



TRƯỜNG ĐẠI HỌC
BÁCH KHOA HÀ NỘI



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Engineering Electromagnetics

Current & Conductors

Contents

- I. Introduction
- II. Vector Analysis
- III. Coulomb's Law & Electric Field Intensity
- IV. Electric Flux Density, Gauss' Law & Divergence
- V. Energy & Potential
- VI. Current & Conductors**
- VII. Dielectrics & Capacitance
- VIII. Poisson's & Laplace's Equations
- IX. The Steady Magnetic Field
- X. Magnetic Forces & Inductance
- XI. Time – Varying Fields & Maxwell's Equations
- XII. Transmission Lines
- XIII. The Uniform Plane Wave
- XIV. Plane Wave Reflection & Dispersion
- XV. Guided Waves & Radiation



Current & Conductors

1. Current & Current Density
2. Metallic Conductors
3. Conductor Properties & Boundary Conditions
4. The Method of Images
5. Semiconductors
6. Applications



Current & Current Density (1)

- Current:

$$I = \frac{dQ}{dt}$$

- Unit A (ampère)
- Current is defined as the motion of positive charges

Current & Current Density (2)

- Current: rate of movement of charge crossing a given reference plane (of one coulomb per second)
- Current density: \mathbf{J} (A/m²)
- The increment of current ΔI crossing an incremental surface ΔS normal to the current density:

$$\Delta I = J_N \Delta S$$

- If the current density is not perpendicular to the surface:

$$\Delta I = \mathbf{J} \cdot \Delta \mathbf{S}$$

- Total current:

$$I = \int_S \mathbf{J} \cdot d\mathbf{S}$$

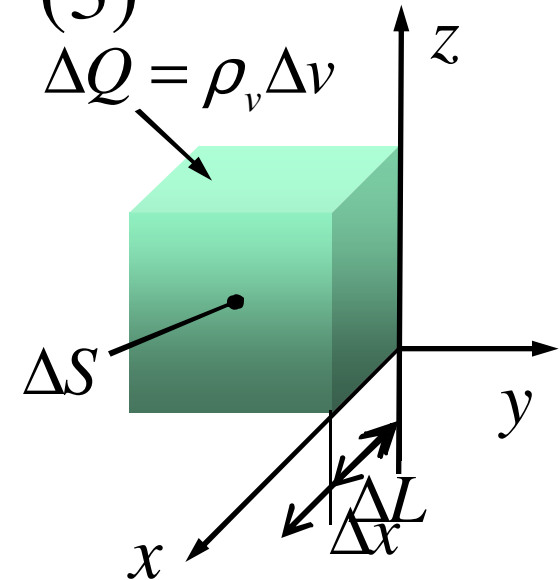
Current & Current Density (3)

$$\Delta Q = \rho_v \Delta v = \rho_v \Delta S \Delta L$$

$$\left. \begin{array}{l} \Delta Q = \rho_v \Delta S \Delta x \\ \Delta I = \frac{\Delta Q}{\Delta t} \end{array} \right\} \rightarrow \Delta I = \rho_v \Delta S \frac{\Delta x}{\Delta t}$$

$$\left. \begin{array}{l} = \rho_v \Delta S v_x \\ \Delta I = J_x \Delta S \end{array} \right\} \rightarrow J_x = \rho_v v_x$$

$$\rightarrow \boxed{\mathbf{J} = \rho_v \mathbf{v}}$$



Current & Current Density (4)

Ex. 1

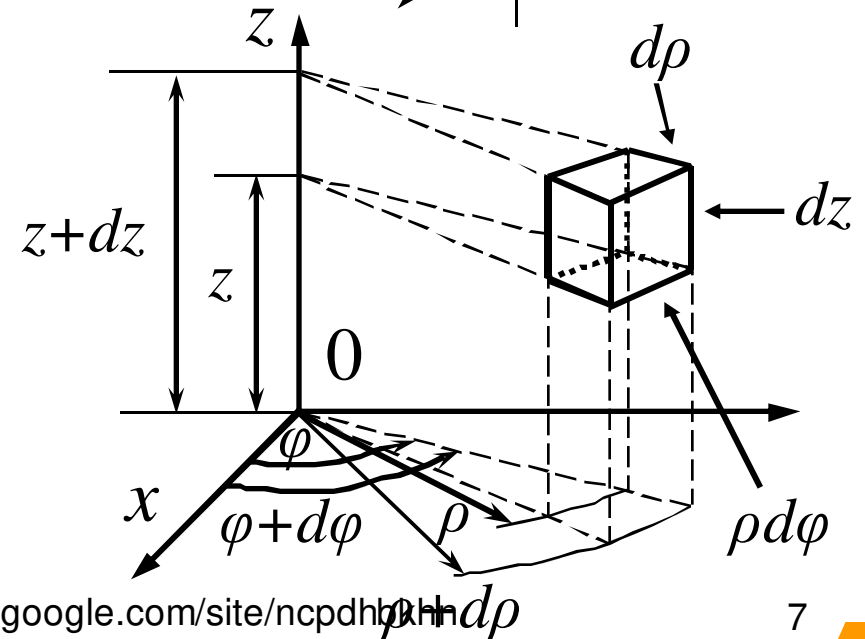
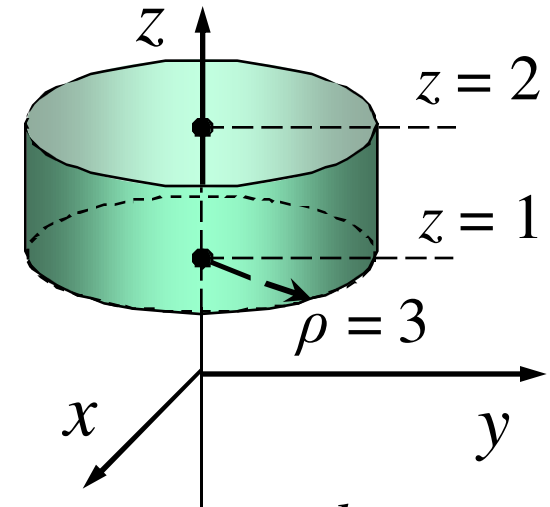
Given $\mathbf{J} = 2\rho z \mathbf{a}_\rho + 7z \sin^2 \varphi \mathbf{a}_\varphi$ mA/m². Find the total current leaving the circular band.

$$I = \int_S \mathbf{J} \cdot d\mathbf{S} = \int_S \mathbf{J}|_{\rho=3} \cdot d\mathbf{S}$$

$$\mathbf{J}|_{\rho=3} = 2 \times 3 z \mathbf{a}_\rho + 7 z \sin^2 \varphi \mathbf{a}_\varphi$$

$$= 6z \mathbf{a}_\rho + 7z \sin^2 \varphi \mathbf{a}_\varphi$$

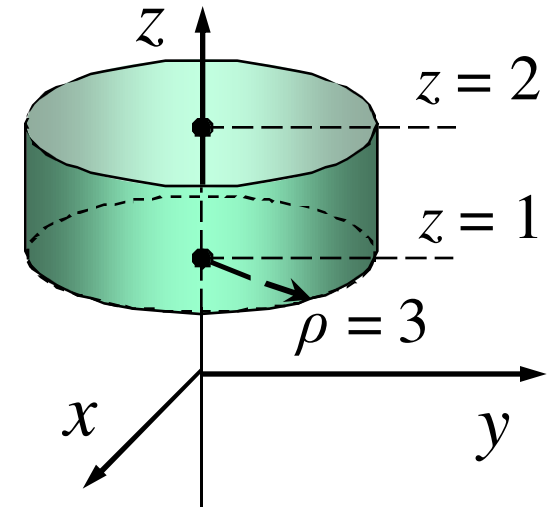
$$d\mathbf{S} = \rho d\varphi dz \mathbf{a}_\rho = 3 d\varphi dz \mathbf{a}_\rho$$



Current & Current Density (5)

Ex. 1

Given $\mathbf{J} = 2\rho z \mathbf{a}_\rho + 7z \sin^2 \varphi \mathbf{a}_\varphi$ mA/m². Find the total current leaving the circular band.



