# **Experimental Methods Exam February 2021**

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MAT. 223450

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I start by evaluating the true value for  $\beta$  and k. These are found by solving:

$$(2\pi
u)^2=\omega_0^2=rac{k}{m}\longrightarrow k=m(2\pi
u)^2$$

$$Q^2=rac{m^2}{eta^2}rac{k}{m}=rac{mk}{eta^2}\longrightarrow eta=rac{\sqrt{mk}}{Q}$$

In [1]: #

# packages used

import scipy.constants # https://docs.scipy.org/doc/scipy/reference/constants.html
from scipy.integrate import quad
import numpy as np

```
import matplotlib.pyplot as plt
from IPython.display import display, Math
# Given parameters
kb = scipy.constants.Boltzmann # J/K
h = scipy.constants.h
c = scipy.constants.c
Nu = 1e9 \# Hz
m = 0.5e-15 \# Kq
Cs = 5e-15 \# Faraday
Im = 100e-6 \# Watt
Lambda = 1e-6 \# m
Rho = 0.1
0 = 10000
Cp = 60e-15 \# Faraday
d = 5e-12 \# m/V
Iref = 100e-3 # Watt
T = 4 \# K
# Extra parameters
Omega0 = 2*np.pi*Nu
k = m*Omega0**2
Beta = np.sqrt(m*k)/Q
# Show value
display(Math(r"k = m (2\pi nu)^2="+"{:.5}".format(k)+r"\ Hz^2\ kg"))
display(Math(r"\beta = \frac{m k}{Q} = + (...5)".format(Beta)+r"\ kg\ Hz"))
```

$$k = m(2\pi\nu)^2 = 1.9739e + 04~Hz^2~kg$$

$$eta = rac{\sqrt{mk}}{Q} = 3.1416e - 10~kg~Hz$$

## **Question 1**

The transfer function between  $V_s$  and x can be obtained by solving the coupled equations given and by noticing that the relation between the input voltage  $V_s$  and the current is (I'm using  $\partial$  instead of d to avoid confusion with the constant):

$$I=C_srac{\partial (V_s-V)}{\partial t}\longrightarrow I(\omega)=i\omega C_s(V_s-V)$$

This relation with the other two given leads to:

$$f(x(\omega)) = rac{C_s d}{C_p - C_S} V_s(\omega) + ig(rac{d^2}{C_p + C_s} - rac{1}{k}ig) F(\omega) \simeq rac{C_s d}{C_p + C_S} V_s - rac{1}{k} F(\omega)$$

Where I dropped the therm  $\frac{d^2}{C_p+C_s}$  as it's a lot smaller than  $\frac{1}{k}$ .

```
In [2]: # Here unit of measurements don't seem to be right, but I'm pretty sure the circuit # analysis is correct. display(Math(r"\frac{d^2}{C_p - C_s}="+"{:.5}".format(d**2 / (Cp - Cs))+ r"\ \frac{m^2}{V^2 F}")) display(Math(r"\frac{1}{k} ="+"{:.5}".format(1/k)+r"\ Hz^{-2}\ kg^{-1}"))
```

$$rac{d^2}{C_p-C_s} = 4.5455e-10 \ rac{m^2}{V^2 F}$$

$$rac{1}{k} = 5.0661e - 05~Hz^{-2}~kg^{-1}$$

This formula can be further expanded by considering the relation between the force and the displacement (harmonic oscillator):

$$\ddot{x} + rac{\omega_0}{Q}\dot{x} + \omega_0^2 x = rac{F}{m} \longrightarrow F(\omega) = (-\omega^2 + rac{i\omega_0}{Q}\omega + \omega_0^2)m \; x(\omega)$$

This implies

$$egin{split} x(\omega)(1+rac{1}{k}(\omega^2-rac{i\omega_0}{Q}\omega-\omega_0^2))&=rac{C_sd}{C_p+C_S}V_s \ h(\omega)&=rac{V_s(\omega)}{x(\omega)}&=rac{C_p+C_S}{C_sd}+rac{C_p+C_s}{C_sdk}(\omega^2-rac{i\omega_0}{Q}\omega-\omega_0^2) \end{split}$$

