

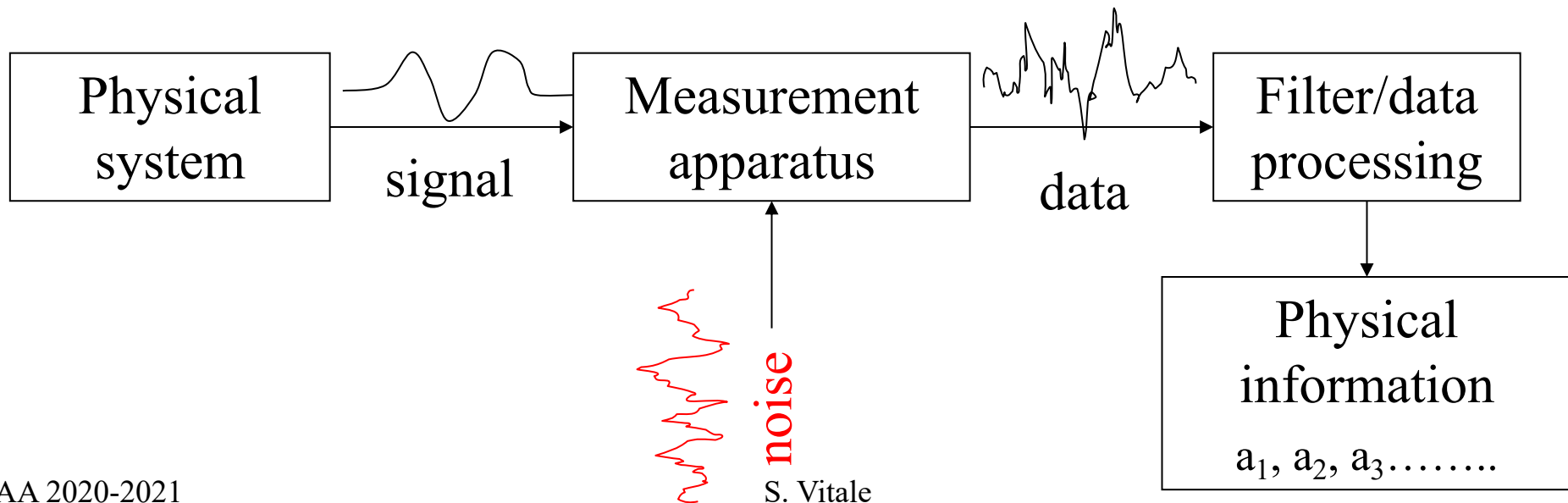
Experimental Methods

Lecture 13

October 19th, 2020

Noise

Noise in physical experiments is described as a random signal $x(t)$:
Independent repetitions of the same experiment produce different functions of time $x(t)$

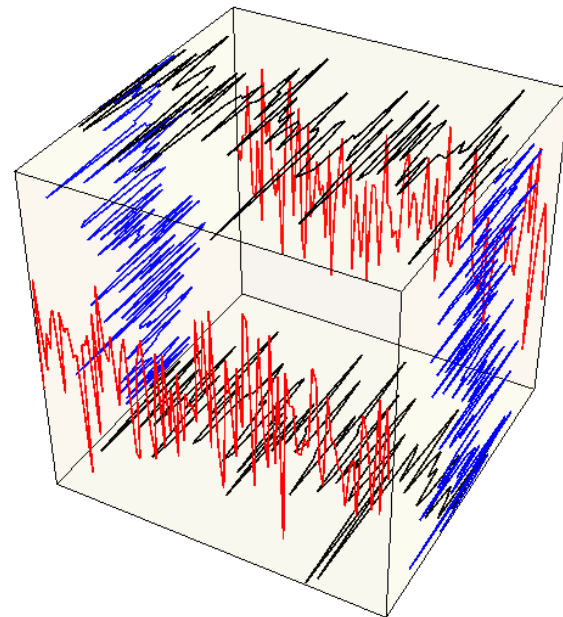


A stochastic process

- At each outcome of an experiment I get a function of time $x(t)$



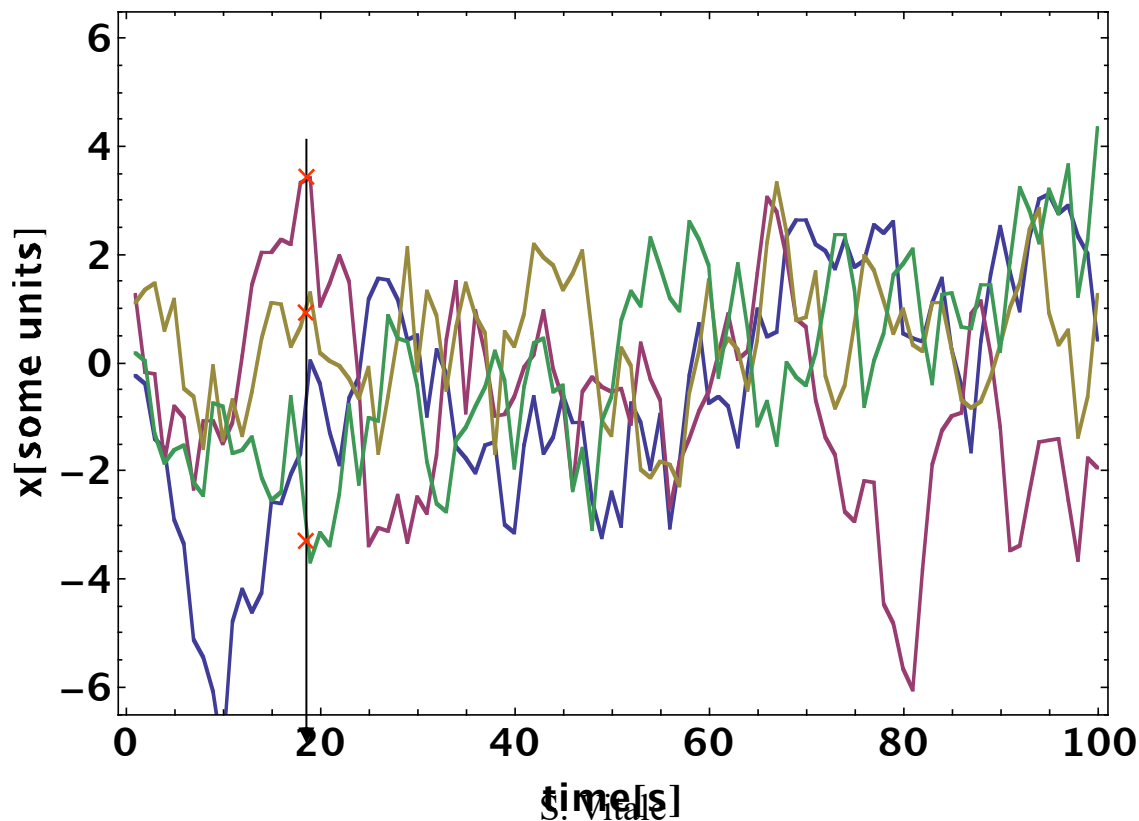
dice



Stochastic dice

A stochastic process

- For each outcome a function of time $x(t)$
- For each time t : a random variable



A stochastic process

- As $x(t_1)$, $x(t_2)$, $x(t_3)$ are all different rv, a stochastic process defines infinite many rv.
- Each rv will have its own density function

$$f_{x(t)}(\chi)$$

- With mean value

$$\eta(t) \equiv \langle x(t) \rangle = \int_{-\infty}^{\infty} \chi f_{x(t)}(\chi) d\chi$$

- This is called the mean value of the process and is in general a function of time.
- Same for the variance of the process

$$\sigma^2(t) \equiv \langle x^2(t) \rangle - \eta(t)^2 = \int_{-\infty}^{\infty} \chi^2 f_{x(t)}(\chi) d\chi - \eta(t)^2$$

