

OPTI 570 LECTURE Th OCT 30

2-level systems

Case 1: $H_0 = E_1 |1\rangle\langle 1| + E_2 |2\rangle\langle 2|$

$$H_{S1,21} = \begin{pmatrix} E_1 & 0 \\ 0 & E_2 \end{pmatrix}$$

$$U_{S1,23} = \begin{pmatrix} e^{-iE_1 t/\hbar} & 0 \\ 0 & e^{-iE_2 t/\hbar} \end{pmatrix}$$

$$|\Psi(0)\rangle = \cos \frac{\theta}{2} |1\rangle + \sin \frac{\theta}{2} e^{i\phi} |2\rangle$$

$$\begin{aligned} |\Psi(t)\rangle &= \cos \frac{\theta}{2} e^{-iE_1 t/\hbar} |1\rangle + \sin \frac{\theta}{2} e^{i\phi} e^{-iE_2 t/\hbar} |2\rangle \\ &= \text{pl. p.} \cdot \left[\cos \frac{\theta}{2} |1\rangle + \sin \frac{\theta}{2} e^{i\phi} e^{-i(E_2 - E_1)t/\hbar} |2\rangle \right] \end{aligned}$$

Bohr frequency

$$\begin{aligned} |\Psi(0)\rangle &= |1\rangle \Rightarrow |\Psi(t)\rangle = |1\rangle \\ &= |2\rangle \Rightarrow |2\rangle \end{aligned}$$

Case 2: Generic time-independent Hamiltonian

$$H = E_1 |1\rangle\langle 1| + E_2 |2\rangle\langle 2|$$

$$W = W_{21} |2\rangle\langle 1| + W_{12}^* \underbrace{|1\rangle\langle 2|}_{\text{coupling}}$$

$$W = \frac{1}{2}\hbar \omega_0 |2\rangle\langle 1| + \frac{1}{2}\hbar \omega_0^* |1\rangle\langle 2|$$

$$\Omega_0 = |\Omega_0| e^{i\phi} \quad |\Omega_0| - \text{strength of } W\text{-coupling.}$$

$$H_{S1,21} = \begin{pmatrix} E_1 & \frac{1}{2}\hbar |\Omega_0| e^{-i\phi} \\ \frac{1}{2}\hbar |\Omega_0| e^{i\phi} & E_2 \end{pmatrix}$$

Any time-independent Hamiltonian for 2-level

#1: $\Delta = \frac{E_1 - E_2}{\hbar}$ "detuning"

#2: $E_m = \frac{1}{2} (E_1 + E_2)$ "mean energy"

Δ and E_m are real

$$E_1 = E_m + \frac{1}{2}\hbar\Delta$$

$$E_2 = E_m - \frac{1}{2}\hbar\Delta$$

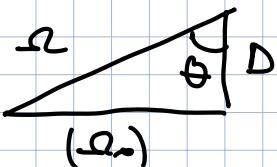
$$H \begin{pmatrix} \zeta_1 \\ \zeta_2 \end{pmatrix} = \begin{pmatrix} E_m + \frac{1}{2}\hbar\Delta & \frac{1}{2}\hbar|\Omega_0|e^{-i\varphi} \\ \frac{1}{2}\hbar|\Omega_0|e^{i\varphi} & E_m - \frac{1}{2}\hbar\Delta \end{pmatrix} =$$

$$= E_m \mathbb{1} + \frac{\hbar}{2} \begin{pmatrix} 0 & |\Omega_0|e^{-i\varphi} \\ |\Omega_0|e^{i\varphi} & -\Delta \end{pmatrix}$$

no phys. conseq.