Which is the transfer function from  $X(\omega)$  to  $V_s(\omega)$ . I can obtain the transfer function from  $V_s(\omega)$  to  $X(\omega)$  by taking the reciprocal. It's possible to lighten the notation by choosing:

$$rac{1}{\omega_1^2} = rac{C_p + C_s}{C_s dk}$$
  $rac{1}{\omega_2} = rac{(C_p + C_s)\omega_0}{QC_s dk}$ 

$$rac{1}{\omega_3} = rac{C_p + C_S}{C_s d} - rac{C_p + C_s}{C_s dk} \omega_0^2$$

In [3]: Omega12 = Cs\*d\*k/(Cp + Cs)
 Omega2 = Q\*Cs\*d\*k/( (Cp+Cs)\*Omega0)
 Omega3 = ( (Cp+Cs)/(Cs\*d) - (Cp+Cs)/(Cs\*d\*k)\*Omega0\*\*2 )\*\*-1

 display(Math(r"\frac{1}{ \omega\_1^2} ="+"{:.5}".format(1/Omega12) ))
 display(Math(r"\frac{1}{ \omega\_2} ="+"{:.5}".format(1/Omega2) ))
 display(Math(r"\frac{1}{ \omega\_3} ="+"{:.5}".format(1/Omega3) ))

$$rac{1}{\omega_1^2} = 1.3172e + 08$$
  $rac{1}{\omega_2} = 8.2761e + 13$   $rac{1}{\omega_3} = -5.2e + 27$ 

So that:

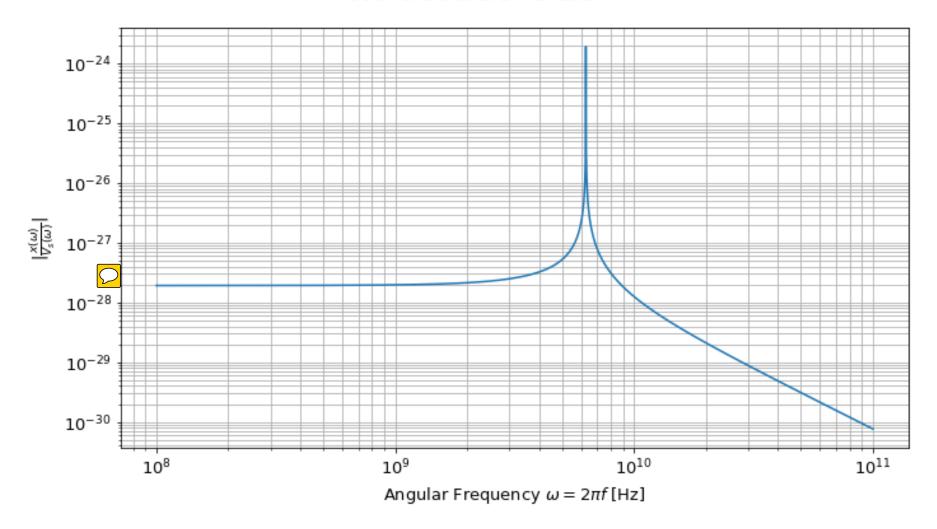
$$h(\omega) = rac{\omega^2}{\omega_1^2} - irac{\omega}{\omega_2} + rac{1}{\omega^3}$$

$$rac{x(\omega)}{V_s(\omega)} = rac{1}{rac{\omega^2}{\omega_1^2} - irac{\omega}{\omega_2} + rac{1}{\omega^3}}$$

It's mandatory to highlight the high magnitude of  $\frac{1}{\omega_3}$  which dominates the frequency response up to GHz.

```
def transfer function(omega):
In [4]:
             return 1/ ( np.power(omega,2)/Omega12 - 1j*omega/Omega2 + 1/Omega3 )
         omega = np.logspace(8, 11, num=int(1e6), endpoint=True, base=10)
         y = np.abs( transfer function(omega) )
         plt.rcParams.update({'font.size': 13})
         fig, axis = plt.subplots(1,1, figsize=(10, 6))
         fig.suptitle("Plot of the transfer function")
         axis.set xscale("log")
         axis.set yscale("log")
         axis.plot(omega, y)
         axis.set xlabel(r"Angular Frequency $\omega=2\pi f$ [Hz]")
         axis.set_ylabel(r"$|\frac{x(\omega)}{V_s(\omega)}|$")
         axis.grid(True, which="both")
         plt.tight_layout()
         plt.show()
```