Multiple random variables

- The statistical properties of multiple random variables like the $x(t_1)$, $x(t_2)$, $x(t_3)$ are described by joint probabilities. Take two rv x and y . Their joint probability density $f_{x,y}$ is defined by:

$$P\{x_0 \leq x \leq x_1 \text{ and } y_0 \leq y \leq y_1\} = \int_{x_0}^{x_1} \int_{y_0}^{y_1} f_{x,y}(\chi, \psi) d\chi d\psi$$

- Also this density needs to be normalized

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f_{x,y}(\chi, \psi) d\chi d\psi = 1$$

- The relation between the joint density $f_{x,y}$ and the marginal densities f_x and f_y is

$$\begin{aligned} \int_{x_0}^{x_1} f_x(\chi) d\chi &= P\{x_0 \leq x \leq x_1\} = \\ &= P\{x_0 \leq x \leq x_1 \text{ and } -\infty \leq y \leq \infty\} = \int_{x_0}^{x_1} \int_{-\infty}^{\infty} f_{x,y}(\chi, \psi) d\chi d\psi \end{aligned}$$

- Then: $f_x(\chi) = \int_{-\infty}^{\infty} f_{x,y}(\chi, \psi) d\psi$ $f_y(\psi) = \int_{-\infty}^{\infty} f_{x,y}(\chi, \psi) d\chi$

Multiple random variables

- Moments $\mu'_{n,m} = \langle x^n y^m \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x^n y^m f_{x,y}(\chi, \psi) d\chi d\psi$
- It follows $\mu'_{n,0} = \langle x^n \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x^n f_{x,y}(\chi, \psi) d\chi d\psi = \int_{-\infty}^{\infty} x^n f_x(\chi) d\chi = \mu'_n$
- Central moments $\mu_{m,n} = \langle (x - \langle x \rangle)^n (y - \langle y \rangle)^m \rangle$
- Specially important: the covariance $\sigma_{x,y} = \mu_{1,1} = \langle (x - \langle x \rangle)(y - \langle y \rangle) \rangle$
- Independent variables. If for any choice of $x_o, x_1, y_o,$ and y_1

$$P\{x_o \leq x \leq x_1 \text{ and } y_o \leq y \leq y_1\} = P\{x_o \leq x \leq x_1\} \times P\{y_o \leq y \leq y_1\}$$
- i.e. if $x \in [x_1, x_2]$ and $y \in [y_1, y_2]$ are independent events x and y are independent rv. It follows

$$\int_{x_o}^{x_1} \int_{y_o}^{y_1} f_{x,y}(\chi, \psi) d\chi d\psi = \int_{x_o}^{x_1} f_x(\chi) d\chi \int_{y_o}^{y_1} f_y(\psi) d\psi \rightarrow f_{x,y}(\chi, \psi) = f_x(\chi) f_y(\psi)$$

Multiple random variables

- For independent random variables $f_{x,y}(\chi, \psi) = f_x(\chi)f_y(\psi)$
- Joint moments

$$\begin{aligned}\mu'_{n,m} = \langle x^n y^m \rangle &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \chi^n \psi^m f_{x,y}(\chi, \psi) d\chi d\psi \\ &= \int_{-\infty}^{\infty} \chi^n f_x(\chi) d\chi \int_{-\infty}^{\infty} \psi^m f_y(\psi) d\psi = \langle x^n \rangle \langle y^m \rangle\end{aligned}$$

- In particular you can calculate that the covariance

$$\sigma_{x,y} = \mu_{1,1} = \langle (x - \langle x \rangle)(y - \langle y \rangle) \rangle = \langle x - \langle x \rangle \rangle \langle y - \langle y \rangle \rangle = 0$$

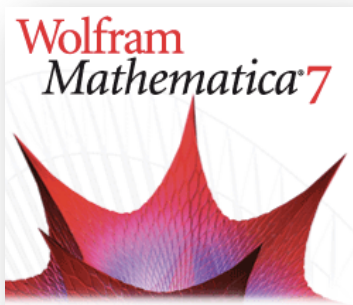
- Independent variables are “uncorrelated”
- The converse is not necessarily true.

