# NaMaster: Scientific Documentation

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## Contents

## 1 Introduction

NaMaster is a C library, python module and standalone program to compute the pseudo-$C_\ell$ estimator of the angular power spectrum between two masked and contaminated fields (this is also the so-called "MASTER" algorithm). The contents of this scientific documentation describe the algorithm, drawing heavily from the methods presented by [1] and [2], extending their results to arbitrary cross-correlations between spin-0 and spin-2 fields (see also [3]).

## 2 Generalities and SHTs

Let $\mathbf{a}(\hat{\mathbf{n}})$ be a spin-$s_a$ quantity defined on the sphere. Then we define its spherical harmonic coefficients as:

$$\mathbf{a}_{\ell m} \equiv \mathrm{SHT}(\mathbf{a}(\hat{\mathbf{n}}))_{\ell m}^{s_a} \equiv \int d\hat{\mathbf{n}}\, \hat{\mathsf{Y}}_{\ell m}^{s_a \dagger}(\hat{\mathbf{n}})\, \mathbf{a}(\hat{\mathbf{n}}), \quad \mathbf{a}(\hat{\mathbf{n}}) = \mathrm{SHT}^{-1}(\mathbf{a}_{\ell m})_{\hat{\mathbf{n}}}^{s_a} \equiv \sum_{\ell m} \hat{\mathsf{Y}}_{\ell m}^{s_a}(\hat{\mathbf{n}})\, \mathbf{a}_{\ell m}. \tag{1}$$

Note that here we will use a vector notation, such that for a complex spin-$s_a$ field $a$ we form the vector $\mathbf{a} \equiv (\mathrm{Re}(a), \mathrm{Im}(a))$. The harmonic coefficients above are decomposed in a similar manner into $E$ and $B$ modes: $\mathbf{a}_{\ell m} \equiv (a_{\ell m}^E, a_{\ell m}^B)$. The spherical harmonic operators $\hat{Y}^s$ are therefore matrix that we define in the following subsection.

## 2.1 Spin-weighted spherical harmonics

Let $\eth$ and $\bar{\eth}$ be the following complex differential operators defined on the sphere when acting on a spin-$s$ quantity $f_s$:

$$\eth f_s \equiv -(\sin\theta)^s \left(\partial_\theta + i\frac{\partial_\varphi}{\sin\theta}\right)(\sin\theta)^{-s} f_s(\theta, \varphi), \quad \bar{\eth} f_s \equiv -(\sin\theta)^{-s}\left(\partial_\theta - i\frac{\partial_\varphi}{\sin\theta}\right)(\sin\theta)^s f_s(\theta, \varphi). \quad (2)$$

The following properties can be easily derived for the action of these operators:

- If $f_s$ is a spin-$s$ quantity, $(f_s)^*$ is a spin-$(-s)$ quantity.

- $\eth f_s$ is a spin-$(s+1)$ quantity, and $\bar{\eth} f_s$ is a spin-$(s-1)$ quantity.

- $(\eth^n f_s)^* = \bar{\eth}^n (f_s)^*$

- $\eth(f\,g) = f\eth g + g\eth f$

- $\eth^2(f\,g) = f\eth^2 g + g\eth^2 f + \eth f \eth g$

We start by defining the spin-weighted spherical harmonics with spin $s \geq 0$:

$$_sY_{\ell m} \equiv \alpha_{\ell,s}\eth^s Y_{\ell m}, \quad _{-s}Y_{\ell m} \equiv \alpha_{\ell,s}(-1)^s\bar{\eth}^s Y_{\ell m}, \quad \alpha_{l,s} \equiv \sqrt{\frac{(\ell-s)!}{(\ell+s)!}}, \quad (3)$$

which have the property: $(_sY_{\ell m})^* = (-1)^{s+m}\,_{-s}Y_{\ell -m}$. We then define the $E$-mode and $B$-mode spherical harmonic vectors as:

$$_s\mathbf{Y}_{\ell m}^E \equiv \mathbf{D}_s^E Y_{\ell m} \equiv -\frac{\alpha_{\ell,s}}{2}\begin{pmatrix} \eth^s + \bar{\eth}^s \\ -i(\eth^s - \bar{\eth}^s) \end{pmatrix} Y_{\ell m} \equiv -\frac{1}{2}\begin{pmatrix} _sY_{\ell m} + (-1)^s\,_{-s}Y_{\ell m} \\ -i(_sY_{\ell m} - (-1)^s\,_{-s}Y_{\ell m}) \end{pmatrix} \quad (4)$$

$$_s\mathbf{Y}_{\ell m}^B \equiv \mathbf{D}_s^B Y_{\ell m} \equiv -\frac{\alpha_{\ell,s}}{2}\begin{pmatrix} i(\eth^s - \bar{\eth}^s) \\ \eth^s + \bar{\eth}^s \end{pmatrix} Y_{\ell m} \equiv -\frac{1}{2}\begin{pmatrix} i(_sY_{\ell m} - (-1)^s\,_{-s}Y_{\ell m}) \\ _sY_{\ell m} + (-1)^s\,_{-s}Y_{\ell m} \end{pmatrix}, \quad (5)$$

which also defines the differential operators $\mathbf{D}_s^{E,B}$. These functions, for $s = 0$ are simply $\mathbf{D}_0^E = (Y_{\ell m}, 0)$ and $\mathbf{D}_0^B = (0, Y_{\ell m})$.

The matrix operator $\hat{Y}_{\ell m}^s$ is then defined as having $_s\mathbf{Y}_{\ell m}^{E,B}$ as columns:

$$\hat{Y}_{\ell m}^s \equiv \left(_s\mathbf{Y}_{\ell m}^E, _s\mathbf{Y}_{\ell m}^B,\right) \equiv -\frac{1}{2}\begin{pmatrix} _sY_{\ell m} + (-1)^s\,_{-s}Y_{\ell m} & i(_sY_{\ell m} - (-1)^s\,_{-s}Y_{\ell m}) \\ -i(_sY_{\ell m} - (-1)^s\,_{-s}Y_{\ell m}) & _sY_{\ell m} + (-1)^s\,_{-s}Y_{\ell m} \end{pmatrix} \quad (6)$$

The matrices $\hat{Y}_{\ell m}^s$ satisfy the following relations:

$$\hat{Y}_{\ell m}^{s\dagger} = (-1)^{m+s}\hat{Y}_{\ell -m}^{-s} \quad (7)$$

$$\int d\hat{\mathbf{n}}\,\hat{Y}_{\ell m}^s \hat{Y}_{\ell' m'}^{s\dagger} = \hat{\mathbf{1}}\delta_{\ell\ell'}\delta_{mm'} \quad (8)$$

$$\hat{D}_{\mathbf{l}\mathbf{l}_1\mathbf{l}_2}^s \equiv \int d\hat{\mathbf{n}}\left(\hat{Y}_{\mathbf{l}}^{s\dagger}(\hat{\mathbf{n}})\,\hat{Y}_{\mathbf{l}_1}^s(\hat{\mathbf{n}})\right)\hat{Y}_{\mathbf{l}_2}^0(\hat{\mathbf{n}}), \quad (9)$$

$$= (-1)^{s+m}\sqrt{\frac{(2\ell+1)(2\ell_1+1)(2\ell_2+1)}{4\pi}}\begin{pmatrix} \ell & \ell_1 & \ell_2 \\ -m & m_1 & m_2 \end{pmatrix}\begin{pmatrix} \ell & \ell_1 & \ell_2 \\ s & -s & 0 \end{pmatrix}\hat{\mathbf{d}}_{\ell+\ell_1+\ell_2}, \quad (10)$$

$$\hat{\mathbf{d}}_n^2 = \frac{1}{2}\begin{pmatrix} 1+(-1)^n & -i[1-(-1)^n] \\ i[1-(-1)^n] & 1+(-1)^n \end{pmatrix}, \quad (11)$$

where we have abbreviated the pair $(\ell, m)$ as $\mathbf{l}$.

Finally, the following orthogonality relation for the Wigner $3j$ symbols is useful:

$$\sum_{mm_1}\begin{pmatrix} \ell & \ell_1 & \ell_2 \\ m & m_1 & m_2 \end{pmatrix}\begin{pmatrix} \ell & \ell_1 & \ell_3 \\ m & m_1 & m_3 \end{pmatrix} = \frac{\delta_{\ell_2\ell_3}\delta_{m_2 m_3}}{2\ell_2+1} \quad (12)$$

## 2.2 $E$ and $B$ mode purification

We define a field $\mathbf{f}$ to be a $B$ ($E$) mode if $(\mathbf{D}_s^{E(B)})^\dagger \mathbf{f} = 0$. At the same time, and under the definition of the dot product:

$$(\mathbf{f}, \mathbf{g}) \equiv \int d\hat{\mathbf{n}} \, \mathbf{f}^\dagger \mathbf{g}, \tag{13}$$

we define a *pure $B$ ($E$) mode* as a field that is orthogonal to all $E$ ($B$) modes.

Since $\mathbf{D}_s^{E\dagger}\mathbf{D}_s^B = 0$, one can always generate a $B$ ($E$) mode by applying $\mathbf{D}_s^{B(E)}$ to a scalar field. It is then possible to show that $E$ and $B$ modes thus defined are orthogonal in the full sky:

$$(\mathbf{D}_s^E \varphi, \mathbf{D}_s^B \psi) = \int d\hat{\mathbf{n}} \, (\mathbf{D}_s^E \varphi)^\dagger \mathbf{D}_s^B \psi = 0 \tag{14}$$

This can be done by integrating by parts and noting that the celestial sphere has no boundaries. On a cut sky, however, and for $s = 2$, this is only true if the fields satisfy Neumann and Dirichlet boundary conditions simultaneously (i.e. vanishing value and first derivative on the boundary of the cut sky region).

