

Pset 2 solutions

$$\boxed{\text{I}} \quad H = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - \alpha \delta(x), \quad \alpha > 0$$

(a.1)

$$\begin{aligned} H \psi(x) &= E \psi(x) \\ \int_{-\varepsilon}^{+\varepsilon} \left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - \alpha \delta(x) \right] \psi(x) dx &= E \int_{-\varepsilon}^{\varepsilon} \psi(x) dx \\ -\frac{\hbar^2}{2m} \frac{d\psi(x)}{dx} \Big|_{-\varepsilon}^{+\varepsilon} - \alpha \psi(0) &= E \int_{-\varepsilon}^{\varepsilon} \psi(x) dx \\ \frac{d\psi(x)}{dx} \Big|_{-\varepsilon}^{+\varepsilon} &= \frac{2m}{\hbar^2} \left[-\alpha \psi(0) - E \int_{-\varepsilon}^{\varepsilon} \psi(x) dx \right] \end{aligned}$$

$$\text{(a.2)} \quad \varepsilon \rightarrow 0 \Rightarrow \int_{-\varepsilon}^{\varepsilon} \psi(x) dx = \int_{-\varepsilon}^{\varepsilon} \psi(0) dx = \psi(0) \cdot 2\varepsilon$$

per hint, define discontinuity as

$$\begin{aligned} \Delta &= \lim_{\varepsilon \rightarrow 0} \left(\frac{d\psi(x)}{dx} \Big|_{\varepsilon} - \frac{d\psi(x)}{dx} \Big|_{-\varepsilon} \right) = \\ &= \lim_{\varepsilon \rightarrow 0} \left\{ \frac{2m}{\hbar^2} \left[-\alpha \psi(0) - 2\varepsilon E \psi(0) \right] \right\} = \end{aligned}$$

$$\boxed{\Delta = -\frac{2m\alpha}{\hbar^2} \psi(0)}$$

(b.1)

$E < 0$ - bound state

$$\begin{aligned} x < 0 \quad \psi(x) &= A_1 e^{px} + A_1' e^{-px} \\ x > 0 \quad \psi(x) &= A_2 e^{ex} + A_2' e^{-px} \end{aligned}$$

$$H \psi(x) = E \psi(x)$$

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi(x)}{dx^2} - \alpha \delta(x) \psi(x) = E \psi(x)$$

$$\left. \begin{array}{l} x < 0 \\ x > 0 \end{array} \right\} \frac{d^2 \psi(x)}{dx^2} = \begin{array}{l} A_1 \rho^2 e^{\rho x} + A_1' \rho^2 e^{-\rho x} \\ A_2 \rho^2 e^{\rho x} + A_2' \rho^2 e^{-\rho x} \end{array} \Rightarrow \frac{d^2 \psi(x)}{dx^2} = \rho^2 \psi(x) \text{ for all } x \neq 0$$

$$\left[-\frac{\hbar^2}{2m} \rho^2 - \alpha \delta(x) \right] \psi(x) = E \psi(x)$$

Evaluate at any x :

$$-\frac{\hbar^2}{2m} \rho^2 = E \Rightarrow \boxed{\rho = \sqrt{\frac{2mE}{\hbar^2}}}$$

(b.2) Find E :

- $\psi(x)$ is Square-integrable $\Rightarrow \psi(x) \rightarrow 0$ as $x \rightarrow \pm \infty$
 $\Rightarrow A_1' = A_2 = 0$

- $\psi(x)$ is continuous every where

$$\psi(0)|_{x<0} = \psi(0)|_{x>0} \Rightarrow A_1 = A_2' = A$$

- Normalize: $\int_{-\infty}^{+\infty} |\psi(x)|^2 dx = 1$ $\psi(x) = A \cdot \rho^2 e^{-\rho|x|}$ for $x \neq 0$

$$|A|^2 \cdot \rho^2 \int_{-\infty}^{+\infty} e^{-2 \cdot \rho |x|} dx = 1$$

$$|A|^2 \cdot \rho^2 \cdot 2 \int_0^{\infty} e^{-2 \rho x} dx = 1$$

$$u = -2 \rho x \quad du = -2 \rho dx$$

$$|A|^2 \cdot \int_{-\infty}^{\infty} e^{-\beta|x|} dx = 1$$

$$-|A|^2 \cdot \int_{-\infty}^{\infty} e^{-\beta|x|} dx = 1$$

$$+|A|^2 \cdot \int_{-\infty}^{\infty} e^{-\beta|x|} dx = 1$$

$A = \sqrt{\beta}$ choosing $A > 0$ solution

So $\boxed{\psi(x) = \sqrt{\beta} e^{-\beta|x|}}$

Now from part (a): $\lim_{\epsilon \rightarrow 0} \left(\frac{d\psi(x)}{dx} \right) \Big|_{-\epsilon}^{\epsilon} = -\frac{2m\alpha}{\hbar^2} \psi(0)$

