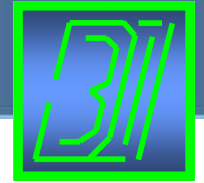




TRƯỜNG ĐẠI HỌC  
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# Engineering Electromagnetics

Time – Varying Fields & Maxwell's Equations

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# Time – Varying Fields & Maxwell's Equations

1. Faraday's Law
2. Displacement Current
3. Maxwell's Equations in Point Form
4. Maxwell's Equations in Integral Form
5. The Retarded Potentials

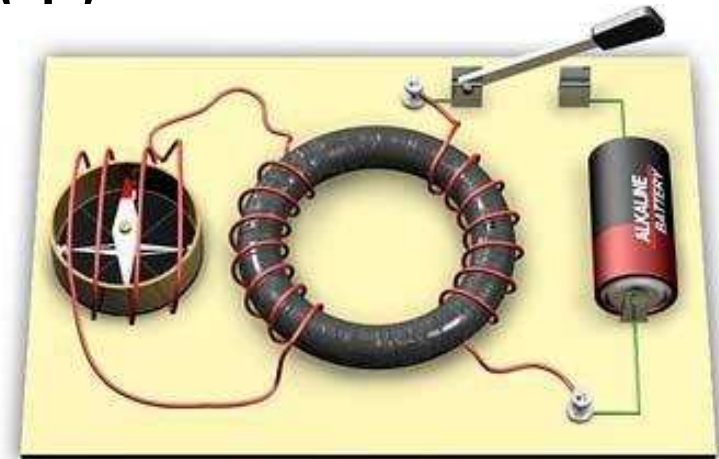


## Faraday's Law (1)

$$\text{emf} = -\frac{d\Phi}{dt} \text{ V}$$

Emf is nonzero if one of any:

- A time-changing flux linking a stationary closed path
- Relative motion between a steady flux and a closed path
- A combination of the two



<http://micro.magnet.fsu.edu/electromag/electricity/inductance.html>



Minus sign ?

Lenz's law

<http://www.engineering-timelines.com/how/electricity/transformer.asp>

## Faraday's Law (2)

$$\left. \begin{aligned} \text{emf} &= -\frac{d\Phi}{dt} \\ \text{emf} &= \oint \mathbf{E} \cdot d\mathbf{L} \\ \Phi &= \int_S \mathbf{B} \cdot d\mathbf{S} \end{aligned} \right\} \rightarrow \left. \begin{aligned} \text{emf} &= \oint \mathbf{E} \cdot d\mathbf{L} = -\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{S} \\ &\mathbf{B} = \mathbf{B}(t) \end{aligned} \right\}$$

$$\rightarrow \left. \begin{aligned} \text{emf} &= \oint \mathbf{E} \cdot d\mathbf{L} = -\int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} \end{aligned} \right\}$$

Stokes' theorem:  $\oint \mathbf{E} \cdot d\mathbf{L} = \int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{S}$

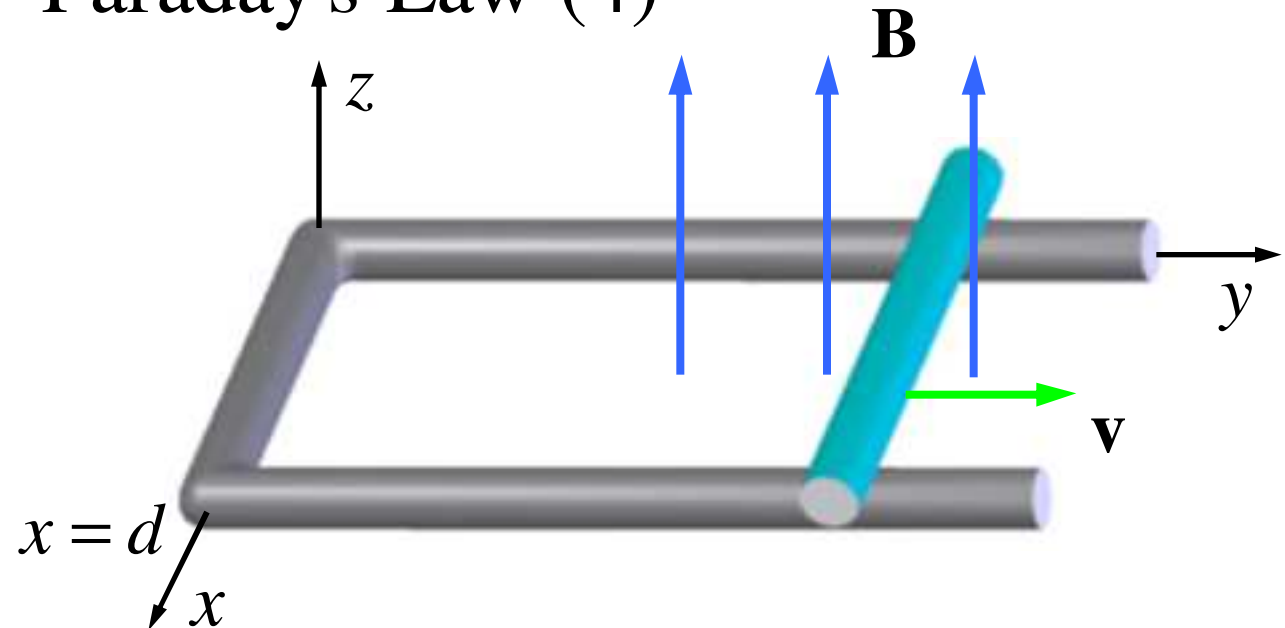
$$\rightarrow \int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{S} = -\int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} \rightarrow (\nabla \times \mathbf{E}) \cdot d\mathbf{S} = -\frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S}$$

$$\rightarrow \boxed{\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}}$$

## Faraday's Law (3)

$$\left. \begin{aligned} \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} \\ \text{emf} = \oint \mathbf{E} \cdot d\mathbf{L} &= -\int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} \\ \frac{\partial \mathbf{B}}{\partial t} &= 0 \quad (\text{steady}) \end{aligned} \right\} \longrightarrow \left\{ \begin{aligned} \oint \mathbf{E} \cdot d\mathbf{L} &= 0 \\ \nabla \times \mathbf{E} &= 0 \end{aligned} \right.$$

