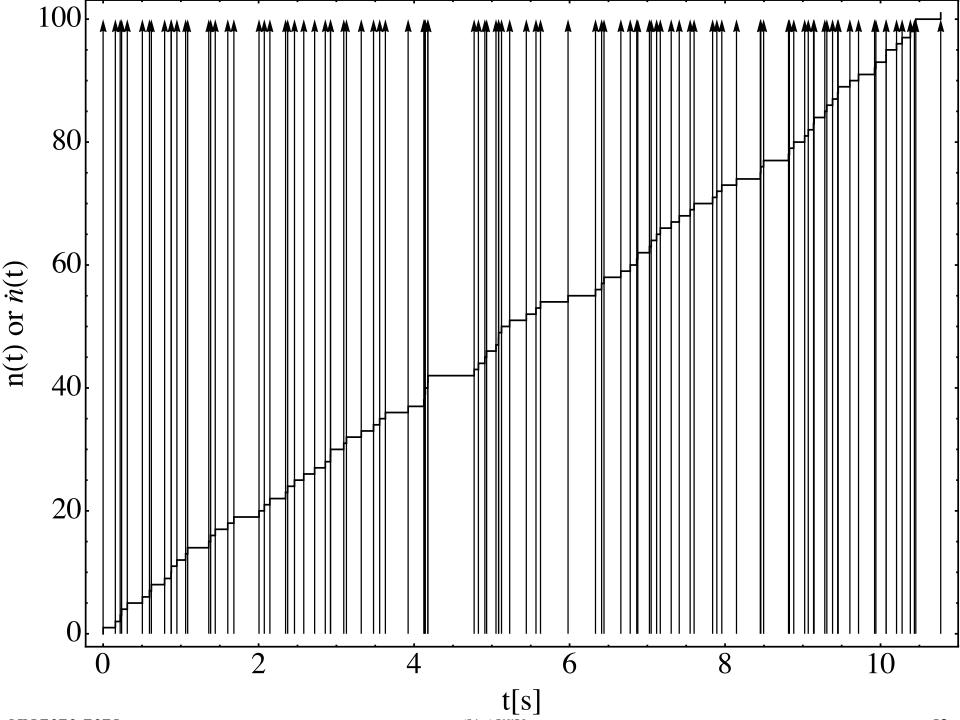
$$\dot{n}(t) = \frac{dn}{dt} \bigg|_{t}$$

- For photon beams, energy is proportional to a Poisson process: $E(t)=\hbar\omega n(t)$, power is proportional to shot noise $P(t)=\hbar\omega \dot{n}(t)$
- For random charge flow, I(t)=e $\dot{n}(t)$ is the current
- From the definition of shot noise

$$\dot{\mathbf{n}}(t) = \frac{\mathrm{d}}{\mathrm{d}t} \left[\sum_{k=0}^{\infty} \Theta(t - t_k) \right] = \sum_{k=0}^{\infty} \delta(t - t_k)$$

Thus shot noise consists of a series of impulses arriving at random times (see next page with $\lambda=10 \text{ s}^{-1}$)

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Mean value and autocorrelation of shot noise

- Mean value of Poisson noise $\langle n(t) \rangle = \lambda t$
- Mean value of shot noise is equal to derivative of mean value:

$$\left\langle \dot{\mathbf{n}}(t)\right\rangle = d\left\langle \mathbf{n}(t)\right\rangle / dt = \lambda$$
Autocorrelation of Poisson noise

- $\langle n(t)n(t')\rangle = \lambda^2 tt' + \lambda \left[t\Theta(t'-t) + t'\Theta(t-t')\right]$
- Using derivative rules, the autocorrelation of shot noise is:

$$\left\langle \dot{\mathbf{n}}(t)\dot{\mathbf{n}}(t')\right\rangle = \frac{\partial^{2}\left\langle \mathbf{n}(t)\mathbf{n}(t')\right\rangle}{\partial t \partial t'} = \frac{\partial}{\partial t'} \left\{ \lambda^{2}t' + \lambda \left[\Theta(t'-t) - t\delta(t-t') + t'\delta(t-t')\right] \right\}$$