$-2\beta \sqrt{\beta} \cdot 1 = \frac{2m\alpha}{\hbar^2} \sqrt{\beta}$

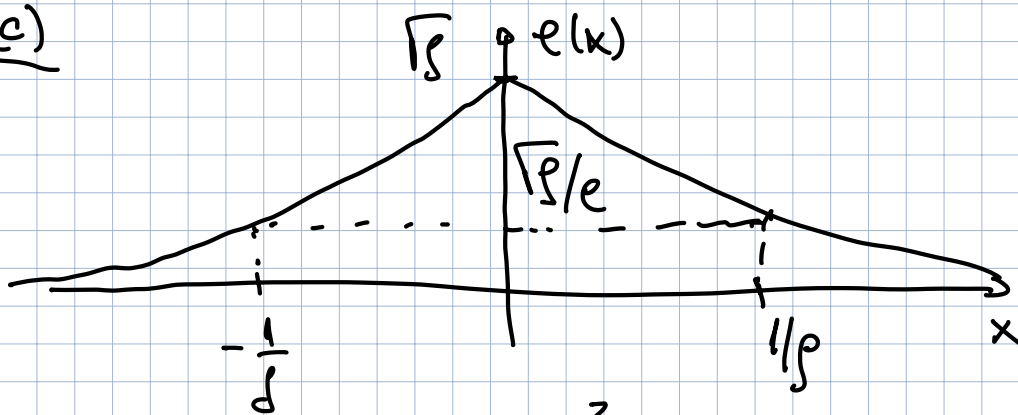
$\beta = \frac{m\alpha}{\hbar^2}$

$\beta = \sqrt{\frac{2mE}{\hbar^2}} \Rightarrow -\frac{m\alpha^2}{\hbar^3} = \frac{2mE}{\hbar^2}$

$\boxed{E = -\frac{m\alpha^2}{2\hbar^2}} < 0$ solution

$\psi(x) = \sqrt{\beta} e^{-\beta|x|}$ where $\beta = \frac{m\alpha}{\hbar^2}$

(c)



$$\Delta x = \frac{2}{p} = \frac{2\hbar^2}{m\omega}$$

(d) $\varphi(x) = \sqrt{p} e^{-p|x|}$

$$\bar{\varphi}(p) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} e^{-\frac{ipx}{\hbar}} \sqrt{p} e^{-p|x|} dx =$$

$$= \frac{\sqrt{p}}{\sqrt{2\pi\hbar}} \left(\int_{-\infty}^0 e^{-\frac{ipx}{\hbar}} e^{+px} dx + \int_0^{\infty} e^{-\frac{ipx}{\hbar}} e^{-px} dx \right) =$$

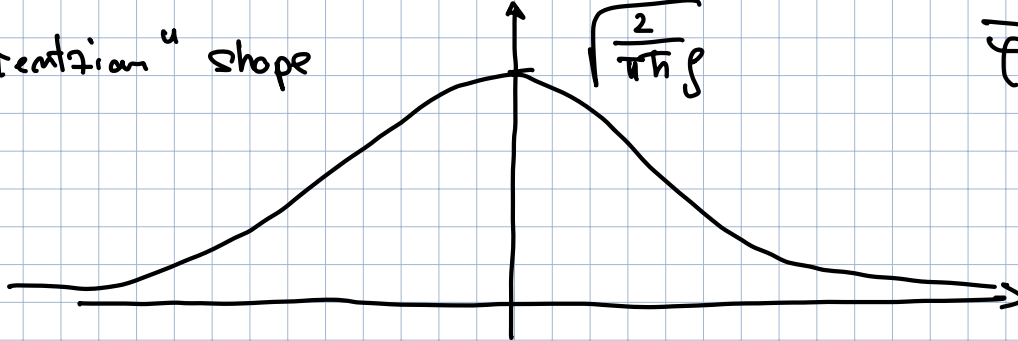
$$= \sqrt{\frac{p}{2\pi\hbar}} \left(\left. \frac{1}{p - \frac{ip}{\hbar}} e^{\left(p - \frac{ip}{\hbar}\right)x} \right|_{-\infty}^0 - \left. \frac{1}{p + \frac{ip}{\hbar}} e^{-\left(p + \frac{ip}{\hbar}\right)x} \right|_0^{\infty} \right) =$$