#### Plot of the transfer function



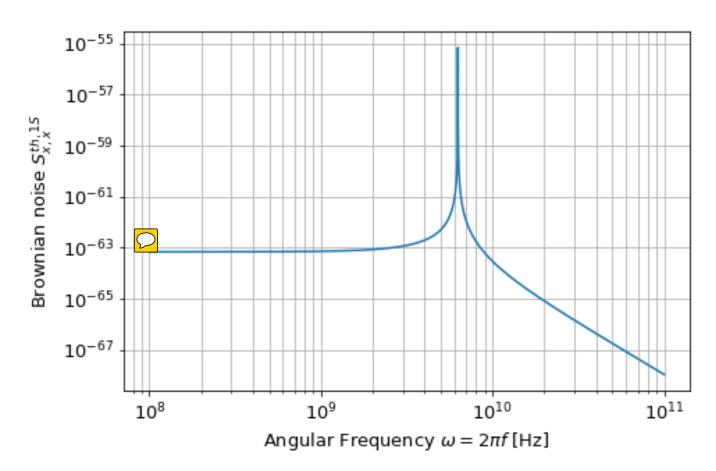
## Question 2

As the mass is very small it's subjected to the effect of the Brownian noise. Its PSD (one sided) can be evaluated via the fluctuation-dissipation theorem:

$$S_{x,x}^{th,1S} = rac{S_{V_s,V_s}^{th,1S}}{\left|h(\omega)
ight|^2} = rac{1}{\left|h(\omega)
ight|^2} 4k_b T rac{\mathcal{I}m[h(\omega)]}{\omega} = rac{4k_b T}{\omega_2 \left|h(\omega)
ight|^2} = rac{4k_b T}{\omega_2} rac{1}{(rac{\omega^2}{\omega_1^2} + rac{1}{\omega^3})^2 + (rac{\omega}{\omega_2})^2}$$

```
def brownian noise(omega):
In [6]:
             return 4*kb*T/Omega2 * 1/( np.power((np.power(omega, 2)/Omega12 + 1/Omega3),2 )
                                       + np.power(omega/Omega2,2))
         omega = np.logspace(8, 11, num=int(1e6), endpoint=True, base=10)
         y = brownian noise(omega)
         plt.rcParams.update({'font.size': 13})
         fig, axis = plt.subplots(1,1, figsize=(7, 5))
         fig.suptitle("Plot of brownian noise")
         axis.set xscale("log")
         axis.set yscale("log")
         axis.plot(omega, y)
         axis.set_xlabel(r"Angular Frequency $\omega=2\pi f$ [Hz]")
         axis.set ylabel(r"Brownian noise $S {x,x}^{th, 1S}$")
         axis.grid(True, which="both")
         plt.tight layout()
         plt.show()
```

#### Plot of brownian noise



## **Question 3**

The noise due to the intrinsic fluctuation of the beam is the shot noise associated to the arrive of the photons itself. The shot noise for the two beams is therefore proportional to their intensity:

$$S_{I_1,I_1}^{1S}(\omega) = 2rac{h\lambda}{c}I_1 = rac{h\lambda}{c}(I_{ref} + I_m - \sqrt{I_mI_ref})cos(\Delta\phi)$$

$$S_{I_2,I_2}^{1S}(\omega) = 2rac{h\lambda}{c}I_2 = rac{h\lambda}{c}(I_{ref} + I_m + \sqrt{I_mI_ref})cos(\Delta\phi)$$

I'm asked for the readout noise which is proportional to  $cos(\Delta\phi)$ . The assignment says that the phase is freely adjustable withing  $2\pi$  so I resolve to set it so that  $cos(\Delta\phi) \to sin(\Delta\phi) \simeq \Delta\phi$ . This is done in order to linearize the expression. This angle is therefore proportional to the difference between the intensities:

$$I_1-I_2=\sqrt{I_mI_{ref}}~\Delta\phi$$

This means that it's PSD is proportional to the sum of the two single PSD:

$$S^{1S}_{\Delta\phi,\Delta\phi}(\omega) = rac{S^{1S}_{I_1,I_1}(\omega) + S^{1S}_{I_2,I_2}(\omega)}{\sqrt{I_m I_{ref}}^2} = rac{2h\lambda}{c}rac{I_{ref} + I_m}{I_{ref}I_m}$$

The relation between the angle and the displacement  $\delta x$  is:

$$\delta x = -rac{1}{2}\delta l = -rac{1}{2}\Delta\phi\lambda$$

The beam PSD can then be expressed as displacement PSD:

$$S_{x,x}^{beam,1S}(\omega) = rac{\lambda^2}{4} S_{\Delta\phi,\Delta\phi}(\omega) = rac{h\lambda^4}{2c^2} rac{I_{ref} + I_m}{I_{ref}I_m}$$



This value is definitely negligible if compared to the Brownian noise which has magnitude between at  $10^{-64}$  and  $10^{-55}$  around resonance.