$$I = \int_S \mathbf{J} \cdot d\mathbf{S} = \int_S \mathbf{J}|_{\rho=3} \cdot d\mathbf{S}$$

$$\mathbf{J}|_{\rho=3} = 2 \times 3z \mathbf{a}_\rho + 7z \sin^2 \varphi \mathbf{a}_\varphi$$

$$\left. \begin{aligned} &= 6z \mathbf{a}_\rho + 7z \sin^2 \varphi \mathbf{a}_\varphi \\ d\mathbf{S} &= \rho d\varphi dz \mathbf{a}_\rho = 3d\varphi dz \mathbf{a}_\rho \end{aligned} \right\} \rightarrow \mathbf{J}|_{\rho=3} \cdot d\mathbf{S} = 18z d\varphi dz$$

$$\rightarrow I = \int_{z=1}^{z=2} \int_{\varphi=0}^{\varphi=2\pi} 18z d\varphi dz = \int_{z=1}^{z=2} 2\pi \times 18z dz = 169 \text{ mA}$$



Current & Current Density (6)

The current leaving a closed surface: $I = \oint_S \mathbf{J} \cdot d\mathbf{S}$

The total charge in the surface: Q_i

The law of conservation of charge

$$\rightarrow I = \oint_S \mathbf{J} \cdot d\mathbf{S} = -\frac{dQ_i}{dt}$$

- in circuit analysis, $I = dQ/dt$ because this is an entering current
- in electromagnetism, $I = -dQ/dt$ because this is a leaving one

Current & Current Density (7)

$$\left. \begin{aligned} I &= \oint_S \mathbf{J} \cdot d\mathbf{S} = -\frac{dQ_i}{dt} \\ \oint_S \mathbf{J} \cdot d\mathbf{S} &= \int_V (\nabla \cdot \mathbf{J}) dv \quad (\text{div. theo.}) \end{aligned} \right\} \rightarrow \left. \begin{aligned} \int_V (\nabla \cdot \mathbf{J}) dv &= -\frac{dQ_i}{dt} \\ Q_i &= \int_V \rho_v dv \end{aligned} \right\}$$

$$\rightarrow \int_V (\nabla \cdot \mathbf{J}) dv = -\frac{d}{dt} \int_V \rho_v dv = \int_V -\frac{\partial \rho_v}{\partial t} dv$$

$$\rightarrow (\nabla \cdot \mathbf{J}) \Delta v = -\frac{\partial \rho_v}{\partial t} \Delta v \rightarrow \boxed{\nabla \cdot \mathbf{J} = -\frac{\partial \rho_v}{\partial t}}$$

Ex. 2 Current & Current Density (9)

Consider the current density $\mathbf{J} = \frac{e^{-t}}{r} \mathbf{a}_r$ A/m².

$$\left. \begin{aligned} -\frac{\partial \rho_v}{\partial t} &= \nabla \cdot \mathbf{J} = \nabla \cdot \left(\frac{e^{-t}}{r} \mathbf{a}_r \right) \\ \nabla \cdot \mathbf{D} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 D_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta D_\theta) + \frac{1}{r \sin \theta} \frac{\partial D_\phi}{\partial \phi} \end{aligned} \right\}$$

$$\rightarrow -\frac{\partial \rho_v}{\partial t} = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{e^{-t}}{r} \right) = \frac{e^{-t}}{r^2} \rightarrow \rho_v = -\int \frac{e^{-t}}{r^2} dt + K(r) = \frac{e^{-t}}{r^2} + K(r) \left\}$$

Suppose $\rho_v \rightarrow 0$ as $t \rightarrow \infty$, then $K(r) = 0$

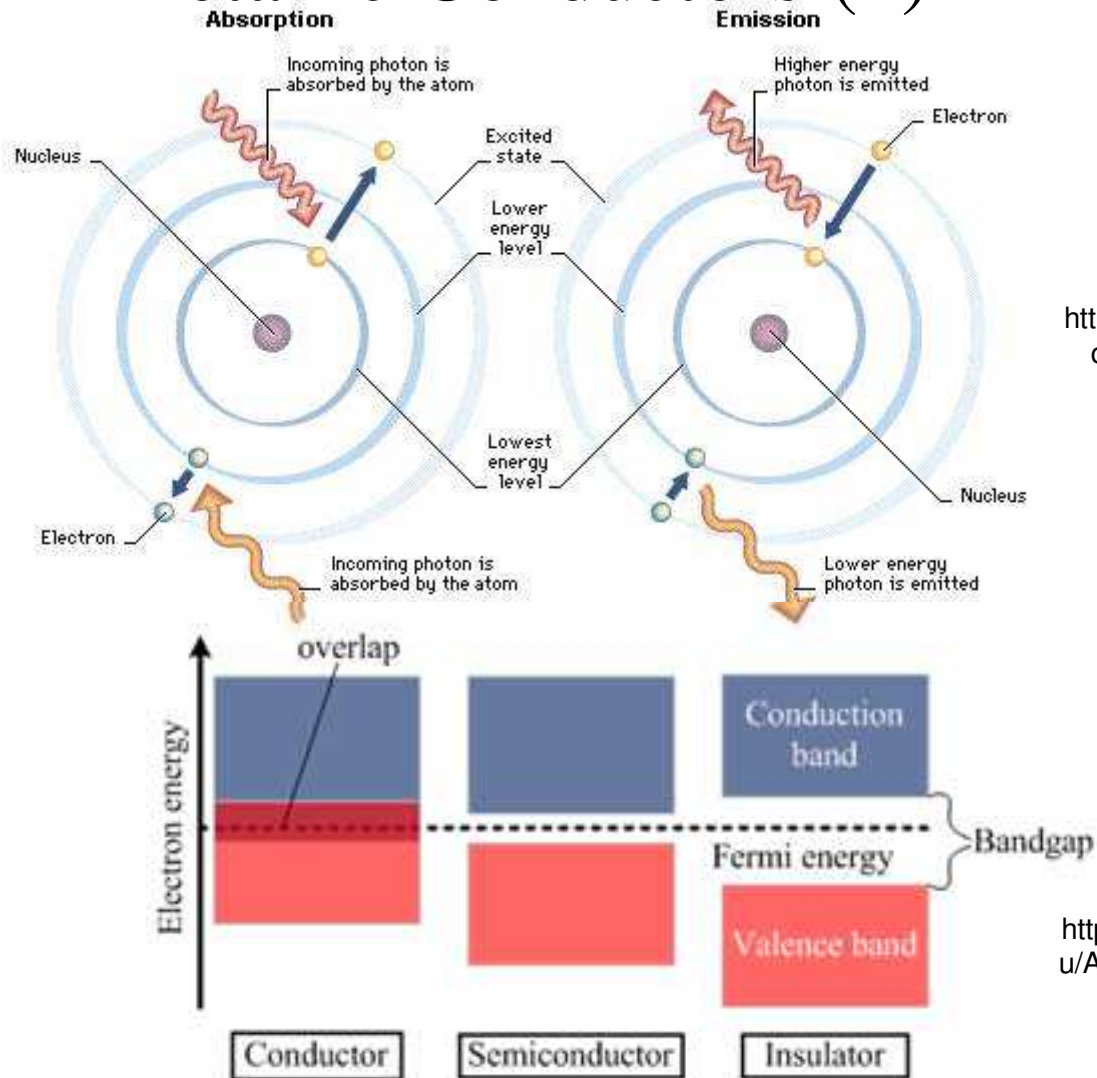
$$\rightarrow \rho_v = \frac{e^{-t}}{r^2} \text{ C/m}^3 \quad \rightarrow v_r = \frac{J_r}{\rho_v} = \left(\frac{e^{-t}}{r} \right) / \left(\frac{e^{-t}}{r^2} \right) = r \text{ m/s}$$

Current & Conductors

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Metallic Conductors (1)



<http://samuelsolema.blogspot.com/2013/03/hw-10-quantum-theory.html>

<http://potential.eecs.utk.edu/About.php?topic=PowerSemiconductors>

Metallic Conductors (2)

$$\mathbf{F} = -e\mathbf{E}$$

- In free space, the electron will accelerate
- In conductors, the electron will soon obtain a constant average velocity:

$$\mathbf{v}_d = -\mu_e \mathbf{E}$$

- μ_e : the mobility of an electron, m^2/Vs , positive
- Ex.: Al: 0.0012; Cu: 0.0032; Ag: 0.0056
- $\mathbf{J} = \rho_v \mathbf{v}$
- $\rightarrow \boxed{\mathbf{J} = -\rho_e \mu_e \mathbf{E}}$

