

Notes for “The Quantum Theory of Fields 1, Foundations” - Weinberg

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abc

I. INTRODUCTION

4. Eq. (2.2.22) (P.54)

abc

II. RELATIVISTIC QUANTUM MECHANICS

$$\begin{aligned} -if_{bc}^a t_a - t_b t_c &\stackrel{(2.2.21)}{=} t_{bc} = t_{cb} \stackrel{(2.2.21)}{=} -if_{cb}^a t_a - t_c t_b \\ \Leftrightarrow [t_b, t_c] &= i(f_{cb}^a - f_{bc}^a) t_a \end{aligned}$$

A. Quantum Mechanics

B. Symmetries

1. “For this to be unitary and linear, t must be Hermitian and linear” (P.51)

Linearity is trivial and hermiticity follow from the following observation:

$$\begin{aligned} \langle U\Psi|U\Phi\rangle &= \langle (1 + i\varepsilon t)\Psi|(1 + i\varepsilon t)\Phi\rangle \\ &= \langle \Psi|\Phi\rangle + \varepsilon i(\langle \Psi|t\Phi\rangle - \langle t\Psi|\Phi\rangle) + \mathcal{O}(\varepsilon^2) \end{aligned}$$

$$\text{Eq. (2.2.2)} \quad \Leftrightarrow \langle \Psi|t\Phi\rangle = \langle t\Psi|\Phi\rangle$$

$$\text{Eq. (2.1.5)} \quad t^\dagger = t$$

2. Eq. (2.2.19) (P.54)

f_{bc}^a and f^a have to be real as θ^a are real.

3. Eq. (2.2.21) (P.54)

From Eq. (2.2.20) we have up to $\mathcal{O}(\theta^2, \bar{\theta}^2)$

$$\begin{aligned} 1 + i(\theta^a + \bar{\theta}^a + f_{bc}^a \bar{\theta}^b \theta^c) t_a + \frac{1}{2}(\theta^b + \bar{\theta}^b)(\theta^c + \bar{\theta}^c) t_{bc} \\ = \left[1 + i\bar{\theta}^a t_a + \frac{1}{2}\bar{\theta}^b \bar{\theta}^c t_{bc} \right] \cdot \left[1 + i\theta^a t_a + \frac{1}{2}\theta^b \theta^c t_{bc} \right] \\ = 1 + i\theta^a t_a + \frac{1}{2}\theta^b \theta^c t_{bc} + i\bar{\theta}^a t_a - \bar{\theta}^b t_b \theta^c t_c + \frac{1}{2}\bar{\theta}^b \bar{\theta}^c t_{bc} \\ \Leftrightarrow if_{bc}^a \bar{\theta}^b \theta^c t_a + \frac{1}{2}(\theta^b \bar{\theta}^c + \bar{\theta}^b \theta^c) t_{bc} = -\bar{\theta}^b \theta^c t_b t_c \\ \Leftrightarrow \bar{\theta}^b \theta^c [t_{bc} + if_{bc}^a t_a + t_b t_c] = 0 \\ \Leftrightarrow t_{bc} = -if_{bc}^a t_a - t_b t_c \end{aligned}$$

C. Quantum Lorentz Transformations

1. “ $\Lambda^\mu{}_\nu$ has an *inverse*” (P.57)

This is true because

$$\det(\Lambda) = \pm 1 \neq 0$$

2. “ $(\bar{\Lambda}\Lambda)^0{}_0 \geq \bar{\Lambda}^0{}_0 \Lambda^0{}_0 - \sqrt{(\Lambda^0{}_0)^2 - 1} \sqrt{(\bar{\Lambda}^0{}_0)^2 - 1} \geq 1$ ” (P.58)

The first inequality is trivial from the preceding considerations. For the second inequality assume the contrary holds, then:

$$\begin{aligned} 0 &\leq \bar{\Lambda}^0{}_0 \Lambda^0{}_0 - 1 < \sqrt{(\Lambda^0{}_0)^2 - 1} \sqrt{(\bar{\Lambda}^0{}_0)^2 - 1} \\ &\Rightarrow (\bar{\Lambda}^0{}_0)^2 (\Lambda^0{}_0)^2 - 2\bar{\Lambda}^0{}_0 \Lambda^0{}_0 + 1 \\ &< (\bar{\Lambda}^0{}_0)^2 (\Lambda^0{}_0)^2 - (\bar{\Lambda}^0{}_0)^2 - (\Lambda^0{}_0)^2 + 1 \\ &\Rightarrow (\Lambda^0{}_0 + \bar{\Lambda}^0{}_0)^2 = (\bar{\Lambda}^0{}_0)^2 - 2\bar{\Lambda}^0{}_0 \Lambda^0{}_0 + (\Lambda^0{}_0)^2 < 0 \end{aligned}$$

Which is a contradiction as $\Lambda^0{}_0 + \bar{\Lambda}^0{}_0 \geq 1 + 1 = 2$ and therefore completes the proof.

D. The Poincaré Algebra

1. “In order for $U(1 + \omega, \varepsilon)$ to be unitary, the operators $J^{\rho\sigma}$ and P^ρ must be **Hermitian**” (P.59)

Analog to IIB 1.

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2. Eqs. (2.4.8/9) (P.60)

$$\begin{aligned}
& \frac{1}{2} \omega_{\rho\sigma} U J^{\rho\sigma} U^{-1} - \varepsilon_\rho U P^\rho U^{-1} \\
& \stackrel{(2.4.7)}{=} \frac{1}{2} (\Lambda \omega \Lambda^{-1})_{\mu\nu} J^{\mu\nu} - (\Lambda \varepsilon - \Lambda \omega \Lambda^{-1} a)_\mu P^\mu \\
& = \frac{1}{2} \Lambda_\mu{}^\rho \omega_{\rho\sigma} (\Lambda^{-1})^\sigma{}_\nu J^{\mu\nu} \\
& - \left(\Lambda_\mu{}^\rho \varepsilon_\rho - \Lambda_\mu{}^\rho \omega_{\rho\sigma} (\Lambda^{-1})^\sigma{}_\nu a^\nu \right) P^\mu \\
& \stackrel{(2.3.10)}{=} \frac{1}{2} \Lambda_\mu{}^\rho \omega_{\rho\sigma} \Lambda_\nu{}^\sigma J^{\mu\nu} \\
& - \left(\Lambda_\mu{}^\rho \varepsilon_\rho - \Lambda_\mu{}^\rho \omega_{\rho\sigma} \Lambda_\nu{}^\sigma a^\nu \right) P^\mu \\
& = \frac{1}{2} \omega_{\rho\sigma} (\Lambda_\mu{}^\rho \Lambda_\nu{}^\sigma J^{\mu\nu} + \Lambda_\mu{}^\rho \Lambda_\nu{}^\sigma a^\nu P^\mu) \\
& - \varepsilon_\rho \Lambda_\mu{}^\rho P^\mu
\end{aligned}$$

In order to be able to compare coefficients in this, the coefficient of $\omega_{\rho\sigma}$ has to be anti symmatrized. With this Eqs. (2.4.8/9) follow immediately.