- $= \frac{\partial}{\partial t'} \left\{ \lambda^2 t' + \lambda \left[\Theta(t' t) + (t' t) \delta(t' t) \right] \right\}$
- As $(t-t')\delta(t-t')=0 \Rightarrow = \frac{\partial}{\partial t'} \{\lambda^2 t' + \lambda \Theta(t'-t)\} = \lambda^2 + \lambda \delta(t-t')$
 - Finally the auto-covariance is $\langle \dot{n}(t)\dot{n}(t')\rangle \langle \dot{n}(t)\rangle \langle \dot{n}(t')\rangle = \lambda\delta(t-t')$

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Mean value and autocorrelation of shot noise

- In summary
- Mean value of shot noise is a constant: the "rate"

$$\eta_{\dot{n}}(t) = \langle \dot{n}(t) \rangle = \lambda$$

Autocorrelation of Poisson noise only depends on delay t-t'

$$R_{\dot{n}(t)\dot{n}(t')} = \langle \dot{n}(t)\dot{n}(t')\rangle = \lambda^2 + \lambda\delta(t-t')$$

And so does the auto-covariance:

$$C_{\dot{\mathbf{n}}(t)\dot{\mathbf{n}}(t')} = \langle \dot{\mathbf{n}}(t)\dot{\mathbf{n}}(t')\rangle - \langle \dot{\mathbf{n}}(t)\rangle \langle \dot{\mathbf{n}}(t')\rangle = \lambda\delta(t-t')$$

- This only holds for t, $t' \ge 0$. However the time origin can be moved $\rightarrow -\infty$, and then shot noise becomes at least wide-sense stationary.
- It is easy to convince yourself that it is also stationary

Stationary noise and linear stationary systems

- Let's consider the case where:
 - the input to a linear system x(t) is a stationary stochastic process
 - the linear system is time-invariant.
- The output signal y(t) is

$$y(t) = \int_{-\infty}^{\infty} h(t')x(t-t')dt'$$

• Then for any N

$$\langle y(t_1 + T)y(t_2 + T)...y(t_N + T)\rangle =$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} h(t')h(t'')h(t'') \langle x(t_1 + T - t')x(t_2 + T - t'')x(t_3 + T - t''') \rangle dt' dt'' dt'''$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} h(t')h(t'')h(t'') \langle x(t_1 - t')x(t_2 - t'')x(t_3 - t''') \rangle dt'dt''dt'''$$

$$= \langle y(t_1)y(t_2)....y(t_N) \rangle$$

That is, y(t) is stationary.

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Stationary noise and linear stationary systems

- A stationary noise at the input of a linear time-invariant system
 - 1. Mean value

$$\left\langle y\!\left(t\right)\right\rangle\!=\int_{-\infty}^{\infty}h\!\left(t'\right)\!\eta_{x}\,dt'\!=\eta_{x}\int_{-\infty}^{\infty}h\!\left(t'\right)\!dt'\!=Constant\equiv\eta_{y}$$

2. Input-output cross correlation

$$R_{y,x}(t,t+\Delta t) = \langle y(t)x(t+\Delta t)\rangle = \int_{-\infty}^{\infty} h(t')\langle x(t-t')x(t+\Delta t)\rangle dt'$$
Using time-invariance of x

$$R_{y,x}(\Delta t) = \int_{-\infty}^{\infty} h(t') R_{x,x}(\Delta t + t') dt'$$

3. Output auto-correlation

$$R_{y,y}(t,t+\Delta t) = \langle y(t)y(t+\Delta t)\rangle =$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dt' dt'' h(t') h(t'') \langle x(t-t')x(t+\Delta t-t'') \rangle$$

That is

$$R_{y,y}(\Delta t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dt' dt'' h(t') h(t'') R_{x,x}(\Delta t + t' - t'')$$



Examples

• Shot noise into a low-pass filter. A good model for charge detection.

