

PERTURBATION THEORY

$$H = \underbrace{\frac{p^2}{2m} + \frac{1}{2} m \omega_0^2 (x^2 + y^2)} + \hbar \omega \left( \frac{x}{G} \right)^4$$

Q: eigenvalues and eigenstates of  $H$ ?

$$H_0 = \frac{p^2}{2m} + \frac{1}{2} m \omega_0^2 (x^2 + y^2)$$

w) eigenstate basis  $\{ | \varphi_m^i \rangle \}$   
 where  $H_0 | \varphi_m^i \rangle = E_m | \varphi_m^i \rangle$

$$E_m = \hbar \omega_0 (m + 1), \quad m = m_x + m_y$$

$\{ | m_x, m_y \rangle \}$  written now  $\{ | \varphi_m^i \rangle \}$

$$| 0, 0 \rangle = | \varphi_0 \rangle \quad - E_0$$

$$\left\{ \begin{array}{l} | 1, 0 \rangle = | \varphi_1^1 \rangle \\ | 0, 1 \rangle = | \varphi_1^2 \rangle \end{array} \right\} \text{ both energy } E_1$$

$$\left\{ \begin{array}{l} | 2, 0 \rangle = | \varphi_2^1 \rangle \\ | 1, 1 \rangle = | \varphi_2^2 \rangle \\ | 0, 2 \rangle = | \varphi_2^3 \rangle \end{array} \right\} \text{ all energy } E_2$$

...

$$H = H_0 + \hbar \omega \left( \frac{x}{G} \right)^4$$

- Options:
- analytic solution, if possible (in general NOT)
  - numerical / computational
  - approximation methods

$$H = H_0 + \underbrace{\hbar \omega \left( \frac{x}{\ell} \right)^4}_W$$

case:  $W$  is small compared to  $H_0$

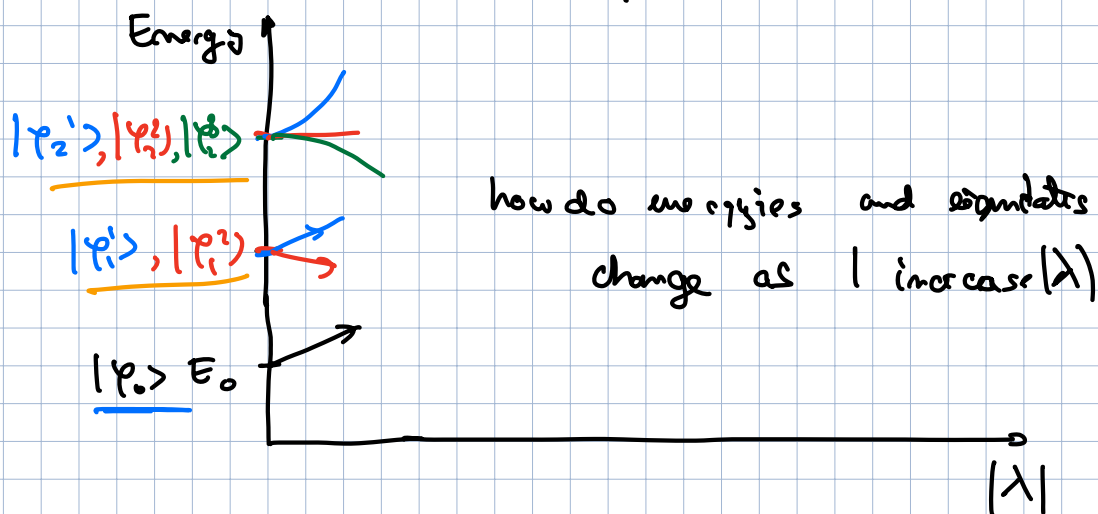
$$H = H_0 + \underbrace{W}_{\text{small perturbation}}$$

$$W = \hbar \omega \left( \frac{x}{\ell} \right)^4$$

$$\begin{aligned} W &= \underbrace{\frac{\omega}{\omega_0}}_{\lambda} \hbar \omega_0 \left( \frac{x}{\ell} \right)^4 = \\ &= \underbrace{\lambda}_{\lambda} \underbrace{\hbar \omega_0 \left( \frac{x}{\ell} \right)^4}_{\hat{W}} \end{aligned}$$

$$W = \sum \hat{W}$$

$$W - \text{small} \Rightarrow |\lambda| \ll 1$$



### Stationary perturbation theory

- (1) Non-degenerate: no degeneracies in  $H_0$   $g_m = 1$
- (2) Degenerate in  $H_0$ :  $g_m > 1$