3. Eqs. (2.4.10/11) (P.60)

Up to $\mathcal{O}(\omega, \varepsilon)$ one can identify

$$U^{-1}(1 + \omega, \varepsilon) = U(1 - \omega, -\varepsilon)$$

,since

$$U(1 + \omega, \varepsilon)U(1 - \omega, -\varepsilon) = U(1 - \omega + \omega, -\varepsilon + \varepsilon) = U(1, 0).$$

With this we have up to $\mathcal{O}(\omega, \varepsilon)$

$$\begin{aligned}
& i \left[\frac{1}{2} \omega_{\mu\nu} J^{\mu\nu} - \varepsilon_\mu P^\mu, J^{\rho\sigma} \right] \\
& = \left(1 + \frac{1}{2} i \omega_{\mu\nu} J^{\mu\nu} - i \varepsilon_\mu P^\mu \right) J^{\rho\sigma} \\
& \cdot \left(1 - \frac{1}{2} i \omega_{\mu\nu} J^{\mu\nu} + i \varepsilon_\mu P^\mu \right) - J^{\rho\sigma} \\
& = U J^{\rho\sigma} U^{-1} - J^{\rho\sigma} \\
& \stackrel{(2.4.8)}{=} (1 + \omega)_\mu{}^\rho (1 + \omega)_\nu{}^\sigma \\
& \cdot (J^{\mu\nu} - \varepsilon^\mu P^\nu + \varepsilon^\nu P^\mu) - J^{\rho\sigma} \\
& = -\varepsilon^\rho P^\sigma + \varepsilon^\sigma P^\rho + \omega_\nu{}^\sigma J^{\rho\nu} + \omega_\mu{}^\rho J^{\mu\sigma}
\end{aligned}$$

and also

$$\begin{aligned}
& i \left[\frac{1}{2} \omega_{\mu\nu} J^{\mu\nu} - \varepsilon_\mu P^\mu, P^\rho \right] \\
& = \left(1 + \frac{1}{2} i \omega_{\mu\nu} J^{\mu\nu} - i \varepsilon_\mu P^\mu \right) P^\rho \\
& \cdot \left(1 - \frac{1}{2} i \omega_{\mu\nu} J^{\mu\nu} + i \varepsilon_\mu P^\mu \right) - P^\rho \\
& = U P^\rho U^{-1} - P^\rho \\
& \stackrel{(2.4.9)}{=} \omega_\mu{}^\rho P^\mu
\end{aligned}$$

4. Eqs. (2.4.12/13/14) (P.60)

The only difficulty lies in the derivation of Eq. (2.4.12) because for this one has to anti symmatrize the coefficient of $\omega_{\mu\nu}$ in Eq. (2.4.10) in order to be able to compare coefficients, similar to II D 2. (Similarly if one derives Eq. (2.4.13) from Eq. (2.4.11))

5. Eqs. (2.4.18 – 24) (P.61)

First observe

$$J^{ij} = \varepsilon_{ijk} J_k$$

and

$$J_i = \frac{1}{2} \varepsilon_{lmi} J^{lm}$$

$$\begin{aligned}
[J_i, J_j] &= \frac{\varepsilon_{lmi} \varepsilon_{kpj}}{4} [J^{lm}, J^{kp}] \\
&\stackrel{(2.4.12)}{=} -i \frac{\varepsilon_{lmi} \varepsilon_{kpj}}{4} [\eta^{mk} J^{lp} - \eta^{lk} J^{mp} - \eta^{pl} J^{km} + \eta^{pm} J^{kl}] \\
&= -\frac{i}{2} [\varepsilon_{kil} \varepsilon_{kmj} J^{lm} + \varepsilon_{kim} \varepsilon_{kjl} J^{lm}] \\
&= -\frac{i}{2} [J^{ji} - J^{ij}] \\
&= i J^{ij} = i \varepsilon_{ijk} J_k \\
[J_i, K_j] &= [J^{lm}, J^{0j}] \frac{\varepsilon_{lmi}}{2} \\
&\stackrel{(2.4.12)}{=} -i \frac{\varepsilon_{lmi}}{2} [\eta^{m0} J^{lj} - \eta^{l0} J^{mj} - \eta^{jl} J^{0m} + \eta^{jm} J^{0l}] \\
&= -i \frac{\varepsilon_{lmi}}{2} [\delta_{jm} K_l - \delta_{jl} K_m] = i \varepsilon_{ijl} K_l \\
[K_i, K_j] &= [J^{0i}, J^{0j}] \\
&\stackrel{(2.4.12)}{=} -i [\eta^{i0} J^{0j} - \eta^{00} J^{ij} - \eta^{j0} J^{0i} + \eta^{ij} J^{00}] \\
&= -i J^{ij} = -i \varepsilon_{ijk} J_k \\
[J_i, P_j] &= \frac{\varepsilon_{lmi}}{2} [J^{lm}, P^j] \\
&\stackrel{(2.4.13)}{=} \frac{i \varepsilon_{lmi}}{2} [\eta^{jl} P^m - \eta^{jm} P^l] \\
&= \frac{i}{2} [\varepsilon_{jmi} P^m - \varepsilon_{mji} P^m] = i \varepsilon_{ijm} P_m \\
[K_i, P_j] &= [J^{0i}, P^j] \\
&\stackrel{(2.4.13)}{=} i [\eta^{j0} P^i - \eta^{ji} P^0] \\
&= -i \delta_{ji} P^0 = -i \delta_{ij} H \\
[J_i, H] &\stackrel{(2.4.13)}{=} \frac{i \varepsilon_{lmi}}{2} [\eta^{0l} P^m - \eta^{0m} P^l] = 0 \\
[P_i, H] &= [P^i, P^0] \stackrel{(2.4.14)}{=} 0 \\
[H, H] &= [P^0, P^0] \stackrel{(2.4.14)}{=} 0 \\
[K_i, H] &= [J^{0i}, P^0] \\
&\stackrel{(2.4.13)}{=} i [\eta^{00} P^i - \eta^{0i} P^0] = -i P_i
\end{aligned}$$

6. Eq. (2.4.27) (P.61)

When trying to check this for e.g. the standard representation in 4 dimensions, in terms of infinitesimal rotations, attention with the sign in front of sin for the different rotation axis (see remark at the end of I E 23).

7. “Inspection of Eqs. (2.4.18 – 24) shows that these commutation relations have a limit for $v \ll 1$ of the form ...” (P.62)

Always equate same orders in v for this.

$$8. \quad \text{“exp}(-i\mathbf{K} \cdot \mathbf{v}) \text{exp}(-i\mathbf{P} \cdot \mathbf{a}) = \text{exp}(iM\mathbf{a} \cdot \mathbf{v}/2) \text{exp}(-i(\mathbf{K} \cdot \mathbf{v} + \mathbf{P} \cdot \mathbf{a})) \text{” (P.62)}$$

Use BCH Formula

$$\begin{aligned} & \exp(-iK_i v_i) \exp(-iP_j a_j) \\ &= \exp\left(-i(K_i v_i + P_j a_j) + \frac{1}{2}(-i)^2 [K_i, P_j] v_i a_j + 0\right) \\ &= \exp\left(-i(K_i v_i + P_j a_j) + \frac{1}{2}iM v_i a_i\right) \end{aligned}$$

E. One-Particle States

1. Eq. (2.5.2) (P.63)

Λ^{-1} shows up here since in Eq. (2.4.9) UPU^{-1} is given but here $U^{-1}PU$ is being used.

