

OPTI 570 FINAL Extra SOLUTIONS

Problem 1

a. Need to keep at least 3×3 , since we are going to x^3 :

$$\hat{X}^2 = \frac{\sigma^2}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & \sqrt{2} \\ 0 & \sqrt{2} & 0 \dots \end{pmatrix} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & \sqrt{2} \\ 0 & \sqrt{2} & 0 \dots \end{pmatrix} = \frac{\sigma^2}{2} \begin{pmatrix} 1 & 0 & \sqrt{2} \\ 0 & 3 & 0 \\ \sqrt{2} & 0 & 2 \dots \end{pmatrix}$$

$$\hat{X}^3 = \frac{\sigma^3}{4} \begin{pmatrix} 1 & 0 & \sqrt{2} \\ 0 & 3 & 0 \\ \sqrt{2} & 0 & 2 \dots \end{pmatrix} \begin{pmatrix} 1 & 0 & \sqrt{2} \\ 0 & 3 & 0 \\ \sqrt{2} & 0 & 2 \dots \end{pmatrix} = \frac{\sigma^3}{4} \begin{pmatrix} 3 & 0 & 3\sqrt{2} \\ 0 & 9 & 0 \\ 3\sqrt{2} & 0 & 6 \dots \end{pmatrix}$$

Alternative:

$$\hat{X} = \frac{\sigma}{\sqrt{2}} (a^\dagger + a) \quad \hat{X}^2 = \frac{\sigma^2}{2} (a^{\dagger 2} + a^2 + a^\dagger a + a a^\dagger)$$

notice these already
need more states

keep only terms that bring $|0\rangle$ back to $|0\rangle$

$$\hat{X}^3 = \frac{\sigma^3}{4} (a^2 a^{\dagger 2} + a^\dagger a a a^\dagger + \dots)$$

$$\langle 0 | \hat{X}^3 | 0 \rangle = \langle 0 | \frac{\sigma^3}{4} (a^2 a^{\dagger 2} + a^\dagger a a a^\dagger) | 0 \rangle = \frac{\sigma^3}{4} (\sqrt{1} \cdot \sqrt{2} \cdot \sqrt{2} \cdot \sqrt{1} + \sqrt{1} \cdot \sqrt{1} \cdot \sqrt{1})$$

$$= \frac{\sigma^3}{4} \cdot 3$$

$$\underline{b.} \quad E_0^{(1)} = \langle 0, 0 | W | 0, 0 \rangle = \lambda \hbar \omega \cdot \frac{\sigma}{\sqrt{2}} \cdot \frac{\sigma}{\sqrt{2}} \cdot 3 = 3 \lambda \hbar \omega$$

$$E_0^{(2)} = \lambda^2 \sum_{P_x \neq 0} \frac{|\langle P_x, 0 | \hbar \omega \left(\frac{\sqrt{2}x}{\sigma} \right)^3 | 0, 0 \rangle|^2}{\hbar \omega - \hbar \omega (P_x + 1)}$$

Summing only non-zero matrix elements

$$E_0^{(2)} = \lambda^2 \hbar^2 \omega^2 \cdot \left(\underbrace{\frac{72}{-2 \hbar \omega}}_{P_x=2} + \underbrace{\frac{24}{-4 \hbar \omega}}_{P_x=4} \right) = \lambda^2 \hbar \omega (-36 - 6) = -42 \lambda^2 \hbar \omega$$

$$E_0 = \hbar\omega + 3\lambda\hbar\omega - 42\lambda^2\hbar\omega + \dots$$

$$\begin{aligned} \underline{\text{c.}} \quad |\Psi_0\rangle &\approx |0,0\rangle + \lambda \sum_{p_x \neq 0} \frac{\langle p_x, 0 | \hbar\omega \left(\frac{\sqrt{2}x}{\sigma}\right)^2 | 0,0 \rangle}{\hbar\omega - p_x} |p_x, 0\rangle \\ &\approx |0,0\rangle + \lambda \cdot \left(\underbrace{\frac{6\sqrt{2}}{-2} |2,0\rangle}_{p_x=2} + \underbrace{\frac{2\sqrt{6}}{-4} |4,0\rangle}_{p_x=4} \right) \\ |\Psi_0\rangle &\approx |0,0\rangle - 3\sqrt{2} \lambda |2,0\rangle - \frac{\sqrt{6}}{2} \lambda |4,0\rangle \end{aligned}$$