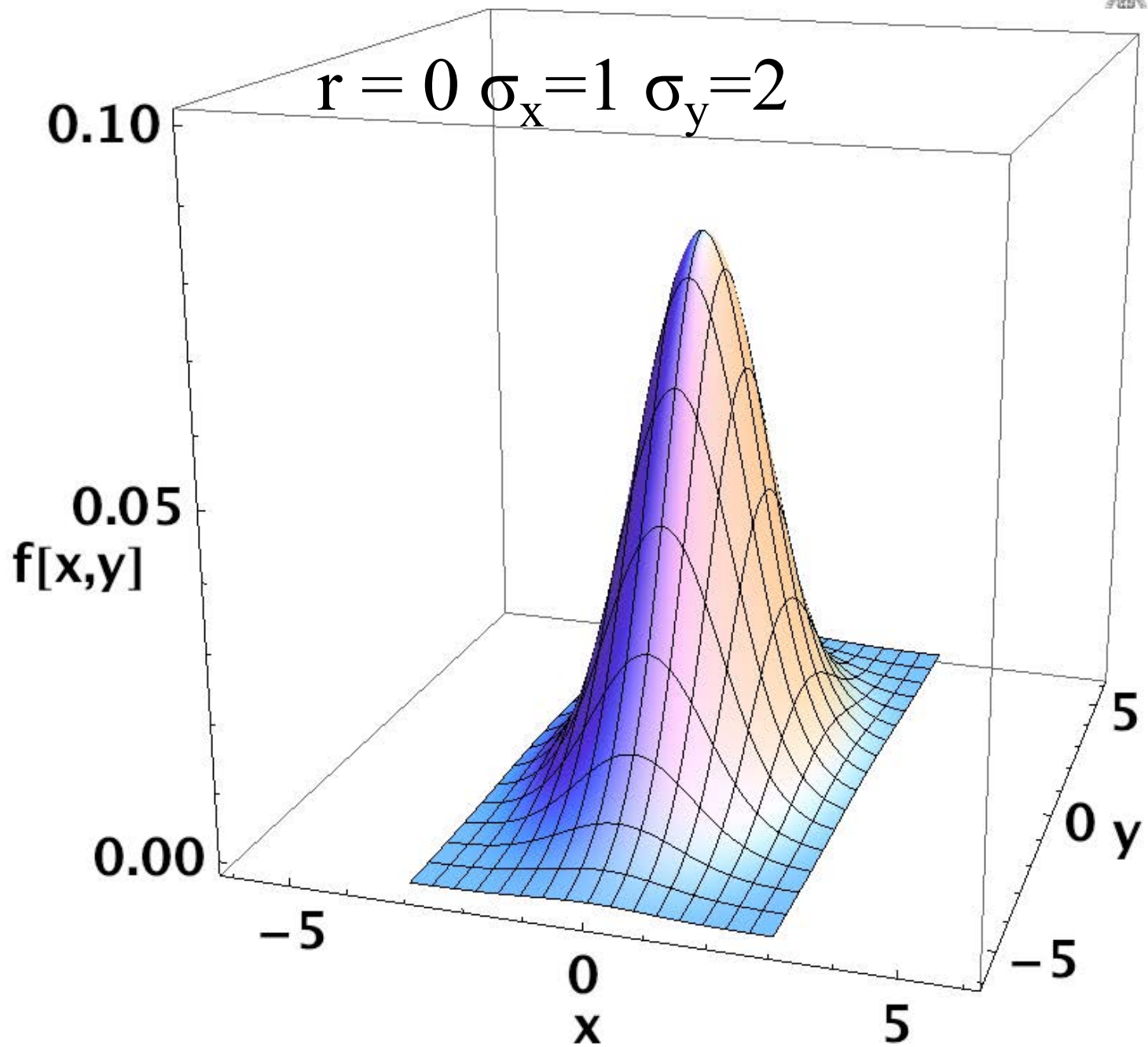
One example: a bidimensional Gaussian

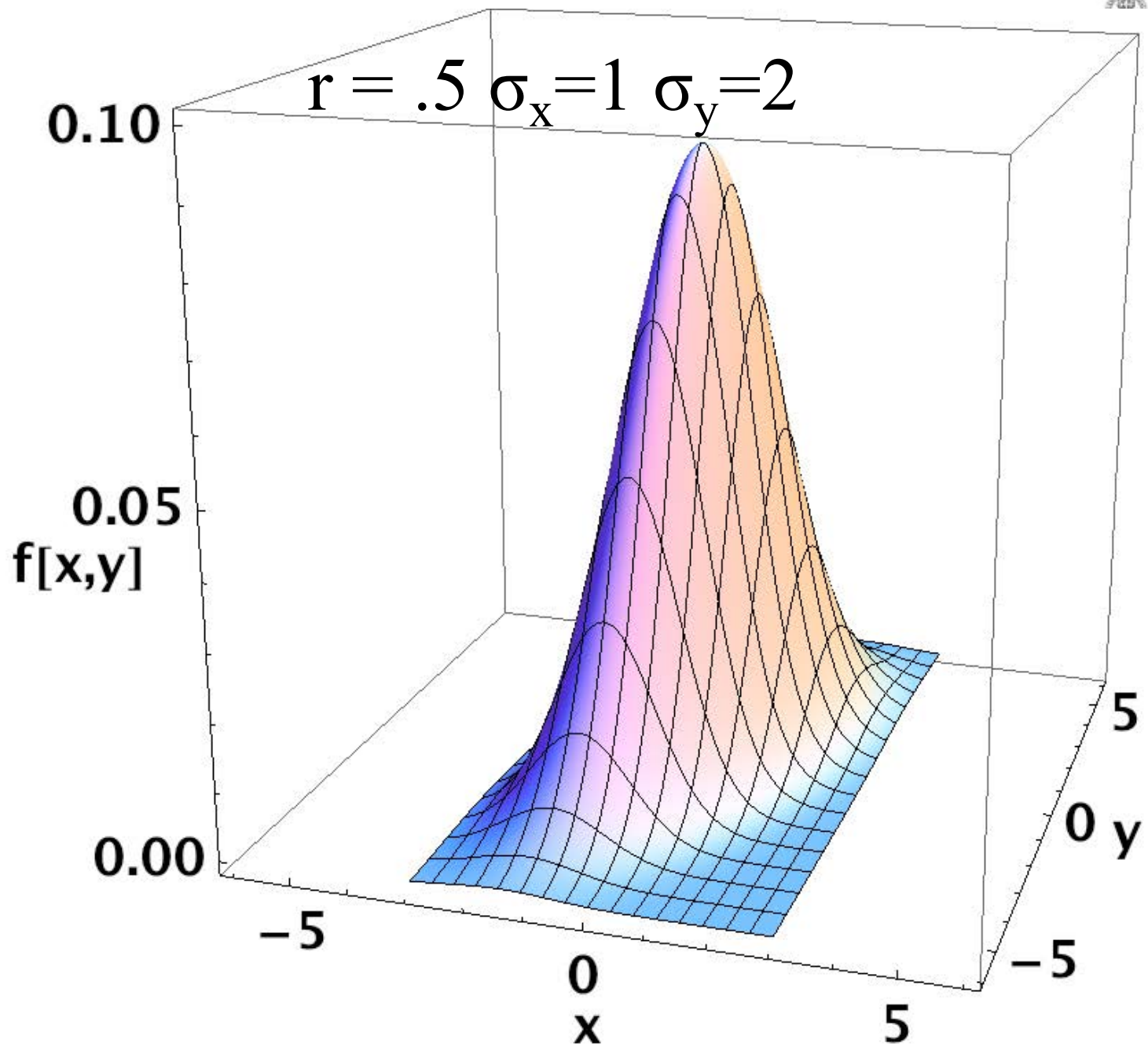
- The Joint Probability Density Function

$$f[x, y] = \frac{e^{-\frac{1}{2(1-r^2)} \left(\frac{(y-y_0)^2}{\sigma_y^2} - \frac{2r}{\sigma_x \sigma_y} (x-x_0)(y-y_0) + \frac{(x-x_0)^2}{\sigma_x^2} \right)}}{2\pi \sqrt{1-r^2} \sigma_x \sigma_y}$$

- For $r=0$ x and y are independent







Auto-correlation of a stochastic process

- Key high order moments of a stochastic processes are

- The auto-correlation

$$R(t, t') = \langle x(t) x(t') \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \chi \psi f_{x(t), x(t')}(\chi, \psi) d\chi d\psi$$

- And the auto-covariance

$$C(t, t') = \langle [x(t) - \eta(t)] [x(t') - \eta(t')] \rangle = R(t, t') - \eta(t) \eta(t')$$

- Notice that a stochastic process can always be written as its mean value, a deterministic function of time, plus the zero-mean process

$$\tilde{x}(t) = x(t) - \eta(t) \rightarrow x(t) = \eta(t) + \tilde{x}(t)$$

- The auto-covariance is the autocorrelation of the “noisy” part of the process
- Auto-covariance expresses the statistical memory of the process. To be further discussed

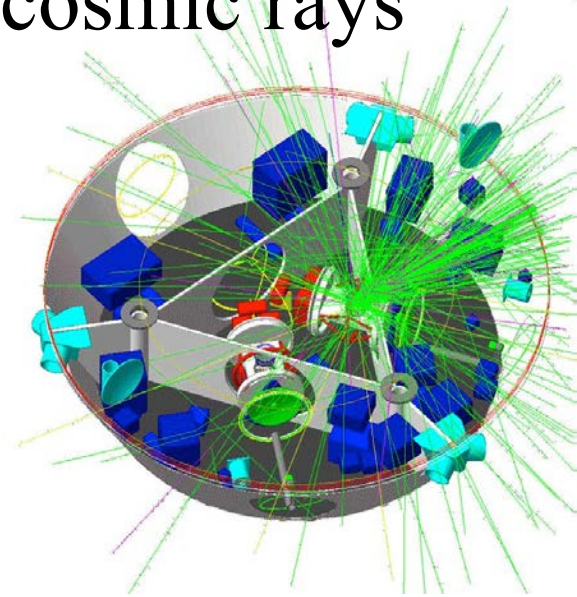
An important example: Poisson process and shot noise

- A stochastic process based model, for all phenomena where carriers arrive at random at a collector/detector:
 - Photons in a laser beam;
 - Charge carriers across a junction;
 - Particles from a radioactive source;

One example: random arrival of charge on LISA Pathfinder test-masses because of cosmic rays