2. “with σ within any one block by themselves furnish a representation of the inhomogeneous Lorentz group” (P.63)

In this case the Blocks do not mix with other blocks.

3. “and for $p^2 \leq 0$, also the sign of p^0 ” (P.64)

For $p^2 \leq 0$ we have

$$\begin{aligned} p^2 &= -(p^0)^2 + \vec{p}^2 \leq 0 \\ \Rightarrow |\vec{p}| &\leq |p^0| \end{aligned}$$

and from Eq. (2.3.13) we know

$$|\Lambda^0_0| \geq |\Lambda^{\vec{0}}_0|.$$

First suppose $p^0 \geq 0$:

$$\begin{aligned} p'^0 &= \Lambda^0_0 p^0 + \Lambda^0_i p^i \\ &\geq \Lambda^0_0 p^0 - |\Lambda^{\vec{0}}_0| |\vec{p}| \\ &= |\Lambda^0_0| |p^0| - |\Lambda^{\vec{0}}_0| |\vec{p}| \\ &\geq |\Lambda^{\vec{0}}_0| (|p^0| - |\vec{p}|) \geq 0 \end{aligned}$$

Now suppose $p^0 \leq 0$:

$$\begin{aligned} p'^0 &= \Lambda^0_0 p^0 + \Lambda^0_i p^i \\ &\leq \Lambda^0_0 p^0 + |\Lambda^{\vec{0}}_0| |\vec{p}| \\ &= -|\Lambda^0_0| |p^0| + |\Lambda^{\vec{0}}_0| |\vec{p}| \\ &\leq |\Lambda^0_0| (-|p^0| + |\vec{p}|) \leq 0 \end{aligned}$$

4. Eq. (2.5.12) “The delta function appears here because $\Psi_{k,\sigma}$ and $\Psi_{k',\sigma'}$ are eigenstates of a Hermitian operator with eigenvalues \mathbf{k} and \mathbf{k}' , respectively.” (P.66)

Note: From this point forward in this section of the book the states considered have standard momentum k^μ ! (see Text in the Book preceding Eq. (2.5.12))

Suppose that for given values of \mathbf{k} and \mathbf{k}'

$$\langle \Psi_{k',\sigma'} | \Psi_{k,\sigma} \rangle \neq 0$$

then from

$$\begin{aligned} & k^i \langle \Psi_{k',\sigma'} | \Psi_{k,\sigma} \rangle \\ &= \langle \Psi_{k',\sigma'} | P^i \Psi_{k,\sigma} \rangle \\ &= \langle P^i \Psi_{k',\sigma'} | \Psi_{k,\sigma} \rangle \\ &= k'^i \langle \Psi_{k',\sigma'} | \Psi_{k,\sigma} \rangle \end{aligned}$$

follows

$$k^i = k'^i.$$

On why there is only a 3 dimensional delta function showing up:

This is because both states are assumed to have standard momentum k^μ and therefore their 0-th component is completely fixed from their spatial components:

$$k'^0 = \pm \sqrt{\mathbf{k}'^2 - k^2}$$

Choice between $+$ and $-$ comes from sign of k^0 , i.e. $+$ for cases (a) and (c) of Table 2.1.

5. Eq. (2.5.13) (P.67)

$$\begin{aligned}
& \langle U(W)\Psi_{k',\sigma'} | U(W)\Psi_{k,\sigma} \rangle \\
& \stackrel{(2.5.8)}{=} \sum_{\sigma''\sigma'''} \langle D_{\sigma''\sigma'} \Psi_{k',\sigma''} | D_{\sigma'''\sigma} \Psi_{k,\sigma'''} \rangle \\
& = \sum_{\sigma''\sigma'''} (D_{\sigma''\sigma'})^* D_{\sigma'''\sigma} \langle \Psi_{k',\sigma''} | \Psi_{k,\sigma'''} \rangle \\
& \stackrel{(2.5.12)}{=} \sum_{\sigma''} (D_{\sigma''\sigma'})^* D_{\sigma''\sigma} \delta^{(3)}(\mathbf{k}' - \mathbf{k}) \\
& \stackrel{(2.2.2)}{=} \langle \Psi_{k',\sigma'} | \Psi_{k,\sigma} \rangle \\
& \stackrel{(2.5.12)}{=} \delta_{\sigma'\sigma} \delta^{(3)}(\mathbf{k}' - \mathbf{k})
\end{aligned}$$

From this it follows

$$D^\dagger(W) = D^{-1}(W)$$

6. “ $\langle \Psi_{p',\sigma'} | \Psi_{p,\sigma} \rangle = \dots$ ” (P.67)

Note: What is meant by “arbitrary momenta” is that these momenta still have the standard momentum k^μ , but now none of the states has exactly k^μ as its momentum.

First define

$$k' := L^{-1}(p)p'$$

with this we get:

$$\begin{aligned}
& \langle \Psi_{p',\sigma'} | \Psi_{p,\sigma} \rangle \\
& \stackrel{(2.5.5)}{=} \langle \Psi_{p',\sigma'} | N(p) U(L(p)) \Psi_{k,\sigma} \rangle \\
& = N(p) \langle U(L^{-1}(p)) \Psi_{p',\sigma'} | \Psi_{k,\sigma} \rangle \\
& \stackrel{(2.5.11)}{=} N(p) \frac{N^*(p')}{N^*(k')} \\
& \cdot \sum_{\sigma''} D_{\sigma''\sigma'}^* (W(L^{-1}(p), p')) \langle \Psi_{k',\sigma''} | \Psi_{k,\sigma} \rangle \\
& \stackrel{(2.5.12)}{=} N(p) N^*(p') D_{\sigma\sigma'}^* (W(L^{-1}(p), p')) \delta^{(3)}(\mathbf{k}' - \mathbf{k})
\end{aligned}$$

Where in the last step we used

$$\frac{\delta^{(3)}(\mathbf{k}' - \mathbf{k})}{N^*(k')} = \frac{\delta^{(3)}(\mathbf{k}' - \mathbf{k})}{N^*(k)} \stackrel{(2.5.5)}{=} \delta^{(3)}(\mathbf{k}' - \mathbf{k})$$

This is valid as from Eq. (2.5.5) we know that $N(p)$ is implicitly dependent on k^μ . This therefor fixes $p^2 = k^2$

and the sign of p^0 (see end of II E 4), s.t.

$$N(p) = N(\mathbf{p}).$$

Further we are allowed to use Eq. (2.5.12) in the last step since k' also has k as its standard momentum:

$$\begin{aligned}
k' &= L^{-1}(p)p' = \underbrace{L^{-1}(p)L(p')}_{=:L(k')} k \\
7. \quad & “W(L^{-1}(p), p) = 1” \text{ (P.67)}
\end{aligned}$$

First note

$$L(k) = 1 = L^{-1}(k)$$

with this we get:

$$\begin{aligned}
W(L^{-1}(p), p) & \stackrel{(2.5.10)}{=} L^{-1}(L^{-1}(p)p) L^{-1}(p)L(p) \\
& = L^{-1}(k) = 1
\end{aligned}$$

8. “So we see that the **invariant delta function** is ...” (P.67)

Invariant in this case means w.r.t. proper orthochronous Lorentz transformations, which can be interpreted as change of variables under the $\int d^4p$ integral.

