

# Assignment 10

## OPTI 570 Quantum Mechanics

### University of Arizona

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Total time: 7 hours

### Problem I

a) The probabilities for x and y directions are:

$$P_x = |\langle \phi_x | \phi_\mu \rangle|^2 = \cos^2 \mu, \quad \text{and} \quad P_y = |\langle \phi_y | \phi_\mu \rangle|^2 = \sin^2 \mu.$$

For circular polarization we do the same:

$$P_{\sigma_+} = |\langle \sigma_+ | \phi_\mu \rangle|^2 = \left| -\frac{1}{\sqrt{2}} [(\langle \phi_x | - i \langle \phi_y |)(\cos \mu | \phi_x \rangle + \sin \mu | \phi_y \rangle)] \right|^2 = \frac{1}{2} (\sin^2 \mu + \cos^2 \mu) = \frac{1}{2}$$

$$P_{\sigma_-} = 1 - P_{\sigma_+} = \frac{1}{2}.$$

b) The state produced by the source is:

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|\phi_x^L\rangle |\phi_y^R\rangle - |\phi_x^R\rangle |\phi_y^L\rangle).$$

The joint measurement is performed through the following projector which is the tensorproduct of the measurement for each photon:

$$P_{jk} = P_j^L \otimes P_k^R = |\phi_j^L\rangle \langle \phi_j^L| |\phi_k^R\rangle \langle \phi_k^R|, \quad j, k \in \{x, y\}.$$

We are let  $j, k$  to vary in either direction so that we will have four types of measurements:

$$P_{xy} = \langle \psi | P_{xy} | \psi \rangle = |\langle \phi_x^L \phi_y^R | \psi \rangle|^2 = \frac{1}{2}$$

$$P_{yx} = \langle \psi | P_{yx} | \psi \rangle = |\langle \phi_y^L \phi_x^R | \psi \rangle|^2 = \frac{1}{2}$$

$$P_{xx} = \langle \psi | P_{xx} | \psi \rangle = |\langle \phi_x^L \phi_x^R | \psi \rangle|^2 = 0$$

$$P_{yy} = \langle \psi | P_{yy} | \psi \rangle = |\langle \phi_y^L \phi_y^R | \psi \rangle|^2 = 0$$

Only the non-zero probabilities are possible state the system may be left after the measurement, which are then:

$$|\psi\rangle \xrightarrow{P_{xy}} \frac{P_{xy} |\psi\rangle}{\sqrt{\langle \psi | P_{xy} | \psi \rangle}} = |\phi_x^L \phi_y^R\rangle, \quad \text{and} \quad |\psi\rangle \xrightarrow{P_{yx}} \frac{P_{yx} |\psi\rangle}{\sqrt{\langle \psi | P_{yx} | \psi \rangle}} = |\phi_y^L \phi_x^R\rangle.$$

On the other hand, measurement of a single photon is kind of a marginal measurement. The probabilities can be constructed from the above probabilities:

$$\begin{aligned} P_x^L &= P_{xy} + P_{xx} = \frac{1}{2} + 0 = \frac{1}{2} \\ P_x^R &= P_{xx} + P_{yx} = 0 + \frac{1}{2} = \frac{1}{2} \\ P_y^L &= P_{yx} + P_{yy} = \frac{1}{2} + 0 = \frac{1}{2} \\ P_y^R &= P_{xy} + P_{yy} = \frac{1}{2} + 0 = \frac{1}{2} \end{aligned}$$

c) We first change of basis, by passing from  $\phi_{x,y}$  to  $\phi_{+,-}$ . The transformation matrix is:

$$M = \frac{-1}{\sqrt{2}} \begin{bmatrix} 1 & i \\ 1 & -i \end{bmatrix} \longrightarrow M^\dagger = \frac{-1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ -i & i \end{bmatrix}$$

. By extracting the equation from  $M^\dagger$  and substitute them in  $\phi_{x,y}$  we find the following equivalent state:

$$|\psi\rangle \propto \frac{1}{\sqrt{2}}(|\phi_+^L \phi_-^R\rangle - |\phi_-^L \phi_+^R\rangle).$$

The measurements are analogous to the previous part, but now in the circular polarized basis, with the exactly same measurement projector, but with  $j, k$  changed:

$$P_{jk} = P_j^L \otimes P_k^R = |\phi_j^L\rangle\langle\phi_j^L| |\phi_k^R\rangle\langle\phi_k^R|, \quad j, k \in \{-, +\}.$$

Therefore,

$$\begin{aligned} P_{++} &= |\langle\phi_+^L \phi_+^R|\psi\rangle|^2 = 0 \\ P_{+-} &= |\langle\phi_+^L \phi_-^R|\psi\rangle|^2 = \frac{1}{2} \\ P_{-+} &= |\langle\phi_-^L \phi_+^R|\psi\rangle|^2 = \frac{1}{2} \\ P_{--} &= |\langle\phi_-^L \phi_-^R|\psi\rangle|^2 = 0. \end{aligned}$$

We see the polarization for either basis is exactly the same, but the states after the measurement will be different.