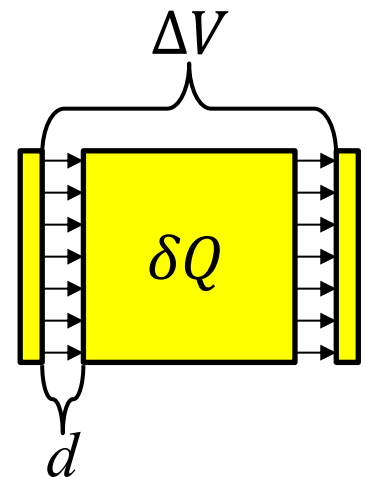
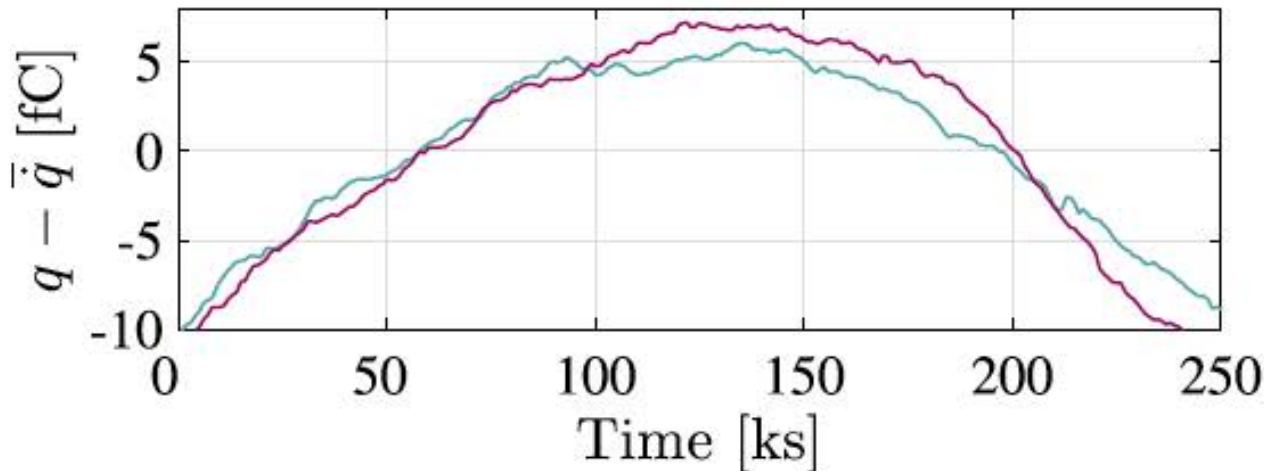
- Cosmic rays keep charging up the test-mass
- Random charge δQ produces force noise

$$\delta F_Q = \frac{\delta Q \Delta V}{d}$$



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PHYSICAL REVIEW LETTERS

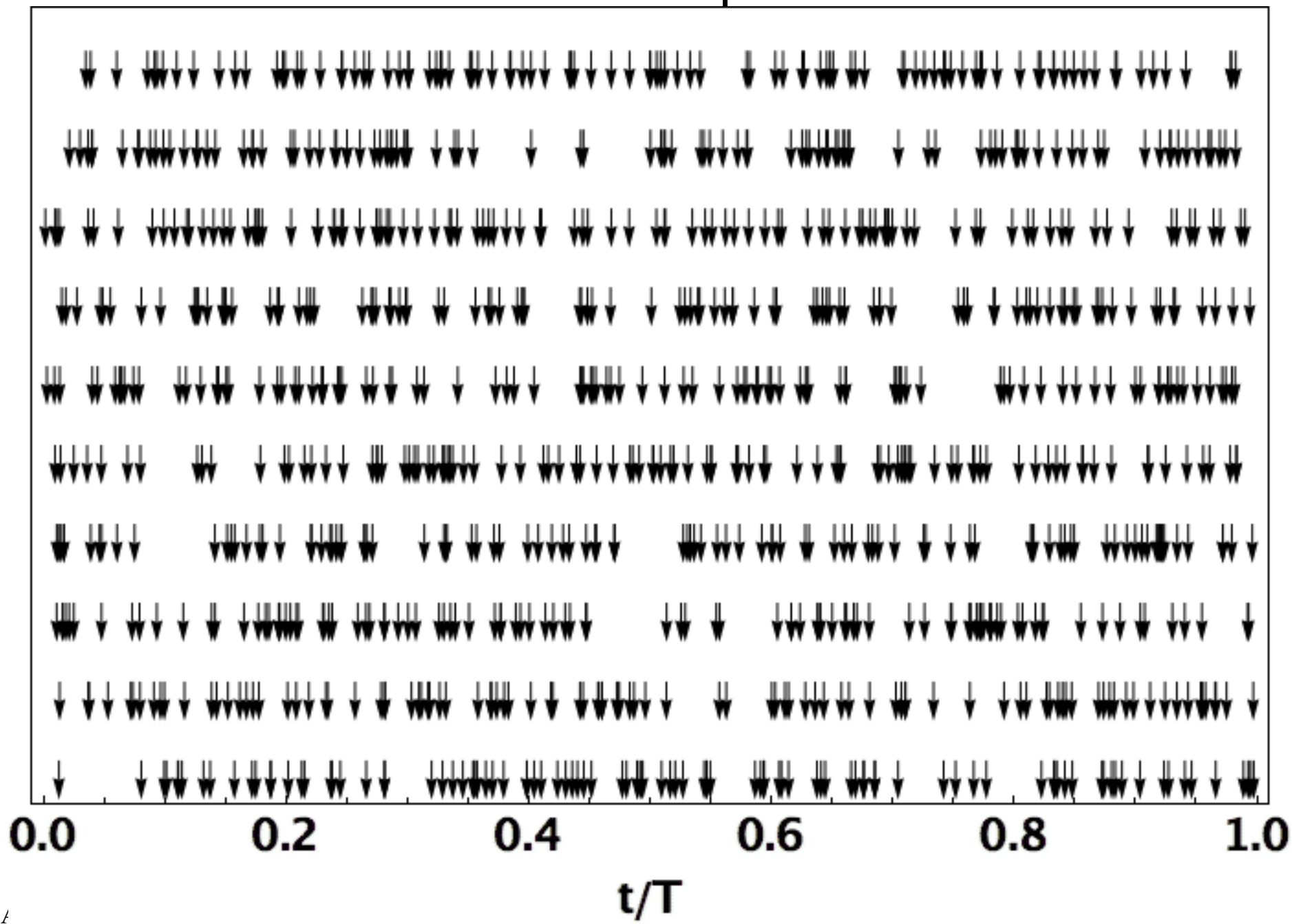
week ending
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Basics of the Poisson process

- Take a time interval with $0 \leq t \leq T$
- Pick N time instants at random, that is:
 - Instant t_i and instant t_k , for $i \neq k$, and $1 \leq i$ and $k \leq N$, are independently picked.
 - The probability density function for the random variable t_i is given by
- Conditions above mean that for $i \neq k$, $1 \leq i$ and $k \leq N$, $t_1 \leq t_2$, and $t_3 \leq t_4$:

$$P\{t_1 \leq t_i \leq t_2\} = \frac{t_2 - t_1}{T} \quad P\{t_1 \leq t_i \leq t_2 \text{ \& } t_3 \leq t_j \leq t_4\} = \frac{t_2 - t_1}{T} \frac{t_4 - t_3}{T}$$

Numerical example N=100



Probability of finding k points in a given interval

- Probability for one point of falling between t_1 and t_2 ($t_1 < t_2$)

$$P\{t_1 \leq t \leq t_2\} = \frac{t_2 - t_1}{T} \equiv p$$

- The events $t \in [t_1, t_2]$ and $t \notin [t_1, t_2]$ are mutually exclusive with probability p and $1-p$ respectively. Thus the probability of k points ($k \in \text{Integers}, 0 \leq k \leq N$) falling between t_1 and t_2 ($t_1 < t_2$) (Binomial distribution) is

$$P(k, t \in [t_1, t_2]) = \frac{N!}{k!(N-k)!} p^k (1-p)^{N-k}$$

- Note: k is a random variable. Its probability density is

$$f_k(x) = \sum_{k=0}^N \frac{N!}{k!(N-k)!} p^k (1-p)^{N-k} \delta(x-k)$$

The Poisson limit

- Binomial distribution

$$P(k, t \in [t_1, t_2]) = \frac{N!}{k!(N-k)!} p^k (1-p)^{N-k}$$

- Now try the following limit: $N \rightarrow \infty$, $T \rightarrow \infty$, $N/T \rightarrow \lambda$. Then

$$p = (t_2 - t_1)/T \rightarrow 0 \quad Np = N(t_2 - t_1)/T \rightarrow \lambda(t_2 - t_1)$$

- In the next slide it is shown that in this limit (Poisson)

$$P(k, t \in [t_1, t_2]) = \frac{[\lambda(t_2 - t_1)]^k}{k!} e^{-\lambda(t_2 - t_1)}$$

- You can check that

$$P\{0 \leq k \leq \infty\} = e^{-\lambda(t_2 - t_1)} \sum_{k=0}^{\infty} \frac{[\lambda(t_2 - t_1)]^k}{k!} = e^{-\lambda(t_2 - t_1)} e^{+\lambda(t_2 - t_1)} = 1$$