## Faraday's Law (4)



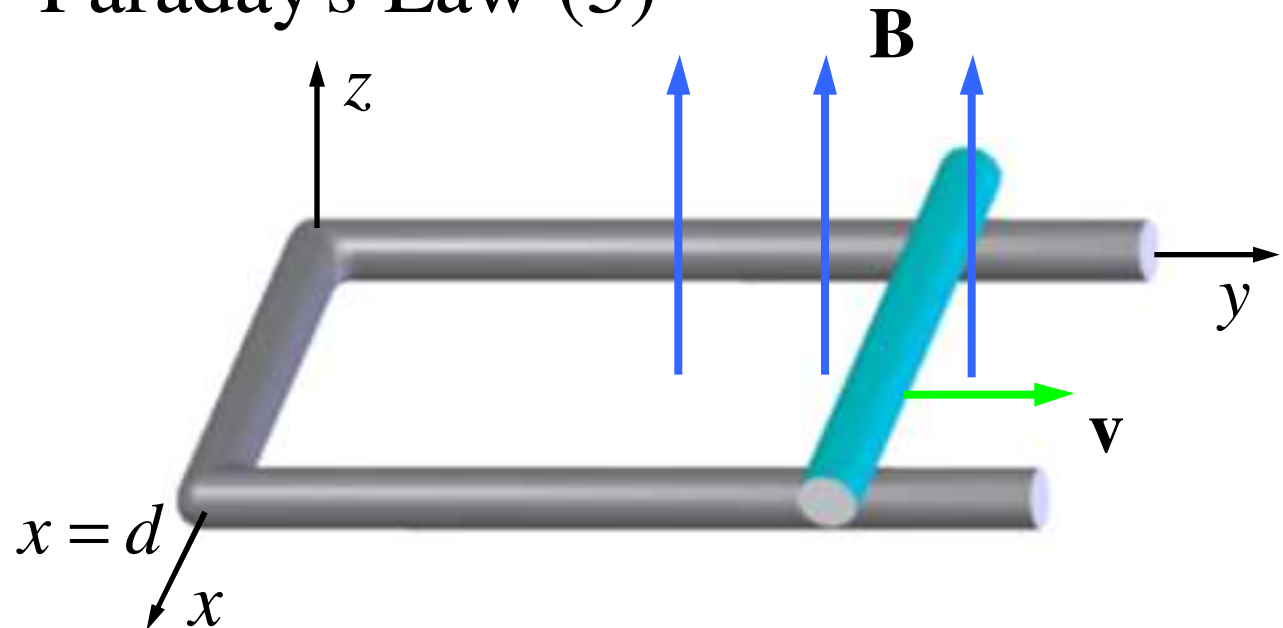
$$\left. \begin{aligned} \Phi &= \int_S \mathbf{B} \cdot d\mathbf{S} = Byd \\ \text{emf} &= -\frac{d\Phi}{dt} \end{aligned} \right\} \rightarrow \text{emf} = -B \frac{dy}{dt} d = -Bvd$$

## Faraday's Law (5)

$$\mathbf{F} = Q\mathbf{v} \times \mathbf{B}$$

$$\rightarrow \frac{\mathbf{F}}{Q} = \mathbf{v} \times \mathbf{B}$$

$$\mathbf{E}_m = \mathbf{v} \times \mathbf{B}$$

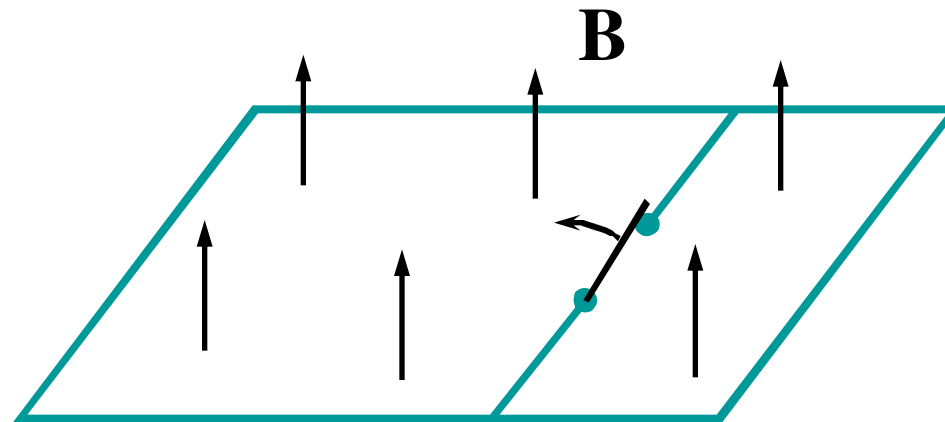


$$\text{emf} = \oint \mathbf{E}_m \cdot d\mathbf{L} = \oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L} = \int_d^0 v B dx = -Bvd$$



## Faraday's Law (6)

$$\text{emf} = \oint \mathbf{E} \cdot d\mathbf{L} = - \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} + \oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L}$$



## Ex. 1

## Faraday's Law (7)

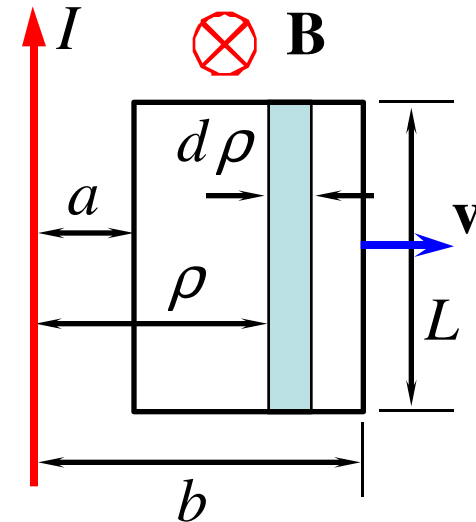
A single turn loop is situated in air, with a uniform magnetic field normal to its plane. The area of the loop is  $10 \text{ m}^2$ . If the rate of change of flux density is  $5 \text{ Wb/m}^2/\text{s}$ , what is the emf appearing at the terminals of the loop?

$$\left. \begin{array}{l} \text{emf} = -N \frac{d\Phi}{dt} \\ \Phi = BS \end{array} \right\} \rightarrow \text{emf} = -\frac{dB}{dt} S = 5 \times 10 = 50 \text{ V}$$

**Ex. 2**

# Faraday's Law (8)

$$\Phi(t) = \frac{\mu_0 IL}{2\pi} \ln \frac{b_0 + vt}{a_0 + vt}$$

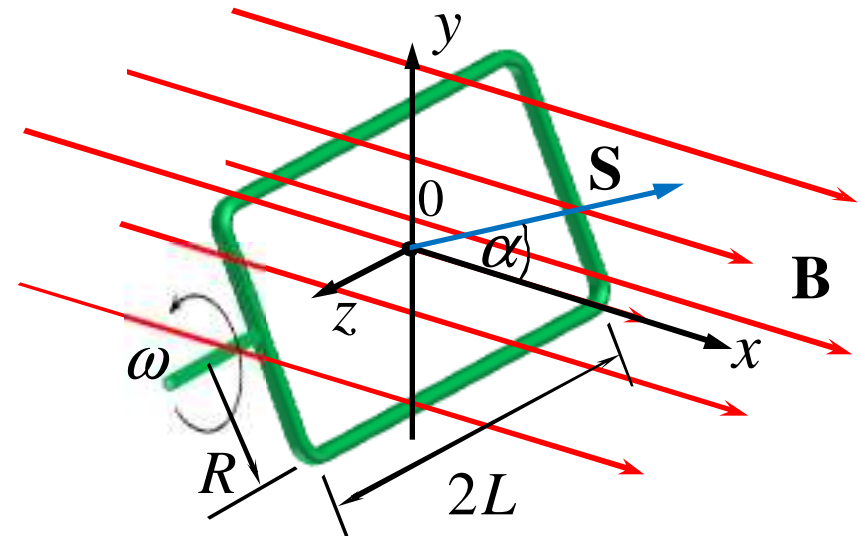
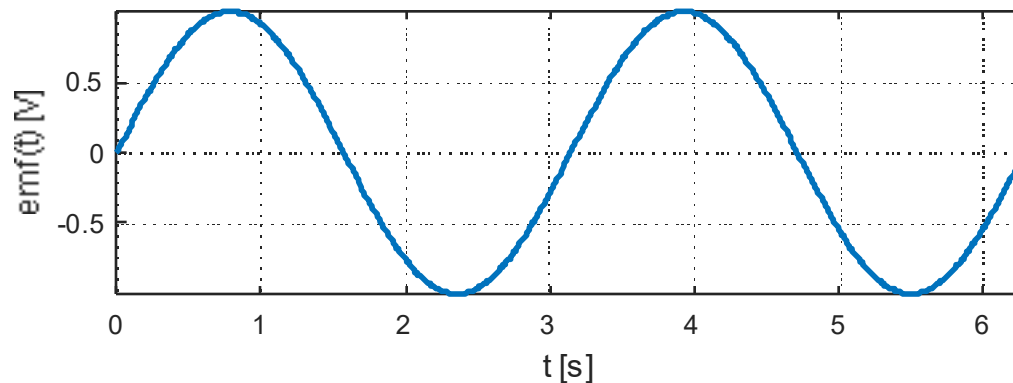
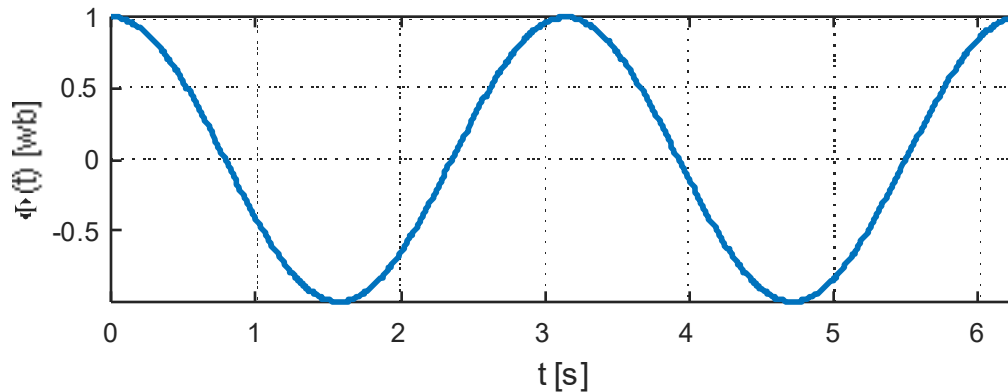


