

Q1:

0.4 We divide the semicylinder into many vertical small infinite line charges: dq

$$dq = \lambda_s ds = \lambda_s dz$$

$$\Rightarrow \int \lambda_s R dq dz = \int \lambda_s dz$$

$$\Rightarrow \int \lambda_s R dq = \lambda_s$$

+ We have the EFI of a infinite line charge:

$$d\vec{E} = \frac{\lambda_s}{2\pi\epsilon_0 R} \vec{a}_p = \frac{\lambda_s dy R}{2\pi\epsilon_0 R} \vec{a}_p = \frac{\lambda_s dy}{2\pi\epsilon_0} \vec{a}_p$$

$$+ d\vec{E} = d\vec{E}_1 + d\vec{E}_2$$

$$\Rightarrow dE = 2dE_1 \cos y = \frac{2\cos y \lambda_s dy}{2\pi\epsilon_0} = \frac{\lambda_s \cos y dy}{\pi\epsilon_0}$$

$$\Rightarrow E = \int_0^{\pi/2} \frac{\lambda_s \cos y dy}{\pi\epsilon_0} = \frac{\lambda_s \sin y}{\pi\epsilon_0} \Big|_0^{\pi/2} = \frac{\lambda_s}{\pi\epsilon_0}$$

From the sketch we can use that \vec{E} is parallel with a axis in opposite direction

$$\Rightarrow \vec{E} = \frac{\lambda_s}{\pi\epsilon_0} (-\vec{a}_x)$$

b, We have: $dq = \lambda_s ds = \lambda_v dv$

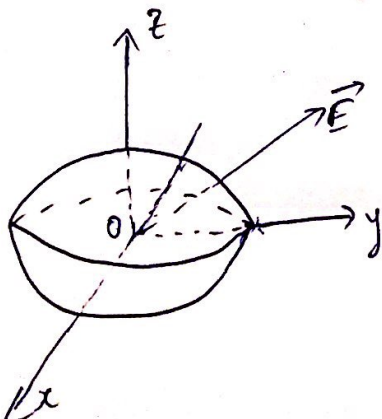
$$\Rightarrow \int \rho dq dz = \int \lambda_v \rho d\rho d\phi dz$$

$$\Rightarrow \lambda_s = \lambda_v \rho$$

$$dE = \frac{\lambda_v \rho}{\pi\epsilon_0} \Rightarrow E = \int_0^R \frac{\lambda_v \rho d\rho}{\pi\epsilon_0} = \frac{R \lambda_v}{\pi\epsilon_0}$$

$$\Rightarrow \vec{E} = (-\vec{a}_x) \frac{R \lambda_v}{\pi\epsilon_0}$$

Q2



+ For the sphere of radius R , the EFI of the surface will be: $\vec{E} = E \vec{a}_r = AR \vec{a}_r$ ($\lambda = R$)

+ Apply Gauss law: $Q = \oint_{\text{sphere}} \vec{D}_s \cdot d\vec{s} \Rightarrow \frac{Q}{\epsilon_0} = \oint_{\text{sphere}} \vec{E}_s \cdot d\vec{s}$

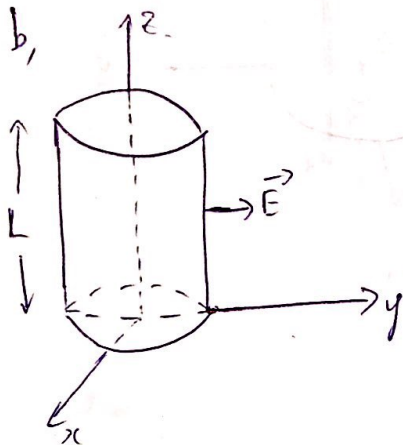
+ Since sphere has \vec{E}_s is everywhere normal to the closed surface.