Metallic Conductors (3)

$$\mathbf{J} = -\rho_e \mu_e \mathbf{E}$$

- ρ_e : free-electron charge density, negative
- \mathbf{J} is in the same direction as \mathbf{E}

$$\mathbf{J} = \sigma \mathbf{E}$$

- σ : conductivity, S/m
- Ex.: Al: 3.82×10^7 ; Cu: 5.80×10^7 ; Ag: 6.17×10^7

$$\sigma = -\rho_e \mu_e$$

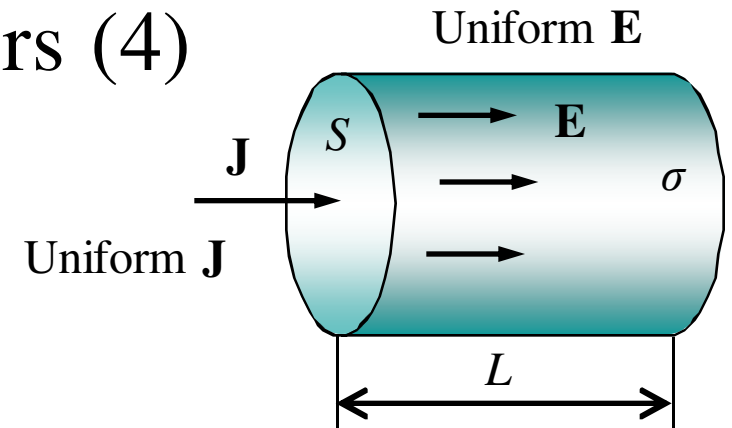
Metallic Conductors (4)

$$\begin{aligned}
 I &= \int_S \mathbf{J} \cdot d\mathbf{S} = JS \rightarrow J = \frac{I}{S} \\
 V_{ab} &= -\int_b^a \mathbf{E} \cdot d\mathbf{L} \\
 &= -\mathbf{E} \cdot \int_b^a d\mathbf{L} \\
 &= -\mathbf{E} \cdot \mathbf{L}_{ba} = \mathbf{E} \cdot \mathbf{L}_{ab} \\
 \rightarrow V &= EL \\
 J &= \sigma E \rightarrow J = \sigma \frac{V}{L}
 \end{aligned}
 \left. \begin{array}{l} \\ \\ \\ \\ \end{array} \right\} \rightarrow \sigma \frac{V}{L} = \frac{I}{S} \rightarrow V = \frac{L}{\sigma S} I$$

$$R = \frac{L}{\sigma S}$$

$\rightarrow V = RI$
 (Ohm's law)

$$R = \frac{V_{ab}}{I} = \frac{-\int_b^a \mathbf{E} \cdot d\mathbf{L}}{\int_S \sigma \mathbf{E} \cdot d\mathbf{S}}$$

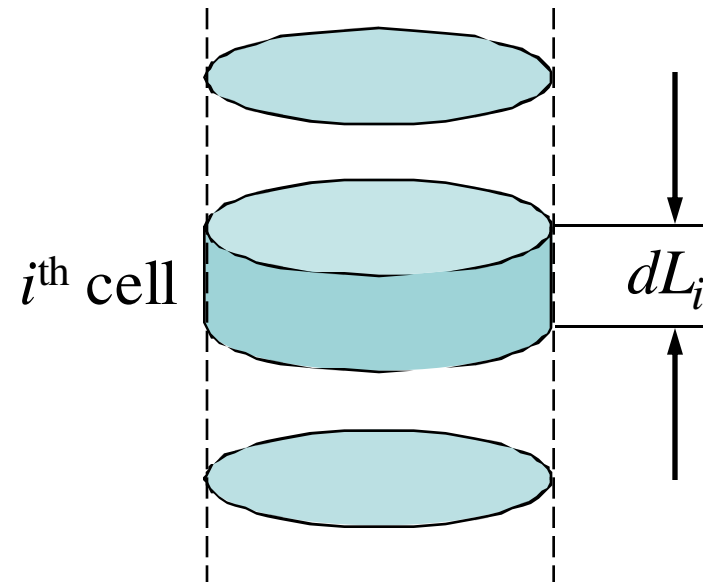


Metallic Conductors (5)

$$R_i = \frac{dL_i}{\sigma_i S_i}$$

$$\rightarrow R = \sum_{i=1}^N R_i = \sum_{i=1}^N \frac{dL_i}{\sigma_i S_i}$$

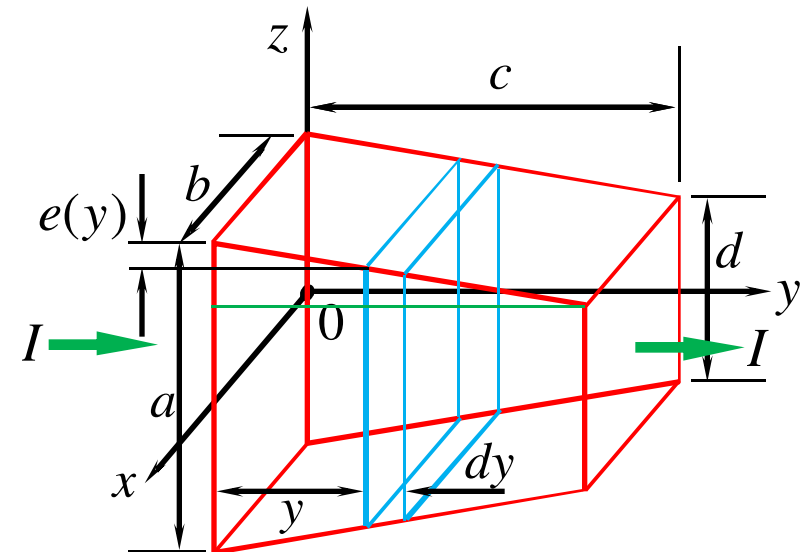
$$\rightarrow R = \int \frac{dL}{\sigma S}$$



Ex. 1

Metallic Conductors (6)

$$\left. \begin{aligned} dR &= \frac{dy}{\sigma S(y)} \\ S(y) &= b[a - 2e(y)] \\ \frac{e(y)}{y} &= \frac{a-d}{2c} \rightarrow e(y) = \frac{a-d}{2c} y \end{aligned} \right\}$$



$$\rightarrow dR = \frac{dy}{\sigma b \left[a - 2 \frac{a-d}{2c} y \right]} = \frac{c}{\sigma b [ac - (a-d)y]} dy$$

$$\rightarrow R = \int_{y=0}^c \frac{c}{\sigma b [ac - (a-d)y]} dy$$

**Ex. 2**

Metallic Conductors (7)

A material with conductivity $\sigma = m/\rho + k$, where m & k are constants, fills the space between two concentric, cylindrical conductors of radii a & b . L is the length of each conductor. Find the resistance of the material?

$$\left. \begin{aligned} R_i &= \frac{dL_i}{\sigma_i S_i} \\ dL_i &= d\rho \\ \sigma_i &= \frac{m}{\rho} + k \\ S_i &= 2\pi\rho L \end{aligned} \right\} \rightarrow R = \int_a^b \frac{d\rho}{(k\rho + m)2\pi L} = \frac{1}{2\pi L k} \ln \frac{kb + m}{ka + m}$$

Ex. 3

Metallic Conductors (8)

$$R = \frac{V}{I} = \frac{-\int_b^a \mathbf{E} \cdot d\mathbf{L}}{\int_S \sigma \mathbf{E} \cdot d\mathbf{S}}$$

$$\mathbf{E} = E(\rho) \mathbf{a}_\phi$$

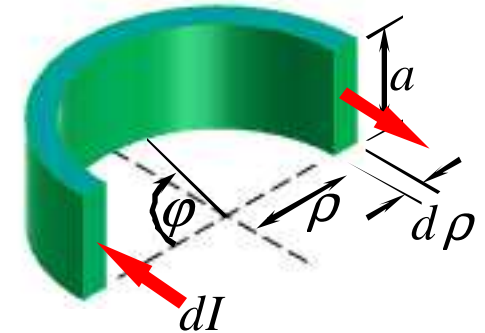
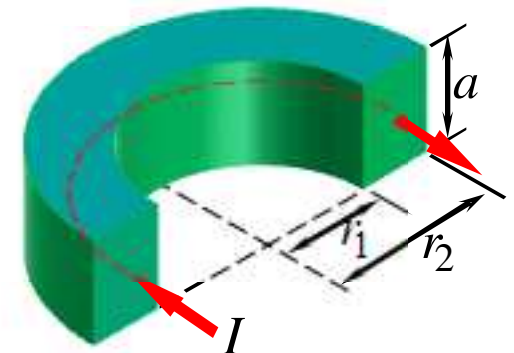
$$V = \int_{\phi=0}^{\pi} [E(\rho) \mathbf{a}_\phi] \cdot (\rho d\phi \mathbf{a}_\phi)$$

$$= \int_{\phi=0}^{\pi} E(\rho) \rho d\phi = E(\rho) \rho \int_0^{\pi} d\phi = E(\rho) \rho \pi \rightarrow E(\rho) = \frac{V}{\pi \rho}$$

$$I = \int_S \sigma \mathbf{E} \cdot d\mathbf{S} = \int_{\rho=r_1}^{r_2} [\sigma E(\rho) \mathbf{a}_\phi] \cdot (a d\rho \mathbf{a}_\phi) = \int_{\rho=r_1}^{r_2} \sigma E(\rho) a d\rho$$

$$= \int_{\rho=r_1}^{r_2} \sigma \frac{V}{\pi \rho} a d\rho = \frac{\sigma V a}{\pi} \ln \frac{r_2}{r_1}$$

$$\rightarrow R = \frac{V}{\frac{\sigma V a}{\pi} \ln \frac{r_2}{r_1}} = \boxed{\frac{\pi}{\sigma a \ln \frac{r_2}{r_1}}}$$