$$\rightarrow \text{emf} = -\frac{d\Phi(t)}{dt} = \frac{\mu_0 IL}{2\pi} \cdot \frac{(b_0 - a_0)v}{(a_0 + vt)^2} \cdot \frac{a_0 + vt}{b_0 + vt} = \frac{\mu_0 IL}{2\pi} \cdot \frac{(b_0 - a_0)v}{(a_0 + vt)b_0 + vt}$$

### Ex. 3

## Faraday's Law (9)

### Method 1



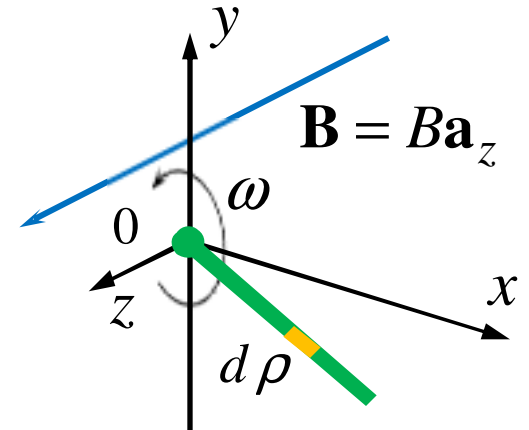
$$\Phi = BS \cos \omega t$$

$$\rightarrow \text{emf} = -\frac{d\Phi(t)}{dt} = BS \omega \sin \omega t$$

## Ex. 4

## Faraday's Law (10)

A conductive strip of length  $L$  pivoted at one end is rotating freely in the  $xy$ -plane with an angular frequency  $\omega$  in a uniform magnetic flux  $\mathbf{B}$ . Find the induced emf between the two ends of the strip?

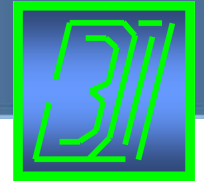


$$\text{emf} = \int (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L}$$

$$\mathbf{v} = \rho\omega\mathbf{a}_\phi$$

$$\mathbf{v} \times \mathbf{B} = (\rho\omega\mathbf{a}_\phi) \times (B\mathbf{a}_z) = \rho\omega B\mathbf{a}_\rho$$

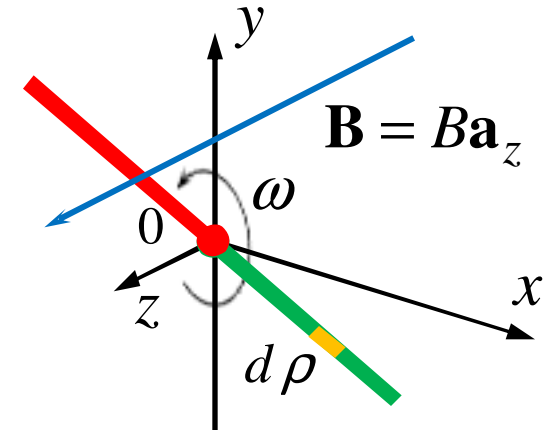
$$\text{emf} = \int_0^L (\rho\omega B\mathbf{a}_\rho) \cdot d\rho\mathbf{a}_\rho = \int_0^L \rho\omega B d\rho = \omega B \int_0^L \rho d\rho = \frac{B\omega L^2}{2}$$



**Ex. 5**

# Faraday's Law (11)

A conductive strip of length  $2L$  pivoted at the midpoint is rotating freely in the  $xy$ -plane with an angular frequency  $\omega$  in a uniform magnetic flux  $\mathbf{B}$ . Find the induced emf between the two ends of the strip?



$$\text{emf}_g = \frac{B\omega L^2}{2}$$

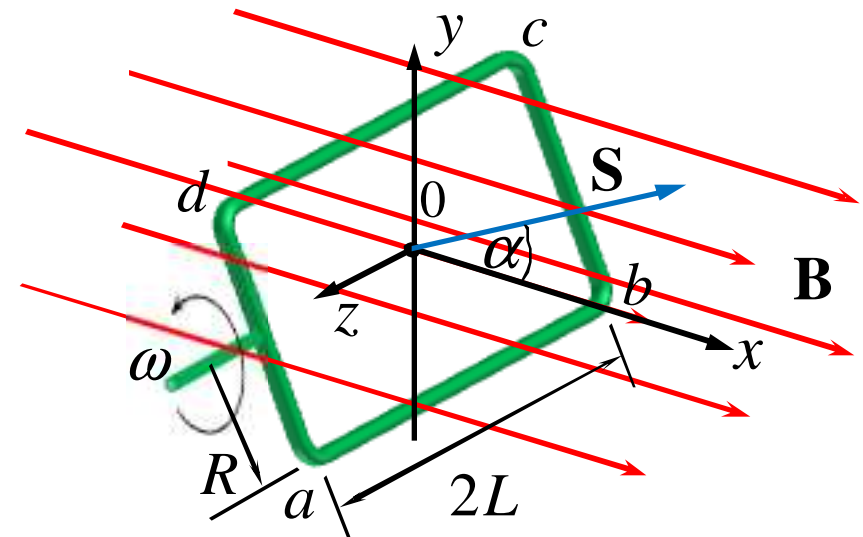
$$\text{emf}_r = \frac{B\omega L^2}{2}$$

$$\text{emf}_t = \text{emf}_g - \text{emf}_r = 0$$

**Ex. 3**

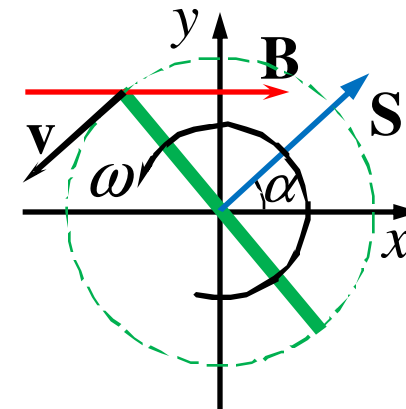
# Faraday's Law (12)