d. degenerate states $|1,0\rangle$ and $|0,1\rangle$ with energy $2\hbar\omega$

$$\begin{aligned} \text{subspace} \quad \begin{pmatrix} \langle 01 | w | 01 \rangle & \langle 01 | w | 10 \rangle \\ \langle 10 | w | 01 \rangle & \langle 10 | w | 10 \rangle \end{pmatrix} &= \begin{pmatrix} \langle 0|w|0\rangle \langle 1|1\rangle & \langle 0|w|1\rangle \langle 1|0\rangle \\ \langle 1|w|0\rangle \langle 0|1\rangle & \langle 1|w|1\rangle \langle 0|0\rangle \end{pmatrix} = \\ \text{matrix} \quad \text{is} : & \begin{pmatrix} 3 & 0 \\ 0 & 15 \end{pmatrix}. \text{ No mixing, so already diagonal} \Rightarrow \end{aligned}$$

$$E_+ = 2\hbar\omega + 15\lambda\hbar\omega$$

$$E_- = 2\hbar\omega + 3\lambda\hbar\omega$$

Problem 2

$$W(+)=i\frac{\hbar\Omega}{2} \left[a^2 e^{2i\omega t} - (a^\dagger)^2 e^{-2i\omega t} \right]$$

a. we need matrix elements that are non-zero

$$\langle 2 | a^{+2} | 0 \rangle = \sqrt{2} \text{ is the only non-zero}$$

$$\Rightarrow \lambda b_2^{(1)}(t) = \frac{1}{i\hbar} \int_0^t dt' e^{i\omega_{20} t'} W_{20}(t') =$$

$$= \frac{1}{i\hbar} \cdot i\frac{\hbar\Omega}{2} \int_0^t dt' e^{i \cdot 2 \omega t'} \cdot (-\sqrt{2}) \cdot e^{-2i\omega t'}$$

$$= -\frac{\Omega}{\sqrt{2}} t \Rightarrow P_2^{(1)}(+)= \frac{\Omega^2 t^2}{2}$$

Second order, we can couple to 2 or 4

$$\langle 4 | a^{+2} | 2 \rangle = \sqrt{3} \cdot \sqrt{5}$$

$$\begin{aligned} \lambda^2 b_4^{(1)}(t) &= \frac{1}{i\hbar} \int_0^t dt' e^{i(\omega_{02} + \omega_{04})t'} \langle 4 | b_2^{(1)} | t' \rangle = \\ &= \frac{1}{i\hbar} \int_0^t dt' e^{i(2\omega_0 + \omega_4)t'} \cdot \sqrt{3} \frac{i\hbar \omega_2}{\hbar} (-e^{-i\omega_0 t'}) \left(-\frac{\omega_2}{\hbar} + \right) = \\ &= \sqrt{2}^2 \cdot \sqrt{\frac{3}{2}} \cdot \frac{t^2}{2} = \\ &= \frac{\sqrt{6}}{2} (2t)^2 \quad P_4^{(1)}(t) = \frac{3}{8} (2t)^4 \end{aligned}$$

b. $\max [P_{0 \rightarrow 2}(\tau)] = 0.005 = \frac{\sqrt{\tau_{\max}}}{2}$

$$\tau_{\max} = \frac{\infty}{10}$$

$$\gamma = \sigma e^{-\frac{1}{10}} = \sigma \left(1 - \frac{1}{10}\right) \sim \frac{9}{10} \sigma. \text{ A little bit of squaring.}$$

Problem 3

$$w(t) = \frac{\lambda}{2} m \omega_0^2 x^2 \exp(-|t|/\tau)$$

$$\underline{a} \quad \lambda b_1^{(1)}(0) = \frac{1}{i\hbar} \int_{-\infty}^0 dt' e^{i(\omega_0 + \omega_1)t'} e^{-|t'|/\tau} \underbrace{\langle 1 | \frac{\lambda}{2} m \omega_0^2 x^2 | 0 \rangle}_{=0} = 0$$

$$= 0$$

$\Rightarrow P_1 = 0$ to all orders since w only couples states 2 away from each other

$$\underline{b} \quad \lambda b_2^{(1)}(0) = \frac{1}{i\hbar} \int_{-\infty}^0 dt' e^{i(2\omega_0 + \omega_2)t'} e^{-|t'|/\tau} \langle 2 | \frac{\lambda}{2} m \omega_0^2 x^2 | 0 \rangle =$$

$$= \frac{\lambda m \omega_0^2}{i\hbar} \int_{-\infty}^0 e^{i(2\omega_0 + \omega_2)t'} e^{-|t'|/\tau} \cdot \frac{\sigma^2}{\sqrt{2}} dt' =$$

$$x^2 = \frac{\sigma^2}{2} (a^{+2} + a^2) \Rightarrow \langle 2 | x^2 | 0 \rangle = \frac{\sigma^2}{2} \sqrt{2}$$