Metallic Conductors (9)

Ex. 4

$$R = \frac{V}{I} = \frac{-\int_b^a \mathbf{E} \cdot d\mathbf{L}}{\int_S \sigma \mathbf{E} \cdot d\mathbf{S}}$$

(Method 1)

$$\nabla \cdot \mathbf{D} = \rho_v = 0 \rightarrow \nabla \cdot \epsilon \mathbf{E} = 0$$

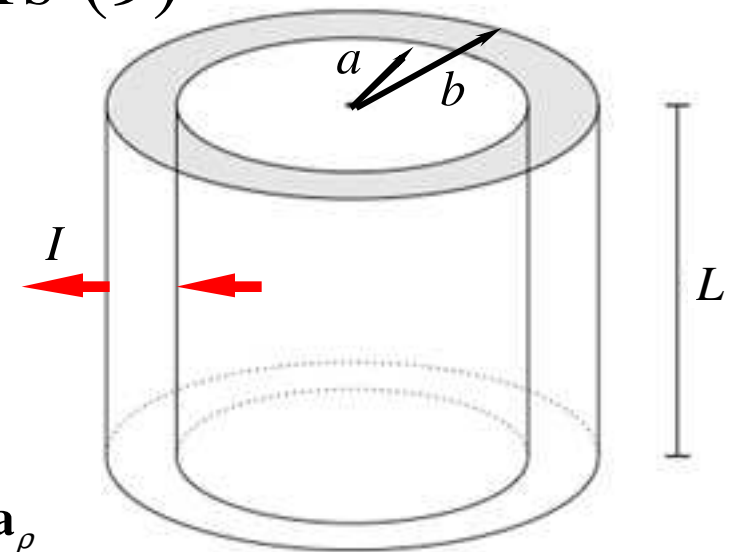
$$\left. \begin{aligned} \nabla \cdot \mathbf{D} &= \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho D_\rho) + \frac{1}{\rho} \frac{\partial D_\phi}{\partial \phi} + \frac{\partial D_z}{\partial z} \\ \mathbf{E} &= E_\rho \mathbf{a}_\rho \end{aligned} \right\} \rightarrow \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho \epsilon E_\rho) = 0$$

$$\rightarrow E_\rho = \frac{C}{\rho} \rightarrow \mathbf{E} = \frac{C}{\rho} \mathbf{a}_\rho$$

$$V = \int_a^b \mathbf{E} \cdot d\mathbf{L} = \int_a^b \frac{C}{\rho} \mathbf{a}_\rho \cdot d\rho \mathbf{a}_\rho = C \ln \frac{b}{a} \rightarrow C = \frac{V}{\ln(b/a)} \rightarrow \mathbf{E} = \frac{V}{\rho \ln(b/a)} \mathbf{a}_\rho$$

$$I = \int_S \sigma \mathbf{E} \cdot d\mathbf{S} = \int_{z=0}^L \int_{\phi=0}^{2\pi} \sigma \frac{V}{\rho \ln(b/a)} \mathbf{a}_\rho \cdot \rho d\phi dz \mathbf{a}_\rho = \frac{\sigma V 2\pi L}{\ln(b/a)}$$

$$\rightarrow R = \frac{V}{I} = \frac{V}{\frac{\sigma V 2\pi L}{\ln(b/a)}} = \frac{\ln(b/a)}{2\pi\sigma L}$$



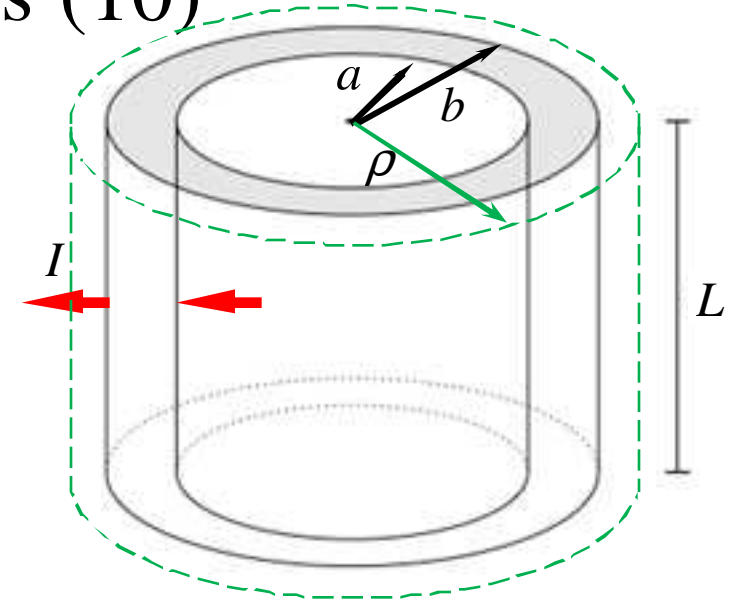
Metallic Conductors (10)

Ex. 4

$$R = \frac{V}{I} = \frac{-\int_b^a \mathbf{E} \cdot d\mathbf{L}}{\int_S \sigma \mathbf{E} \cdot d\mathbf{S}} \quad (\text{Method 2})$$

$$\left. \begin{aligned} Q &= \oint_S \mathbf{D} \cdot d\mathbf{S} \\ \mathbf{E} &= E_\rho \mathbf{a}_\rho \end{aligned} \right\} \rightarrow Q = \oint_S \epsilon E_\rho \mathbf{a}_\rho \cdot d\mathbf{S} \\ = \epsilon E_\rho (2\pi\rho L) \\ \rightarrow \mathbf{E} = \frac{Q}{2\pi\epsilon\rho L} \mathbf{a}_\rho$$

$$\left. \begin{aligned} V &= \int_a^b \mathbf{E} \cdot d\mathbf{L} = \int_a^b \frac{Q}{2\pi\epsilon\rho L} \mathbf{a}_\rho \cdot d\rho \mathbf{a}_\rho = \frac{Q}{2\pi\epsilon L} \ln \frac{b}{a} \\ I &= \int_S \sigma \mathbf{E} \cdot d\mathbf{S} = \int_{z=0}^L \int_{\phi=0}^{2\pi} \sigma \frac{Q}{2\pi\epsilon\rho L} \mathbf{a}_\rho \cdot \rho d\phi dz \mathbf{a}_\rho = \frac{\sigma Q}{\epsilon} \end{aligned} \right\} \rightarrow R = \frac{V}{I} = \frac{\frac{Q}{2\pi\epsilon L} \ln \frac{b}{a}}{\frac{\sigma Q}{\epsilon}} = \boxed{\frac{\ln(b/a)}{2\pi\sigma L}}$$