Let $w(\hat{\mathbf{n}})$ be a sky window function defining the sky region to be analyzed (and the weight to be applied in each pixel). The standard pseudo $B$-mode of a field $\mathbf{P}$ is then given by

$$\tilde{B}_{\ell m} = \int d\hat{\mathbf{n}} \, w(\hat{\mathbf{n}}) \left( {}_s\mathbf{Y}_{\ell m}^B(\hat{\mathbf{n}}) \right)^\dagger \mathbf{P} = \int d\hat{\mathbf{n}} \, w(\hat{\mathbf{n}})(\mathbf{D}_s^B Y_{\ell m})^\dagger \mathbf{P}(\hat{\mathbf{n}}), \tag{15}$$

Now, since $\mathbf{D}_s^B Y_{\ell m}$ is a $B$-mode, in the absence of $w$ this expresion would correspond to a projection that filters out all the $E$-modes from $\mathbf{P}$. However, $w(\hat{\mathbf{n}})\mathbf{D}_s^B Y_{\ell m}$ is not a $B$-mode, and therefore $\tilde{B}_{\ell m}$ receives contributions from ambiguous $E$ modes (which then propagate into the variance of the pseudo-$C_\ell$ estimator of the power spectrum).

The idea behind $B$-mode purification is to move $w$ to the right of $\mathbf{D}_s^B$, defining the field:

$$B_{\ell m}^p = \int d\hat{\mathbf{n}} \left( \mathbf{D}_2^B (w \, Y_{\ell m}) \right)^\dagger \mathbf{P}(\hat{\mathbf{n}}). \tag{16}$$

Since $\mathbf{D}_2^B (w Y_{\ell m})$ is a $B$-mode quantity, $B_{\ell m}^p$ should receive contributions only from $B$-modes.

Expanding $\mathbf{D}_2^B (w Y_{\ell m})$, we can write $B_{\ell m}^p$ as:

$$B_{\ell m}^p = \left( \tilde{P}_2 \right)_{\ell m}^B + 2 \frac{\alpha_{\ell,2}}{\alpha_{\ell,1}} \left( \tilde{P}_1 \right)_{\ell m}^B + \alpha_{\ell,2} \left( \tilde{P}_2 \right)_{\ell m}^B, \tag{17}$$

where $(f)_{\ell m}^B$ stands for the $B$-mode of field $f$, and we have defined the fields $\tilde{P}_n = (\eth^n w)^* (Q + iU)$, where $Q$ and $U$ are the real and imaginary parts of the field $P$.

Note that the derivatives of $w$ can be computed as:

$$w \to w_{\ell m} = \text{SHT}(w) \to \left\{ {}_n w_{\ell m}^E = (-1)^H w_{\ell m}/\alpha_{\ell,n}, \, {}_n w_{\ell m}^B = 0 \right\} \to \eth^n w = \text{SHT}^{-1}(\{ {}_n w_{\ell m}^E, \, {}_n w_{\ell m}^B \}) \tag{18}$$

# 3 Contaminant cleaning

Let $\mathbf{a}$ be a random field defined on the sphere, let $v(\hat{\mathbf{n}})$ a mask for $\mathbf{a}$ and let $\mathbf{f}^i$ be a set of $N_a$ contaminants of $\mathbf{a}$ such that the observed version of $\mathbf{a}$ be:

$$\mathbf{d}_a(\hat{\mathbf{n}}) = \mathbf{a}^v(\hat{\mathbf{n}}) + \sum_{i=1}^{N_a} \alpha_i \mathbf{f}^i(\hat{\mathbf{n}}), \tag{19}$$

where $\mathbf{a}^v(\hat{\mathbf{n}}) = v(\hat{\mathbf{n}}) \mathbf{a}(\hat{\mathbf{n}})$ (note that we have implicitly applied the same mask to $\mathbf{f}^i$). The best-fit value for the coefficients $\alpha_i$ assuming the same weights for all points in $\mathbf{d}_a$ can be found as

$$\tilde{\alpha}_i = \sum_j M_{ij} \int d\hat{\mathbf{n}} \, \mathbf{f}^{j\dagger}(\hat{\mathbf{n}}) \, \mathbf{d}_a(\hat{\mathbf{n}}), \quad (\hat{\mathsf{M}}^{-1})_{ij} \equiv \int d\hat{\mathbf{n}} \, \mathbf{f}^{i\dagger}(\hat{\mathbf{n}}) \, \mathbf{f}^j(\hat{\mathbf{n}}). \tag{20}$$

Thus we can find a cleaned version of $\mathbf{a}$ as:

$$\tilde{\mathbf{a}}(\hat{\mathbf{n}}) \equiv \mathbf{d}_a(\hat{\mathbf{n}}) - \mathbf{f}^i(\hat{\mathbf{n}}) \, M_{ij} \int d\hat{\mathbf{n}}' \mathbf{f}^{j\dagger}(\hat{\mathbf{n}}') \mathbf{d}_a(\hat{\mathbf{n}}')$$

$$= \mathbf{a}^v(\hat{\mathbf{n}}) - \mathbf{f}^i(\hat{\mathbf{n}}) \, M_{ij} \int d\hat{\mathbf{n}}' \mathbf{f}^{j\dagger}(\hat{\mathbf{n}}') \mathbf{a}^v(\hat{\mathbf{n}}'), \tag{21}$$

where there is an implicit summation sign over $i$ and $j$ (we will omit these from now on).

The harmonic coefficients of the cleaned and masked field are:

$$\tilde{\mathbf{a}}_{\ell m} = \mathbf{a}^v_{\ell m} - \mathbf{f}^i_{\ell m} M_{ij} \sum_{\ell' m'} \mathbf{f}^{j\dagger}_{\ell' m'} \mathbf{a}^v_{\ell' m'}. \tag{22}$$

From now on we will simplify the notation by abbreviating the pair $\ell m$ as $\mathbf{l}$, so that the previous equation reads:

$$\tilde{\mathbf{a}}_\mathbf{l} = \mathbf{a}^v_\mathbf{l} - \mathbf{f}^i_\mathbf{l} M_{ij} \sum_{\mathbf{l}'} \mathbf{f}^{j\dagger}_{\mathbf{l}'} \mathbf{a}^v_{\mathbf{l}'}. \tag{23}$$

The harmonic coefficients for the masked field can be related to those of the unmasked one and the mask $v$ (understood as a spin-0 field) as:

$$\mathbf{a}^v_\mathbf{l} = \sum_{\mathbf{l}_1 \mathbf{l}_2} \hat{\mathsf{D}}^{s_a}_{\mathbf{l}\mathbf{l}_1 \mathbf{l}_2} \mathbf{a}_{\mathbf{l}_1} v_{\mathbf{l}_2}. \tag{24}$$

# 4  Pseudo-$C_\ell$ estimators with mode deprojection

In what follows, for two fields $\mathbf{a}$ and $\mathbf{b}$ we will define their observed power spectrum as:

$$\tilde{\mathsf{C}}^{ab}_\ell \equiv \frac{1}{2\ell + 1} \sum_m \mathbf{a}_{\ell m} \mathbf{b}^\dagger_{\ell m}. \tag{25}$$

This must not be confused with the true power spectrum defined as an ensemble average for isotropic fields:

$$\langle \mathbf{a}_\mathbf{l} \mathbf{b}^\dagger_{\mathbf{l}'} \rangle \equiv \hat{\mathsf{C}}^{ab}_\ell \delta_{\ell\ell'} \delta_{mm'}. \tag{26}$$

Now, let $\tilde{\mathbf{a}}$ and $\tilde{\mathbf{b}}$ be the contaminant-cleaned versions of two random fields $\mathbf{a}$ and $\mathbf{b}$ with contaminants $\mathbf{f}^i$ and $\mathbf{g}^j$ and masks $v$ and $w$ respectively, and let us define

$$(\hat{\mathsf{N}}^{-1})_{ij} \equiv \int d\hat{\mathbf{n}} \, \mathbf{g}^{i\dagger}(\hat{\mathbf{n}}) \, \mathbf{g}^j(\hat{\mathbf{n}}). \tag{27}$$

The observed power spectrum of the contaminant-cleaned maps can be written as:

$$\tilde{\mathsf{C}}^{\tilde{a}\tilde{b}}_\ell = \frac{1}{2\ell + 1} \sum_m \mathbf{a}^v_{\ell m} \mathbf{b}^{w\dagger}_{\ell m} - \frac{N^*_{ij}}{2\ell + 1} \sum_m \sum_{\mathbf{l}'} \mathbf{a}^v_\mathbf{l} \mathbf{b}^{w\dagger}_{\mathbf{l}'} \mathbf{g}^j_{\mathbf{l}'} \mathbf{g}^{i\dagger}_\mathbf{l} - $$

$$- \frac{M_{ij}}{2\ell + 1} \sum_m \sum_{\mathbf{l}'} \mathbf{f}^i_\mathbf{l} \mathbf{f}^{j\dagger}_{\mathbf{l}'} \mathbf{a}^v_{\mathbf{l}'} \mathbf{b}^{w\dagger}_\mathbf{l} + \frac{M_{ij} N^*_{pq}}{2\ell + 1} \sum_m \sum_{\mathbf{l}'\mathbf{l}''} \mathbf{f}^i_\mathbf{l} \mathbf{f}^{j\dagger}_{\mathbf{l}'} \mathbf{a}^v_{\mathbf{l}'} \mathbf{b}^{w\dagger}_{\mathbf{l}''} \mathbf{g}^q_{\mathbf{l}''} \mathbf{g}^{p\dagger}_\mathbf{l}. \tag{28}$$