9. Eq. (2.5.24) (P.68)

$$\begin{aligned}
p^0 &= L^0_0 k^0 + L^0_i k^i \\
&= \frac{\sqrt{\mathbf{p}^2 + M^2}}{M} M = \sqrt{\mathbf{p}^2 + M^2} \\
p^i &= L^i_0 k^0 + L^i_j k^j \\
&= \frac{p_i}{|\mathbf{p}|} \sqrt{\frac{\mathbf{p}^2 + M^2}{M^2}} - 1M \\
&= \frac{p_i}{|\mathbf{p}|} \sqrt{\frac{\mathbf{p}^2}{M^2}} M = p_i
\end{aligned}$$

10. “To see this, note that the boost Eq. (2.5.24) may be expressed as $L(p) = R(\hat{\mathbf{p}})B(|\mathbf{p}|)R^{-1}(\hat{\mathbf{p}})$ ” (P.68)

First note that the columns and rows of the matrix $B(|\mathbf{p}|)$ are counted in the order 1,2,3,0.

Let

$$\hat{\mathbf{p}} = \begin{pmatrix} \sin(\theta) \cos(\phi) \\ \sin(\theta) \sin(\phi) \\ \cos(\theta) \end{pmatrix}$$

then we have

$$R(-\hat{\mathbf{p}}) = R_3(-\phi)R_2(\theta) = \begin{pmatrix} \cos(\phi)\cos(\theta) & -\sin(\phi) & \cos(\phi)\sin(\theta) \\ \sin(\phi)\cos(\theta) & \cos(\phi) & \sin(\phi)\sin(\theta) \\ -\sin(\theta) & 0 & \cos(\theta) \end{pmatrix}$$

(see IIE 23) and

$$\begin{aligned} R^{-1}(\hat{\mathbf{p}}) &= R_2(\theta)R_3(\phi) = \begin{pmatrix} \cos(\theta) & 0 & -\sin(\theta) \\ 0 & 1 & 0 \\ \sin(\theta) & 0 & \cos(\theta) \end{pmatrix} \begin{pmatrix} \cos(\phi) & \sin(\phi) & 0 \\ -\sin(\phi) & \cos(\phi) & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} \cos(\theta)\cos(\phi) & \cos(\theta)\sin(\phi) & -\sin(\theta) \\ -\sin(\phi) & \cos(\phi) & 0 \\ \sin(\theta)\cos(\phi) & \sin(\theta)\sin(\phi) & \cos(\theta) \end{pmatrix} = R^\top(\hat{\mathbf{p}}) \end{aligned}$$

From Eq. (2.5.24) we know

$$L(p) = \begin{pmatrix} 1 + (\gamma - 1)s_\theta^2 c_\phi^2 & (\gamma - 1)s_\theta c_\phi s_\theta s_\phi & (\gamma - 1)s_\theta c_\phi c_\theta & s_\theta c_\phi \sqrt{\gamma^2 - 1} \\ (\gamma - 1)s_\theta c_\phi s_\theta s_\phi & 1 + (\gamma - 1)s_\theta^2 s_\phi^2 & (\gamma - 1)s_\theta s_\phi c_\theta & s_\theta s_\phi \sqrt{\gamma^2 - 1} \\ (\gamma - 1)s_\theta c_\phi c_\theta & (\gamma - 1)s_\theta s_\phi c_\theta & 1 + (\gamma - 1)c_\theta^2 & c_\theta \sqrt{\gamma^2 - 1} \\ s_\theta c_\phi \sqrt{\gamma^2 - 1} & s_\theta s_\phi \sqrt{\gamma^2 - 1} & c_\theta \sqrt{\gamma^2 - 1} & \gamma \end{pmatrix}$$

where s and c are shorthand for \sin and \cos . With this we can now check:

$$\begin{aligned} R(\hat{\mathbf{p}})B(|\mathbf{p}|)R^{-1}(\hat{\mathbf{p}}) &= \begin{pmatrix} c_\phi c_\theta & -s_\phi & c_\phi s_\theta & 0 \\ s_\phi c_\theta & c_\phi & s_\phi s_\theta & 0 \\ -s_\theta & 0 & c_\theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \gamma & \sqrt{\gamma^2 - 1} \\ 0 & 0 & \sqrt{\gamma^2 - 1} & \gamma \end{pmatrix} \begin{pmatrix} c_\phi c_\theta & s_\phi c_\theta & -s_\theta & 0 \\ -s_\phi & c_\phi & 0 & 0 \\ c_\phi s_\theta & s_\phi s_\theta & c_\theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} c_\phi c_\theta & -s_\phi & c_\phi s_\theta & 0 \\ s_\phi c_\theta & c_\phi & s_\phi s_\theta & 0 \\ -s_\theta & 0 & c_\theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c_\phi c_\theta & s_\phi c_\theta & -s_\theta & 0 \\ -s_\phi & c_\phi & 0 & 0 \\ \gamma c_\phi s_\theta & \gamma s_\phi s_\theta & \gamma c_\theta & \sqrt{\gamma^2 - 1} \\ \sqrt{\gamma^2 - 1} c_\phi s_\theta & \sqrt{\gamma^2 - 1} s_\phi s_\theta & \sqrt{\gamma^2 - 1} c_\theta & \gamma \end{pmatrix} = L(p) \end{aligned}$$

11. “ $W(\mathbf{R}, p) = \mathbf{R}$ ” (P.69)

13. Eq. (2.5.26) (P.70)

To see this just substitute $R(\theta)$ back into the previous result.

This is a Lorentz transformation, since

12. Eq. (2.5.25) (P.70)

$$\begin{aligned} S^\top \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} S &= \begin{pmatrix} 1 & 0 & \alpha & \alpha \\ 0 & 1 & \beta & \beta \\ -\alpha & -\beta & 1 - \zeta & -\zeta \\ \alpha & \beta & \zeta & 1 + \zeta \end{pmatrix} \\ &\cdot \begin{pmatrix} 1 & 0 & -\alpha & \alpha \\ 0 & 1 & -\beta & \beta \\ \alpha & \beta & 1 - \zeta & \zeta \\ -\alpha & -\beta & \zeta & -1 - \zeta \end{pmatrix} \\ &= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \end{aligned}$$

$$\begin{aligned} -1 &= (Wt)^\mu (Wt)_\mu = \alpha^2 + \beta^2 + \zeta^2 - (1 + \zeta)^2 \\ &\Leftrightarrow \alpha^2 + \beta^2 = 2\zeta \end{aligned}$$

14. Eqs. (2.5.29/30) (P.70)

Eq. (2.5.30) is trivial since the rotations around the three axis satisfy this group multiplication and

$$S(0,0) = \mathbf{1}.$$

For Eq. (2.5.29) explicit calculation shows it (see **TODO**) together with

$$R(0) = \mathbf{1}.$$

15. Eq. (2.5.31) (P.70)

Explicit calculation shows Eq. (2.5.31) (see **TODO**) and the invariance follows then immediately, since:

$$W' S W'^{-1} = S' \underbrace{R' S R'^{-1}}_{=S''} S'^{-1} = S'''$$

16. “ $W(\theta, \alpha, \beta) = 1 + \omega$ ” (P.71)