## Method 2



$$\begin{aligned} \text{emf} &= \oint \mathbf{E}_m \cdot d\mathbf{L} = \oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L} \\ &= \int_a^b + \int_b^c + \int_c^d + \int_d^a \\ &\quad \int_b^c = \int_d^a = 0 \end{aligned}$$

$$\begin{aligned} \text{emf}_{ab} &= \int_a^b = \int_L^{-L} [(R\omega \mathbf{a}_\phi) \times (B\mathbf{a}_x)] \cdot (dz\mathbf{a}_z) \\ &= \int_L^{-L} [BR\omega \sin(\omega t)(-\mathbf{a}_z)] \cdot (dz\mathbf{a}_z) \\ &= -BR\omega \sin \omega t \int_L^{-L} dz = 2LBR\omega \sin \omega t \end{aligned}$$



**Ex. 3**

# Faraday's Law (13)

## Method 2

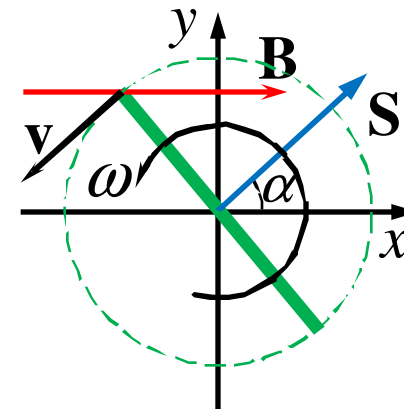
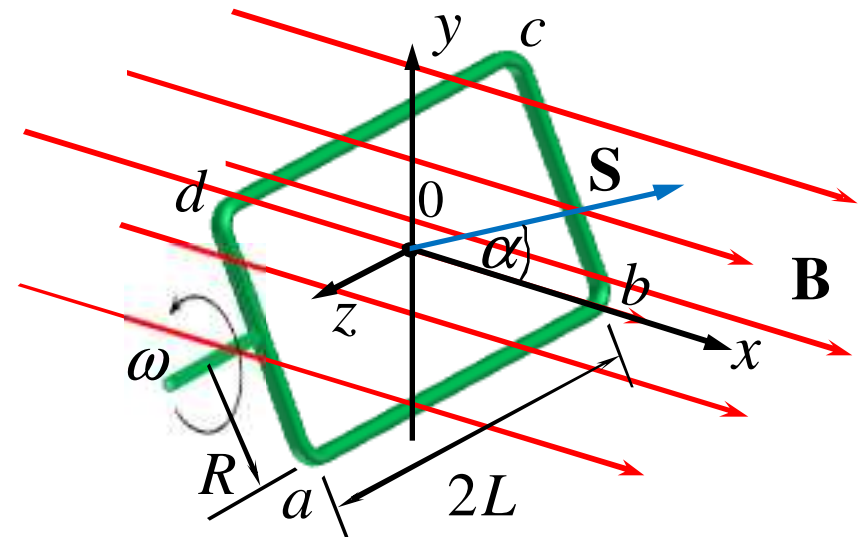
$$\text{emf} = \oint \mathbf{E}_m \cdot d\mathbf{L} = \oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L}$$

$$\text{emf}_{bc} = \text{emf}_{da} = 0$$

$$\text{emf}_{ab} = 2LBR\omega \sin \omega t$$

$$\text{emf}_{cd} = 2LBR\omega \sin \omega t$$

$$\begin{aligned} \text{emf} &= \text{emf}_{ab} + \text{emf}_{cd} \\ &= 4LBR\omega \sin \omega t = \boxed{BS\omega \sin \omega t} \end{aligned}$$





**Ex. 6**

# Faraday's Law (14)

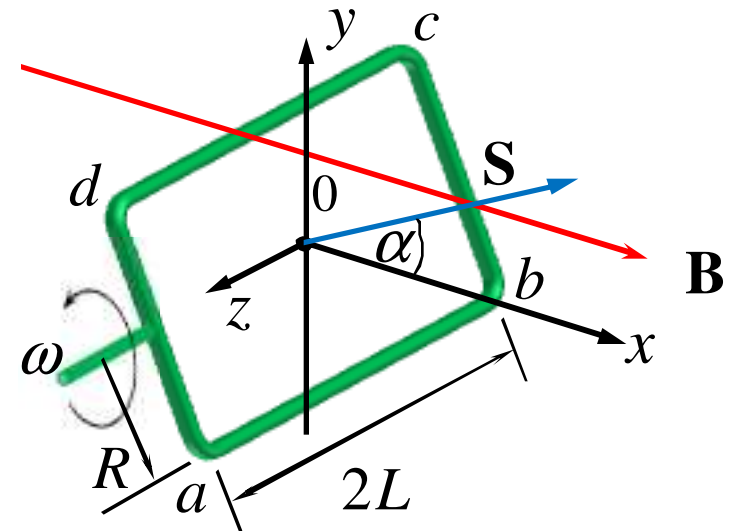
$$\mathbf{B} = B_m \sin \omega t \mathbf{a}_\rho.$$

$$\begin{aligned} \Phi &= \int_S \mathbf{B} \cdot d\mathbf{S} = \int_S (B_m \sin \omega t \mathbf{a}_\rho) \cdot (d\mathbf{S} \mathbf{a}_S) \\ &= \int_S B_m \sin \omega t \cos \omega t dS \end{aligned}$$

$$= B_m \sin \omega t \cos \omega t \int_S dS = B_m S \sin \omega t \cos \omega t = \frac{1}{2} B_m S \sin 2\omega t$$

$$\text{emf} = -\frac{d\Phi}{dt} = -B_m S \omega \cos 2\omega t$$

## Method 1



**Ex. 6**

# Faraday's Law (15)

$$\mathbf{B} = B_m \sin \omega t \mathbf{a}_\rho$$

$$\text{emf} = - \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} + \oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L}$$

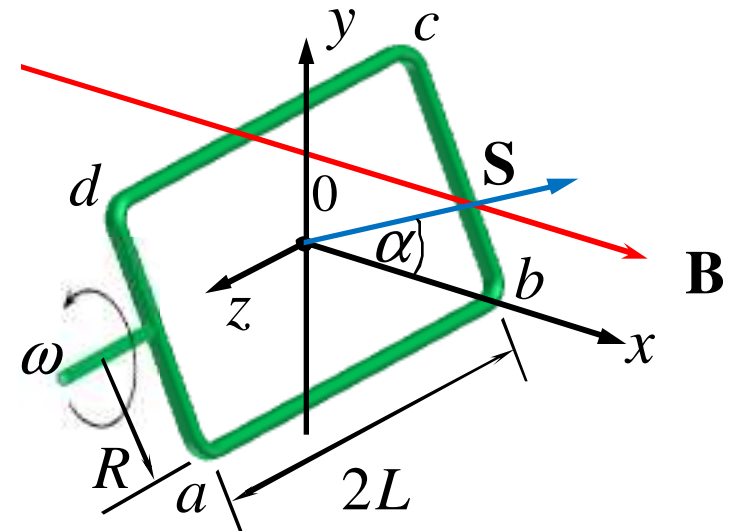
$$- \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} = - \int_S \left( \frac{\partial B_m \sin \omega t}{\partial t} \mathbf{a}_\rho \right) \cdot (dS \mathbf{a}_S)$$