#3. $\tan \theta = \frac{|\Omega_0|}{\Delta} \quad 0 \leq \theta \leq \frac{\pi}{2}$

#4. $\Omega = \sqrt{|\Omega_0|^2 + \Delta^2} \quad - \text{real, pos#}$



$$\Omega = \frac{2}{\pi} \sqrt{(E_1 - E_2)^2 + |W_{21}|^2}$$

$$H \begin{pmatrix} \zeta_1 \\ \zeta_2 \end{pmatrix} = E_m \mathbb{1} + \frac{\hbar \Omega}{2} \begin{pmatrix} \cos \theta & \sin \theta e^{-i\varphi} \\ \sin \theta e^{i\varphi} & -\cos \theta \end{pmatrix}$$

$$H = E_m \mathbb{1} + \frac{\hbar \Omega}{2} \vec{\sigma}_v$$

$$G_u = \vec{G} \cdot \hat{u}$$

$$\hat{u} = \begin{pmatrix} \sin \theta & \cos \varphi \\ \sin \theta & \sin \varphi \\ \cos \theta \end{pmatrix} = \begin{pmatrix} \frac{|\Omega_0|}{\Omega} \operatorname{Re}[e^{i\varphi}] \\ \frac{|\Omega_0|}{\Omega} \operatorname{Im}[e^{i\varphi}] \\ \frac{\Delta}{\Omega} \end{pmatrix} =$$

$$\hat{u} = \frac{1}{\Omega} \begin{pmatrix} \operatorname{Re}[\Omega_0] \\ \operatorname{Im}[\Omega_0] \\ \Delta \end{pmatrix} = \frac{1}{\hbar \Omega} \begin{pmatrix} 2 \operatorname{Re}\{W_{21}\} \\ 2 \operatorname{Im}\{W_{21}\} \\ (E_1 - E_2) / \hbar \end{pmatrix}$$

Energy eigenvalues and eigenstates

$$H |E_{\pm}\rangle = E_{\pm} |E_{\pm}\rangle$$

$$E_+ = E_m + \frac{\hbar\omega}{2}$$

$$\psi |E_+\rangle = \cos \frac{\theta}{2} |1\rangle + \sin \frac{\theta}{2} e^{i\phi} |2\rangle$$

$$E_- = E_m - \frac{\hbar\omega}{2}$$

$$\psi |E_-\rangle = \cos \frac{\theta}{2} |1\rangle - \sin \frac{\theta}{2} e^{-i\phi} |2\rangle$$

Time evolution:

$$U \{E_+, E_-\} = E_m \mathbb{1} + \frac{\hbar\omega}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$U \{E_-, E_+\} = e^{-i E_m t / \hbar} \cdot \begin{pmatrix} e^{-i \frac{\omega t}{2}} & 0 \\ 0 & e^{i \frac{\omega t}{2}} \end{pmatrix}$$

$$U \{S_{1,2}\} = \alpha^+ U M$$

$$U(t) = \underbrace{e^{-i E_m t / \hbar}}_{\dots} \underbrace{e^{-i \frac{\omega t}{2}} \hat{G}_u}_{\dots} =$$

$$= \cos \left(\frac{\omega t}{2} \right) \mathbb{1} + i \hat{G}_u \sin \left(\frac{\omega t}{2} \right) =$$

$$U \{S_{1,2}\} = \begin{pmatrix} e^{\frac{i \omega t}{2}} \sin^2 \frac{\theta}{2} + e^{-\frac{i \omega t}{2}} \cos^2 \frac{\theta}{2} & -i \sin \left(\frac{\omega t}{2} \right) \sin \theta e^{-i\phi} \\ -i \sin \left(\frac{\omega t}{2} \right) \sin \theta e^{i\phi} & e^{\frac{i \omega t}{2}} \cos^2 \frac{\theta}{2} + e^{-\frac{i \omega t}{2}} \sin^2 \frac{\theta}{2} \end{pmatrix}$$