Current & Conductors

1. Current & Current Density
2. Metallic Conductors
- 3. Conductor Properties & Boundary Conditions**
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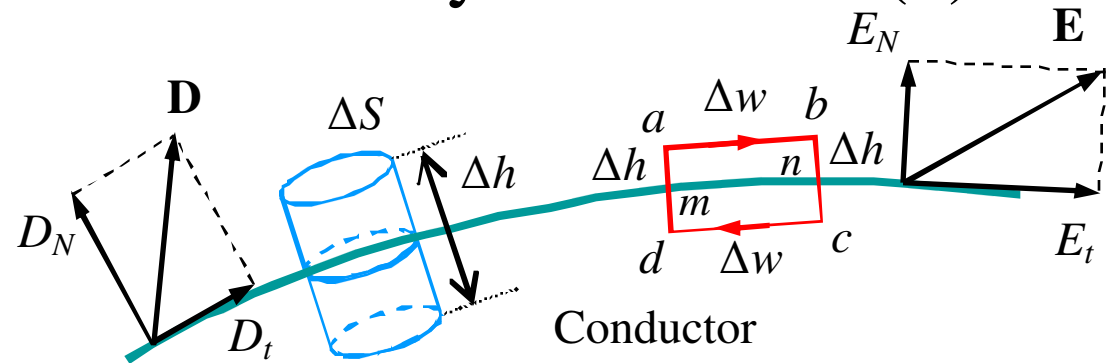
Conductor Properties & Boundary Conditions (1)

- Given some electrons in the interior of a conductor
- They will begin to accelerate away from each other, until they reach the surface of the conductor
- Characteristic 1: the charge density inside a conductor is zero, the exterior surface has a surface charge density
- Within a conductor: no charge \rightarrow no current \rightarrow no electric field intensity (Ohm)
- Characteristic 2: the electric field intensity within the conductor is zero



Conductor Properties & Boundary Conditions (2)

$$\oint \mathbf{E} \cdot d\mathbf{L} = 0$$



$$\rightarrow \int_a^b + \int_b^c + \int_c^d + \int_d^a = 0$$

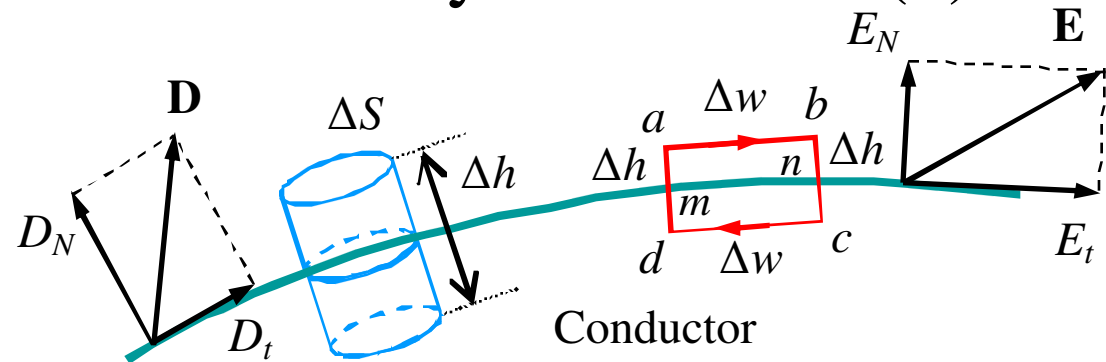
$$\left. \begin{array}{l} \int_a^b \mathbf{E}_{ab} \cdot d\mathbf{L}_{ab} \\ \mathbf{E}_{ab} = \mathbf{E}_N + \mathbf{E}_t \end{array} \right\} \rightarrow \int_a^b \mathbf{E}_{ab} \cdot d\mathbf{L}_{ab} = \int_a^b (\mathbf{E}_N + \mathbf{E}_t) \cdot d\mathbf{L}_{ab} = \int_a^b E_t dL_{ab} \left. \begin{array}{l} \\ E_t \approx \text{const} \end{array} \right\}$$

$$\rightarrow \int_a^b \mathbf{E}_{ab} \cdot d\mathbf{L}_{ab} = E_t \int_a^b dL_{ab} = E_t \Delta w$$

$$\left. \begin{array}{l} \int_c^d \mathbf{E}_{cd} \cdot d\mathbf{L}_{cd} \\ \mathbf{E}_{\text{within conductor}} = 0 \end{array} \right\} \rightarrow \mathbf{E}_{cd} = 0 \rightarrow \int_c^d \mathbf{E}_{cd} \cdot d\mathbf{L}_{cd} = 0$$

Conductor Properties & Boundary Conditions (3)

$$\oint \mathbf{E} \cdot d\mathbf{L} = 0$$



$$\rightarrow \int_a^b + \int_b^c + \int_c^d + \int_d^a = 0$$

$$\left. \begin{aligned} \int_b^c \mathbf{E}_{bc} \cdot d\mathbf{L}_{bc} &= \int_b^n \mathbf{E}_{bn} \cdot d\mathbf{L}_{bn} + \int_n^c \mathbf{E}_{nc} \cdot d\mathbf{L}_{nc} \\ \mathbf{E}_{\text{within conductor}} &= 0 \rightarrow \mathbf{E}_{nc} = 0 \\ \mathbf{E}_{bn} &= \mathbf{E}_{N,b} + \mathbf{E}_{tt} \end{aligned} \right\} \rightarrow \left. \begin{aligned} \int_b^c \mathbf{E}_{bc} \cdot d\mathbf{L}_{bc} &= - \int_b^n E_{N,b} dL_{bn} \\ E_{N,b} &\approx \text{const} \end{aligned} \right\}$$

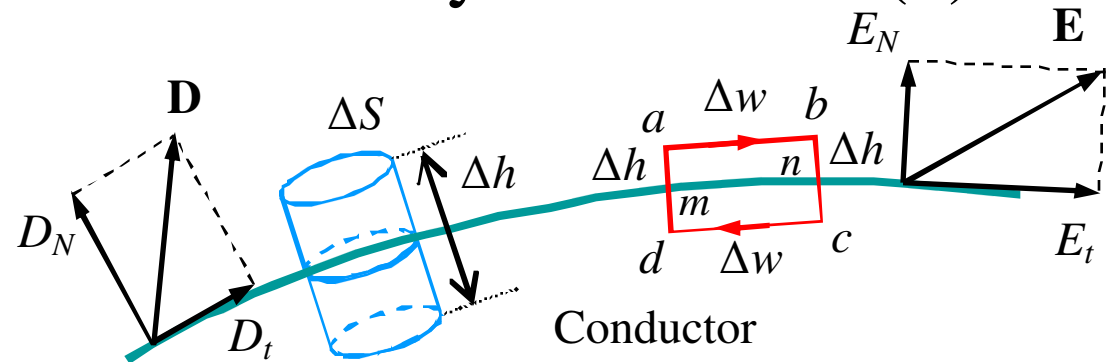
$$\rightarrow \int_b^c \mathbf{E}_{bc} \cdot d\mathbf{L}_{bc} = -E_{N,b} \int_b^n dL_{bn} = -\frac{E_{N,b} \Delta h}{2}$$

$$\int_d^a \mathbf{E}_{da} \cdot d\mathbf{L}_{da} = \frac{E_{N,a} \Delta h}{2}$$



Conductor Properties & Boundary Conditions (4)

$$\oint \mathbf{E} \cdot d\mathbf{L} = 0$$



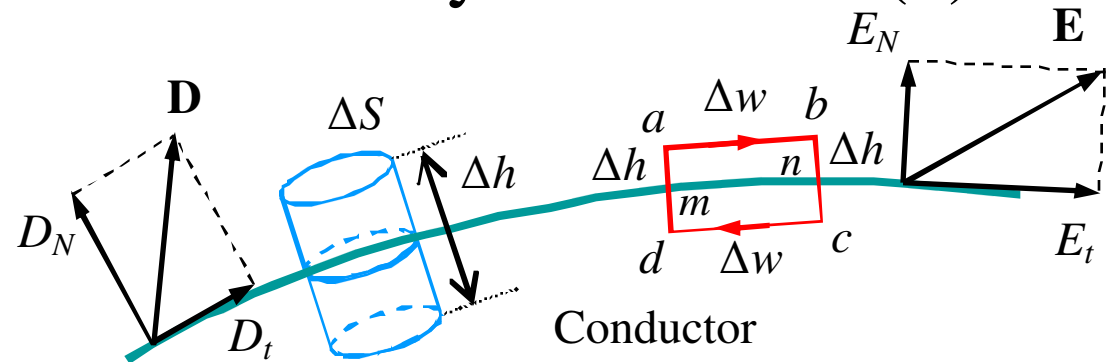
$$\left. \begin{aligned} \rightarrow \int_a^b + \int_b^c + \int_c^d + \int_d^a &= 0 \\ \int_a^b &= E_{tt} \Delta w \\ \int_b^c &= -\frac{E_{N,b} \Delta h}{2} \\ \int_c^d &= 0 \\ \int_d^a &= \frac{E_{N,a} \Delta h}{2} \end{aligned} \right\} \rightarrow E_t \Delta w - \frac{E_{N,b} \Delta h}{2} + 0 + \frac{E_{N,a} \Delta h}{2} = 0$$

$$\Delta h \rightarrow 0 \left\{ \begin{aligned} \rightarrow E_t \Delta w &= 0 \rightarrow E_t = 0 \\ \rightarrow D_t = \epsilon_0 E_t &= 0 \rightarrow \boxed{D_t = E_t = 0} \end{aligned} \right.$$