$$= - \int_S (B_m \omega \cos \omega t \mathbf{a}_\rho) \cdot (dS \mathbf{a}_S)$$

$$= -B_m \omega \cos \omega t \int_S \mathbf{a}_\rho \cdot d\mathbf{S} \mathbf{a}_S = -B_m \omega \cos \omega t \int_S \cos \omega t dS$$

$$= -B_m \omega \cos^2 \omega t \int_S dS = -B_m S \omega \cos^2 \omega t$$

## Method 2



**Ex. 6**

# Faraday's Law (16)

$$\mathbf{B} = B_m \sin \omega t \mathbf{a}_\rho$$

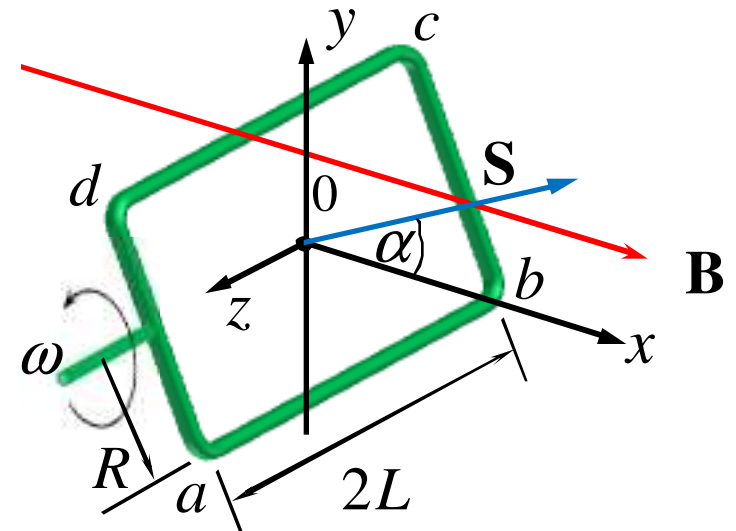
$$\text{emf} = - \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} + \oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L}$$

$$\oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L} = 2 \int_a^b (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L}$$

$$= 2 \int_L^{-L} [(R\omega \mathbf{a}_\varphi) \times (B_m \sin \omega t \mathbf{a}_\rho)] \cdot (dz \mathbf{a}_z)$$

$$= 2 \int_L^{-L} [(R\omega B_m \sin^2(-\mathbf{a}_z)] \cdot (dz \mathbf{a}_z) = B_m S \omega \sin^2 \omega t$$

## Method 2



**Ex. 6**

# Faraday's Law (17)

$$\mathbf{B} = B_m \sin \omega t \mathbf{a}_\rho$$

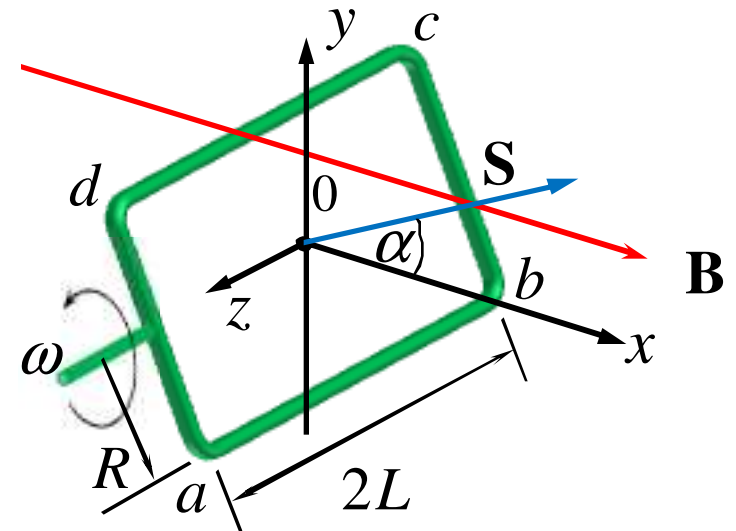
$$\text{emf} = - \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} + \oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L}$$

$$- \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} = -B_m S \omega \cos^2 \omega t$$

$$\oint (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{L} = B_m S \omega \sin^2 \omega t$$

$$\text{emf} = -B_m S \omega \cos^2 \omega t + B_m S \omega \sin^2 \omega t = \boxed{-B_m S \omega \cos 2\omega t}$$

## Method 2



**Ex. 7****Faraday's Law (18)**

A circular conducting loop of radius  $R$  lies in the  $xy$  plane. Find the emf of the loop if  $\mathbf{B} = 0.5\sin 500t\mathbf{a}_x + 0.3\sin 400t\mathbf{a}_y + 0.9\cos 314t\mathbf{a}_z$  T?

$$\Phi = \int_S \mathbf{B} \cdot d\mathbf{S}$$

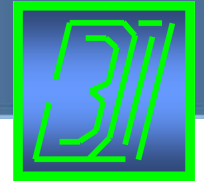
$$= \int_S (0.5 \sin 500t\mathbf{a}_x + 0.3 \sin 400t\mathbf{a}_y + 0.9 \cos 314t\mathbf{a}_z) \cdot (dS\mathbf{a}_z)$$

$$= \int_S (0.9 \cos 314t) dS$$

$$= 0.9 \cos 314t \int_S dS$$

$$= 0.9 \cos 314t (\pi R^2)$$

$$\text{emf} = -\frac{d\Phi}{dt} = -\frac{d}{dt} (0.9\pi R^2 \cos 314t) = 888R^2 \sin 314t \text{ V}$$



# Time – Varying Fields & Maxwell's Equations

1. Faraday's Law
- 2. Displacement Current**
3. Maxwell's Equations in Point Form
4. Maxwell's Equations in Integral Form
5. The Retarded Potentials



## Displacement Current (1)

$$\left. \begin{aligned} \nabla \times \mathbf{H} = \mathbf{J} &\rightarrow \nabla \cdot \nabla \times \mathbf{H} = \nabla \cdot \mathbf{J} \\ \nabla \cdot \nabla \times \mathbf{H} &= 0 \\ \nabla \cdot \mathbf{J} &= -\frac{\partial \rho_v}{\partial t} \end{aligned} \right\} \rightarrow \frac{\partial \rho_v}{\partial t} = 0 \text{ (unreasonable)}$$

$$\left. \begin{aligned} \nabla \times \mathbf{H} = \mathbf{J} + \mathbf{G} &\rightarrow 0 = \nabla \cdot \mathbf{J} + \nabla \cdot \mathbf{G} \\ \nabla \cdot \mathbf{J} &= -\frac{\partial \rho_v}{\partial t} \end{aligned} \right\} \rightarrow \left. \begin{aligned} \nabla \cdot \mathbf{G} &= \frac{\partial \rho_v}{\partial t} \\ \nabla \cdot \mathbf{D} &= \rho_v \end{aligned} \right\}$$

$$\left. \begin{aligned} \rightarrow \nabla \cdot \mathbf{G} &= \frac{\partial}{\partial t} (\nabla \cdot \mathbf{D}) = \nabla \cdot \frac{\partial \mathbf{D}}{\partial t} \rightarrow \mathbf{G} = \frac{\partial \mathbf{D}}{\partial t} \\ \nabla \times \mathbf{H} &= \mathbf{J} + \mathbf{G} \end{aligned} \right\} \rightarrow \boxed{\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}}$$