From Eq. (2.5.28) we immediately get for infinitesimal θ, α, β :

$$(W(\theta, \alpha, \beta))^\mu{}_\nu = \delta^\mu_\nu + \begin{pmatrix} 0 & \theta & -\alpha & \alpha \\ -\theta & 0 & -\beta & \beta \\ \alpha & \beta & 0 & 0 \\ \alpha & \beta & 0 & 0 \end{pmatrix}^\mu{}_\nu$$

$$\Rightarrow \omega_{\mu\nu} = \begin{pmatrix} 0 & \theta & -\alpha & \alpha \\ -\theta & 0 & -\beta & \beta \\ \alpha & \beta & 0 & 0 \\ -\alpha & -\beta & 0 & 0 \end{pmatrix}_{\mu\nu}$$

17. Eq. (2.5.32) (P.71)

$$\begin{aligned} U(1 + \omega) &\stackrel{(2.4.3)}{=} 1 + \frac{1}{2} i \omega_{\rho\sigma} J^{\rho\sigma} \\ &= 1 + i\theta J^{12} + i\alpha(-J^{13} + J^{10}) + i\beta(-J^{23} + J^{20}) \\ &= 1 + i\theta J_3 + i\alpha(J_2 - K_1) + i\beta(-J_1 - K_2) \end{aligned}$$

18. Eqs. (2.5.35/36/37) (P.71)

From Eqs. (2.4.18/19/20) we get:

$$\begin{aligned} [J_3, A] &= [J_3, J_2 - K_1] \\ &= [J_3, J_2] - [J_3, K_1] \\ &= -iJ_1 - iK_2 = iB \\ [J_3, B] &= -[J_3, J_1] - [J_3, K_2] \\ &= -iJ_2 - i(-K_1) = -iA \\ [A, B] &= -[J_2, J_1] - [J_2, K_2] + [K_1, J_1] + [K_1, K_2] \\ &= iJ_3 - iJ_3 = 0 \end{aligned}$$

19. “Since A and B are commuting Hermitian operators they can be simultaneously diagonalized by states $\Psi_{k,a,b}$ ” (P.71)

This is valid, including the label k , since

$$\begin{aligned} [A, P_1] \Psi_{k,a,b} &= ([J_2, P_1] - [K_1, P_1]) \Psi_{k,a,b} \\ &= (-iP_3 + iP^0) \Psi_{k,a,b} = 0 \\ [A, P_2] \Psi_{k,a,b} &= ([J_2, P_2] - [K_1, P_2]) \Psi_{k,a,b} \\ &= 0 \\ [A, P_3] \Psi_{k,a,b} &= ([J_2, P_3] - [K_1, P_3]) \Psi_{k,a,b} \\ &= (iP_1) \Psi_{k,a,b} = 0 \\ [A, P^0] \Psi_{k,a,b} &= ([J_2, P^0] - [K_1, P^0]) \Psi_{k,a,b} \\ &= (-iP_1) \Psi_{k,a,b} = 0 \\ [B, P_1] \Psi_{k,a,b} &= (-[J_1, P_1] - [K_2, P_1]) \Psi_{k,a,b} \\ &= 0 \\ [B, P_2] \Psi_{k,a,b} &= ([J_1, P_2] - [K_2, P_2]) \Psi_{k,a,b} \\ &= (iP_3 - iP^0) \Psi_{k,a,b} = 0 \\ [B, P_3] \Psi_{k,a,b} &= ([J_1, P_3] - [K_2, P_3]) \Psi_{k,a,b} \\ &= (-iP_2) \Psi_{k,a,b} = 0 \\ [B, P^0] \Psi_{k,a,b} &= ([J_1, P^0] - [K_2, P^0]) \Psi_{k,a,b} \\ &= (-iP_2) \Psi_{k,a,b} = 0 \end{aligned}$$

20. “ σ gives the component of angular momentum in the direction of motion, or helicity” (P.72)

The derivation of Eq. (2.5.26) starts among other conditions from the explicit form of k . Which results in

$$\text{Eqs. (2.5.27/28)} \rightarrow \text{Eq. (2.5.32)} \rightarrow 2.5.39,$$

s.t. this is really connected to the direction of motion.

$$21. \quad “\mathcal{U}(W) \Psi_{k,\sigma} = \exp(i\theta\sigma) \Psi_{k,\sigma}” \text{ (P.72)}$$

Use Eqs. (2.5.38/39).

22. Eq. (2.5.44) (P.73)

With $B(u)$ from Eq. (2.5.45) we get:

$$B\left(\frac{|\mathbf{p}|}{\kappa}\right)k = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \frac{(\frac{|\mathbf{p}|}{\kappa})^2+1}{2\frac{|\mathbf{p}|}{\kappa}} & \frac{(\frac{|\mathbf{p}|}{\kappa})^2-1}{2\frac{|\mathbf{p}|}{\kappa}} \\ 0 & 0 & \frac{(\frac{|\mathbf{p}|}{\kappa})^2-1}{2\frac{|\mathbf{p}|}{\kappa}} & \frac{(\frac{|\mathbf{p}|}{\kappa})^2+1}{2\frac{|\mathbf{p}|}{\kappa}} \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ \kappa \\ \kappa \end{pmatrix} \\ = |\mathbf{p}| \begin{pmatrix} 0 \\ 0 \\ 1 \\ 1 \end{pmatrix}$$

23. “take $R(\hat{\mathbf{p}})$ as a rotation by angle θ around the two-axis followed by a rotation by angle ϕ around the three-axis” (P.73)

Let

$$\hat{\mathbf{p}} = \begin{pmatrix} \sin(\theta) \cos(\phi) \\ \sin(\theta) \sin(\phi) \\ \cos(\theta) \end{pmatrix}$$

then

$$R(\hat{\mathbf{p}}) = R_3(-\phi)R_2(-\theta) \\ = \begin{pmatrix} \cos(\phi) & -\sin(\phi) & 0 \\ \sin(\phi) & \cos(\phi) & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \cos(\theta) & 0 & \sin(\theta) \\ 0 & 1 & 0 \\ -\sin(\theta) & 0 & \cos(\theta) \end{pmatrix} \\ = \begin{pmatrix} \cos(\phi) \cos(\theta) & -\sin(\phi) \cos(\theta) & \sin(\theta) \\ \sin(\phi) \cos(\theta) & \cos(\phi) \cos(\theta) & \sin(\theta) \\ -\sin(\theta) & 0 & \cos(\theta) \end{pmatrix},$$

(With this sign convention everything is consistent, see definition of $R(\theta)$ after Eq. (2.5.27) and compare Eq. (2.5.47) to Eq. (2.4.27)) with this we have

$$\hat{\mathbf{p}} = R(\hat{\mathbf{p}})\hat{\mathbf{e}}_3.$$

F. Space Inversion and Time-Reversal

1. Eqs. (2.6.7 – 12) (P.76)

$$\begin{aligned} \mathbf{P}J_i\mathbf{P}^{-1} &= \frac{1}{2}\varepsilon_{ijk}\mathbf{P}J^{jk}\mathbf{P}^{-1} \\ &= \frac{1}{2}\varepsilon_{ijk}\left(-\delta_l^j\right)\left(-\delta_m^k\right)J^{lm} \\ &= \frac{1}{2}\varepsilon_{ijk}J^{jk} = J_i \\ \mathbf{P}K_i\mathbf{P}^{-1} &= \mathbf{P}J^{0i}\mathbf{P}^{-1} \\ &= \delta_\mu^0\left(-\delta_\nu^j\right)J^{\mu\nu} \\ &= -J^{0i} = -K_i \\ \mathbf{P}P_i\mathbf{P}^{-1} &= \left(-\delta_\nu^i\right)P^\mu = -P_i \\ \mathbf{T}J_i\mathbf{T}^{-1} &= \frac{1}{2}\varepsilon_{ijk}\mathbf{T}J^{jk}\mathbf{T}^{-1} \\ &= -\frac{1}{2}\varepsilon_{ijk}\delta_l^j\delta_m^k J^{lm} \\ &= -\frac{1}{2}\varepsilon_{ijk}J^{jk} = -J_i \\ \mathbf{T}K_i\mathbf{T}^{-1} &= \mathbf{T}J^{0i}\mathbf{T}^{-1} \\ &= -\left(-\delta_\mu^0\right)\delta_\nu^j J^{\mu\nu} \\ &= J^{0i} = K_i \\ \mathbf{T}P_i\mathbf{T}^{-1} &= -\delta_\nu^i P^\mu = -P_i \end{aligned}$$