In order to compute the bias of $\tilde{\mathsf{C}}^{\tilde{a}\tilde{b}}_\ell$ with respect to $\hat{\mathsf{C}}^{ab}_\ell$, we need to compute the ensemble average of the former, which we will write as:

$$\left\langle \tilde{\mathsf{C}}^{\tilde{a}\tilde{b}}_\ell \right\rangle = \hat{\mathsf{F}}^1_\ell - \hat{\mathsf{F}}^2_\ell - \hat{\mathsf{F}}^3_\ell + \hat{\mathsf{F}}^4_\ell, \tag{29}$$

where:

$$\hat{\mathsf{F}}^1_\ell \equiv \frac{1}{2\ell + 1} \sum_m \langle \mathbf{a}^v_{\ell m} \mathbf{b}^{w\dagger}_{\ell m} \rangle, \quad \hat{\mathsf{F}}^2_\ell \equiv \frac{N^*_{ij}}{2\ell + 1} \sum_m \sum_{\mathbf{l}'} \langle \mathbf{a}^v_\mathbf{l} \mathbf{b}^{w\dagger}_{\mathbf{l}'} \mathbf{g}^j_{\mathbf{l}'} \mathbf{g}^{i\dagger}_\mathbf{l} \rangle \tag{30}$$

$$\hat{\mathsf{F}}^3_\ell \equiv \frac{M_{ij}}{2\ell + 1} \sum_m \sum_{\mathbf{l}'} \langle \mathbf{f}^i_\mathbf{l} \mathbf{f}^{j\dagger}_{\mathbf{l}'} \mathbf{a}^v_{\mathbf{l}'} \mathbf{b}^{w\dagger}_\mathbf{l} \rangle, \quad \hat{\mathsf{F}}^4_\ell \equiv \frac{M_{ij} N^*_{pq}}{2\ell + 1} \sum_m \sum_{\mathbf{l}'\mathbf{l}''} \langle \mathbf{f}^i_\mathbf{l} \mathbf{f}^{j\dagger}_{\mathbf{l}'} \mathbf{a}^v_{\mathbf{l}'} \mathbf{b}^{w\dagger}_{\mathbf{l}''} \mathbf{g}^q_{\mathbf{l}''} \mathbf{g}^{p\dagger}_\mathbf{l} \rangle. \tag{31}$$

We will now compute the ensemble average of each of these terms.

## 4.1 $\hat{\mathsf{F}}_\ell^1$

$$\hat{\mathsf{F}}_\ell^1 = \frac{1}{2\ell+1} \sum_{m\mathbf{l}_{1,2,3,4}} v_{\mathbf{l}_2} w_{\mathbf{l}_4}^* \hat{\mathsf{D}}_{\mathbf{ll}_1\mathbf{l}_2}^{s_a} \langle \mathbf{a}_{\mathbf{l}_1} \mathbf{b}_{\mathbf{l}_3}^\dagger \rangle \hat{\mathsf{D}}_{\mathbf{ll}_3\mathbf{l}_4}^{s_b\dagger}$$

$$= \frac{1}{2\ell+1} \sum_{m\mathbf{l}_{1,2,3}} v_{\mathbf{l}_2} w_{\mathbf{l}_3}^* \hat{\mathsf{D}}_{\mathbf{ll}_1\mathbf{l}_2}^{s_a} \hat{\mathsf{C}}_{\mathbf{l}_1}^{ab} \hat{\mathsf{D}}_{\mathbf{ll}_1\mathbf{l}_3}^{s_b\dagger}$$

$$= \frac{1}{2\ell+1} \sum_{\ell_1\mathbf{l}_{2,3}} v_{\mathbf{l}_2} w_{\mathbf{l}_3}^* \frac{(2\ell+1)(2\ell_1+1)}{4\pi} \sqrt{(2\ell_2+1)(2\ell_3+1)} \begin{pmatrix} \ell & \ell_1 & \ell_2 \\ s_a & -s_a & 0 \end{pmatrix} \begin{pmatrix} \ell & \ell_1 & \ell_3 \\ s_b & -s_b & 0 \end{pmatrix}$$

$$\hat{\mathsf{d}}_{\ell+\ell_1+\ell_2}^{s_a} \hat{\mathsf{C}}_{\ell_1}^{ab} \hat{\mathsf{d}}_{\ell+\ell_1+\ell_3}^{s_b\dagger} \sum_{mm_1} \begin{pmatrix} \ell & \ell_1 & \ell_2 \\ -m & m_1 & m_2 \end{pmatrix} \begin{pmatrix} \ell & \ell_1 & \ell_3 \\ -m & m_1 & m_3 \end{pmatrix}$$

$$= \sum_{\ell_1\ell_2} \frac{(2\ell_1+1)(2\ell_2+1)}{4\pi} \tilde{C}_{\ell_2}^{vw} \begin{pmatrix} \ell & \ell_1 & \ell_2 \\ s_a & -s_a & 0 \end{pmatrix} \begin{pmatrix} \ell & \ell_1 & \ell_2 \\ s_b & -s_b & 0 \end{pmatrix} \hat{\mathsf{d}}_{\ell+\ell_1+\ell_2}^{s_a} \hat{\mathsf{C}}_{\ell_1}^{ab} \hat{\mathsf{d}}_{\ell+\ell_1+\ell_2}^{s_b\dagger} \quad (32)$$

For $v = w = 1$ this reduces to $\tilde{C}_\ell^{vw} = 4\pi\delta_{\ell 0}$ and:

$$\hat{\mathsf{F}}_\ell^1 = \sum_{\ell_1\ell_2} (2\ell_1+1)(2\ell_2+1)\delta_{\ell_2 0} \begin{pmatrix} \ell & \ell_1 & \ell_2 \\ s_a & -s_a & 0 \end{pmatrix} \begin{pmatrix} \ell & \ell_1 & \ell_2 \\ s_b & -s_b & 0 \end{pmatrix} \hat{\mathsf{d}}_{\ell+\ell_1+\ell_2}^{s_a} \hat{\mathsf{C}}_{\ell_1}^{ab} \hat{\mathsf{d}}_{\ell+\ell_1+\ell_2}^{s_b\dagger}$$

$$= \sum_{\ell_1} (2\ell_1+1) \begin{pmatrix} \ell & \ell_1 & 0 \\ s_a & -s_a & 0 \end{pmatrix} \begin{pmatrix} \ell & \ell_1 & 0 \\ s_b & -s_b & 0 \end{pmatrix} \hat{\mathsf{d}}_{\ell+\ell_1}^{s_a} \hat{\mathsf{C}}_{\ell_1}^{ab} \hat{\mathsf{d}}_{\ell+\ell_1}^{s_b\dagger}$$

$$= \sum_{\ell_1} (2\ell_1+1)\delta_{\ell\ell_1} \frac{(-1)^{\ell-s_a}}{\sqrt{2\ell+1}} \delta_{\ell\ell_1} \frac{(-1)^{\ell-s_b}}{\sqrt{2\ell+1}} \hat{\mathsf{d}}_{\ell+\ell_1}^{s_a} \hat{\mathsf{C}}_{\ell_1}^{ab} \hat{\mathsf{d}}_{\ell+\ell_1}^{s_b\dagger}$$

$$= \hat{\mathsf{d}}_{2\ell}^{s_a} \hat{\mathsf{C}}_\ell^{ab} \hat{\mathsf{d}}_{2\ell}^{s_b\dagger}$$

$$= \hat{\mathsf{C}}_\ell^{ab}$$

## 4.2 $\hat{\mathsf{F}}_\ell^2$

$$\hat{\mathsf{F}}_\ell^2 = N_{ij}^* \int d\hat{\mathbf{n}}\, d\hat{\mathbf{n}}' \frac{\sum_{m\mathbf{l}'\mathbf{l}_{1,2,3,4}}}{2\ell+1} \hat{\mathsf{Y}}_{\mathbf{l}}^{s_a\dagger}(\hat{\mathbf{n}}) \hat{\mathsf{Y}}_{\mathbf{l}_1}^{s_a}(\hat{\mathbf{n}}) \langle \mathbf{a}_{\mathbf{l}_1} \mathbf{b}_{\mathbf{l}_3}^\dagger \rangle \hat{\mathsf{Y}}_{\mathbf{l}_3}^{s_b\dagger}(\hat{\mathbf{n}}') \hat{\mathsf{Y}}_{\mathbf{l}'}^{s_b}(\hat{\mathbf{n}}') \mathbf{g}_{\mathbf{l}'}^j \mathbf{g}_{\mathbf{l}}^{i\dagger} v_{\mathbf{l}_2} w_{\mathbf{l}_4}^* Y_{\mathbf{l}_2}(\hat{\mathbf{n}}) Y_{\mathbf{l}_4}^*(\hat{\mathbf{n}}')$$