Conductor Properties & Boundary Conditions (5)

$$E_{tt} = 0$$

$$\oint_S \mathbf{D} \cdot d\mathbf{S} = Q$$



$$\rightarrow \int_{\text{top}} + \int_{\text{bottom}} + \int_{\text{sides}} = \rho_S \Delta S$$

$$\int_{\text{top}} = \int_{\text{top}} \mathbf{D}_N \cdot d\mathbf{S}_{\text{top}} = \int_{\text{top}} D_N dS_{\text{top}} = D_N \Delta S$$

$$\int_{\text{bottom}} = \int_{\text{bottom}} 0 \cdot d\mathbf{S}_{\text{bottom}} = 0$$

$$\int_{\text{sides}} = \int_{\text{s, top}} \mathbf{D}_N \cdot d\mathbf{S}_{\text{s, top}} + \int_{\text{s, bottom}} 0 \cdot d\mathbf{S}_{\text{s, bottom}} = 0$$

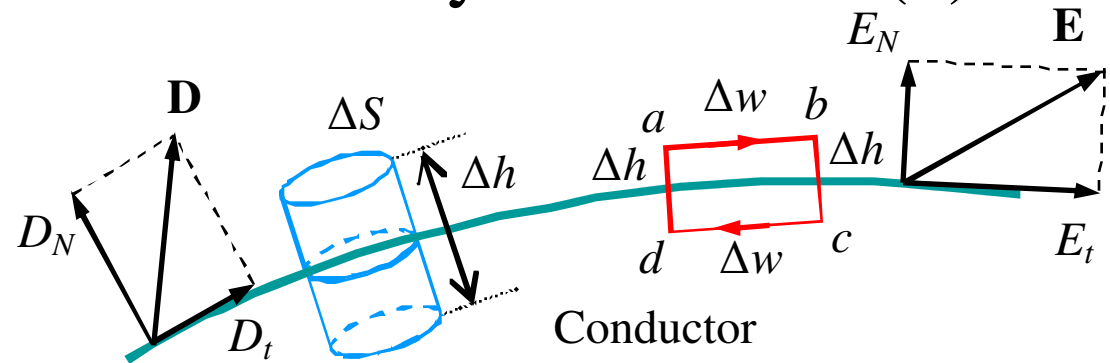
$$\rightarrow D_N \Delta S = \rho_S \Delta S \rightarrow \boxed{D_N = \rho_S = \epsilon_0 E_N}$$

Conductor Properties & Boundary Conditions (6)

$$D_t = E_t = 0$$

$$D_N = \epsilon_0 E_N = \rho_S$$

$$V_{xy} = -\int_y^x \mathbf{E} \cdot d\mathbf{L} = 0$$



Characteristics of conductors in static field:

1. The static EFI inside a conductor is zero
2. The static EFI at the surface of a conductor is everywhere directed normal to that surface
3. The conductor surface is an equipotential surface

Conductor Properties & Boundary Conditions (7)

Ex.

Given $V = x^2 - 10yz$ V & $P(2, 1, 2)$ lies on a conductor – free space boundary. Find V , \mathbf{E} , \mathbf{D} , ρ_S at P , & the equation of the conductor surface.

$$V_P = 2^2 - 10 \times 1 \times 2 = -16 \text{ V} \quad \rightarrow -16 = x^2 - 10yz$$

$$\mathbf{E} = -\nabla V = -\nabla(x^2 - 10yz) = -2x\mathbf{a}_x + 10z\mathbf{a}_y + 10y\mathbf{a}_z \text{ V/m}$$

$$\rightarrow \mathbf{E}_P = (-2x\mathbf{a}_x + 10z\mathbf{a}_y + 10y\mathbf{a}_z) \Big|_{x=2, y=1, z=2}$$

$$= -40\mathbf{a}_x + 20\mathbf{a}_y + 10\mathbf{a}_z \text{ V/m}$$

$$\mathbf{D}_P = \epsilon_0 \mathbf{E}_P = 8.854 \times 10^{-12} (-40\mathbf{a}_x + 20\mathbf{a}_y + 10\mathbf{a}_z) \text{ C/m}^2$$

$$\rho_{S,P} = D_N$$

$$D_{N,P} = |\mathbf{D}_P| = 8.854 \times 10^{-12} \sqrt{40^2 + 20^2 + 10^2} = 406 \text{ pC/m}^2 \left. \vphantom{\begin{matrix} \rho_{S,P} = D_N \\ D_{N,P} = |\mathbf{D}_P| \end{matrix}} \right\} \rightarrow \rho_{S,P} = 406 \text{ pC/m}^2$$

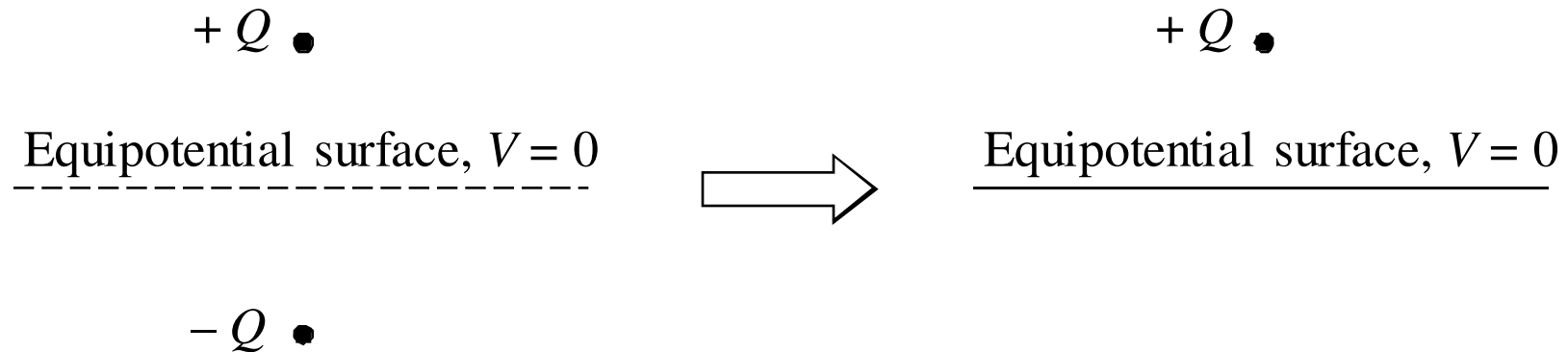
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3. Conductor Properties & Boundary Conditions
- 4. The Method of Images**
5. Semiconductors
6. Applications