2. “ $\mathcal{P}L(p)\mathcal{P}^{-1} = L(\mathcal{P}p)$ ” (P.77)

From Eq. (2.5.24) we immediately see:

$$\begin{aligned} (\mathcal{P}L(p)\mathcal{P}^{-1})^i_k &= (L(p))^i_k = (L(\mathcal{P}p))^i_k \\ (\mathcal{P}L(p)\mathcal{P}^{-1})^i_0 &= -(L(p))^i_0 = (L(\mathcal{P}p))^i_0 \\ (\mathcal{P}L(p)\mathcal{P}^{-1})^0_0 &= (L(p))^0_0 = (L(\mathcal{P}p))^0_0 \end{aligned}$$

3. “Using Eq. (2.6.14) again on the left, we see that the square-root factors cancel...” (P.77)

$$\begin{aligned} (-J_1 \pm iJ_2)\zeta_\sigma\Psi_{k,-\sigma} &= -(J_1 \mp iJ_2)\zeta_\sigma\Psi_{k,-\sigma} \\ &\stackrel{(2.6.14)}{=} -\sqrt{(j \pm (-\sigma))(j \mp (-\sigma) + 1)}\zeta_\sigma\Psi_{k,-\sigma \mp 1} \\ &= -\sqrt{(j \mp \sigma)(j \pm \sigma + 1)}\zeta_\sigma\Psi_{k,-\sigma \mp 1} \end{aligned}$$

4. “The time-reversal phase ζ has no physical significance.” (P.78)

This redefinition only works because \mathbf{T} is *anti-linear*.

5. “ $\mathcal{T}L(p)\mathcal{T}^{-1} = L(\mathcal{P}p)$ ” (P.78)

From Eq. (2.5.24) we immediately see:

$$\begin{aligned} (\mathcal{T}L(p)\mathcal{T}^{-1})^i_k &= (L(p))^i_k = (L(\mathcal{P}p))^i_k \\ (\mathcal{T}L(p)\mathcal{T}^{-1})^i_0 &= -(L(p))^i_0 = (L(\mathcal{P}p))^i_0 \\ (\mathcal{T}L(p)\mathcal{T}^{-1})^0_0 &= (L(p))^0_0 = (L(\mathcal{P}p))^0_0 \end{aligned}$$

6. “ \mathcal{P} yields a state with four-momentum ...” (P.78)

What is meant by this is:

$$\begin{aligned} P^i \mathcal{P} \Psi_{k,\sigma} &\stackrel{(2.6.9)}{=} \mathcal{P} (-P^i) \Psi_{k,\sigma} = (-\delta^i_3 \kappa) \mathcal{P} \Psi_{k,\sigma} \\ H \mathcal{P} \Psi_{k,\sigma} &\stackrel{(2.6.13)}{=} \mathcal{P} H \Psi_{k,\sigma} = \kappa \mathcal{P} \Psi_{k,\sigma} \\ J_3 \mathcal{P} \Psi_{k,\sigma} &\stackrel{(2.6.7)}{=} \mathcal{P} J_3 \Psi_{k,\sigma} = \sigma \mathcal{P} \Psi_{k,\sigma} \end{aligned}$$

7. Eq. (2.6.20) (P.79)

We have

$$R_2 = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} = R_2^{-1}$$

from which we get (Is $U(R_2^{-1}) = U^{-1}(R_2)$? **TODO**)

$$\begin{aligned} U^{-1}(R_2) J_3 U(R_2) &= U^{-1}(R_2) J^{12} U(R_2) \\ &\stackrel{(2.4.8)}{=} (R_2^{-1})^1_\mu (R_2^{-1})^2_\nu J^{\mu\nu} \\ &= -J^{12} = -J_3 \\ U^{-1}(R_2) P^\nu U(R_2) &\stackrel{(2.4.9)}{=} (R_2^{-1})^\nu_\mu P^\mu \end{aligned}$$

such that

$$\begin{aligned} J_3 U(R_2^{-1}) \mathcal{P} \Psi_{k,\sigma} &= -U(R_2^{-1}) J_3 \mathcal{P} \Psi_{k,\sigma} \\ &\stackrel{II F 6}{=} -\sigma U(R_2^{-1}) \mathcal{P} \Psi_{k,\sigma} \\ P^i U(R_2^{-1}) \mathcal{P} \Psi_{k,\sigma} &\stackrel{II F 6}{=} U(R_2^{-1}) (-(-\delta^i_3 \kappa)) \mathcal{P} \Psi_{k,\sigma} \\ &= \kappa^i U(R_2^{-1}) \mathcal{P} \Psi_{k,\sigma} \\ H U(R_2^{-1}) \mathcal{P} \Psi_{k,\sigma} &\stackrel{II F 6}{=} \kappa U(R_2^{-1}) \mathcal{P} \Psi_{k,\sigma} \\ &= \kappa^0 U(R_2^{-1}) \mathcal{P} \Psi_{k,\sigma}. \end{aligned}$$

Further we get

$$R_2^{-1} \mathcal{P} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

8. “ \mathcal{P} commutes with the rotation $R(\hat{\mathbf{p}})$ ” (P.79)

This is true because $R(\hat{\mathbf{p}})$ only acts on space components non trivially, which all just get a “-” sign from \mathcal{P} .

9. “ $\mathcal{P} \Psi_{p,\sigma} = \sqrt{\frac{\kappa}{p^0}} \eta_\sigma U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,-\sigma}$ ” (P.79)

$$\begin{aligned} \mathcal{P} \Psi_{p,\sigma} &\stackrel{(2.5.5)}{=} N(p) \mathcal{P} U(L(p)) \Psi_{k,\sigma} \\ &\stackrel{(2.5.44)}{=} \sqrt{\frac{\kappa}{p^0}} \mathcal{P} U \left(R(\hat{\mathbf{p}}) B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\ &= \sqrt{\frac{\kappa}{p^0}} U \left(\mathcal{P} R(\hat{\mathbf{p}}) B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\ &= \sqrt{\frac{\kappa}{p^0}} U \left(R(\hat{\mathbf{p}}) \mathcal{P} B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\ &= \sqrt{\frac{\kappa}{p^0}} U \left(R(\hat{\mathbf{p}}) R_2 R_2^{-1} \mathcal{P} B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\ &= \sqrt{\frac{\kappa}{p^0}} U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) R_2^{-1} \mathcal{P} \right) \Psi_{k,\sigma} \\ &= \sqrt{\frac{\kappa}{p^0}} U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) U(R_2^{-1}) \mathcal{P} \Psi_{k,\sigma} \\ &\stackrel{(2.6.20)}{=} \sqrt{\frac{\kappa}{p^0}} \eta_\sigma U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,-\sigma} \end{aligned}$$

10. “But a rotation of $\pm 180^\circ$ around the three-axis reverses the sign of J_2, \dots ” (P.79)

Analogously to II F 6 ($2 \leftrightarrow 3$).