The Poisson Limit

Let's now calculate the limit for $N \rightarrow \infty$, $T \rightarrow \infty$, $N/T \rightarrow \lambda$ of the binomial formula

$$P_k = \frac{N!}{k! (N-k)!} p^k (1-p)^{N-k}$$

First rewrite probability as $p = (N (t_2 - t_1)/T)/N = \lambda(t_2 - t_1)/N$ then:

$$\begin{aligned} P_k &= \frac{N!}{k! (N-k)!} \left[\frac{\lambda(t_2 - t_1)}{N} \right]^k \left[1 - \frac{\lambda(t_2 - t_1)}{N} \right]^{N-k} = \\ &= \frac{N!}{N^k (N-k)!} \frac{[\lambda(t_2 - t_1)]^k}{k!} \left[1 - \frac{\lambda(t_2 - t_1)}{N} \right]^{N-k} \end{aligned}$$

Now consider the limits of the various parts

$$\begin{aligned} \lim_{N \rightarrow \infty} \left[1 - \frac{\lambda(t_2 - t_1)}{N} \right]^N &= e^{-\lambda(t_2 - t_1)} \quad \lim_{N \rightarrow \infty} \left[1 - \frac{\lambda(t_2 - t_1)}{N} \right]^{-k} = 1 \\ \lim_{n \rightarrow \infty} \frac{N!}{N^k (N-k)!} &= \lim_{N \rightarrow \infty} \frac{N^k + O[N^{k-1}]}{N^k} = 1 \end{aligned}$$

As a consequence

$$\lim_{N \rightarrow \infty, T \rightarrow \infty, N/T \rightarrow \lambda} P_k = \frac{[\lambda(t_2 - t_1)]^k}{k!} e^{-\lambda(t_2 - t_1)}$$

The Poisson limit

- Thus in the limit: $N \rightarrow \infty$, $T \rightarrow \infty$, $N/T \rightarrow \lambda$, k has probability density:

$$f_k(x) = e^{-\lambda(t_2 - t_1)} \sum_{m=0}^{\infty} \frac{[\lambda(t_2 - t_1)]^m}{m!} \delta(x - m)$$

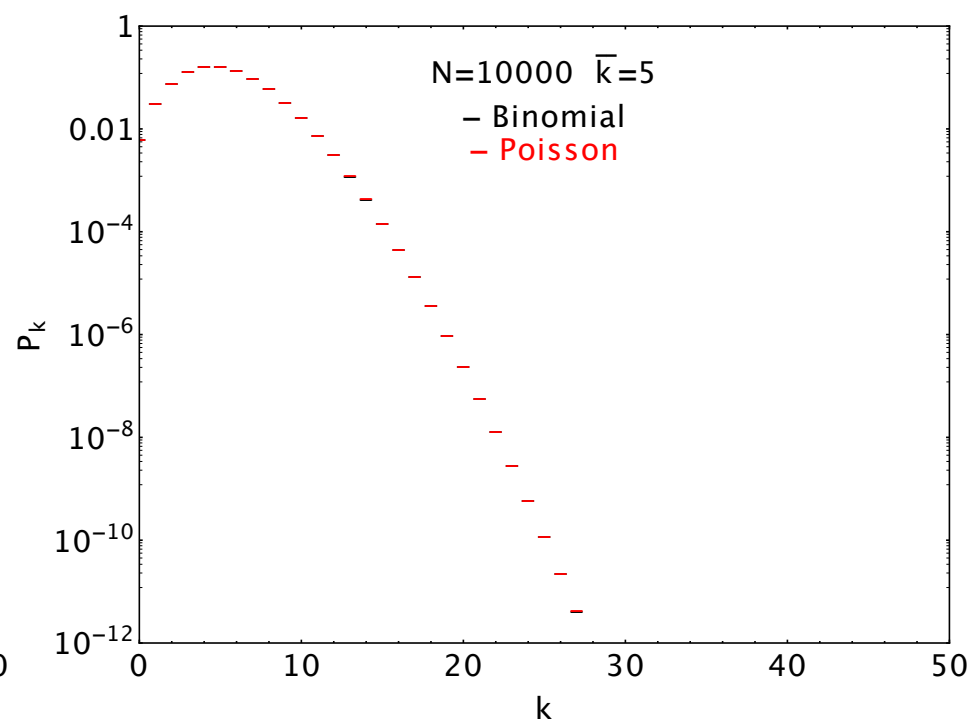
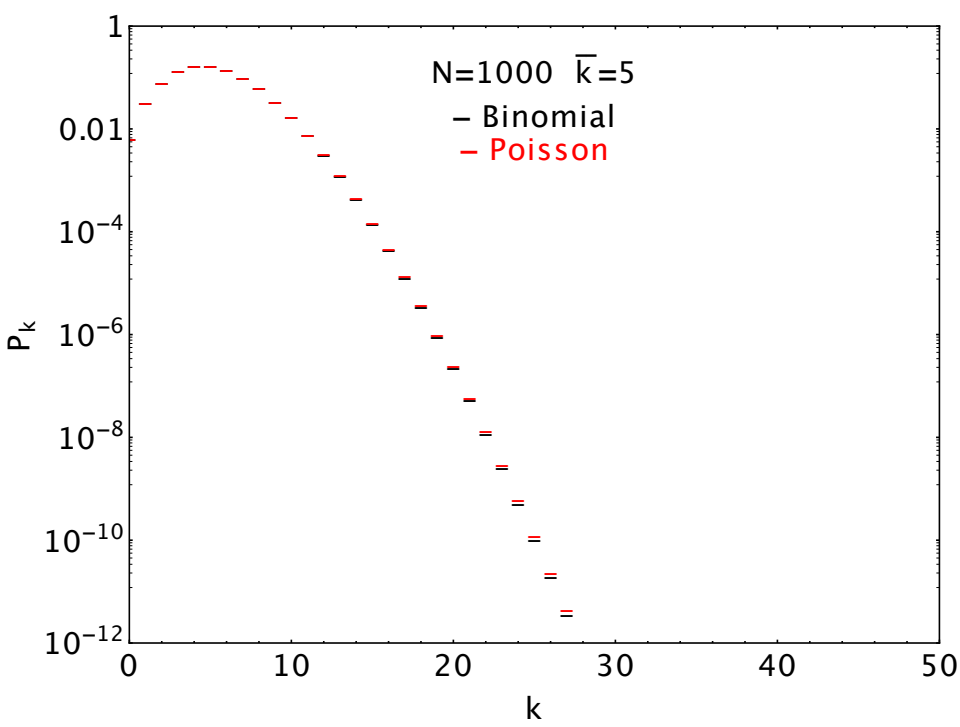
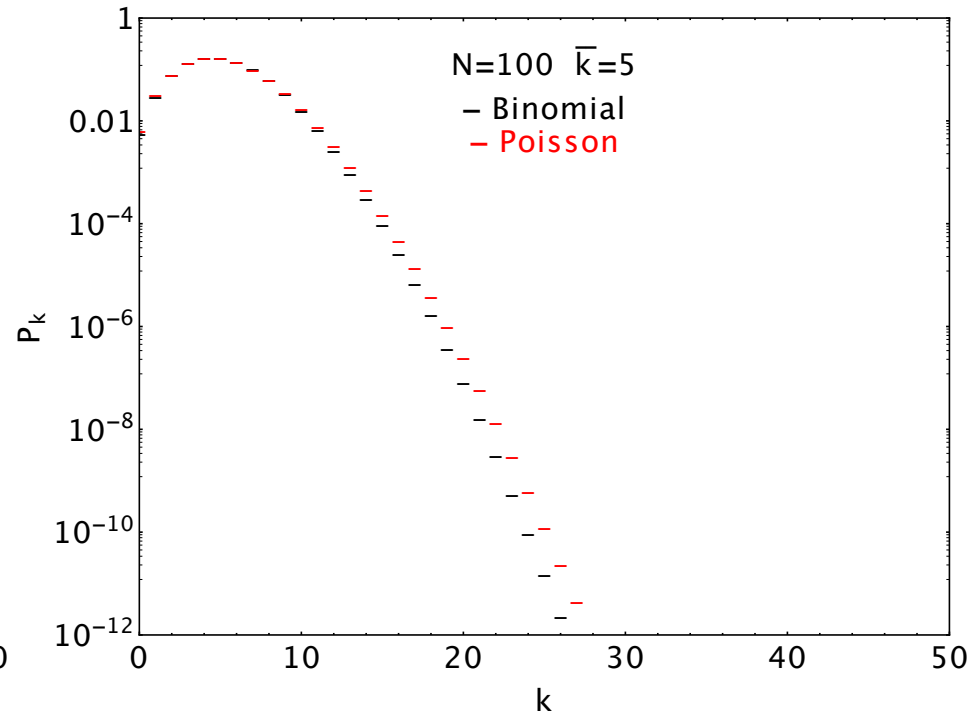
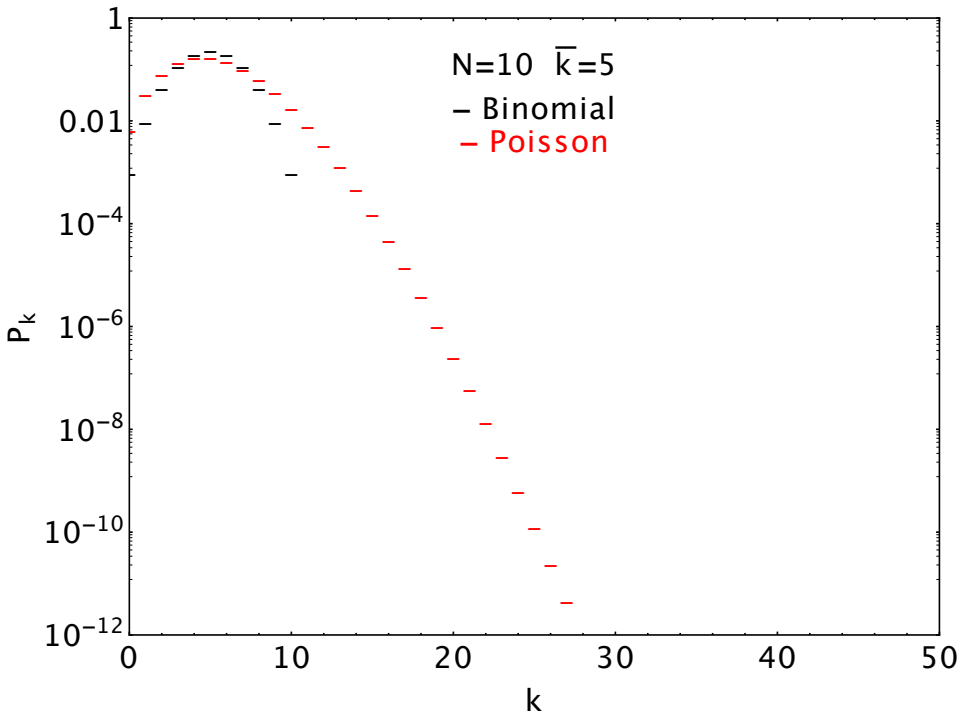
- Mean value of which is