Define: $P_{j \rightarrow k}(t)$ - transition probability

$$P_{1 \rightarrow 1}(0) = 1$$

$$P_{1 \rightarrow 2}(0) = 0$$

$$P_{2 \rightarrow 2}(0) = 1$$

$$P_{2 \rightarrow 1}(0) = 0$$

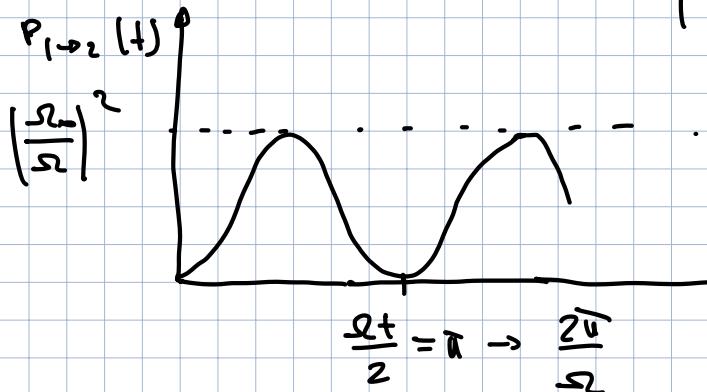
$$P_{1 \rightarrow 2}(t)$$

$$P_{1 \rightarrow 1}(t) + P_{1 \rightarrow 2}(t) = 1$$

$$|\Psi(0)\rangle = |1\rangle, \text{ what is } P_{1 \rightarrow 2}(t) = ?$$

$$|\Psi(t)\rangle = \left(\cos^2 \frac{\theta}{2} e^{-i \frac{\omega t}{2}} + \sin^2 \frac{\theta}{2} e^{i \frac{\omega t}{2}} \right) |1\rangle - i \sin \theta e^{i\phi} \sin \left(\frac{\omega t}{2} \right) |2\rangle$$

$$P_{1 \rightarrow 2}(t) = |\langle \Psi | z \rangle|^2 = \sin^2 \theta \sin^2 \left(\frac{\Omega t}{2} \right) = \left| \frac{\Omega_0}{\Omega} \right|^2 \sin^2 \left(\frac{\Omega t}{2} \right)$$



"Rabi Oscillations"

Ω_0 : "Bare Rabi Frequency"

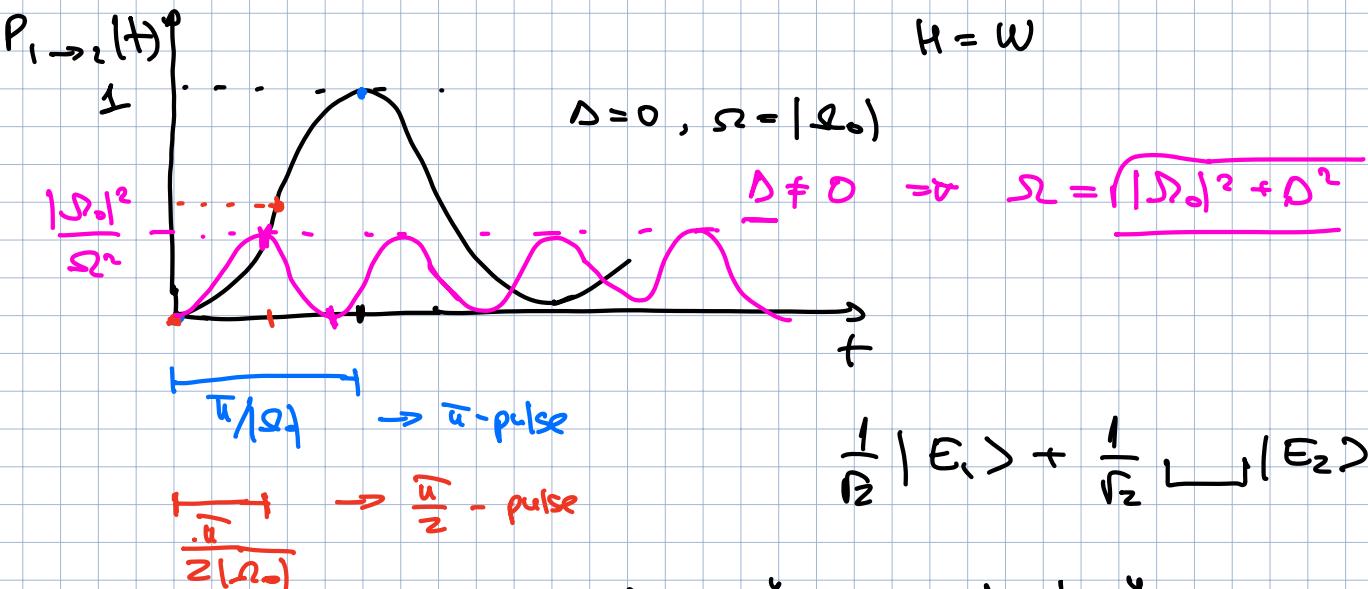
Ω : "Rabi Frequency"