• White noise into a low-pass filter. The basic colored noise.

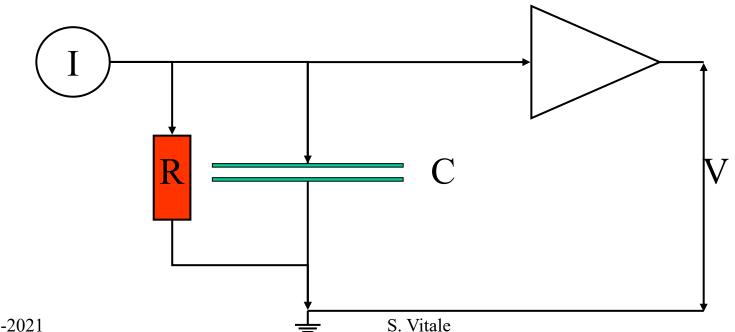
Random Charging



- Assume that a flow of particles, each of charge *e*, hits a detector represented by capacitor of capacitance C.
- Particles arrive as Poisson points with a rate λ . They are represented by a current generator generating a current

$$I(t) = \sum_{i=-\infty}^{\infty} e\delta(t - t_i)$$

• Losses in the capacitor are represented by a parallel resistor of Resistance R. Voltage V across C is read out by an ideal amplifier



Random Charging



- The circuit responds to the stochastic signal as to any other signal
- The impulse response of the circuit can be easily obtained from

$$C(dV(t)/dt) = -(V(t)/R) + I(t)$$

• Fourier transform gives the frequency response

$$V(\omega) = (I(\omega)/C)(i\omega + 1/RC)^{-1} \equiv h(\omega)I(\omega)$$

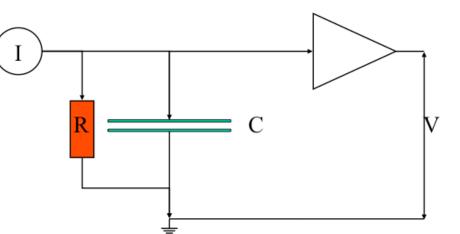
• The inverse Fourier transform of which is

$$h(t) = \Theta(t)e^{-t/RC}/C$$
As the current is just formed by impulses
$$I(t) = \sum_{i=1}^{\infty} e\delta(t - t_i)$$

• and the system is linear, then

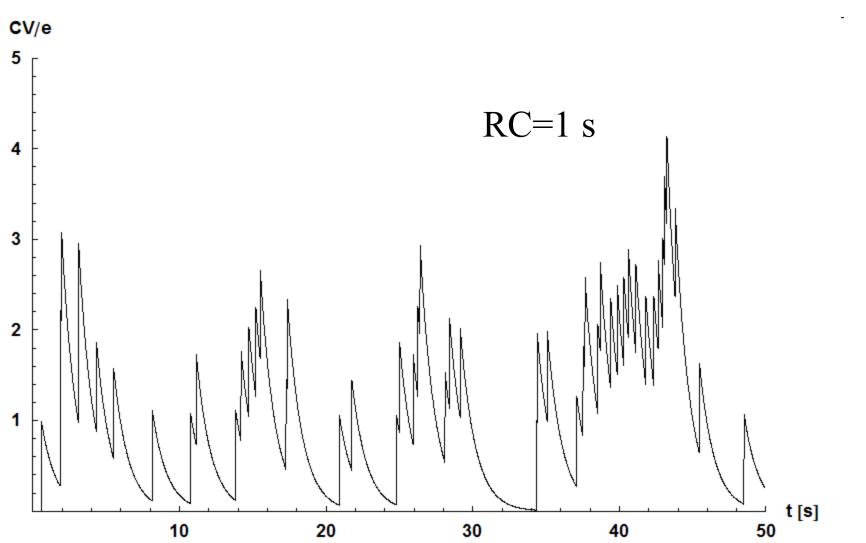
$$V(t) = \frac{e}{C} \sum_{i=-\infty}^{\infty} \Theta(t - t_i) e^{-(t - t_i)/RC}$$