$$\bar{k} = \int_{-\infty}^{\infty} x f_k(x) dx = e^{-\lambda(t_2 - t_1)} \sum_{m=0}^{\infty} m \frac{[\lambda(t_2 - t_1)]^m}{m!} = e^{-\lambda(t_2 - t_1)} \sum_{m=1}^{\infty} m \frac{[\lambda(t_2 - t_1)]^m}{m!}$$

- By expanding

$$\begin{aligned} \bar{k} &= e^{-\lambda(t_2 - t_1)} \sum_{m=1}^{\infty} \frac{[\lambda(t_2 - t_1)]^m}{(m-1)!} = e^{-\lambda(t_2 - t_1)} \lambda(t_2 - t_1) \sum_{m=1}^{\infty} \frac{[\lambda(t_2 - t_1)]^{m-1}}{(m-1)!} \\ &= e^{-\lambda(t_2 - t_1)} \lambda(t_2 - t_1) \sum_{n=0}^{\infty} \frac{[\lambda(t_2 - t_1)]^n}{n!} = e^{-\lambda(t_2 - t_1)} \lambda(t_2 - t_1) e^{\lambda(t_2 - t_1)} = \lambda(t_2 - t_1) \end{aligned}$$

- Then
- $$P_k = e^{-\bar{k}} \frac{\bar{k}^m}{m!} \quad f_k(x) = e^{-\bar{k}} \sum_{m=0}^{\infty} \frac{\bar{k}^m}{m!} \delta(x - m)$$



The Poisson limit

- k has probability density function:

$$f_k(x) = e^{-\bar{k}} \sum_{m=0}^{\infty} \frac{\bar{k}^m}{m!} \delta(x - m)$$

- With mean value

$$\bar{k} = \lambda(t_2 - t_1)$$

- Variance. Begin by calculating

$$\begin{aligned} \langle k^2 \rangle &= \int_{-\infty}^{\infty} x^2 f_k(x) dx = e^{-\bar{k}} \sum_{m=0}^{\infty} m^2 \frac{\bar{k}^m}{m!} = e^{-\bar{k}} \sum_{m=1}^{\infty} m^2 \frac{\bar{k}^m}{m!} = e^{-\bar{k}} \bar{k} \sum_{m=1}^{\infty} m \frac{\bar{k}^{m-1}}{(m-1)!} \\ &= \bar{k} \left[e^{-\bar{k}} \sum_{n=0}^{\infty} (n+1) \frac{\bar{k}^n}{n!} \right] = \bar{k}(\bar{k} + 1) \end{aligned}$$

- Then the variance is

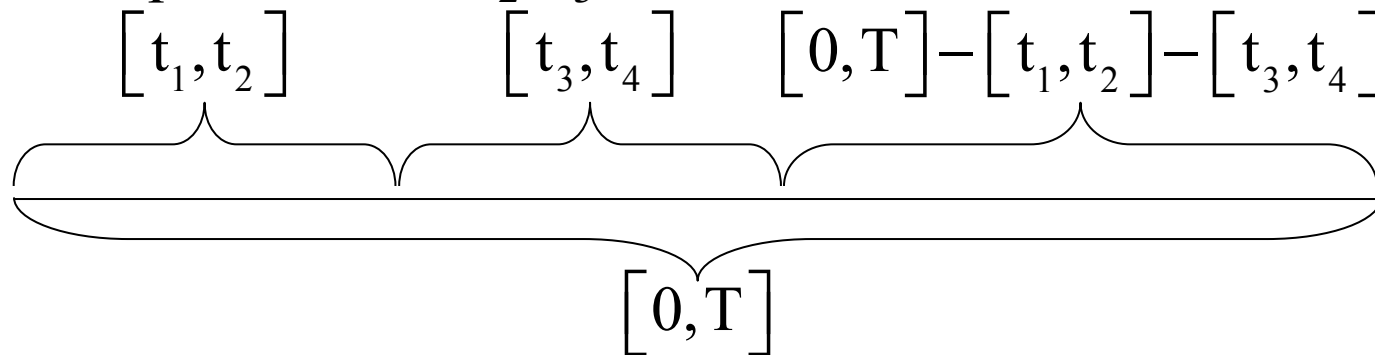
$$\sigma_k^2 = \langle k^2 \rangle - \langle k \rangle^2 = \bar{k}(\bar{k} + 1) - \bar{k}^2 = \bar{k}$$

- And the standard deviation is $\sigma_k = \sqrt{\bar{k}}$

- Notice $\sigma_k / \bar{k} = 1 / \sqrt{\bar{k}}$

Numbers of instants falling in non-overlapping time intervals are not independent rv