Given:  $H_0 | \varphi_m^i \rangle = E_m^0 | \varphi_m^i \rangle$ ,  $i \in \{1, 2, \dots, g_m\}$   
• solutions known exactly for  $m$

$$H = H_0 + W$$

Question: find the eigenstates + eigenvalues of  $H$

Challenge: no exact solutions

SPT:  $W = \lambda \hat{W}$        $\lambda$  - real #, dimensionless

Goal:  $H(\lambda) | \Psi_{m,j}(\lambda) \rangle = E_{m,j}(\lambda) | \Psi_{m,j}(\lambda) \rangle$   
 $j \in \{1, \dots, g_m\}$

Approach: write  $| \Psi_{m,j}(\lambda) \rangle$ ,  $E_{m,j}(\lambda)$  as power series in  $\lambda$

Main results

Non-degenerate:  $g_m = 1$

$$E_m \approx \underbrace{E_m^0}_{\text{0th order}} + \underbrace{\lambda \langle \varphi_m | \hat{W} | \varphi_m \rangle}_{\text{1st order}} + \underbrace{\lambda^2 \sum_{p \neq m} \sum_{i=1}^{g_p} \frac{|\langle \varphi_p^i | \hat{W} | \varphi_m \rangle|^2}{E_m^0 - E_p^0}}_{\text{2nd order}}$$

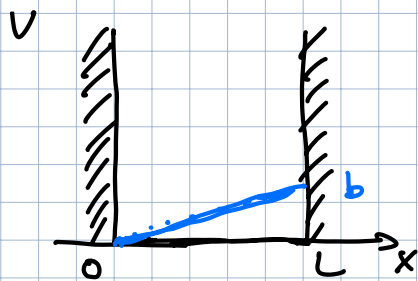
$$\approx E_m^0 \langle \varphi_m | W | \varphi_m \rangle + \sum_{p \neq m} \sum_{i=1}^{g_p} \frac{|\langle \varphi_p^i | W | \varphi_m \rangle|^2}{E_m^0 - E_p^0}$$

$$| \Psi_m \rangle \approx | \varphi_m \rangle + \lambda \sum_{p \neq m} \sum_{i=1}^{g_p} \frac{\langle \varphi_p^i | \hat{W} | \varphi_m \rangle}{E_m^0 - E_p^0} | \varphi_p^i \rangle$$

$$\approx | \varphi_m \rangle + \sum_{p \neq m} \sum_{i=1}^{g_p} \frac{\langle \varphi_p^i | W | \varphi_m \rangle}{E_m^0 - E_p^0} | \varphi_p^i \rangle$$

$| \Psi_m \rangle$  - normalized? - No

Example: non-degenerate perturbation theory



infinite square well

$$\varphi_n(x) = \langle x | \varphi_n \rangle = \sqrt{\frac{2}{L}} \sin \left( \frac{n\pi x}{L} \right)$$

$$n \geq 1$$

$$E_n = n^2 \frac{\hbar^2 \pi^2}{2mL^2} = n^2 \cdot E_1^0, E_1^0 = \frac{\hbar^2 \pi^2}{2mL^2}$$

Add perturbation:

$$W = b \frac{x}{L}, \quad b \ll E_1^0$$

$$\lambda = \frac{b}{E_1^0}$$

$$W = \frac{b}{E_1^0} \cdot E_1^0 \frac{x}{L}$$

$$W = \lambda \cdot \hat{W}, \quad \text{with } \lambda \ll 1$$

How does  $W$  change the ground state energy and eigenstate?