$$= N_{ij}^* \frac{\sum_m}{2\ell+1} \left\{ \int d\hat{\mathbf{n}}\, v(\hat{\mathbf{n}}) \hat{\mathsf{Y}}_{\mathbf{l}}^{s_a\dagger}(\hat{\mathbf{n}}) \left[ \sum_{\ell_1 m_1} \hat{\mathsf{Y}}_{\mathbf{l}_1}^{s_a}(\hat{\mathbf{n}}) \hat{\mathsf{C}}_{\ell_1}^{ab} \left( \int d\hat{\mathbf{n}}' \hat{\mathsf{Y}}_{\mathbf{l}_1}^{s_b\dagger}(\hat{\mathbf{n}}') \mathbf{g}^j(\hat{\mathbf{n}}') w(\hat{\mathbf{n}}') \right) \right] \mathbf{g}_{\mathbf{l}}^{i\dagger} \right\}$$

$$= N_{ij}^* \frac{\sum_m}{2\ell+1} \mathrm{SHT} \left\{ v(\hat{\mathbf{n}}) \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab} \mathrm{SHT} \left( w\, \mathbf{g}^j \right)_{\mathbf{l}_1}^{s_b} \right]_{\hat{\mathbf{n}}}^{s_a} \right\}_{\mathbf{l}}^{s_a} \mathbf{g}_{\mathbf{l}}^{i\dagger} \quad (33)$$

For $v = w = 1$ this reduces to:

$$\hat{\mathsf{F}}_\ell^2 = N_{ij}^* \frac{\sum_m}{2\ell+1} \mathrm{SHT} \left\{ \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab} \mathrm{SHT} \left( \mathbf{g}^j \right)_{\mathbf{l}_1}^{s_b} \right]_{\hat{\mathbf{n}}}^{s_a} \right\}_{\mathbf{l}}^{s_a} \mathbf{g}_{\mathbf{l}}^{i\dagger}$$

$$= N_{ij}^* \hat{\mathsf{C}}_\ell^{ab} \frac{\sum_m \mathbf{g}_{\ell m}^j \mathbf{g}_{\ell m}^{i\dagger}}{2\ell+1}$$

$$= N_{ij}^* \hat{\mathsf{C}}_\ell^{ab} \tilde{\mathsf{C}}_\ell^{g^j g^i}$$

### 4.3 $\hat{\mathsf{F}}_\ell^3$

$$\hat{\mathsf{F}}_\ell^3 = M_{ij} \int d\hat{\mathbf{n}}\, d\hat{\mathbf{n}}' \frac{\sum_{m\mathbf{l}'\mathbf{l}_{1,2,3,4}}}{2\ell+1} \mathbf{f}_\mathbf{l}^i \mathbf{f}_{\mathbf{l}'}^{j\dagger} \hat{Y}_{\mathbf{l}'}^{s_a\dagger}(\hat{\mathbf{n}}') \hat{Y}_{\mathbf{l}_3}^{s_a}(\hat{\mathbf{n}}') \langle \mathbf{a}_{\mathbf{l}_3} \mathbf{b}_{\mathbf{l}_1}^\dagger \rangle \hat{Y}_{\mathbf{l}_1}^{s_b\dagger}(\hat{\mathbf{n}}) \hat{Y}_\mathbf{l}^{s_b}(\hat{\mathbf{n}}) v_{\mathbf{l}_4} w_{\mathbf{l}_2}^* Y_{\mathbf{l}_2}(\hat{\mathbf{n}}) Y_{\mathbf{l}_4}^*(\hat{\mathbf{n}}')$$

$$= M_{ij} \frac{\sum_m}{2\ell+1} \mathbf{f}_\mathbf{l}^i \left\{ \int d\hat{\mathbf{n}}\, w(\hat{\mathbf{n}}) \left[ \sum_{\ell_1 m_1} \left( \int d\hat{\mathbf{n}}'\, v(\hat{\mathbf{n}}') \mathbf{f}^{j\dagger}(\hat{\mathbf{n}}') \hat{Y}_{\mathbf{l}_1}^{s_a}(\hat{\mathbf{n}}') \right) \hat{\mathsf{C}}_{\ell_1}^{ab} \hat{Y}_{\mathbf{l}_1}^{s_b\dagger}(\hat{\mathbf{n}}) \right] \hat{Y}_\mathbf{l}^{s_b}(\hat{\mathbf{n}}) \right\}$$

$$= M_{ij} \frac{\sum_m}{2\ell+1} \mathbf{f}_\mathbf{l}^i\, \mathrm{SHT} \left\{ w(\hat{\mathbf{n}})\, \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab\dagger} \mathrm{SHT} \left( v\, \mathbf{f}^j \right)_{\mathbf{l}_1}^{s_a} \right]_{\hat{\mathbf{n}}}^{s_b} \right\}_\mathbf{l}^{s_b\dagger} \tag{34}$$

For $v = w = 1$ this reduces to:

$$\hat{\mathsf{F}}_\ell^3 = M_{ij} \frac{\sum_m}{2\ell+1} \mathbf{f}_\mathbf{l}^i\, \mathrm{SHT} \left\{ \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab\dagger} \mathrm{SHT} \left( \mathbf{f}^j \right)_{\mathbf{l}_1}^{s_a} \right]_{\hat{\mathbf{n}}}^{s_b} \right\}_\mathbf{l}^{s_b\dagger}$$

$$= M_{ij} \frac{\sum_m \mathbf{f}_{\ell m}^i \mathbf{f}_{\ell m}^{j\dagger}}{2\ell+1} \hat{\mathsf{C}}_\ell^{ab}$$

$$= M_{ij} \tilde{\mathsf{C}}_\ell^{f^i f^j} \hat{\mathsf{C}}_\ell^{ab}$$

### 4.4 $\hat{\mathsf{F}}_\ell^4$

$$\hat{\mathsf{F}}_\ell^4 = \frac{M_{ij} N_{pq}^*}{2\ell+1} \int d\hat{\mathbf{n}}\, d\hat{\mathbf{n}}' \sum_{m\mathbf{l}'\mathbf{l}''\mathbf{l}_{1,2,3,4}} \mathbf{f}_\mathbf{l}^i \mathbf{f}_{\mathbf{l}'}^{j\dagger} \hat{Y}_{\mathbf{l}'}^{s_a\dagger}(\hat{\mathbf{n}}) \hat{Y}_{\mathbf{l}_1}^{s_a}(\hat{\mathbf{n}}) \langle \mathbf{a}_{\mathbf{l}_1} \mathbf{b}_{\mathbf{l}_3}^\dagger \rangle \hat{Y}_{\mathbf{l}_3}^{s_b\dagger}(\hat{\mathbf{n}}') \hat{Y}_{\mathbf{l}''}^{s_b}(\hat{\mathbf{n}}') \mathbf{g}_{\mathbf{l}''}^q \mathbf{g}_\mathbf{l}^{p\dagger} v_{\mathbf{l}_2} w_{\mathbf{l}_4}^* Y_{\mathbf{l}_4}(\hat{\mathbf{n}}) Y_{\mathbf{l}_4}^*(\hat{\mathbf{n}}')$$

$$= \frac{M_{ij} N_{pq}^*}{2\ell+1} \sum_m \mathbf{f}_\mathbf{l}^i \left\{ \int d\hat{\mathbf{n}}\, v(\hat{\mathbf{n}}) \mathbf{f}^{j\dagger}(\hat{\mathbf{n}}) \left[ \sum_{\ell_1 m_1} \hat{Y}_{\mathbf{l}_1}^{s_a}(\hat{\mathbf{n}}) \hat{\mathsf{C}}_{\ell_1}^{ab} \left( \int d\hat{\mathbf{n}}'\, \hat{Y}_{\mathbf{l}_1}^{s_b\dagger}(\hat{\mathbf{n}}') \mathbf{g}^q(\hat{\mathbf{n}}') w(\hat{\mathbf{n}}') \right) \right] \right\} \mathbf{g}_\mathbf{l}^{p\dagger}$$

$$= M_{ij} N_{pq}^* \left\{ \int d\hat{\mathbf{n}}\, v(\hat{\mathbf{n}}) \mathbf{f}^{j\dagger}(\hat{\mathbf{n}})\, \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab} \mathrm{SHT} \left( w\, \mathbf{g}^q \right)_{\mathbf{l}_1}^{s_b} \right]_{\hat{\mathbf{n}}}^{s_a} \right\} \tilde{\mathsf{C}}_\ell^{f^i g^p} \tag{35}$$