- Consider 2 time intervals $[t_1, t_2]$ and $[t_3, t_4]$ between 0 and T that have no common points. Take $t_2 < t_3$

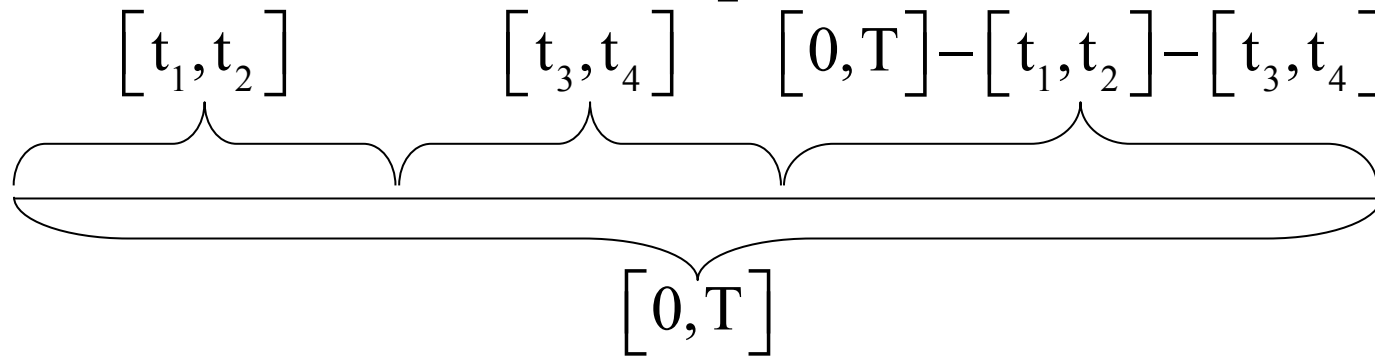


- The events $t \in [t_1, t_2]$, $t \in [t_3, t_4]$, and events $t \in [0, T] - [t_1, t_2] - [t_3, t_4]$ are mutually exclusive. The probability of k_1 time-instants falling into $[t_1, t_2]$, k_2 instants falling into $[t_3, t_4]$ and $N - k_1 - k_2$ falling into $[0, T] - [t_1, t_2] - [t_3, t_4]$ is given by:

$$P[k_1 \in [t_1, t_2] \text{ and } k_2 \in [t_3, t_4]] =$$

$$= \frac{N!}{k_1! k_2! (N - k_1 - k_2)!} \left(\frac{t_2 - t_1}{T} \right)^{k_1} \left(\frac{t_4 - t_3}{T} \right)^{k_2} \left(1 - \frac{t_2 - t_1}{T} - \frac{t_4 - t_3}{T} \right)^{N - k_1 - k_2}$$

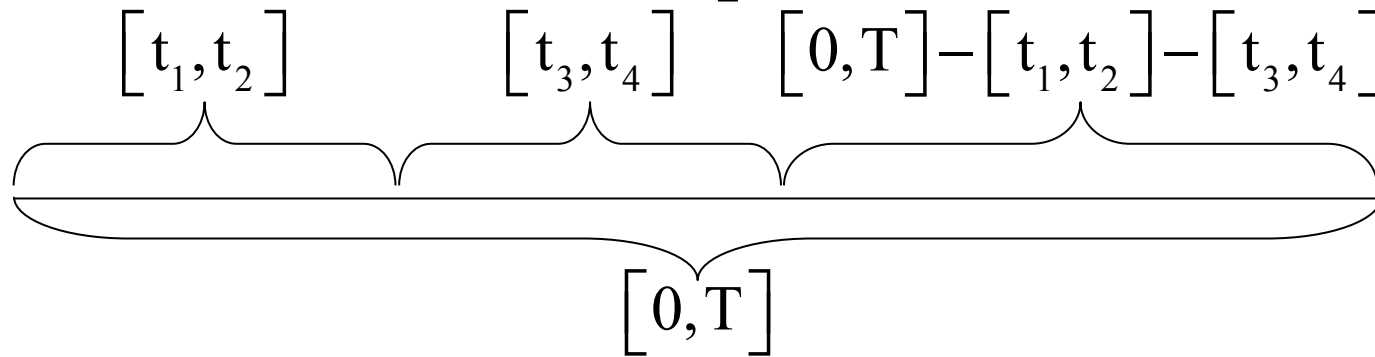
Numbers of instants in non-overlapping time intervals are not independent rv



- Notice

$$\begin{aligned}
 &P[k_1 \in [t_1, t_2] \text{ and } k_2 \in [t_3, t_4]] = \\
 &= \frac{N!}{k_1! k_2! (N - k_1 - k_2)!} \left(\frac{t_2 - t_1}{T} \right)^{k_1} \left(\frac{t_4 - t_3}{T} \right)^{k_2} \left(1 - \frac{t_2 - t_1}{T} - \frac{t_4 - t_3}{T} \right)^{N - k_1 - k_2} \neq \\
 &\neq \frac{N!}{k_1! (N - k_1)!} \left(\frac{t_2 - t_1}{T} \right)^{k_1} \left(1 - \frac{t_2 - t_1}{T} \right)^{N - k_1} \times \\
 &\times \frac{N!}{k_2! (N - k_2)!} \left(\frac{t_4 - t_3}{T} \right)^{k_2} \left(1 - \frac{t_4 - t_3}{T} \right)^{N - k_2} = P[k_1 \in [t_1, t_2]] P[k_2 \in [t_3, t_4]]
 \end{aligned}$$

Numbers of instants in non-overlapping time intervals are not independent rv



- As

$$P[k_1 \in [t_1, t_2] \text{ and } k_2 \in [t_3, t_4]] \neq P[k_1 \in [t_1, t_2]]P[k_2 \in [t_3, t_4]]$$

- k_1 and k_2 are not independent.
- They only *become independent* in the now famous limit $N \rightarrow \infty$, $T \rightarrow \infty$, $N/T \rightarrow \lambda$ (see next 2 pages)

The limit $N \rightarrow \infty, T \rightarrow \infty, N/T \rightarrow \lambda$

Start with the formula:

$$\frac{N!}{k_1! k_2! (N - k_1 - k_2)!} \left(\frac{t_2 - t_1}{T} \right)^{k_1} \left(\frac{t_4 - t_3}{T} \right)^{k_2} \left(1 - \frac{t_2 - t_1}{T} - \frac{t_4 - t_3}{T} \right)^{N - k_1 - k_2}$$

Use the same trick as before $(t_2 - t_1)/T = [(N/T)(t_2 - t_1)/N] = \lambda(t_2 - t_1)$ and same for t_3 and t_4 . Then

$$= \frac{N!}{k_1! k_2! (N - k_1 - k_2)!} (\lambda(t_2 - t_1) / N)^{k_1} (\lambda(t_4 - t_3) / N)^{k_2} \times \\ \times (1 - \lambda(t_2 - t_1) / N - \lambda(t_4 - t_3) / N)^{N - k_1 - k_2}$$

You can proceed as before

$$= \frac{N!}{N^{k_1} N^{k_2} (N - k_1 - k_2)!} \left[(\lambda(t_2 - t_1))^{k_1} / k_1! \right] \left[(\lambda(t_4 - t_3))^{k_2} / k_2! \right] \times \\ \times (1 - \lambda(t_2 - t_1) / N - \lambda(t_4 - t_3) / N)^N (1 - \lambda(t_2 - t_1) / N - \lambda(t_4 - t_3) / N)^{-k_1 - k_2}$$