$$= \frac{\lambda m \omega_0^2 \sigma^2}{2\sqrt{2} i\hbar} \int_{-\infty}^0 e^{(2i\omega_0 - \frac{1}{\tau})t} dt$$

$$= \frac{\lambda m \omega_0^2 \sigma^2}{2\sqrt{2} i\hbar} \left[\frac{e^{2i\omega_0 t - \frac{1}{\tau} t}}{2i\omega_0 + \frac{1}{\tau}} \right]_0^{-\infty}$$

$$\sigma = \sqrt{\frac{\hbar}{m\omega}}$$

$$= \frac{\lambda m \omega_0^2 \sigma^2}{2\sqrt{2} i\hbar} \cdot \frac{\tau}{1 + 2i\omega_0 \tau} = \frac{\lambda m \omega_0^2 \frac{\tau}{2\sqrt{2} i\hbar}}{1 + 2i\omega_0 \tau} \frac{\tau}{1 + 2i\omega_0 \tau}$$

$$P_2^{(1)}(\infty) = \frac{\lambda^2 \omega_0^2 \tau^2}{8 \cdot (1 + 4\omega_0^2 \tau^2)}$$

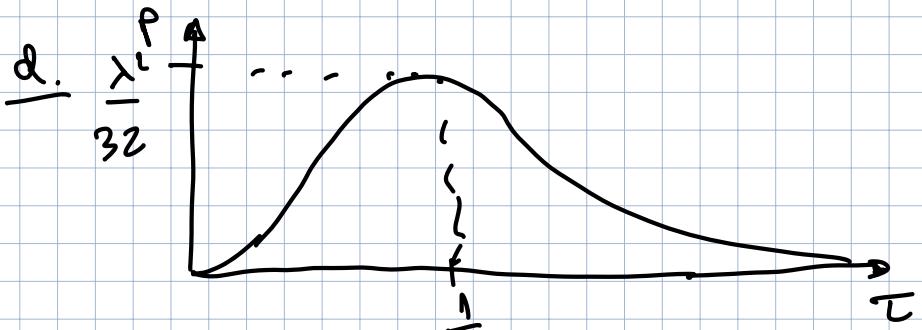
$$\underline{c.} \quad \lambda b_i^{(1)}(\infty) = \frac{1}{i\hbar} \int_{-\infty}^{+\infty} dt' e^{i\omega_0 t'} e^{-|t'|/\tau} \frac{\lambda m \omega_0^2}{2} \cdot \frac{\sigma^2}{\sqrt{2}} =$$

$$= -i \frac{\lambda m \omega_0^2 \sigma^2}{2\sqrt{2} \hbar} \int_{-\infty}^{+\infty} e^{-|t'|/\tau} e^{2i\omega_0 t'} dt' =$$

$$= -i \frac{\lambda m \omega_0^2 \frac{\tau}{2\sqrt{2} i\hbar}}{2\sqrt{2} \cdot \frac{\tau}{2}} \left[\frac{2\tau}{1 + \sqrt{2}^2 \tau^2} \right] \Big|_{\Omega = 2\omega_0}$$

$$= -i \frac{\lambda \omega_0 \tau}{2\sqrt{2}} \frac{2\tau}{1 + 4\omega_0^2 \tau^2} = -i \frac{\lambda \omega_0 \tau}{\sqrt{2} (1 + 4\omega_0^2 \tau^2)}$$

$$P_2^{(1)}(\infty) = \frac{\lambda^2 \omega_0^2 \tau^2}{2 (1 + 4\omega_0^2 \tau^2)^2}$$



$$\frac{d P_2''(\infty)}{d x} = \frac{\lambda^2}{2 \omega_0^2} = 0 \Rightarrow$$

$$\frac{2x}{(1 + \zeta \omega_0^2 x^2)^2} - \frac{2x^2}{(1 + \zeta \omega_0^2 x^2)^3} \cdot 4\omega_0^2 \cdot 2 \cdot x = 0$$

$$2x(1 + \zeta \omega_0^2 x^2) = 8\omega_0^2 x^2$$

$$\zeta \omega_0^2 x^2 = 1 \quad \omega_0^2 x_{\max}^2 = \frac{1}{\zeta} \quad x_{\max} = \frac{1}{\sqrt{\zeta} \omega_0}$$

$$P_{\max} = \frac{\lambda^2 \cdot \frac{1}{\zeta}}{2(1 + \zeta \cdot \frac{1}{\zeta})^2} = \frac{\lambda^2 \cdot \frac{1}{\zeta}}{2 \cdot \zeta} = \frac{\lambda^2}{32}$$