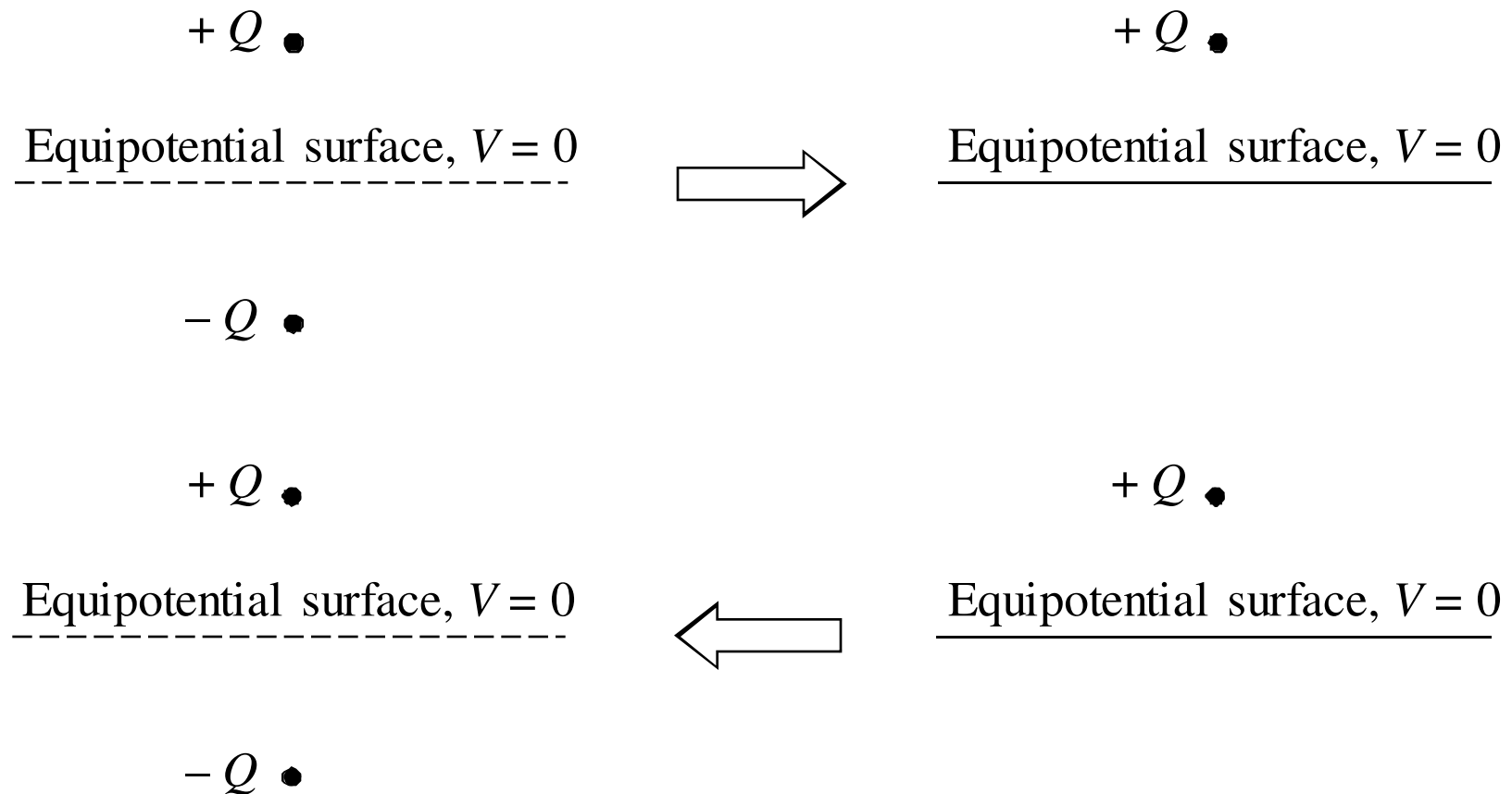


The Method of Images (1)



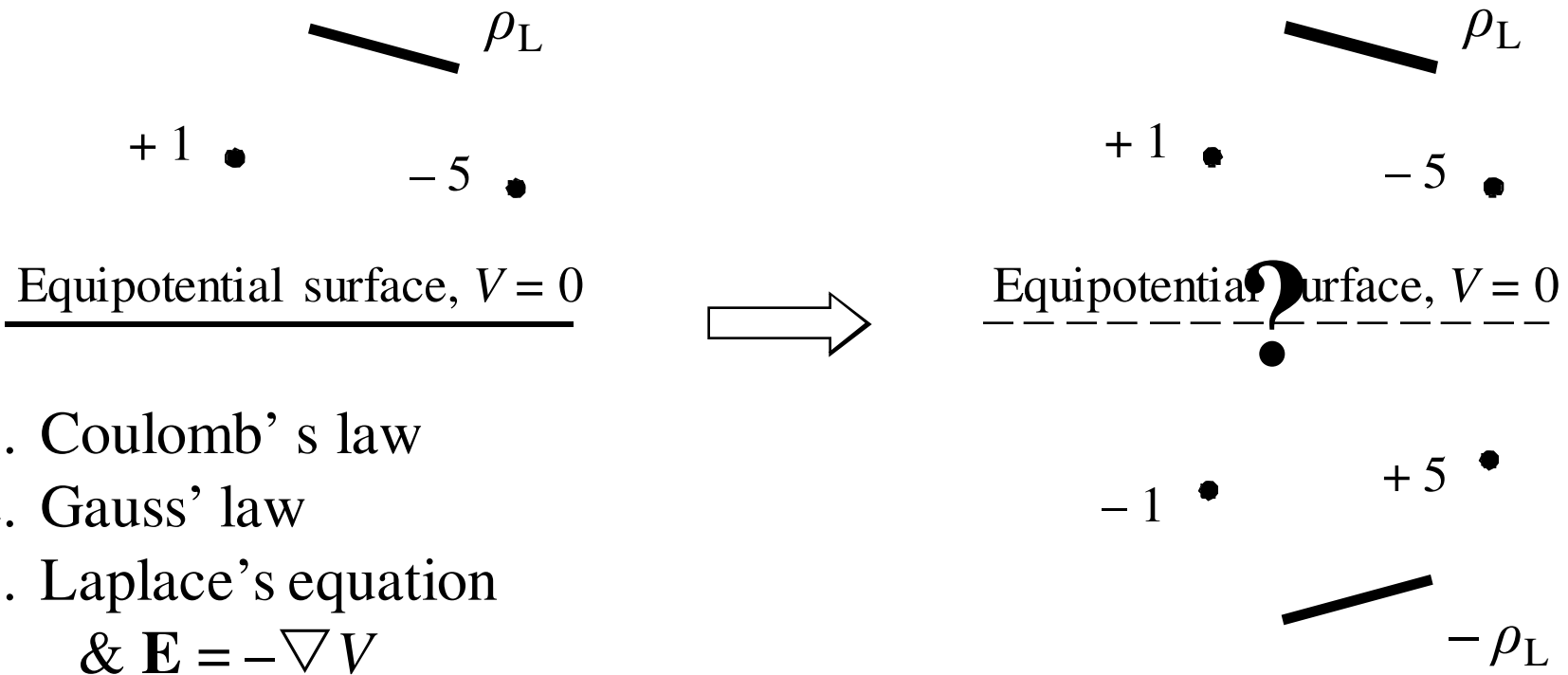
- Dipole: the plane between the 2 charges is zero potential
- That plane can be represented by a vanishingly thin conducting plane, infinite in extent
- \rightarrow the dipole can be substituted for a system of a charge and a conducting plane, & then the fields above the conducting plane obtain equivalence

The Method of Images (2)



Ex. 1

The Method of Images (3)



1. Coulomb's law
2. Gauss' law
3. Laplace's equation
& $\mathbf{E} = -\nabla V$

Ex. 2

The Method of Images (4)

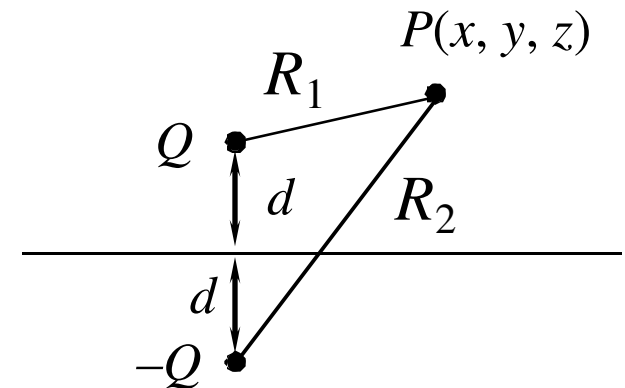
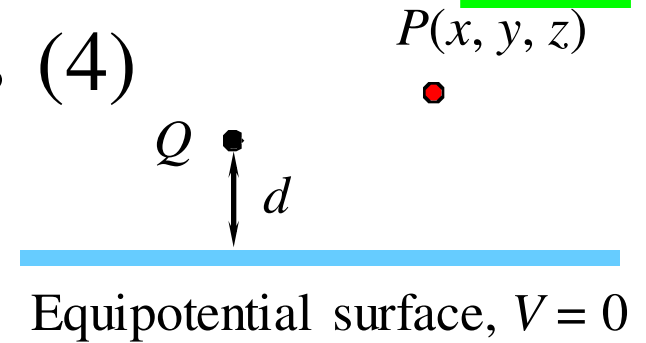
Given Q at $(0, 0, d)$. Find the potential & EFI at P ?

$$V_{+Q} = \frac{Q}{4\pi\epsilon_0 R_1} = \frac{Q}{4\pi\epsilon_0 \sqrt{x^2 + y^2 + (z-d)^2}}$$

$$V_{-Q} = \frac{-Q}{4\pi\epsilon_0 R_2} = \frac{-Q}{4\pi\epsilon_0 \sqrt{x^2 + y^2 + (z+d)^2}}$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{\sqrt{x^2 + y^2 + (z-d)^2}} - \frac{1}{\sqrt{x^2 + y^2 + (z+d)^2}} \right]$$

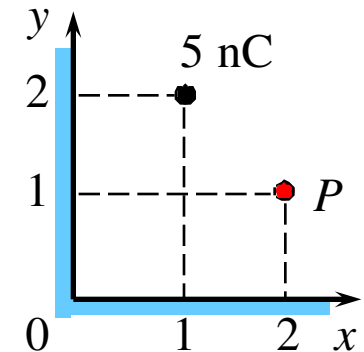
$$\mathbf{E} = -\nabla V = -\frac{Q}{4\pi\epsilon_0} \left[\left(\frac{x}{R_2^3} - \frac{x}{R_1^3} \right) \mathbf{a}_x + \left(\frac{y}{R_2^3} - \frac{y}{R_1^3} \right) \mathbf{a}_y + \left(\frac{z+d}{R_2^3} - \frac{z-d}{R_1^3} \right) \mathbf{a}_z \right]$$



Ex. 3

The Method of Images (5)

Find the potential at P ?