The limit $N \rightarrow \infty, T \rightarrow \infty, N/T \rightarrow \lambda$

Now take the limits of the various parts

$$\frac{N!}{N^{k_1} N^{k_2} (N - k_1 - k_2)!} \rightarrow \frac{N^{k_1+k_2} + O[N^{k_1+k_2-1}]}{N^{k_1+k_2}} \rightarrow 1$$

$$(1 - \lambda(t_2 - t_1) / N - \lambda(t_4 - t_3) / N)^{-k_1-k_2} \rightarrow 1$$

$$(1 - \lambda(t_2 - t_1) / N - \lambda(t_4 - t_3) / N)^N \rightarrow e^{-\lambda(t_2-t_1)-\lambda(t_4-t_3)}$$

then

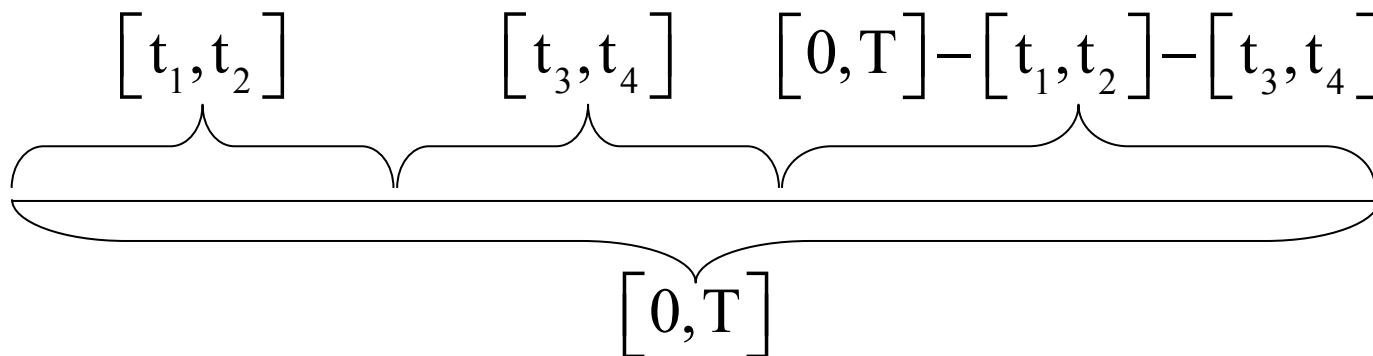
$$\begin{aligned} &= \frac{N!}{N^{k_1} N^{k_2} (N-k_1-k_2)!} \left[(\lambda(t_2 - t_1))^{k_1} / k_1! \right] \left[(\lambda(t_4 - t_3))^{k_2} / k_2! \right] \times \\ &\times (1 - \lambda(t_2 - t_1) / N - \lambda(t_4 - t_3) / N)^N (1 - \lambda(t_2 - t_1) / N - \lambda(t_4 - t_3) / N)^{-k_1-k_2} \\ &\rightarrow \left[(\lambda(t_2 - t_1))^{k_1} / k_1! \right] \left[(\lambda(t_4 - t_3))^{k_2} / k_2! \right] e^{-\lambda(t_2-t_1)} e^{-\lambda(t_4-t_3)} \end{aligned}$$

That is

$$P[k_1 \in [t_1, t_2] \text{ and } k_2 \in [t_3, t_4]] = P[k_1 \in [t_1, t_2]] P[k_2 \in [t_3, t_4]]$$

Basics of the Poisson process

- In the limit $N \rightarrow \infty$, $T \rightarrow \infty$, $N/T \rightarrow \lambda$ k_1 and k_2 , numbers of points within disjoint intervals $[t_1, t_2]$ and $[t_3, t_4]$, are independent random variable



$$f_{k_1, k_2}(x, y) = \left\{ e^{-\bar{k}_1} \sum_{m=0}^{\infty} \frac{\bar{k}_1^m}{m!} \delta(x - m) \right\} \left\{ e^{-\bar{k}_2} \sum_{m=0}^{\infty} \frac{\bar{k}_2^m}{m!} \delta(y - m) \right\} = f_{k_1}(x) f_{k_2}(y)$$

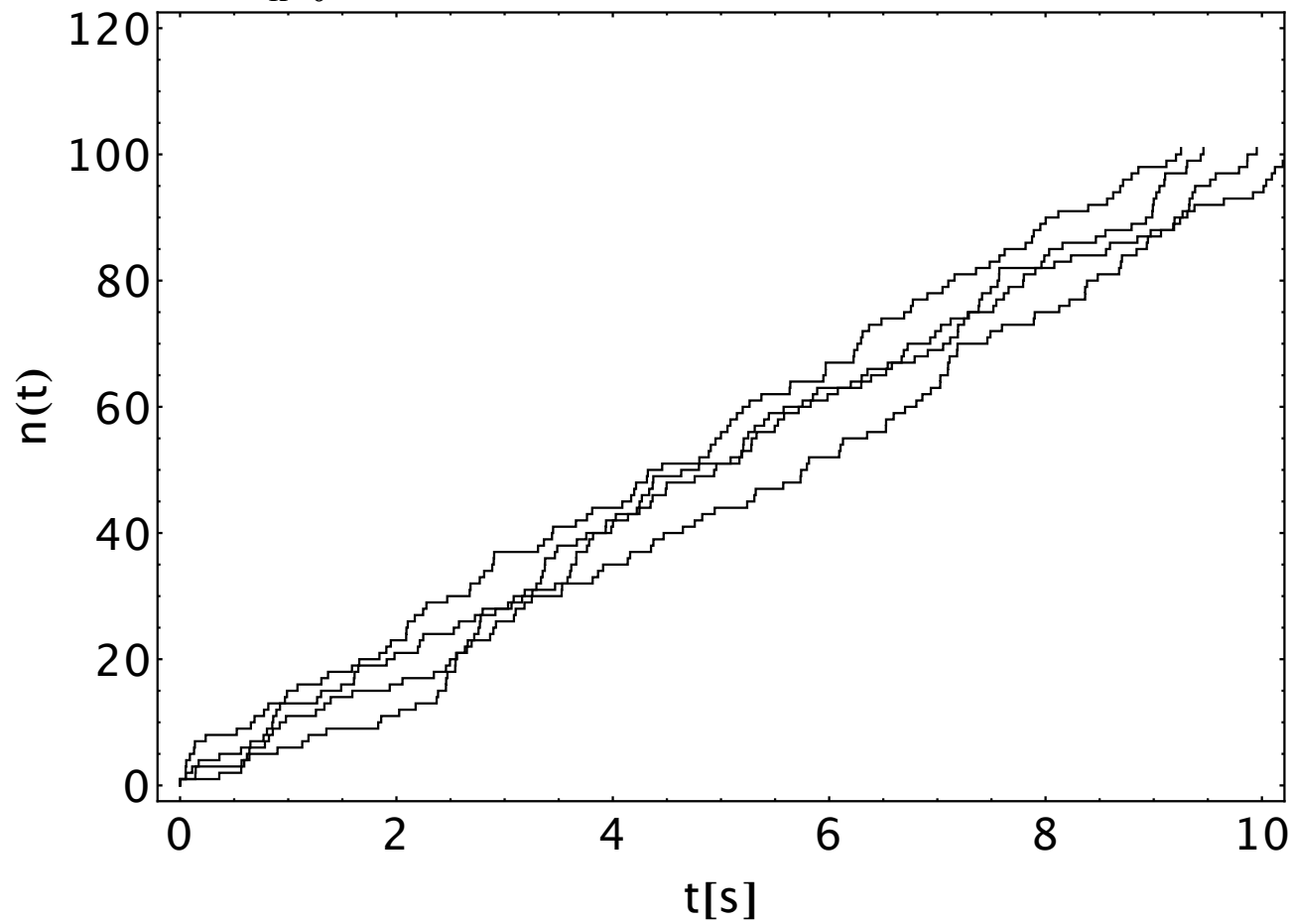
The Poisson stochastic process

- Assume that the time instants above represent the time of arrival of some carrier at a given point. We define the process $n(t)$ as the number of carriers that have already arrived at time t :

$$n(t) = \sum_{k=0} \Theta(t - t_k)$$

Example with $\lambda = 10 \text{ s}^{-1}$.

Different lines corresponds to different outcomes of the experiment



Geiger counting