"Generalized Rabi Frequency"

units:

$$\frac{\text{rad}}{\text{s}} \quad \Omega : \text{"Detuning"} \quad \Omega = \sqrt{\Omega_0^2 + D^2}$$

Rabi oscillations $\Delta = 0 \rightarrow \tan \theta = \infty, \theta = \frac{\pi}{2}$



$\Delta = 0 - \text{"special angle phase"}$

O|-resonance:

2-level dynamics on the Bloch sphere

$$|\Psi(+)\rangle = a_1(+)|1\rangle + a_2(+)|2\rangle$$

$$|\Psi(+)\rangle_{S_{1,21}} = \begin{pmatrix} a_1(+) \\ a_2(+) \end{pmatrix}$$

Bloch Vector:

$$\langle \hat{\sigma}_z \rangle = \langle \Psi | \hat{\sigma}_z | \Psi \rangle = (a_1^* a_1) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = |a_1|^2 - |a_2|^2$$

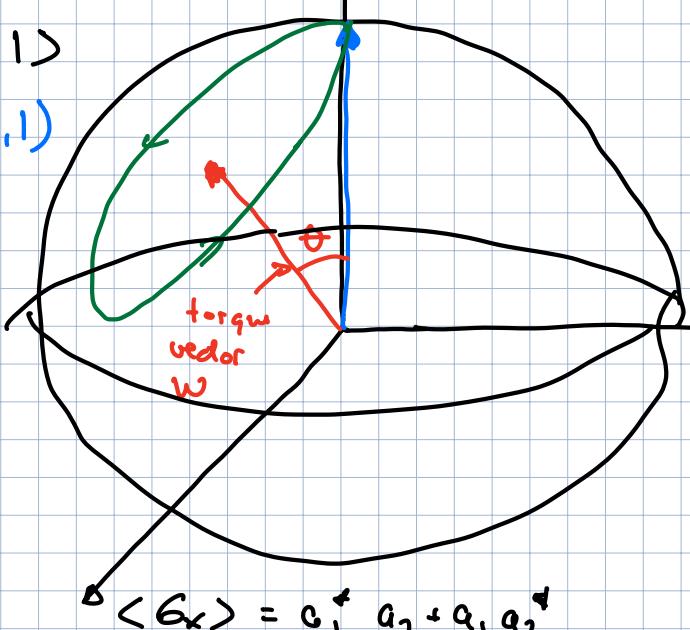
$$\langle \hat{\sigma}_x \rangle = a_1^* a_2 + a_1 a_2^*$$

$$\langle \hat{\sigma}_y \rangle = -i a_1^* a_2 + i a_1 a_2^*$$

$$\langle \hat{\sigma}_z \rangle = |a_1|^2 - |a_2|^2 = P(|+\rangle_z) - P(|-\rangle_z)$$

$$|\Psi(0)\rangle = |1\rangle$$

$$\langle \hat{\sigma} \rangle(0) = (0, 0, 1)$$



$$\begin{aligned} \langle \hat{\sigma}_y \rangle &= -i a_1^* a_2 + i a_1 a_2^* \\ &= P(|+\rangle_y) - P(|-\rangle_y) \end{aligned}$$