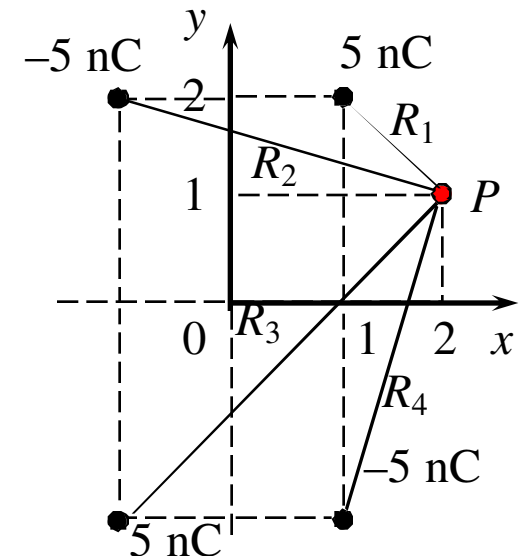
$$R_1 = \sqrt{1^2 + 1^2} = 1.41$$

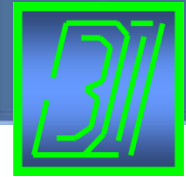
$$R_2 = \sqrt{3^2 + 1^2} = 3.16$$

$$R_3 = \sqrt{3^2 + 3^2} = 4.24$$

$$R_4 = \sqrt{1^2 + 3^2} = 3.16$$

$$V_P = \frac{5 \times 10^{-9}}{4\pi\epsilon_0} \left(\frac{1}{R_1} - \frac{1}{R_2} + \frac{1}{R_3} - \frac{1}{R_4} \right) = 14.03 \text{ V}$$





Ex. 4

The Method of Images (6)

A point charge Q at a distance d from a center of a grounded conducting sphere of radius a . Find the image charge?

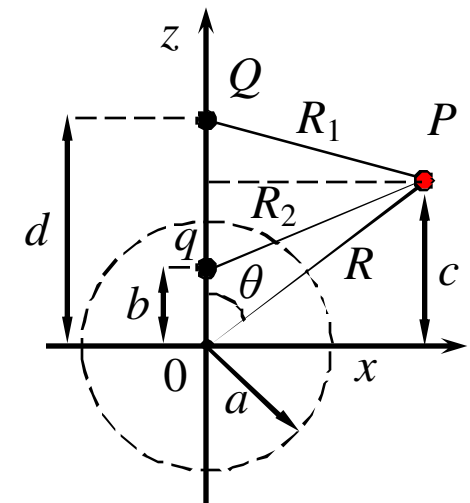
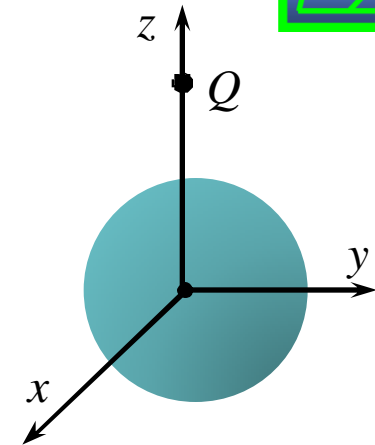
Problem: find q & b

$$R_1 = \sqrt{(d - R \cos \theta)^2 + (R \sin \theta)^2} = \sqrt{R^2 + d^2 - 2Rd \cos \theta}$$

$$R_2 = \sqrt{(R \cos \theta - b)^2 + (R \sin \theta)^2} = \sqrt{R^2 + b^2 - 2Rb \cos \theta}$$

$$V_P = \frac{Q}{4\pi\epsilon R_1} - \frac{q}{4\pi\epsilon R_2} = \frac{Q}{4\pi\epsilon R_1} - \frac{mQ}{4\pi\epsilon R_2} = \frac{Q}{4\pi\epsilon} \left(\frac{1}{R_1} - \frac{m}{R_2} \right)$$

$$R = a \rightarrow V_P = 0 \rightarrow \frac{1}{R_1} - \frac{m}{R_2} = 0 \rightarrow \begin{cases} m = \frac{a}{d} \rightarrow q = -\frac{a}{d}Q \\ b = \frac{a^2}{d} \end{cases}$$



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Semiconductors

- Germani, silicon
- Conductivity of conductors:

$$\sigma = -\rho_e \mu_e$$

- Conductivity of semiconductors:

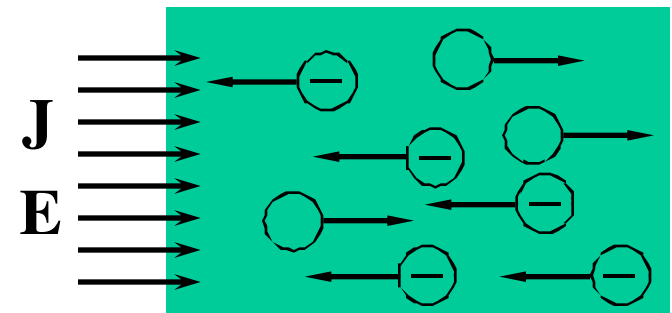
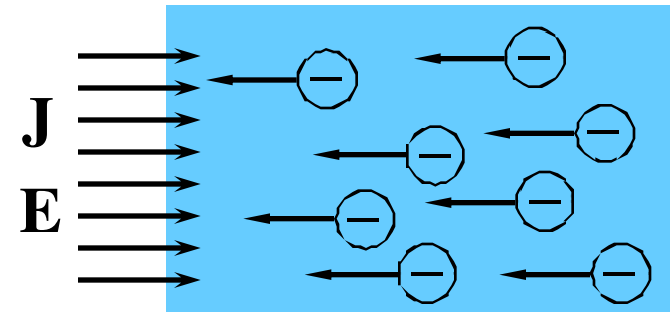
$$\sigma = -\rho_e \mu_e + \rho_h \mu_h$$

- h : hole

- At 300K:

– $\mu_{e, \text{Germani}}: 0.36 \text{ m}^2/\text{Vs}; \quad \mu_{h, \text{Germani}}: 0.17 \text{ m}^2/\text{Vs}$

– $\mu_{e, \text{Silicon}}: 0.12 \text{ m}^2/\text{Vs}; \quad \mu_{h, \text{Silicon}}: 0.025 \text{ m}^2/\text{Vs}$



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Applications – Faraday's Cage



<https://lifeonthebluehighways.com/2013/04/20/faradays-cage/>

Current & Conductors - sites.google.com/site/ncpdhbkhn

$$\begin{array}{ccccccc}
 Q & \longrightarrow & \mathbf{F} = \frac{Q_1 Q_2}{4\pi\epsilon R^2} \mathbf{a}_R & \longrightarrow & \mathbf{E} = \frac{Q}{4\pi\epsilon R^2} \mathbf{a}_R & \longrightarrow & \mathbf{D} = \epsilon \mathbf{E} \\
 \downarrow & & \downarrow & & & & \\
 & & W = -Q \int \mathbf{E} \cdot d\mathbf{L} & \longrightarrow & V = -\int \mathbf{E} \cdot d\mathbf{L} & \longrightarrow & C = \frac{Q}{V} \\
 \downarrow & & & & & & \\
 I = \frac{dQ}{dt} & \longrightarrow & R = \frac{V}{I} & & & &
 \end{array}$$