$$= \sqrt{\frac{p}{2\pi\hbar}} \left(\frac{1}{p - \frac{ip}{\hbar}} + \frac{1}{p + \frac{ip}{\hbar}} \right) =$$

$$= \sqrt{\frac{p}{2\pi\hbar}} \frac{p + \frac{ip}{\hbar} + p - \frac{ip}{\hbar}}{p^2 + \frac{p^2}{\hbar^2}} =$$

$$\bar{\varphi}(p) = \sqrt{\frac{2p}{\pi\hbar}} \cdot \frac{p}{p^2 + \frac{p^2}{\hbar^2}}$$

"Lorentzian" shape



$$\begin{aligned}\overline{\varphi}(0) &= \sqrt{\frac{2g}{\pi \hbar}} \cdot \frac{1}{g} = \\ &= \sqrt{\frac{2}{\pi \hbar g}}\end{aligned}$$

At half maximum: $g^2 = \frac{P_{1/2}}{\hbar^2} \Rightarrow \boxed{P_{1/2} = \hbar g}$

Full width at half maximum: $\Delta p = 2\hbar g$

$$\Delta x \cdot \Delta p = \frac{2}{g} \cdot 2\hbar g = 4\hbar \sim \text{few } \hbar \text{ depending on how we define width}$$

II

$$U(m, n) = |\varphi_m\rangle \langle \varphi_n|$$

$$H |\varphi_m\rangle = E_m |\varphi_m\rangle$$

a.
$$U^\dagger(m, n) = (|\varphi_m\rangle \langle \varphi_n|)^\dagger =$$
$$= |\varphi_n\rangle \langle \varphi_m|$$

b.
$$[H, U(m, n)] = H U(m, n) - U(m, n) H =$$
$$= H |\varphi_m\rangle \langle \varphi_n| - |\varphi_m\rangle \langle \varphi_n| H =$$

$$H |\varphi_m\rangle = E_m |\varphi_m\rangle \quad |^\dagger$$

$$\langle \varphi_n| H^\dagger = E_m^* \langle \varphi_n|$$

$$H \text{ is Hermitian} \Rightarrow H^\dagger = H \quad E_m^* = E$$

$$\langle \varphi_n| H = E_m \langle \varphi_n|$$

$$[H, U(m, n)] = E_m |\varphi_m\rangle \langle \varphi_n| - E_n |\varphi_m\rangle \langle \varphi_n| =$$
$$= (E_m - E_n) U(m, n)$$

c.
$$U(m, n) U^\dagger(p, q) = |\varphi_m\rangle \langle \varphi_n| (|\varphi_p\rangle \langle \varphi_q|)^\dagger =$$
$$= |\varphi_m\rangle \underbrace{\langle \varphi_n| \varphi_q\rangle}_{\delta_{n, q}} \langle \varphi_p| =$$

$$= \delta_{n, q} |\varphi_m\rangle \langle \varphi_p| = \delta_{n, q} U(m, p)$$

d.
$$\text{Tr} \{ U(m, n) \} \equiv \sum_p \langle \varphi_p | \varphi_m \rangle \langle \varphi_n | \varphi_p \rangle =$$
$$= \sum_p \delta_{p, m} \delta_{n, p} =$$
$$= \delta_{m, n}$$

e. $A_{mn} = \langle \varphi_m | A | \varphi_n \rangle$

$$\begin{aligned} \sum_{m,n} A_{mn} U(m,n) &= \sum_{m,n} \underbrace{\langle \varphi_m | A | \varphi_n \rangle}_{\text{scalar so can move around}} | \varphi_m \rangle \langle \varphi_n | \\ &= \sum_{m,n} | \varphi_m \rangle \langle \varphi_m | A | \varphi_n \rangle \langle \varphi_n | = \\ &\quad \swarrow \text{use } \sum_m | \varphi_m \rangle \langle \varphi_m | = \mathbb{I} \\ &= \mathbb{I} \cdot A \cdot \mathbb{I} = \text{closure relation} \\ &= A \end{aligned}$$

f. $\text{Tr} \{ A U^\dagger(p,q) \} = \sum_m \langle \varphi_m | A | \varphi_q \rangle \langle \varphi_p | \varphi_m \rangle =$

$$\begin{aligned} &= \sum_m A_{mq} \cdot \delta_{p,m} = \\ &= A_{pq} \end{aligned}$$