$$E_1 = E_1^0 + \lambda \langle \varphi_1 | \hat{W} | \varphi_1 \rangle + \lambda^2 \sum_{p \neq 1} \frac{|\langle \varphi_p | \hat{W} | \varphi_1 \rangle|^2}{E_1^0 - E_p^0}$$

$g_p = 1$  for all states

$$E_1 = E_1^0 + \lambda \langle \varphi_1 | \hat{W} | \varphi_1 \rangle + \lambda^2 \sum_{p=2}^{\infty} \frac{|\langle \varphi_p | \hat{W} | \varphi_1 \rangle|^2}{E_1^0 - E_p^0}$$

$$\langle \varphi_p | \hat{W} | \varphi_p \rangle =$$

$$= \int_0^L dx \frac{2}{L} \sin\left(\frac{\pi x}{L}\right) \cdot \sin\left(\frac{p\pi x}{L}\right) \cdot \frac{x}{L}$$

$$= \begin{cases} \frac{1}{2} & p=1 \\ -\frac{8p}{\pi^2(p^2-1)^2} & \text{for } p \geq 2 \text{ even} \\ 0 & \text{for } p \geq 3 \text{ odd} \end{cases}$$

$$\text{Then: } E_1 = E_1^0 \left[ 1 + \lambda \cdot \frac{1}{2} - \lambda^2 \cdot \sum_{\substack{p=2 \\ \text{even}}}^{\infty} \frac{1}{(p^2-1)} \frac{\delta \psi p^2}{\pi^4 (p^2-1)^4} \right]$$

$$= E_1^0 \left[ 1 + \frac{\lambda}{2} - \frac{\lambda^2}{\pi^4} \sum_{\substack{p=2 \\ \text{even}}}^{\infty} \frac{p^2}{(p^2-1)^5} \right]$$

$\approx 0.004$

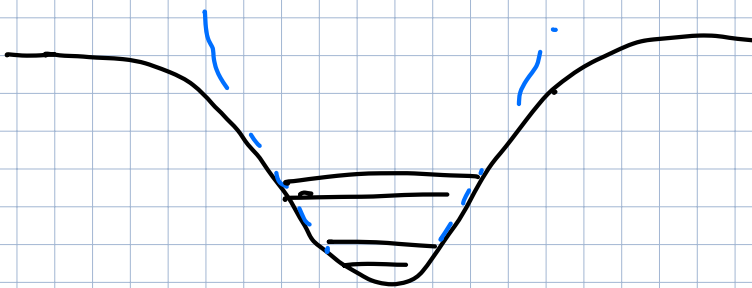
$$= E_1^0 \left( 1 + \frac{\lambda}{2} - 0.011 \lambda^2 \right) \quad \text{w/ } \lambda = \frac{b}{E_1}$$



$$|\psi_1\rangle = |\varphi_1\rangle - \lambda \sum_{\substack{p=2 \\ \text{even}}}^{\infty} \frac{1}{p^2-1} \left( -\frac{\delta p}{\pi^2 (p^2-1)^2} \right) |\varphi_p\rangle$$

$$= |\varphi_1\rangle + \lambda \sum_{\substack{p=2 \\ \text{even}}}^{\infty} \frac{8}{\pi^2} \frac{p}{(p^2-1)^3} |\varphi_p\rangle =$$

$$= |\varphi_1\rangle + \lambda \left[ \underbrace{0.06}_{\text{blue}} |\varphi_2\rangle + \underbrace{0.001}_{\text{blue}} |\varphi_4\rangle + \underbrace{0.0001}_{\text{blue}} |\varphi_6\rangle + \dots \right]$$



Degenerate stationary perturbation theory

CT. , main results

$$\sum_{i'=1}^{l_n} \langle \varphi_n^i | \hat{W} | \varphi_n^{i'} \rangle \langle \varphi_n^{i'} | 0_j \rangle = \epsilon_{ij} \langle \varphi_n^i | 0_j \rangle$$