## Displacement Current (2)

$$\left. \begin{array}{l} \nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \\ \text{Define } \mathbf{J}_d = \frac{\partial \mathbf{D}}{\partial t} \end{array} \right\} \rightarrow \nabla \times \mathbf{H} = \mathbf{J} + \mathbf{J}_d$$

In nonconducting medium  $\mathbf{J} = 0 \rightarrow \nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t}$

$$\left. \begin{array}{l} I_d = \int_S \mathbf{J}_d \cdot d\mathbf{S} = \int_S \frac{\partial \mathbf{D}}{\partial t} \cdot d\mathbf{S} \\ \int_S (\nabla \times \mathbf{H}) \cdot d\mathbf{S} = \int_S \mathbf{J} \cdot d\mathbf{S} + \int_S \frac{\partial \mathbf{D}}{\partial t} \cdot d\mathbf{S} \\ \oint \mathbf{H} \cdot d\mathbf{L} = \int_S (\nabla \times \mathbf{H}) \cdot d\mathbf{S} \end{array} \right\} \rightarrow \boxed{\oint \mathbf{H} \cdot d\mathbf{L} = I + I_d = I + \int_S \frac{\partial \mathbf{D}}{\partial t} \cdot d\mathbf{S}}$$



## Displacement Current (3)

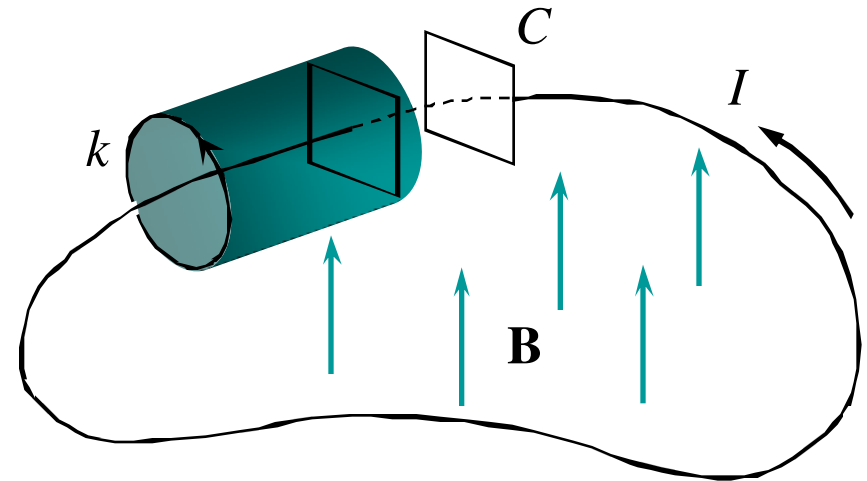
$$\text{emf} = V_0 \cos \omega t$$

$$\rightarrow I = -\omega C V_0 \sin \omega t$$

$$= -\omega \frac{\epsilon S}{d} V_0 \sin \omega t$$

$$\oint_k \mathbf{H} \cdot d\mathbf{L} = I_k$$

$$\left. \begin{aligned} D &= \epsilon E = \epsilon \left( \frac{V_0}{d} \cos \omega t \right) \\ I_d &= \int_S \frac{\partial \mathbf{D}}{\partial t} \cdot d\mathbf{S} = \frac{\partial D}{\partial t} S \end{aligned} \right\} \rightarrow I_d = -\omega \frac{\epsilon S}{d} V_0 \sin \omega t$$



**Ex. 1****Displacement Current (4)**

Given a magnetic field in free space as  $\mathbf{H} = H_0 \sin(\omega t - \beta z) \mathbf{a}_y$  A/m. Determine the displacement current density.

$$\left. \begin{aligned} \mathbf{J}_d &= \frac{\partial \mathbf{D}}{\partial t} \\ \nabla \times \mathbf{H} &= \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \\ \mathbf{J} &= 0 \end{aligned} \right\}$$

$$\begin{aligned} \rightarrow \mathbf{J}_d &= \nabla \times \mathbf{H} = \left( \frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} \right) \mathbf{a}_x + \left( \frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} \right) \mathbf{a}_y + \left( \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right) \mathbf{a}_z \\ &= H_0 \beta \cos(\omega t - \beta z) \mathbf{a}_x \end{aligned}$$

## Ex. 2

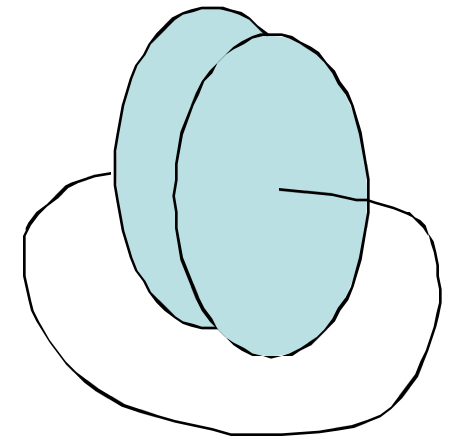
## Displacement Current (5)

A parallel – plate capacitor consists of two circular plates of radius  $R$ . Suppose that the capacitor is being charged at a uniform rate so that the  $\mathbf{E}$  between the plates changes at a constant rate  $dE/dt = 10^{12}$  V/m/s. Find the displacement current for the capacitor?

$$\mathbf{J}_d = \frac{\partial \mathbf{D}}{\partial t} = \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

$$I_d = \int_S \mathbf{J}_d \cdot d\mathbf{S} = \int_S \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \cdot d\mathbf{S} = \int_S \epsilon_0 \frac{\partial E}{\partial t} dS$$

$$= \epsilon_0 \frac{\partial E}{\partial t} \int_S dS = \epsilon_0 \frac{dE}{dt} \pi R^2$$



# Time – Varying Fields & Maxwell's Equations

1. Faraday's Law
2. Displacement Current
- 3. Maxwell's Equations in Point Form**
4. Maxwell's Equations in Integral Form
5. The Retarded Potentials