For $v = w = 1$ this reduces to:

$$\hat{\mathsf{F}}_\ell^4 = M_{ij} N_{pq}^* \left\{ \int d\hat{\mathbf{n}}\, \mathbf{f}^{j\dagger}(\hat{\mathbf{n}})\, \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab} \mathrm{SHT} \left( \mathbf{g}^q \right)_{\mathbf{l}_1}^{s_b} \right]_{\hat{\mathbf{n}}}^{s_a} \right\} \tilde{\mathsf{C}}_\ell^{f^i g^p}$$

$$= M_{ij} N_{pq}^* \left\{ \int d\hat{\mathbf{n}} \sum_{\mathbf{l}_2} \mathbf{f}_{\mathbf{l}_2}^{j\dagger} \hat{Y}_{\mathbf{l}_2}^{s_a\dagger}(\hat{\mathbf{n}}) \sum_{\mathbf{l}_1} \hat{Y}_{\mathbf{l}_1}^{s_a}(\hat{\mathbf{n}}) \hat{\mathsf{C}}_{\ell_1}^{ab} \mathbf{g}_{\mathbf{l}_1}^q \right\} \tilde{\mathsf{C}}_\ell^{f^i g^p}$$

$$= M_{ij} N_{pq}^* \mathbf{f}_{\mathbf{l}_1}^{j\dagger} \hat{\mathsf{C}}_{\ell_1}^{ab} \mathbf{g}_{\mathbf{l}_1}^q \tilde{\mathsf{C}}_\ell^{f^i g^p}$$

$$= M_{ij} N_{pq}^* \left[ \sum_{\ell_1} (2\ell_1+1) \mathrm{Tr} \left( \hat{\mathsf{C}}_{\ell_1}^{ab} \tilde{\mathsf{C}}_{\ell_1}^{g^q f^j} \right) \right] \tilde{\mathsf{C}}_\ell^{f^i g^p}$$

### 4.5 Final form of the estimator

Putting together the results from Equations 32, 33, 34 and 35, we can write down an unbiased estimator for the pseudo-$C_\ell$ of the cut-sky maps free from contamination from $\mathbf{f}$ and $\mathbf{g}$:

$$\tilde{\mathsf{C}}_\ell^{ab} = \tilde{\mathsf{C}}_\ell^{\tilde{a}\tilde{b}} + N_{ij}^* \frac{\sum_m}{2\ell+1} \mathrm{SHT} \left\{ v(\hat{\mathbf{n}}) \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab} \mathrm{SHT} \left( w\, \mathbf{g}^j \right)_{\mathbf{l}_1}^{s_b} \right]_{\hat{\mathbf{n}}}^{s_a} \right\}_\mathbf{l}^{s_a} \mathbf{g}_\mathbf{l}^{i\dagger} +$$

$$+ M_{ij} \frac{\sum_m}{2\ell+1} \mathbf{f}_\mathbf{l}^i\, \mathrm{SHT} \left\{ w(\hat{\mathbf{n}})\, \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab\dagger} \mathrm{SHT} \left( v\, \mathbf{f}^j \right)_{\mathbf{l}_1}^{s_a} \right]_{\hat{\mathbf{n}}}^{s_b} \right\}_\mathbf{l}^{s_b\dagger} -$$

$$- M_{ij} N_{pq}^* \left\{ \int d\hat{\mathbf{n}}\, v(\hat{\mathbf{n}}) \mathbf{f}^{j\dagger}(\hat{\mathbf{n}})\, \mathrm{SHT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab} \mathrm{SHT} \left( w\, \mathbf{g}^q \right)_{\mathbf{l}_1}^{s_b} \right]_{\hat{\mathbf{n}}}^{s_a} \right\} \tilde{\mathsf{C}}_\ell^{f^i g^p} \tag{36}$$

Once $\tilde{\mathsf{C}}^{ab}$ is calculated, it can be corrected for the effects of masking by inverting the linear transformation in Eq 32. This transformation can be written explicitly by first transforming the power spectrum

matrices into vectors $_v\mathsf{C}$. E.g. for $s_a = s_b = 2$ we transform:

$$\hat{\mathsf{C}}_\ell^{ab} \equiv \left( \begin{array}{cc} C_\ell^{E_a E_b} & C_\ell^{E_a B_b} \\ C_\ell^{B_a E_b} & C_\ell^{B_a B_b} \end{array} \right) \quad \text{into} \quad _v\hat{\mathsf{C}}_\ell^{ab} \equiv \left( \begin{array}{c} C_\ell^{E_a E_b} \\ C_\ell^{E_a B_b} \\ C_\ell^{B_a E_b} \\ C_\ell^{B_a B_b} \end{array} \right). \tag{37}$$

We can then write, in general:

$$_v\tilde{\mathsf{C}}_\ell^{ab} = \sum_{\ell'} \mathsf{M}_{\ell\ell'}^{s_a s_b} \cdot {}_v\hat{\mathsf{C}}_{\ell'}^{ab}, \tag{38}$$

where:

$$\mathsf{M}_{\ell\ell'}^{00} = \frac{2\ell'+1}{4\pi} \sum_{\ell''} (2\ell''+1) C_{\ell''}^{vw} \left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 0 & 0 & 0 \end{array} \right)^2 \tag{39}$$

$$\mathsf{M}_{\ell\ell'}^{02} = M_{\ell\ell'}^{0+} \hat{\mathbb{1}}, \quad M_{\ell\ell'}^{0+} = \frac{2\ell'+1}{4\pi} \sum_{\ell''} (2\ell''+1) C_{\ell''}^{vw} \left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 0 & 0 & 0 \end{array} \right) \left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 2 & -2 & 0 \end{array} \right) \tag{40}$$

$$\mathsf{M}_{\ell\ell'}^{22} = \left( \begin{array}{cccc} M_{\ell\ell'}^{++} & 0 & 0 & M_{\ell\ell'}^{--} \\ 0 & M_{\ell\ell'}^{++} & -M_{\ell\ell'}^{--} & 0 \\ 0 & -M_{\ell\ell'}^{--} & M_{\ell\ell'}^{++} & 0 \\ M_{\ell\ell'}^{--} & 0 & 0 & M_{\ell\ell'}^{++} \end{array} \right) \tag{41}$$

$$M_{\ell\ell'}^{\pm\pm} = \frac{2\ell'+1}{4\pi} \sum_{\ell''} (2\ell''+1) C_{\ell''}^{vw} \left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 2 & -2 & 0 \end{array} \right)^2 \frac{1 \pm (-1)^{\ell+\ell'+\ell''}}{2} \tag{42}$$

If either the $B$ or $E$ modes of a spin-2 field has been purified, the equations above must be modified by carrying out the following modification in the equations above:

$$\left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 2 & -2 & 0 \end{array} \right) \longrightarrow \tag{43}$$

$$\left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 2 & -2 & 0 \end{array} \right) + 2\sqrt{\frac{(\ell+1)!(\ell-2)!(\ell''+1)!}{(\ell-1)!(\ell+2)!(\ell''-1)!}} \left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 1 & -2 & 1 \end{array} \right) + \sqrt{\frac{(\ell-2)!(\ell''+2)!}{(\ell+2)!(\ell''-2)!}} \left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 0 & -2 & 2 \end{array} \right).$$

This change must be applied to the corresponding factors of $\left( \begin{array}{ccc} \ell & \ell' & \ell'' \\ 2 & -2 & 0 \end{array} \right)$.

### 4.6  Beam

Adding the effect of a beam amounts to redefining:

$$\mathsf{M}_{\ell_1\ell_2}^{s_a s_b} \rightarrow \mathsf{M}_{\ell_1\ell_2}^{s_a s_b} b_{\ell_2}^{ab}, \tag{44}$$

where $b_\ell^{ab}$ is the product of the harmonic transform of the beams for maps $a$ and $b$.

### 4.7  Binning into bandpowers

Given the loss of information implicit in masking the originally full-sky field, it is in general not possible to invert Eq. 38 directly. De usual approach to doing so is by binning the pseudo-$C_\ell$ into bandpowers. A bandpower $b$ is defined by a set of $N_b$ multipoles $\vec{\ell}_b \equiv (\ell_b^1, ..., \ell_b^{N_k})$ and a set of weights $\vec{w}_b \equiv (w_b^1, ..., w_b^{N_b})$ normalized such that $\sum_{i=1}^{N_b} w_b^i = 1$. The $b$−th bandpower for the coupled pseudo-$C_\ell$ is then defined as:

$$_v\tilde{\mathsf{B}}_b \equiv \sum_{i=1}^{N_b} w_b^i \, {}_v\tilde{\mathsf{C}}_{\ell_b^i} = \sum_{i=1}^{N_b} w_b^i \sum_{\ell'} \mathsf{M}_{\ell_b^i \ell'}^{s_a s_b} \, {}_v\hat{\mathsf{C}}_\ell. \tag{45}$$

One then proceeds by assuming that the true power spectrum is a step-wise function, taking constant values over the multipoles corresponding to each bandpower: $_v\hat{\mathsf{C}}_\ell = \sum_b {}_v\hat{\mathsf{B}}_b \Theta(\ell \in \vec{\ell}_b)$ (where $\Theta$ is a

binary step function). The previous equation then reads:

$$_v\tilde{\mathsf{B}}_b = \sum_{b'} \mathcal{M}_{bb'} \, {}_v\hat{\mathsf{B}}_{bb'} \equiv \sum_{b'} \left( \sum_{\ell \in \vec{\ell}_b} \sum_{\ell' \in \vec{\ell}_{b'}} w_b^\ell \mathsf{M}_{\ell\ell'} \right) {}_v\hat{\mathsf{B}}_{b'}, \tag{46}$$

which defines the binned coupling matrix $\mathcal{M}_{bb'}$. The decoupled bandpowers are then estimated by inverting $\mathcal{M}$:

$$_v\hat{\mathsf{B}}_b = \sum_{b'} (\mathcal{M})_{bb'}^{-1} \, {}_v\tilde{\mathsf{B}}_{b'}. \tag{47}$$

Note that, since this procedure is based on the assumption that the true power spectrum is stepwise constant, the bandpowers computed this way should be compared with the theoretical prediction subjected to the same type of transformation. I.e. the theoretical prediction for the bandpowers is:

$$_v\bar{\mathsf{B}}_b = \sum_{b'} (\mathcal{M})_{bb'}^{-1} \sum_{\ell' \in \vec{\ell}_{b'}} w_{b'}^{\ell'} \sum_{\ell''} \mathsf{M}_{\ell'\ell''} \, {}_v\bar{\mathsf{C}}_{\ell''}, \tag{48}$$

where we the overline $^-$ denotes theoretical predictions.