11. Eq. (2.6.22) (P.79)

First note that

$$B \left(\frac{|\mathbf{p}|}{\kappa} \right)$$

and rotations around the three-axis commute (see Eq. (2.5.45)).

$$\begin{aligned}
P\Psi_{p,\sigma} &\stackrel{II F 9}{=} \sqrt{\frac{\kappa}{p^0}} \eta_\sigma U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,-\sigma} \\
&\stackrel{Eq. (2.6.21)}{=} \sqrt{\frac{\kappa}{p^0}} \eta_\sigma U (R(-\hat{\mathbf{p}})) \\
&\quad \cdot \exp(\pm i\pi J_3) U \left(B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,-\sigma} \\
&= \sqrt{\frac{\kappa}{p^0}} \eta_\sigma U (R(-\hat{\mathbf{p}})) \\
&\quad \cdot U \left(B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \exp(\pm i\pi J_3) \Psi_{k,-\sigma} \\
&= \eta_\sigma \exp(\pm i\pi(-\sigma)) \sqrt{\frac{\kappa}{p^0}} U (L(\mathcal{P}p)) \Psi_{k,-\sigma} \\
&= \eta_\sigma \exp(\mp i\pi\sigma) \Psi_{\mathcal{P}p,-\sigma}
\end{aligned}$$

12. “This peculiar change of sign in the operation of parity for mass-less particles of half-integer spin is due to the convention adopted in Eq. (2.5.47) for the rotation used to define mass-less particles states of arbitrary momentum.” (P.79)

This is based on the fact that the choice of $R(\hat{\mathbf{p}})$, which transforms the three axis into the unit vector $\hat{\mathbf{p}}$, is *not* unique. As mentioned in the text right after Eq. (2.5.47), one could always add an initial rotation around the three axis. This is also why the factor of

$$\exp(\pm i\pi J_3)$$

shows up in Eq. (2.6.21), despite $R(\hat{\mathbf{p}}) R_2$ and $R(-\hat{\mathbf{p}})$ both transforming the three axis into the unit vector $-\hat{\mathbf{p}}$.

13. “ \mathbb{T} yields a state which has values ...” (P.79)

What is meant by this is:

$$\begin{aligned}
P^i \mathbb{T} \Psi_{k,\sigma} &\stackrel{(2.6.12)}{=} \mathbb{T} (-P^i) \Psi_{k,\sigma} = (-\delta_3^i \kappa) \mathbb{T} \Psi_{k,\sigma} \\
H \mathbb{T} \Psi_{k,\sigma} &\stackrel{(2.6.13)}{=} \mathbb{T} H \Psi_{k,\sigma} = \kappa \mathbb{T} \Psi_{k,\sigma} \\
J_3 \mathbb{T} \Psi_{k,\sigma} &\stackrel{(2.6.10)}{=} \mathbb{T} (-J_3) \Psi_{k,\sigma} = -\sigma \mathbb{T} \Psi_{k,\sigma}
\end{aligned}$$

14. Eq. (2.6.23) (P.80)

This is completely analog to Eq. (2.6.20) when using IIF 13 (see IIF 7). Further we get

$$R_2^{-1} \mathbb{T} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}.$$

15. Eq. (2.6.24) (P.80)

$$\begin{aligned}
\mathbb{T} \Psi_{p,\sigma} &\stackrel{(2.5.5)}{=} N(p) \mathbb{T} U (L(p)) \Psi_{k,\sigma} \\
&\stackrel{(2.5.44)}{=} \sqrt{\frac{\kappa}{p^0}} \mathbb{T} U \left(R(\hat{\mathbf{p}}) B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\
&= \sqrt{\frac{\kappa}{p^0}} U \left(\mathbb{T} R(\hat{\mathbf{p}}) B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\
&= \sqrt{\frac{\kappa}{p^0}} U \left(R(\hat{\mathbf{p}}) \mathbb{T} B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\
&= \sqrt{\frac{\kappa}{p^0}} U \left(R(\hat{\mathbf{p}}) R_2 R_2^{-1} \mathbb{T} B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\
&= \sqrt{\frac{\kappa}{p^0}} U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) R_2^{-1} \mathbb{T} \right) \Psi_{k,\sigma} \\
&= \sqrt{\frac{\kappa}{p^0}} U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) U (R_2^{-1}) \mathbb{T} \Psi_{k,\sigma} \\
&\stackrel{(2.6.23)}{=} \sqrt{\frac{\kappa}{p^0}} \zeta_\sigma U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma}
\end{aligned}$$

16. Eq. (2.6.25) (P.80)

First note that

$$B \left(\frac{|\mathbf{p}|}{\kappa} \right)$$

and rotations around the three-axis commute (see Eq. (2.5.45)).

$$\begin{aligned}
\mathbb{T} \Psi_{p,\sigma} &\stackrel{II F 15}{=} \sqrt{\frac{\kappa}{p^0}} \zeta_\sigma U \left(R(\hat{\mathbf{p}}) R_2 B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\
&\stackrel{Eq. (2.6.21)}{=} \sqrt{\frac{\kappa}{p^0}} \zeta_\sigma U (R(-\hat{\mathbf{p}})) \\
&\quad \cdot \exp(\pm i\pi J_3) U \left(B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \Psi_{k,\sigma} \\
&= \sqrt{\frac{\kappa}{p^0}} \zeta_\sigma U (R(-\hat{\mathbf{p}})) \\
&\quad \cdot U \left(B \left(\frac{|\mathbf{p}|}{\kappa} \right) \right) \exp(\pm i\pi J_3) \Psi_{k,\sigma} \\
&= \zeta_\sigma \exp(\pm i\pi\sigma) \sqrt{\frac{\kappa}{p^0}} U (L(\mathcal{P}p)) \Psi_{k,\sigma} \\
&= \zeta_\sigma \exp(\pm i\pi\sigma) \Psi_{\mathcal{P}p,\sigma}
\end{aligned}$$

17. “...the total angular momentum j of any state of this system would have to be a half-integer...” (P.81)

This is true as all spins and helicities except for one half integer spin/helicity would couple to a an integer

angular momentum. And this would then couple with the remaining half integer spin/helicity to a half-integer total angular momentum j .

G. Projective Representations

H. The Symmetry Representation Theorem

1. “But $\langle \Psi'_k | \Psi'_k \rangle$ is automatically **real and positive**” (P.91)

This follows immediately from Eq. (2.1.1).