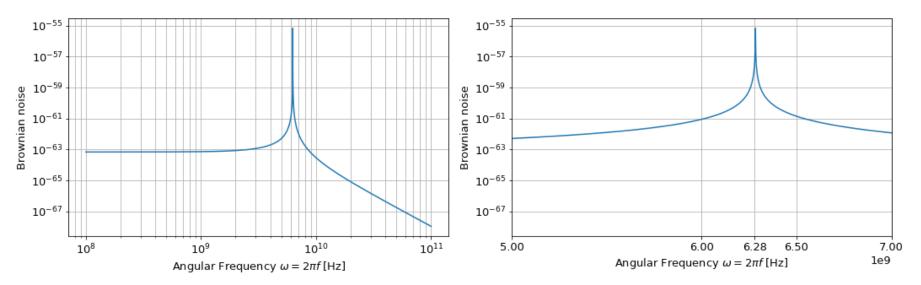
```
In [7]: Sshot = h*Lambda**4 /(2*c**2) * (Iref + Im)/(Iref * Im) display(Math(r"S_{x, x}^{beam, 1S}(\omega) ="+"{:.5}".format(Sshot) )) S_{x,x}^{beam,1S}(\omega) = 3.6899e - 71
```

### Question 4

I plot again the PSD graph around resonance

```
omega = np.logspace(8, 11, num=int(1e6), endpoint=True, base=10)
In [25]:
          y = brownian noise(omega)
          plt.rcParams.update({'font.size': 13})
          fig, axis = plt.subplots(1,2, figsize=(15, 5))
          fig.suptitle("Plot of brownian noise")
          axis[0].set xscale("log")
          axis[0].set_yscale("log")
          axis[0].plot(omega, y)
          axis[0].set xlabel(r"Angular Frequency $\omega=2\pi f$ [Hz]")
          axis[0].set ylabel(r"Brownian noise")
          axis[0].grid(True, which="both")
          axis[1].set yscale("log")
          axis[1].set xlim([0.5e10, 0.7e10])
```

#### Plot of brownian noise



Angular Frequency corresponding to the peak: = 6.2832

The peak is located at  $\omega=6.28*10^9 Hz \to f=10^9$  Hz while the narrow line has an approximately width of  $\Delta f\simeq (6.5-6)*10^9 Hz=0.5 GHz$ . As the shot noise is much smaller than the thermal

noise the time needed to evaluated the PSD of the noise can be obtained trough Dicke's radiometric formula:

$$ho = rac{1}{\sqrt{\Delta f T}} 
ightarrow T = rac{\Delta f}{
ho^2}$$

Minimum measurement time:  $= 5e + 10s \simeq 1585$ years

## **Question 5**

The minimum uncertainty can be evaluated thanks to Wiener's filter theory. As the input is a short pulse (template:  $\mathcal{FT}[\delta(t)] = 1$ ) the measurement time can be approximate as going from  $-\infty$  to  $+\infty$ . The minimum uncertainty is therefore:

$$\sigma^2 = rac{1}{rac{1}{2\pi} \int_0^{+\infty} SNR(\omega) d\omega}$$

where SNR is the signal to noise ratio defined as:

$$SNR(\omega) = rac{\left|f(\omega)
ight|^2}{S_{x,x}(\omega)} = rac{1}{S_{x,x}^{beam,1S}(\omega) + S_{x,x}^{th,1S}}$$

This integral can be performed numerically:

Value of the integral: l=9.4400e-61 Absolute error: 1.75e-60

$$\sigma^2 = 6.656e + 60$$

Which is a very high error, there is probably a mistake somewhere, but I can't find a way to find it out. My hypothesis is that is something related to the noise PSD as the integral heavily depends on it.