## Maxwell's Equations in Point Form (1)

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}$$

$$\nabla \cdot \mathbf{D} = \rho_v$$

$$\nabla \cdot \mathbf{B} = 0$$

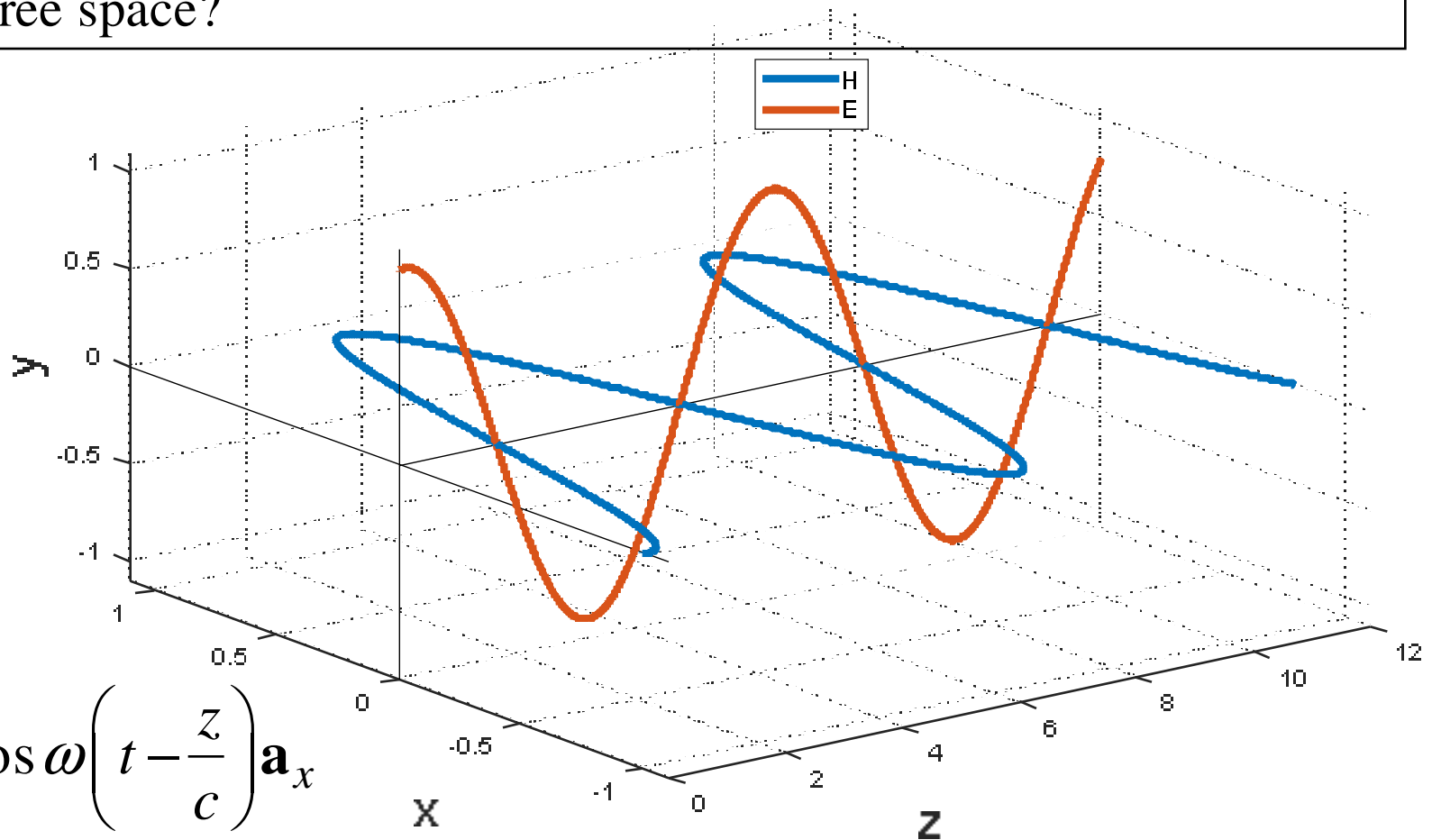
## Ex. 1 Maxwell's Equations in Point Form (2)

Given an electric field  $\mathbf{E} = A \cos \omega(t - z/c) \mathbf{a}_y$ . Determine the time-dependent MFI  $\mathbf{H}$  in free space?

$$\left. \begin{aligned} \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} = -\mu_0 \frac{\partial \mathbf{H}}{\partial t} \\ \nabla \times \mathbf{E} &= -\frac{\partial E_y}{\partial z} \mathbf{a}_x = -\frac{\omega}{c} A \sin \omega \left( t - \frac{z}{c} \right) \mathbf{a}_x \end{aligned} \right\}$$
$$\begin{aligned} \rightarrow \mathbf{H} &= \frac{\omega A}{c \mu_0} \int \sin \omega \left( t - \frac{z}{c} \right) \mathbf{a}_x \\ &= -\frac{A}{c \mu_0} \cos \omega \left( t - \frac{z}{c} \right) \mathbf{a}_x \end{aligned}$$

## Ex. 1 Maxwell's Equations in Point Form (3)

Given an electric field  $\mathbf{E} = A \cos \omega(t - z/c) \mathbf{a}_y$ . Determine the time-dependent MFI  $\mathbf{H}$  in free space?



$$\mathbf{H} = \frac{-A}{c\mu_0} \cos \omega \left( t - \frac{z}{c} \right) \mathbf{a}_x$$

## Ex. 2 Maxwell's Equations in Point Form (4)

$$\text{Find } \mathbf{E} \text{ if } \mathbf{B} = \begin{cases} B_0 \cos(\omega t + \alpha) \mathbf{a}_z & (\rho \leq a) \\ 0 & (\rho > a) \end{cases}$$

$$\left. \begin{aligned} \oint \mathbf{E} \cdot d\mathbf{L} &= \int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{S} \\ \mathbf{E} &= E(\rho) \mathbf{a}_\varphi \end{aligned} \right\} \rightarrow \oint \mathbf{E} \cdot d\mathbf{L} = 2\pi\rho E$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} = \omega B \sin(\omega t + \alpha) \mathbf{a}_z$$

$$\rightarrow \int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{S} = \omega B \sin(\omega t + \alpha) \int_S dS$$

$$\rho \leq a \rightarrow \omega B \sin(\omega t + \alpha) \int_S dS = \omega B \sin(\omega t + \alpha) (\pi \rho^2)$$

$$\begin{aligned} \rho > a \rightarrow \omega B \sin(\omega t + \alpha) \int_S dS &= \omega B \sin(\omega t + \alpha) \int_{S, \rho \leq a} dS + 0 \int_{S, \rho > a} dS \\ &= \omega B \sin(\omega t + \alpha) (\pi a^2) \end{aligned}$$



## Ex. 2 Maxwell's Equations in Point Form (5)

$$\left. \begin{aligned} \text{Find } \mathbf{E} \text{ if } \mathbf{B} = \begin{cases} B_0 \cos(\omega t + \alpha) \mathbf{a}_z & (\rho \leq a) \\ 0 & (\rho > a) \end{cases} \quad \left. \begin{aligned} \oint \mathbf{E} \cdot d\mathbf{L} &= \int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{S} \\ \oint \mathbf{E} \cdot d\mathbf{L} &= 2\pi\rho E \end{aligned} \right\} \\ \int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{S} = \begin{cases} \omega B \sin(\omega t + \alpha)(\pi\rho^2) & (\rho \leq a) \\ \omega B \sin(\omega t + \alpha)(\pi a^2) & (\rho > a) \end{cases} \\ \rightarrow 2\pi\rho E = \begin{cases} \omega B \sin(\omega t + \alpha)(\pi\rho^2) & (\rho \leq a) \\ \omega B \sin(\omega t + \alpha)(\pi a^2) & (\rho > a) \end{cases} \\ \rightarrow \mathbf{E} = \begin{cases} \frac{1}{2} \omega B \rho \sin(\omega t + \alpha) \mathbf{a}_\phi & (\rho \leq a) \\ \frac{1}{2} \omega B \frac{a^2}{\rho} \sin(\omega t + \alpha) \mathbf{a}_\phi & (\rho > a) \end{cases} \end{aligned} \right.$$

### Ex. 3 Maxwell's Equations in Point Form (6)

Given a magnetic field in free space where there is neither current density nor charge,  $\mathbf{B} = A \sin(\omega t - nx) \mathbf{a}_x + Ank \cos(\omega t - nx) \mathbf{a}_y$  (T) where  $A$ ,  $n$ , &  $\omega$  are constants. Use a Maxwell equation to derive the time-dependent part of  $\mathbf{E}$

$$\left. \begin{aligned} \nabla \times \mathbf{H} &= \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \\ \mathbf{J} &= 0 \end{aligned} \right\} \rightarrow \nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} \rightarrow \frac{1}{\mu_0} \nabla \times \mathbf{B} = \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \rightarrow \nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

$$\left. \begin{aligned} \nabla \times \mathbf{B} &= \left( \frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} \right) \mathbf{a}_x + \left( \frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x} \right) \mathbf{a}_y + \left( \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} \right) \mathbf{a}_z \\ &= An^2 k \sin(\omega t - nx) \mathbf{a}_z \end{aligned} \right\}$$

$$\rightarrow An^2 k \sin(\omega t - nx) \mathbf{a}_z = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

$$\rightarrow \mathbf{E} = \frac{An^2 k}{\mu_0 \epsilon_0} \mathbf{a}_z \int_0^t \sin(\omega t - nx) dt = \frac{An^2 k}{\mu_0 \epsilon_0 \omega} \cos(\omega t - nx) \mathbf{a}_z$$