$$\Rightarrow \vec{E}_s \cdot d\vec{s} = E_s ds$$

$$\Rightarrow Q = \oint_{\text{sphere}} \vec{E}_s \cdot d\vec{s} = \oint_{\text{sphere}} E_s ds$$

$$= E_s \oint_{\text{sphere}} ds$$

$$\Rightarrow Q = E_s 4\pi R^2$$



Apply Gauss's law:

$$Q = \oint \vec{D} \cdot d\vec{s} = \epsilon_0 \oint \vec{E} \cdot d\vec{s}$$

$$= \epsilon_0 \int_{\text{bottom}} \vec{E} \cdot d\vec{s} + \epsilon_0 \int_{\text{side}} \vec{E} \cdot d\vec{s} + \epsilon_0 \int_{\text{top}} \vec{E} \cdot d\vec{s}$$

$$\textcircled{1} Q_{\text{bottom}} = \epsilon_0 \int_{\text{bottom}} \vec{E} \cdot d\vec{s} = \epsilon_0 \int_{\text{bot}} R^2 (\vec{a}_p) \rho dp d\phi (-\vec{a}_z) = 0 \quad \textcircled{1}$$

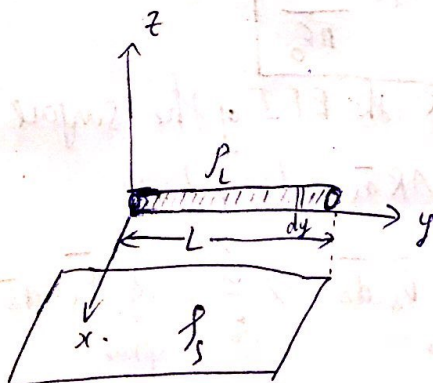
$(\vec{a}_p \cdot \vec{a}_z = 0)$

$$\textcircled{2} Q_{\text{top}} = \epsilon_0 \int_{\text{top}} \vec{E} \cdot d\vec{s} = \epsilon_0 \int_{\text{top}} R^2 (\vec{a}_p) \rho dp d\phi (\vec{a}_z) = 0 \quad (\vec{a}_p \cdot \vec{a}_z = 0) \quad \textcircled{2}$$

$$\textcircled{3} Q_{\text{side}} = \epsilon_0 \int_{\text{side}} R^2 (\vec{a}_p) \rho dp d\phi dz (\vec{a}_p) = \epsilon_0 \int_0^{2\pi} \int_0^L R^2 dz d\phi = \epsilon_0 \int_0^{2\pi} R^2 L d\phi = \epsilon_0 2\pi L R^2 \rho \quad \textcircled{3}$$

$$\textcircled{1} \textcircled{2} \textcircled{3} \Rightarrow Q = \epsilon_0 2\pi L R^2 \rho$$

Q3



① The electric field due to a infinite sheet with surface charge density ρ_s is:

$$\vec{E} = \frac{\rho_s}{2\epsilon_0} \vec{a}_z$$

We divide the line charge into many small pieces of charge dq , length dy

$$\Rightarrow dq = \lambda dy$$

We have work needed to rotate the pieces of charge dq is:

$$\begin{cases} dW = -dq \int_a^y \vec{E} \cdot d\vec{L} \\ d\vec{L} = dx \vec{a}_x + dy \vec{a}_y + dz \vec{a}_z \\ \vec{E} = \frac{\lambda}{2\epsilon_0} \vec{a}_z \end{cases}$$

$$\begin{aligned} \Rightarrow dW &= -dq \int_a^y \frac{\lambda}{2\epsilon_0} dz = -dq \int_0^y \frac{\lambda}{2\epsilon_0} dz \\ &= \frac{-\lambda \cdot dy \cdot y \cdot \lambda}{2\epsilon_0} \quad (dq = dy \lambda) \end{aligned}$$

$$\Rightarrow dW = \frac{-\lambda^2 y dy}{2\epsilon_0}$$

$$\Rightarrow W = \int_0^L \frac{-\lambda^2 y dy}{2\epsilon_0} = \boxed{\frac{-\lambda^2 L^2}{4\epsilon_0}}$$