$$\begin{aligned} \langle \hat{\sigma}_x \rangle &= a_1^* a_2 + a_1 a_2^* \\ &= P(|+\rangle_x) - P(|-\rangle_x) \end{aligned}$$

$$\langle \hat{\sigma}_z \rangle = |a_1|^2 - |a_2|^2 = P(|1\rangle) - P(|2\rangle)$$

$$P(1) = \frac{1}{2} (1 + \langle \hat{\sigma}_z \rangle)$$

$$P(2) = \frac{1}{2} (1 - \langle \hat{\sigma}_z \rangle)$$

Bloch sphere - tells you the exact state, not just probabilities

Example: Spin $\frac{1}{2}$ particle in a magnetic field $\gamma < 0$

$$\vec{B} = (0, 0, B_z)$$

$$\hbar = \frac{1}{2} \hbar \omega_z \hat{\sigma}_z, \quad \omega_z = -\gamma B_z > 0$$

$$\text{at } t=0, \quad \vec{B} = (0, B_y, B_z) \quad t>0$$

$$\omega = \frac{1}{2} \hbar \omega_z \hat{\sigma}_y, \quad \omega_y = -\gamma B_y > 0$$

$$t>0, \quad \hat{\mu} = \frac{1}{2} \hbar (\omega_y \hat{\sigma}_y + \omega_z \hat{\sigma}_z)$$

$$K \{ | \pm \rangle \} = \frac{\hbar}{2} \begin{pmatrix} \omega_z & -i\omega_y \\ i\omega_y & -\omega_z \end{pmatrix}$$

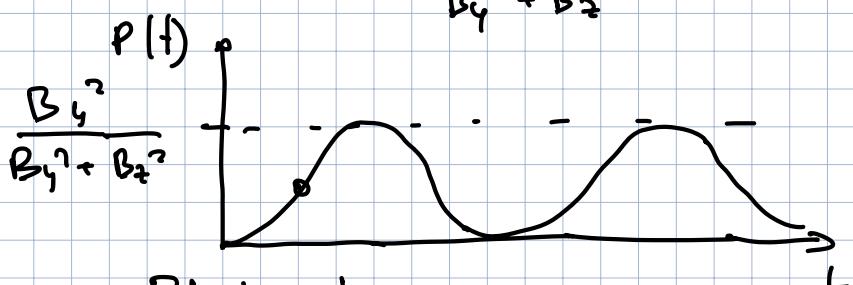
$$= E_m \underline{1} + \frac{\hbar}{2} \begin{pmatrix} \Delta & |\Omega_0| e^{-ip} \\ |\Omega_0| e^{ip} & -\Delta \end{pmatrix}$$

$$\Delta = \omega_z, \quad |\Omega_0| = \omega_y, \quad p = \frac{\pi}{2}, \quad E_m = 0$$

$$\Omega = \sqrt{\omega_y^2 + \omega_z^2} \quad \tan \Theta = \frac{|\Omega_0|}{\Delta} = \frac{\omega_y}{\omega_z} = \frac{B_y}{B_z}$$

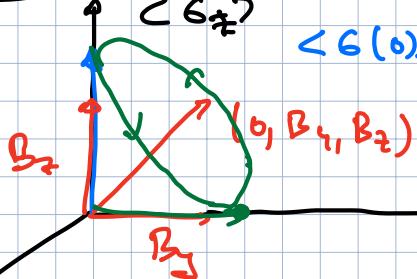
$$P_{|+\rangle_z} \Rightarrow P_{|-\rangle_z}(t) = \frac{|\Omega_0|^2}{\Omega^2} \sin^2 \left(\frac{\Omega t}{2} \right) =$$

$$= \frac{B_y^2}{B_y^2 + B_z^2} \cdot \sin^2 \left(\frac{\pm \sqrt{B_y^2 + B_z^2}}{2} t \right)$$



Bloch sphere:

$$<\sigma_z> \quad <\sigma(0)> = (0, 0, 1)$$



$$<\sigma_y> \quad |\Psi(\frac{t}{\hbar})> = |+\rangle_y$$