# 5 Flat-sky

## 5.1 Fourier transforms

In the flat sky we will write the directional vector $\hat{\mathbf{n}}$ as $\mathbf{x}$. Let $\mathbf{a}(\mathbf{x})$ be a spin-$s_a$ quantity. Under the approximation $\sin\theta \sim 1$, $\Delta\theta \to \Delta x$, $\Delta\varphi \to \Delta y$, the differential operator $\eth$ now takes the form:

$$\eth = -(\partial_x + i\partial_y), \quad \bar{\eth} = -(\partial_x - i\partial_y), \tag{49}$$

and acts on a plane wave $e^{i\mathbf{k}\mathbf{x}}$ as:

$$\eth^s e^{i\mathbf{k}\mathbf{x}} = (-ik)^s e^{i\,s\varphi_k} e^{i\mathbf{k}\mathbf{x}}, \quad \bar{\eth} e^{i\mathbf{k}\mathbf{x}} = (-ik)^s e^{-i\,s\varphi_k} e^{i\mathbf{k}\mathbf{x}} \tag{50}$$

Let us define the basis functions:

$$_s\mathcal{Y}_{\mathbf{k}}(\mathbf{x}) \equiv k^{-s}\eth^s e^{i\mathbf{k}\mathbf{x}} = (-i)^s e^{is\varphi_k} e^{i\mathbf{k}\mathbf{x}}, \quad _{-s}\mathcal{Y}_{\mathbf{k}}(\mathbf{x}) \equiv (-k)^{-s}\bar{\eth}^s e^{i\mathbf{k}\mathbf{x}} = i^s e^{-is\varphi_k} e^{i\mathbf{k}\mathbf{x}}, \tag{51}$$

and the spin-$s$ Fourier coefficients:

$$a(\mathbf{x}) = \int \frac{d\mathbf{l}^2}{2\pi} \, {}_s\mathcal{Y}_{\mathbf{l}}(\mathbf{x}) {}_s a_{\mathbf{l}}, \quad a^*(\mathbf{x}) = \int \frac{d\mathbf{l}^2}{2\pi} \, {}_{-s}\mathcal{Y}_{\mathbf{l}}(\mathbf{x}) {}_{-s} a_{\mathbf{l}}, \tag{52}$$

$$_s a_{\mathbf{l}} \equiv \int \frac{d\mathbf{x}^2}{2\pi} \, {}_s\mathcal{Y}_{\mathbf{l}}^*(\mathbf{x}) a(\mathbf{x}), \quad _{-s} a_{\mathbf{l}} \equiv \int \frac{d\mathbf{x}^2}{2\pi} \, {}_{-s}\mathcal{Y}_{\mathbf{l}}^*(\mathbf{x}) a^*(\mathbf{x}). \tag{53}$$

The $E$ and $B$-mode coefficients are then defined as:

$$_s E_{\mathbf{l}} \equiv -\frac{1}{2} \left[ {}_s a_{\mathbf{l}} + (-1)^s \, {}_{-s} a_{\mathbf{l}} \right], \quad i \, {}_s B_{\mathbf{l}} \equiv -\frac{1}{2} \left[ {}_s a_{\mathbf{l}} - (-1)^s \, {}_{-s} a_{\mathbf{l}} \right] \tag{54}$$

Let us now write $a$ as a vector such that in real space $\mathbf{a}(\mathbf{x}) \equiv (\mathrm{Re}(a), \mathrm{Im}(a))$, and in Fourier space $\mathbf{a}_{\mathbf{l}} \equiv ( {}_s E_{\mathbf{l}}, {}_s B_{\mathbf{l}})$. We can rewrite the equations above in vectorial form:

$$\mathbf{a}(\mathbf{x}) \equiv \int \frac{d\mathbf{l}^2}{2\pi} \, {}_s\mathsf{E}_{\mathbf{l}}(\mathbf{x}) \mathbf{a}_{\mathbf{l}}, \quad \mathbf{a}_{\mathbf{l}} \equiv \int \frac{d\mathbf{x}^2}{2\pi} \, {}_s\mathsf{E}_{\mathbf{l}}^\dagger(\mathbf{x}) \mathbf{a}(\mathbf{x}), \tag{55}$$

where we have defined the matrix basis functions:

$$_s\mathsf{E}_{\mathbf{l}}(\mathbf{x}) \equiv -\frac{1}{2} \begin{pmatrix} {}_s\mathcal{Y}_{\mathbf{l}} + (-1)^s \, {}_{-s}\mathcal{Y}_{\mathbf{l}} & i( {}_s\mathcal{Y}_{\mathbf{l}} - (-1)^s \, {}_{-s}\mathcal{Y}_{\mathbf{l}}) \\ -i( {}_s\mathcal{Y}_{\mathbf{l}} - (-1)^s \, {}_{-s}\mathcal{Y}_{\mathbf{l}}) & {}_s\mathcal{Y}_{\mathbf{l}} + (-1)^s \, {}_{-s}\mathcal{Y}_{\mathbf{l}} \end{pmatrix} \tag{56}$$

$$= -\frac{1}{2l^s} \begin{pmatrix} \eth^s + \bar{\eth}^s & i(\eth^s - \bar{\eth}^s) \\ -i(\eth^s - \bar{\eth}^s) & \eth^s + \bar{\eth}^s \end{pmatrix} e^{i\mathbf{l}\mathbf{x}} \tag{57}$$

$$= -(-i)^s \begin{pmatrix} \cos(s\varphi_l) & \sin(s\varphi_l) \\ -\sin(s\varphi_l) & \cos(s\varphi_l) \end{pmatrix} e^{i\mathbf{l}\mathbf{x}} \tag{58}$$

$$= -(-i)^s \mathsf{R}(s\varphi_l) e^{i\mathbf{l}\mathbf{x}}, \tag{59}$$

where $\mathsf{R}(\varphi)$ is a rotation matrix.

Thus:

$$\begin{pmatrix} Q(\mathbf{x}) \\ U(\mathbf{x}) \end{pmatrix} = -(-i)^s \int \frac{d\mathbf{l}2}{2\pi} \begin{pmatrix} \cos(s\varphi_l) & \sin(s\varphi_l) \\ -\sin(s\varphi_l) & \cos(s\varphi_l) \end{pmatrix} \begin{pmatrix} {}_sE_{\mathbf{l}} \\ {}_sB_{\mathbf{l}} \end{pmatrix} e^{i\mathbf{l}\mathbf{x}} \tag{60}$$

$$\begin{pmatrix} {}_sE_{\mathbf{l}} \\ {}_sB_{\mathbf{l}} \end{pmatrix} = -i^s \begin{pmatrix} \cos(s\varphi_l) & -\sin(s\varphi_l) \\ \sin(s\varphi_l) & \cos(s\varphi_l) \end{pmatrix} \begin{pmatrix} Q_{\mathbf{l}} \\ U_{\mathbf{l}} \end{pmatrix}, \tag{61}$$

where $Q_{\mathbf{l}}$ and $U_{\mathbf{l}}$ are the standard Fourier transforms of $Q$ and $U$.

The functions ${}_s\mathsf{E}_{\mathbf{l}}(\mathbf{x})$ satisfy the following orthogonality and completeness relations:

$$\int \frac{d\mathbf{x}^2}{(2\pi)^2} \, {}_s\mathsf{E}_{\mathbf{l}}(\mathbf{x}) {}_s\mathsf{E}_{\mathbf{l}'}^{\dagger}(\mathbf{x}) = 1 \, \delta(\mathbf{l}-\mathbf{l}'), \quad \int \frac{d\mathbf{l}^2}{(2\pi)^2} \, {}_s\mathsf{E}_{\mathbf{l}}(\mathbf{x}) {}_s\mathsf{E}_{\mathbf{l}}^{\dagger}(\mathbf{x}') = 1 \, \delta(\mathbf{x}-\mathbf{x}'). \tag{62}$$

### 5.2 Pseudo-$C_\ell$ estimator - continuum limit

The Fourier coefficients of the masked field are:

$$\mathbf{a}_{\mathbf{l}}^v = \int\int \frac{d\mathbf{k}^2 d\mathbf{q}^2}{2\pi} \left[ \int \frac{d\mathbf{x}^2}{(2\pi)^2} \, {}_{s_a}\mathsf{E}_{\mathbf{l}}^{\dagger}(\mathbf{x}) \, {}_{s_a}\mathsf{E}_{\mathbf{k}}(\mathbf{x}) \, {}_0E_{\mathbf{q}}(\mathbf{x}) \right] \mathbf{a}_{\mathbf{k}} \, v_{\mathbf{q}} \tag{63}$$

$$= \int \frac{d\mathbf{k}^2}{2\pi} \mathsf{R}^{\dagger}\left(s_a(\varphi_\ell - \varphi_k)\right) \mathbf{a}_{\mathbf{k}} v_{\mathbf{l}-\mathbf{k}} \tag{64}$$

Then:

$$\left\langle \tilde{\mathbf{a}}_{\mathbf{l}} \, \tilde{\mathbf{b}}_{\mathbf{l}}^{\dagger} \right\rangle = \int\int \frac{d\mathbf{k}^2 d\mathbf{q}^2}{(2\pi)^2} \mathsf{R}^{\dagger}(s_a(\varphi_\ell - \varphi_k)) \left\langle \mathbf{a}_{\mathbf{k}} \mathbf{b}_{\mathbf{q}}^{\dagger} \right\rangle \mathsf{R}(s_b(\varphi_\ell - \varphi_q)) v_{\mathbf{l}-\mathbf{k}} w_{\mathbf{l}-\mathbf{q}}^* \tag{65}$$

$$= \int \frac{d\mathbf{k}^2}{(2\pi)^2} \mathsf{R}^{\dagger}(s_a\Delta\varphi) \mathsf{C}_k^{ab} \mathsf{R}(s_b\Delta\varphi) v_{\mathbf{l}-\mathbf{k}} w_{\mathbf{l}-\mathbf{k}}^* \tag{66}$$

where $\Delta\varphi \equiv \varphi_\ell - \varphi_k$.