# Time – Varying Fields & Maxwell's Equations

1. Faraday's Law
2. Displacement Current
3. Maxwell's Equations in Point Form
- 4. Maxwell's Equations in Integral Form**
5. The Retarded Potentials



## Maxwell's Equations in Integral Form (1)

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}$$

$$\nabla \cdot \mathbf{D} = \rho_v$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\oint \mathbf{E} \cdot d\mathbf{L} = -\int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S}$$

$$\oint \mathbf{H} \cdot d\mathbf{L} = I + \int_S \frac{\partial \mathbf{D}}{\partial t} \cdot d\mathbf{S}$$

$$\oint_S \mathbf{D} \cdot d\mathbf{S} = \int_V \rho_v dv$$

$$\oint_S \mathbf{B} \cdot d\mathbf{S} = 0$$

$$E_{t1} = E_{t2}$$

$$H_{t1} = H_{t2}$$

$$D_{N1} - D_{N2} = \rho_S$$

$$B_{N1} = B_{N2}$$

## **Ex.** Maxwell's Equations in Integral Form (2)

Find  $\mathbf{E}$  given an magnetic field  $\mathbf{B} = B_0 e^{bt} \mathbf{a}_z$ ?

$$\left. \begin{aligned} \oint \mathbf{E} \cdot d\mathbf{L} &= - \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} \\ \mathbf{E} &= E(\rho) \mathbf{a}_\phi \end{aligned} \right\}$$

$$\rightarrow \mathbf{E} \cdot \oint d\mathbf{L} = - \frac{\partial \mathbf{B}}{\partial t} \cdot \int_S d\mathbf{S}$$

$$\rightarrow E(2\pi\rho) = -bB_0 e^{bt} (\pi\rho^2)$$

$$\rightarrow E = -\frac{1}{2} b B_0 e^{bt} \pi \rho$$

$$\rightarrow \boxed{\mathbf{E} = -\frac{1}{2} b B_0 e^{bt} \pi \rho \mathbf{a}_\phi}$$

# Time – Varying Fields & Maxwell's Equations

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## The Retarded Potentials (1)

$$\left. \begin{aligned} \mathbf{E} = -\nabla V \rightarrow \nabla \times \mathbf{E} &= -\nabla \times (\nabla V) \\ 0 &= \nabla \times (\nabla V) \end{aligned} \right\} \rightarrow \left. \begin{aligned} \nabla \times \mathbf{E} &= 0 \\ \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} \end{aligned} \right\} \rightarrow \frac{\partial \mathbf{B}}{\partial t} = 0 \quad (\text{unreasonable})$$

$$\left. \begin{aligned} \mathbf{E} = -\nabla V + \mathbf{N} \rightarrow \nabla \times \mathbf{E} &= -\nabla \times (\nabla V) + \nabla \times \mathbf{N} \\ \nabla \times (\nabla V) &= 0 \\ \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} \end{aligned} \right\} \rightarrow \left. \begin{aligned} \nabla \times \mathbf{N} &= -\frac{\partial \mathbf{B}}{\partial t} \\ \mathbf{B} &= \nabla \times \mathbf{A} \end{aligned} \right\}$$

$$\rightarrow \nabla \times \mathbf{N} = -\frac{\partial}{\partial t} (\nabla \times \mathbf{A}) \rightarrow \nabla \times \mathbf{N} = -\nabla \times \frac{\partial \mathbf{A}}{\partial t} \rightarrow \mathbf{N} = -\frac{\partial \mathbf{A}}{\partial t}$$

$$\rightarrow \boxed{\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t}}$$

## The Retarded Potentials (2)

$$\left. \begin{aligned} \mathbf{B} &= \nabla \times \mathbf{A} \\ \mathbf{E} &= -\nabla V - \frac{\partial \mathbf{A}}{\partial t} \\ \nabla \times \mathbf{H} &= \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \\ \nabla \cdot \mathbf{D} &= \rho_v \end{aligned} \right\}$$

$$\rightarrow \left\{ \begin{aligned} \frac{1}{\mu} \nabla \times \nabla \times \mathbf{A} &= \mathbf{J} + \varepsilon \left( -\nabla \frac{\partial V}{\partial t} - \frac{\partial^2 \mathbf{A}}{\partial t^2} \right) \\ \varepsilon \left( -\nabla \cdot \nabla V - \frac{\partial}{\partial t} \nabla \cdot \mathbf{A} \right) &= \rho_v \end{aligned} \right.$$



## The Retarded Potentials (3)

$$\begin{cases} \frac{1}{\mu} \nabla \times \nabla \times \mathbf{A} = \mathbf{J} + \varepsilon \left( -\nabla \frac{\partial V}{\partial t} - \frac{\partial^2 \mathbf{A}}{\partial t^2} \right) \\ \varepsilon \left( -\nabla \cdot \nabla V - \frac{\partial}{\partial t} \nabla \cdot \mathbf{A} \right) = \rho_v \end{cases}$$

$$\rightarrow \begin{cases} \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A} = \mu \mathbf{J} - \mu \varepsilon \left( \nabla \frac{\partial V}{\partial t} + \frac{\partial^2 \mathbf{A}}{\partial t^2} \right) \\ \nabla^2 V + \frac{\partial}{\partial t} (\nabla \cdot \mathbf{A}) = -\frac{\rho_v}{\varepsilon} \end{cases}$$

## The Retarded Potentials (4)

$$\left\{ \begin{array}{l} \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A} = \mu \mathbf{J} - \mu \epsilon \left( \nabla \frac{\partial V}{\partial t} + \frac{\partial^2 \mathbf{A}}{\partial t^2} \right) \\ \nabla^2 V + \frac{\partial}{\partial t} (\nabla \cdot \mathbf{A}) = -\frac{\rho_v}{\epsilon} \\ \text{Define } \nabla \cdot \mathbf{A} = -\mu \epsilon \frac{\partial V}{\partial t} \end{array} \right\}$$

$$\rightarrow \left\{ \begin{array}{l} \nabla^2 \mathbf{A} = -\mu \mathbf{J} + \mu \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} \\ \nabla^2 V = -\frac{\rho_v}{\epsilon} + \mu \epsilon \frac{\partial^2 V}{\partial t^2} \end{array} \right.$$

## The Retarded Potentials (5)

$$\left. \begin{aligned} V &= \int_v \frac{\rho_v dv}{4\pi\epsilon R} \\ t' &= t - \frac{R}{v} \end{aligned} \right\} \rightarrow V = \int_v \frac{[\rho_v] dv}{4\pi\epsilon R}$$

$$\text{Ex : } \rho_v = e^{-r} \cos \omega t \rightarrow [\rho_v] = e^{-r} \cos \left[ \omega \left( t - \frac{R}{v} \right) \right]$$

$$\mathbf{A} = \int_v \frac{\mu \mathbf{J}}{4\pi R} dv \rightarrow \boxed{\mathbf{A} = \int_v \frac{\mu [\mathbf{J}]}{4\pi R} dv}$$