See next page for an example





Numerical example for charged particle detection



Random Charging



- Voltage mean value $\langle V(t) \rangle = \int_{-\infty}^{\infty} \Theta(t') (e^{-t'/RC}/C) \langle I(t-t') \rangle dt'$ • As $\langle I(t) \rangle = e \langle \dot{n}(t) \rangle = e \lambda$
- We obtain (rather obviously) $\langle V(t) \rangle = \frac{e\lambda}{C} \int_0^\infty e^{-t'/RC} dt' = Re\lambda$
- Autocorrelation
 - $R_{VV}(\Delta t) = \int_{0}^{\infty} \int_{0}^{\infty} \left(e^{-t'/RC}/C \right) \left(e^{-t''/RC}/C \right) R_{II}(\Delta t + t' t'') dt' dt''$
- As $R_{II}(\Delta t + t' t'') = e^2 R_{nn}(\Delta t + t' t'') = e^2 \lambda^2 + e^2 \lambda \delta(\Delta t + t' t'')$
- $R_{VV}(\Delta t) = e^2 \lambda^2 \left[\int_0^\infty \frac{e^{-t'/RC}}{C} dt' \right]^2 + \frac{e^2 \lambda}{C^2} \int_0^\infty \int_0^\infty e^{-t'/RC} e^{-t''/RC} \delta(\Delta t + t' t'') dt' dt'' \right]$
- That is $R_{VV}(\Delta t) = (Re\lambda)^{2} + \frac{e^{2}\lambda}{C^{2}} \int_{\text{Max}[-\Delta t,0]}^{\infty} e^{-t'/RC} e^{-(\Delta t+t')/RC} dt'$ $= (Re\lambda)^{2} + \frac{Re^{2}\lambda}{C} e^{-\Delta t/RC} \left[\Theta(\Delta t) \int_{0}^{\infty} e^{-2x} dx + \Theta(-\Delta t) \int_{-\Delta t/RC}^{\infty} e^{-2x} dx \right] =$

Random Charging



In summary

- Mean value
$$\eta_{v} = Re\lambda$$

- Autocorrelation
$$R_{VV}(\Delta t) = (Re\lambda)^2 + \frac{Re^2\lambda}{2C}e^{-|\Delta t|/RC}$$

Then auto-covariance is

$$C_{VV}(\Delta t) = R_{VV}(\Delta t) - \eta_{V}^{2} = \frac{Re^{2}\lambda}{2C}e^{-|\Delta t|/RC}$$

Thus voltage is the sum of a static dc voltage and of a noisy part $\tilde{V}(t)$

$$V(t) = Re\lambda + V(t)$$

$$Re^{2\lambda} = \frac{|\Delta t|}{2} \left(\frac{\Delta t}{2} \right) - \frac{|\Delta t|}{2} \left(\frac{\Delta t}{2} \right) = \frac{|\Delta$$

- $V(t) = Re\lambda + \tilde{V}(t)$ With autocorrelation $R_{\tilde{V}\tilde{V}}(\Delta t) = C_{VV}(\Delta t) = \frac{Re^2\lambda}{2C}e^{-|\Delta t|/RC}$
- Notice that variance of both V(t) and $\tilde{V}(t)$ is:

$$\sigma_{V}^{2} = \sigma_{\tilde{V}}^{2} = R_{\tilde{V}\tilde{V}}(0) = C_{VV}(0) = Re^{2}\lambda/2C$$
The then
$$\sigma_{V}^{2}/\eta_{V} = \frac{Re^{2}\lambda}{2CRe\lambda} = \frac{e}{2C}$$

- The then
- A well known way of measuring electron charge independently of λ !