The pseudo-$C_\ell$ is defined as the nomalized angular average of $\mathbf{a}_{\mathbf{l}} \mathbf{b}_{\mathbf{l}}^{\dagger}$:

$$\tilde{\mathsf{C}}_\ell^{ab} \equiv \frac{(2\pi)^2}{S} \int \frac{d\varphi_\ell}{2\pi} \left\langle \tilde{\mathbf{a}}_{\mathbf{l}} \, \tilde{\mathbf{b}}_{\mathbf{l}}^{\dagger} \right\rangle \tag{67}$$

$$= \int \frac{k\,dk\,q\,dq}{(2\pi)^2} \left[ \frac{(2\pi)^2}{S} \int \frac{d\varphi_q}{2\pi} v_{\mathbf{q}} w_{\mathbf{q}}^* \int d\varphi_\ell d\varphi_k \mathsf{R}^{\dagger}(s_a\Delta\varphi) \mathsf{C}_k^{ab} \mathsf{R}(s_b\Delta\varphi) \delta(\mathbf{q}-\mathbf{l}+\mathbf{k}) \right], \tag{68}$$

where $S$ is the observed sky area (in sterad), and where we have eliminated the dependence on $\mathbf{l}-\mathbf{k}$ by introducing an additional integral over $d\mathbf{q}^2\delta(\mathbf{q}-\mathbf{l}+\mathbf{k})$.

As shown in [4], these expressions can be simplified through the following steps:

1. Substitute

$$\delta(\mathbf{q}-\mathbf{l}+\mathbf{k}) \to \int \frac{d\mathbf{r}^2}{(2\pi)^2} e^{i(\mathbf{q}-\mathbf{l}+\mathbf{k})\mathbf{r}} \tag{69}$$

2. Integrate over the angular parts of $\mathbf{r}$, $\mathbf{l}$ and $\mathbf{k}$ using the following relation:

$$\int_0^{2\pi} d\varphi \, e^{i\,x\,\cos\varphi} e^{i\,n\varphi} = 2\pi \, i^n \, J_n(x), \tag{70}$$

where $J_n$ is the cylindrical Bessel function of order $n$.

3. Integrate over the angular part of $\mathbf{q}$, defining the pseudo-power spectrum of the masks:

$$\tilde{C}_q^{vw} \equiv \frac{(2\pi)^2}{S} \int d\varphi_q v_{\mathbf{q}} w_{\mathbf{q}}^* \tag{71}$$

9

4. Solve the last isolated integral over the radial part of **r** by using the following relation:

$$\int_0^\infty dr\, r\, J_0(qr)\, J_n(kr)\, J_n(\ell r) = \frac{\cos n\theta}{\pi k \ell \sin\theta},$$

(72)

where $\theta$ is the angle between sides $\ell$ and $k$ of the triangle formed by three sides of length $q$, $\ell$ and $k$.

5. Make the change of variables $q^2(\ell, k, \theta) \equiv \ell^2 + k^2 - 2k\ell\cos\theta$ to simplify the integral over the radial part of **q**.

This yields the following relation analogous to Eq. 38:

$$_v\tilde{\mathsf{C}}_\ell^{ab} = \int_0^\infty dk\, \mathsf{M}_{\ell k}^{s_a s_b} \cdot\, _v\hat{\mathsf{C}}_k^{ab},$$

(73)

where:

$$\mathsf{M}_{\ell k}^{00} = \frac{k}{2\pi} \int_0^\pi \frac{d\theta}{\pi} \tilde{C}_{q(\ell,k,\theta)}^{vw}$$

(74)

$$\mathsf{M}_{\ell\ell'}^{02} = M_{\ell\ell'}^{0+}\,\hat{\mathbb{1}}, \qquad M_{\ell k}^{0+} = \frac{k}{2\pi} \int_0^\pi \frac{d\theta}{\pi} \tilde{C}_{q(\ell,k,\theta)}^{vw} \cos 2\theta$$

(75)

$$\mathsf{M}_{\ell\ell'}^{22} = \begin{pmatrix} M_{\ell\ell'}^{++} & 0 & 0 & M_{\ell\ell'}^{--} \\ 0 & M_{\ell\ell'}^{++} & -M_{\ell\ell'}^{--} & 0 \\ 0 & -M_{\ell\ell'}^{--} & M_{\ell\ell'}^{++} & 0 \\ M_{\ell\ell'}^{--} & 0 & 0 & M_{\ell\ell'}^{++} \end{pmatrix}, \qquad M_{\ell\ell'}^{\pm\pm} = \frac{k}{2\pi} \int_0^\pi \frac{d\theta}{\pi} \tilde{C}_{q(\ell,k,\theta)}^{vw} \frac{1 \pm \cos 4\theta}{2}.$$

(76)

These can also be expressed as integrals over $q$:

$$M_{\ell k}^{00} = \frac{k}{2\pi} \int \frac{q\, dq}{2\pi} \tilde{C}_q^{vw} F(\ell, k, q)$$

(77)

$$M_{\ell k}^{0+} = \frac{k}{2\pi} \int \frac{q\, dq}{2\pi} \tilde{C}_q^{vw} F(\ell, k, q) \frac{\ell^4 + k^4 + q^4 - 2k^2 q^2 - 2\ell^2 q^2}{2k^2\ell^2}$$

(78)

$$M_{\ell k}^{++} = \frac{k}{2\pi} \int \frac{q\, dq}{2\pi} \tilde{C}_q^{vw} F(\ell, k, q) \left[ \frac{\ell^4 + k^4 + q^4 - 2k^2 q^2 - 2\ell^2 q^2}{2k^2\ell^2} \right]^2$$

(79)

$$M_{\ell k}^{--} = \frac{k}{2\pi} \int \frac{q\, dq}{2\pi} \tilde{C}_q^{vw} F(\ell, k, q) \left( 1 - \left[ \frac{\ell^4 + k^4 + q^4 - 2k^2 q^2 - 2\ell^2 q^2}{2k^2\ell^2} \right]^2 \right),$$

(80)

where

$$F(\ell, k, q) \equiv \begin{cases} 4\left(2\ell^2 k^2 + 2k^2 q^2 + 2q^2\ell^2 - \ell^4 - k^4 - q^4\right)^{-1/2} & \text{if } |k - q| < \ell < k + q \\ 0 & \text{otherwise} \end{cases}$$

(81)

Since in the flat-sky limit $\ell$ is a continuous variable, we define bandpowers $_v\mathsf{B}_b$ as averages over a given interval in $\ell$, $[\ell_b^i, \ell_b^0]$:

$$_v\tilde{\mathsf{B}}_b \equiv N_b \int_{\ell_b^i}^{\ell_b^f} \frac{d\ell}{\ell_b^f - \ell_b^i}\, _v\tilde{\mathsf{C}}_\ell.$$

(82)

The binned coupling matrix is therefore given by:

$$\mathcal{M}_{bb'} \equiv \int_{\ell_b^i}^{\ell_b^f} \frac{d\ell}{\ell_b^f - \ell_b^i} \int_{\ell_b'^i}^{\ell_b'^f} d\ell'\, \mathsf{M}_{\ell\ell'}$$

(83)

## 5.3 Pseudo-$C_\ell$ estimator - discrete formalism

The main complication in using the pseudo-$C_\ell$ formalism in the continuum limit is the need to compute the angle-averaged mask pseudo power spectrum $C_\ell^{vw}$. Flat fields are most easily analyzed when pixelized in a Cartesian grid, and under this setup the mask power spectrum is not a well defined quantity for

infinitesitally small intervals of $\ell$. This leads to non-negligible biases and poor performance [4] due to the need to use highly resolved finite intervals to compute the integrals presented in the previous section. This motivates the discrete formalism described here, which connects directy with the data storage format of flat-sky maps.

We will discuss the flat-sky pseudo-$C_\ell$ algorithm starting from a pixelized representation of the sky map. Let the patch of sky under inspection be contained by a rectangle of sides $L_x$ and $L_y$ (in units of radians), and let us discretize this rectangle by dividing it into an $N_x \times N_y$ grid with pixels of area $\Delta \mathbf{x}^2 \equiv \Delta x \Delta y \equiv (L_x/N_x)(L_y/N_y)$. Each pixel in this grid is then labelled by a pair of integers $\mathbf{n} \equiv (n_x, n_y)$, and is assigned coordinates $\mathbf{x_n} \equiv (n_x \Delta x, n_y \Delta y)$. The each component of the pixelized map $\mathbf{a}(\mathbf{x})$ is therefore defined for $n_x \in [0, N_x - 1]$, $n_y \in [0, N_y - 1]$.