2. “From Eq. (2.A.1) we have $|c_{kk}| = |c_{k1}| = \frac{1}{\sqrt{2}}$ and for $l \neq k$ and $l \neq 1$: $c_{kl} = 0$ ” (P.91)

$$|c_{kl'}|^2 \stackrel{(2.A.3)}{=} \left| \sum_l c_{kl}^* \langle \Psi'_l | \Psi'_{l'} \rangle \right|^2 = |\langle \Upsilon'_k | \Psi'_{l'} \rangle|^2$$

$$\stackrel{(2.A.1)}{=} |\langle \Upsilon_k | \Psi_{l'} \rangle|^2 = \begin{cases} \frac{1}{2} & \text{for } l' = 1, k \\ 0 & \text{for } l' \neq 1, k \end{cases}$$

3. Eq. (2.A.10) (P.92)

$$|C_k|^2 + |C_1|^2 + 2 \operatorname{Re}(C_k C_1^*) = |C_k + C_1|^2$$

$$\stackrel{(2.A.9)}{=} |C'_k + C'_1|^2 = |C'_k|^2 + |C'_1|^2 + 2 \operatorname{Re}(C'_k C'^*_1)$$

$$\stackrel{\text{Eq. (2.A.8)}}{\Rightarrow} \operatorname{Re}(C_k C_1^*) = \operatorname{Re}(C'_k C'^*_1)$$

$$\stackrel{\text{Eq. (2.A.8)}}{\Rightarrow} \operatorname{Re}\left(\frac{C_k}{C_1}\right) = \operatorname{Re}\left(\frac{C'_k}{C'_1}\right)$$

4. Eq. (2.A.11) (P.92)

$$\left\{ \operatorname{Re}\left(\frac{C_k}{C_1}\right) \right\}^2 + \left\{ \operatorname{Im}\left(\frac{C_k}{C_1}\right) \right\}^2 = \left| \frac{C_k}{C_1} \right|^2$$

$$\stackrel{\text{Eq. (2.A.8)}}{=} \left| \frac{C'_k}{C'_1} \right|^2 = \left\{ \operatorname{Re}\left(\frac{C'_k}{C'_1}\right) \right\}^2 + \left\{ \operatorname{Im}\left(\frac{C'_k}{C'_1}\right) \right\}^2$$

$$\stackrel{\text{Eq. (2.A.10)}}{\Rightarrow} \operatorname{Im}\left(\frac{C_k}{C_1}\right) = \pm \operatorname{Im}\left(\frac{C'_k}{C'_1}\right)$$

5. “This is only possible if $\operatorname{Re}\left(\frac{C_k}{C_1} \frac{C_l^*}{C_1^*}\right) = \operatorname{Re}\left(\frac{C_k}{C_1} \frac{C_l}{C_1}\right)$ or, in other words, if $\operatorname{Im}\left(\frac{C_k}{C_1}\right) \operatorname{Im}\left(\frac{C_l}{C_1}\right) = 0$ ” (P.93)

Define

$$a := \frac{C_k}{C_1}$$

$$b := \frac{C_l}{C_1}$$

With this we have

$$|1 + a + b^*|^2 = |1 + a + b|^2 \quad (1)$$

$$\Leftrightarrow 1 + a^* + b + a + |a|^2 + ab + b^* + b^* a^* + |b|^2 \quad (2)$$

$$= 1 + a^* + b^* + a + |a|^2 + ab^* + b + ba^* + |b|^2 \quad (3)$$

$$\Leftrightarrow ab + b^* a^* = ab^* + ba^* \quad (4)$$

And further rewriting yields

$$\operatorname{Re}\left(\frac{C_k}{C_1} \frac{C_l}{C_1}\right) = \operatorname{Re}(ab) \stackrel{4}{=} \operatorname{Re}(ab^*) = \operatorname{Re}\left(\frac{C_k}{C_1} \frac{C_l^*}{C_1^*}\right)$$

$$\operatorname{Im}\left(\frac{C_k}{C_1}\right) \operatorname{Im}\left(\frac{C_l}{C_1}\right) = \operatorname{Im}(a) \operatorname{Im}(b)$$

$$= -\frac{1}{4}(ab - ab^* - a^* b + a^* b^*) \stackrel{4}{=} 0$$

6. “Then the invariance of transition probabilities requires

$$\text{that } \left| \sum_k B_k^* A_k \right|^2 = \left| \sum_k B_k A_k \right|^2 \text{” (P.93)}$$

$$\left| \sum_k B_k^* A_k \right|^2 \stackrel{\text{Eq. (2.A.2)}}{=} \left| \sum_{kl} B_k^* A_l \langle \Psi_k | \Psi_l \rangle \right|^2$$

$$= \left| \left\langle \sum_k B_k \Psi_k \left| \sum_l A_l \Psi_l \right. \right\rangle \right|^2$$

$$\stackrel{\text{Eq. (2.A.1)}}{=} \left| \left\langle U \left(\sum_k B_k \Psi_k \right) \left| U \left(\sum_l A_l \Psi_l \right) \right. \right\rangle \right|^2$$

$$= \left| \left\langle \sum_k B_k^* U \Psi_k \left| \sum_l A_l U \Psi_l \right. \right\rangle \right|^2$$

$$= \left| \sum_{kl} B_k A_l \langle U \Psi_k | U \Psi_l \rangle \right|^2$$

$$\stackrel{\text{Eq. (2.A.3)}}{=} \left| \sum_k B_k A_k \right|^2$$

7. Eq. (2.A.16)(P.94)

8. “However, for any pair of such state-vectors, with neither A_k nor B_k **all of the same phase**” (P.94)

If they were all of the same phase then

$$\forall k, l : \text{Im}\{A_k^* A_l\} = 0$$

or

$$\forall k, l : \text{Im}\{B_k^* B_l\} = 0$$

See Footnote j, for why this is relevant.

9. “We have thus shown that for a given symmetry transformation T either all state-vectors satisfy Eq. (2.A.14) or else they all satisfy Eq. (2.A.15)” (P.94)

Preceding this a contradiction was derived from the assumption that Eq. (2.A.14) applies for a state-vector

$$\sum_k A_k \Psi_k$$

while Eq. (2.A.15) applies for a state vector

$$\sum_k B_k \Psi_k.$$

Such that the statement is obvious.

I. Group Operators and Homotopy Classes

J. Inversions and Degenerate Multiplets

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$$\begin{aligned} & \sum_{kl} \text{Im}(B_k^* B_l) \text{Im}(A_k^* A_l) \\ &= \text{Im} \left(\sum_{kl} \text{Im}(B_k^* B_l) A_k^* A_l \right) \\ &= \frac{1}{2i} \left[\sum_{kl} \text{Im}(B_k^* B_l) A_k^* A_l - \text{c.c.} \right] \\ &= \frac{1}{2i} \left[\sum_{kl} \frac{1}{2i} (B_k^* B_l - B_k B_l^*) A_k^* A_l - \text{c.c.} \right] \\ &= \frac{1}{2i} \left[\frac{1}{2i} \left(\sum_{kl} B_k^* B_l A_k^* A_l - \sum_{kl} B_k B_l^* A_k^* A_l \right) - \text{c.c.} \right] \\ &= \frac{1}{2i} \left[\frac{1}{2i} \left(\left| \sum_k B_k A_k \right|^2 - \left| \sum_k B_k^* A_k \right|^2 \right) - \text{c.c.} \right] \\ &= \frac{1}{2i} \left[\frac{1}{i} \left(\left| \sum_k B_k A_k \right|^2 - \left| \sum_k B_k^* A_k \right|^2 \right) \right] \\ &\stackrel{II}{=} 0 \end{aligned}$$