In this scenario, the spin Fourier transform of the pixelized field can be computed as its discrete Fourier transform:

$$\mathbf{a_k} \equiv \mathrm{DFT}\,(\mathbf{a})_{\mathbf{k}}^{s_a} \equiv \sum_{\mathbf{x}} \frac{\Delta \mathbf{x}^2}{2\pi} \,{}_{s_a}\mathsf{E_k}(\mathbf{x})\mathbf{a_x}, \tag{84}$$

where the wavenumber $\mathbf{k}$ is now discretized as $\mathbf{k} = (j_x \Delta k_x, j_y \Delta k_y)$, the integers $j_{(x,y)}$ run from $-N_{(x,y)}/2$ tof $N_{(x,y)}/2 - 1$[1] and the pixel size is $\Delta k_{(x,y)} \equiv 2\pi/L_{(x,y)}$.

The following properties of the DFT are worth remembering:

- Periodicity: $\mathbf{a}_{(k_x,k_y)} = \mathbf{a}_{(k_x+N_x\Delta k_x, k_y)} = \mathbf{a}_{(k_x, k_y+N_y\Delta k_y)} = \mathbf{a}_{(k_x+N_x\Delta k_x, k_y+N_y\Delta k_y)}$

- The power spectrum of the DFT is defined as:

$$\langle \mathbf{a_l b_k^\dagger} \rangle \equiv \frac{\delta_{\mathbf{l,k}}}{\Delta \mathbf{k}^2} \mathsf{C}_{\mathbf{l}}^{ab} \tag{85}$$

- For a real-valued scalar $a$, its DFT satisfies $a_{\mathbf{k_j}}^* = a_{\mathbf{k_{N-j}}}$ (where $\mathbf{N} \equiv (N_x, N_y)$).

- The orthogonality relation of the basis functions now takes the form:

$$\sum_{\mathbf{x}} {}_s E_{\mathbf{l}}^\dagger(\mathbf{x}) \,{}_s E_{\mathbf{k}}(\mathbf{x}) = N_x N_y \delta_{\mathbf{l,k}} \hat{1} \tag{86}$$

Using the properties above one can prove that the Fourier coefficients of a masked field are:

$$\mathbf{a_l^v} \equiv (v(\mathbf{x})\mathbf{a}(\mathbf{x}))_{\mathbf{l}} = \sum_{\mathbf{k}} \frac{\Delta \mathbf{k}^2}{2\pi} \mathsf{R}^\dagger(s_a(\varphi_{\mathbf{l}} - \varphi_{\mathbf{k}}))\mathbf{a_k} v_{\mathbf{l-k}}, \tag{87}$$

and we can easily find the following relation for the covariance of two masked fields at the same wavenumber:

$$\left\langle \mathbf{a_l^v b_l^{w\dagger}} \right\rangle = \sum_{\mathbf{k}} \frac{\Delta \mathbf{k}^2}{(2\pi)^2} \mathsf{R}^\dagger(s_a \Delta\varphi)\mathsf{C}_{\mathbf{k}}^{ab}\mathsf{R}^\dagger(s_b\Delta\varphi)v_{\mathbf{l-k}} w_{\mathbf{l-k}}^* \tag{88}$$

At this stage it is natural to connect directly with the final bandpowers (since there is no natural minimal binning as in the case of the full-sky multipoles). In this case we will define a bandpower $\mathsf{B}_b$, indexed by an integer $b$ as the average of the covariance above over a set of $\mathbf{l}$ values $S_b$:

$$_v\mathsf{B}_b^{ab} \equiv \frac{\Delta \mathbf{k}^2}{N_b} \sum_{\mathbf{l} \in S_b} \mathbf{a_l^v b_l^{w\dagger}}, \tag{89}$$

where $N_b$ is the number of Fourier-space pixels covered by $S_b$. The mode-coupling matrix for the bandpowers is therefore given by:

$$\mathcal{M}_{bb'}^{s_a s_b} = \sum_{\mathbf{l} \in S_b} N_b^{-1} \sum_{\mathbf{k} \in S_{b'}} \mathsf{M}_{\mathbf{lk}}^{s_a s_b}, \tag{90}$$

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[1]This is the valid domain when $N_{(x,y)}$ is even. For odd $N_{(x,y)}$ the interval becomes $[-(N_{(x,y)} - 1)/2, (N_{(x,y)} - 1)/2]$

where:

$$M_{\mathbf{l\,k}}^{00} \equiv \frac{2\pi^2}{L_x^2 L_y^2} v_{\mathbf{l-k}} w_{\mathbf{l-k}}^*, \tag{91}$$

$$\mathsf{M}_{\mathbf{l\,k}}^{02} \equiv \frac{2\pi^2}{L_x^2 L_y^2} v_{\mathbf{l-k}} w_{\mathbf{l-k}}^* \begin{pmatrix} \cos 2\Delta\varphi & -\sin 2\Delta\varphi \\ \sin 2\Delta\varphi & \cos 2\Delta\varphi \end{pmatrix} \tag{92}$$

$$\mathsf{M}_{\mathbf{l\,k}}^{02} \equiv \frac{2\pi^2}{L_x^2 L_y^2} v_{\mathbf{l-k}} w_{\mathbf{l-k}}^* \begin{pmatrix} \cos^2 2\Delta\varphi & -\cos 2\Delta\varphi \sin 2\Delta\varphi & -\cos 2\Delta\varphi \sin 2\Delta\varphi & \sin^2 2\Delta\varphi \\ \cos 2\Delta\varphi \sin 2\Delta\varphi & \cos^2 2\Delta\varphi & -\sin^2 2\Delta\varphi & -\cos 2\Delta\varphi \sin 2\Delta\varphi \\ \cos 2\Delta\varphi \sin 2\Delta\varphi & -\sin^2 2\Delta\varphi & \cos^2 2\Delta\varphi & -\cos 2\Delta\varphi \sin 2\Delta\varphi \\ \sin^2 2\Delta\varphi & \cos 2\Delta\varphi \sin 2\Delta\varphi & \cos 2\Delta\varphi \sin 2\Delta\varphi & \cos^2 2\Delta\varphi \end{pmatrix} \tag{93}$$

### 5.4 Contaminant cleaning

Using the same notation as in Section 3, the contaminant-cleaned version of **a** is (c.f. Eqs. 21 23):

$$\tilde{\mathbf{a}}(\mathbf{x}) = \mathbf{a}^v(\mathbf{x}) - \mathbf{f}^i(\mathbf{x}) M_{ij} \sum_{\mathbf{x}} \Delta\mathbf{x}^2\, \mathbf{f}^{j\dagger}(\mathbf{x}') \mathbf{a}^v(\mathbf{x}'), \quad \tilde{\mathbf{a}}_{\mathbf{l}} = \mathbf{a}_{\mathbf{l}}^v - \mathbf{f}_{\mathbf{l}}^i M_{ij} \sum_{\mathbf{k}} \Delta\mathbf{k}^2\, \mathbf{f}_{\mathbf{k}}^{j\dagger} \mathbf{a}_{\mathbf{k}}^v, \tag{94}$$

and the unbiased pseudo-$C_\ell$ estimator takes the form (c.f. Eq. 36):

$$\begin{aligned}
\left\langle \tilde{\mathbf{a}}_{\mathbf{l}} \tilde{\mathbf{b}}_{\mathbf{l}}^\dagger \right\rangle &= \left\langle \mathbf{a}_{\mathbf{l}}^v \mathbf{b}_{\mathbf{l}}^{w\dagger} \right\rangle + N_{ij}^* \operatorname{DFT} \left\{ v(\mathbf{x}) \operatorname{DFT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab} \operatorname{DFT} \left( w\, \mathbf{g}^j \right)_{\mathbf{l}_1}^{s_b} \right]_{\mathbf{x}}^{s_a} \right\}_{\mathbf{l}}^{s_a} \mathbf{g}_{\mathbf{l}}^{i\dagger} + \\
&\quad + M_{ij} \mathbf{f}_{\mathbf{l}}^i \operatorname{DFT} \left\{ w(\mathbf{x}) \operatorname{DFT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab\dagger} \operatorname{DFT} \left( v\, \mathbf{f}^j \right)_{\mathbf{l}_1}^{s_a} \right]_{\mathbf{x}}^{s_b} \right\}_{\mathbf{l}}^{s_b\dagger} - \\
&\quad - M_{ij} N_{pq}^* \left\{ \sum_{\mathbf{x}} \Delta\mathbf{x}^D\, v(\mathbf{x})\, \mathbf{f}^{j\dagger}(\mathbf{x}) \operatorname{DFT}^{-1} \left[ \hat{\mathsf{C}}_{\ell_1}^{ab} \operatorname{DFT} \left( w\, \mathbf{g}^q \right)_{\mathbf{l}_1}^{s_b} \right]_{\mathbf{x}}^{s_a} \right\} \mathbf{f}_{\mathbf{l}}^i \mathbf{g}_{\mathbf{l}}^{p\dagger}
\end{aligned} \tag{95}$$

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