## Problem II

In the deuterium atom we have that  $I = 1$  and  $S = 1/2$ .

a) In the state  $1s$ , we have  $n = 1$  and  $l = 0$ . The quantum number  $J$  is:

$$J \in \{|L - S|, |L - S| + 1, \dots, L + S - 1, L + S\} = \left\{\frac{1}{2}\right\}.$$

For this unique value, we have

$$J = \frac{1}{2} : \quad F \in \{|J - I|, |J - I| + 1, \dots, J + I - 1, J + I\} = \left\{\frac{1}{2}, \frac{3}{2}\right\}.$$

b) For the state  $2p$ , we have  $n = 2$  and  $l = 1$ . The quantum number  $J$  is:

$$J \in \left\{ \left| 1 - \frac{1}{2} \right|, 1 + \frac{1}{2} \right\} = \left\{ \frac{1}{2}, \frac{3}{2} \right\}.$$

For each  $J$ , we have  $F$ :

$$\begin{aligned} J = \frac{1}{2}: \quad F &\in \left\{ \left| \frac{1}{2} - 1 \right|, \frac{1}{2} + 1 \right\} = \left\{ \frac{1}{2}, \frac{3}{2} \right\} \\ J = \frac{3}{2}: \quad F &\in \left\{ \left| \frac{3}{2} - 1 \right|, \left| \frac{3}{2} - 1 \right| + 1, \frac{3}{2} + 1 \right\} = \left\{ \frac{1}{2}, \frac{3}{2}, \frac{5}{2} \right\}. \end{aligned}$$

### Problem III

We are asked to compute the matrix element of  $A$ , but the initial and final states are in the TAM basis, which is not suitable due to the action of  $A$  only on  $|L, M_L\rangle$ ; in the TAM basis those are encoded. We need therefore to convert from TAM basis to TP basis each  $|\phi\rangle$  and  $|\phi'\rangle$ . We need to do the following for each state:

$$|F, M_F\rangle \xrightarrow{CG} |J, M_J\rangle |I, M_I\rangle \xrightarrow{CG} |L, M_L\rangle |S, M_S\rangle |I, M_I\rangle.$$

We do that for each state then,

a)  $|\phi\rangle$  has  $L = 0$ ,  $S = 1/2$ ,  $J = 1/2$ ,  $I = 3/2$ : we look at the  $\frac{3}{2} \times \frac{1}{2}$  table:

$$|\phi\rangle = |F = 1, M_F = 1\rangle = \frac{1}{2} |J = \frac{1}{2}, m_J = \frac{1}{2}\rangle |I = \frac{3}{2}, M_I = \frac{1}{2}\rangle - \frac{\sqrt{3}}{2} |J = \frac{1}{2}, m_J = -\frac{1}{2}\rangle |I = \frac{3}{2}, M_I = \frac{3}{2}\rangle.$$

Now, the second decomposition:

$$|J = \frac{1}{2}, m_J = \pm \frac{1}{2}\rangle = |L = 0, M_L = 0\rangle |S = \frac{1}{2}, M_S = \pm \frac{1}{2}\rangle.$$

Inserting it into the first equation:

$$|\phi\rangle = \frac{1}{2} |0, 0; \frac{1}{2}, \frac{1}{2}; \frac{3}{2}, \frac{1}{2}\rangle - \frac{\sqrt{3}}{2} |0, 0; \frac{1}{2}, -\frac{1}{2}; \frac{3}{2}, \frac{3}{2}\rangle.$$

Doing the same for  $|\phi'\rangle$ :

$$|\phi'\rangle = \frac{\sqrt{3}}{6} |1, 0; \frac{1}{2}, \frac{1}{2}; \frac{3}{2}, \frac{1}{2}\rangle - \frac{1}{2} |1, 0; \frac{1}{2}, -\frac{1}{2}; \frac{3}{2}, \frac{3}{2}\rangle.$$

We now do the inner product:

$$\langle \phi' | A | \phi \rangle = A_{1000} \left[ -\frac{\sqrt{3}}{6} \cdot \frac{1}{2} + \frac{1}{2} \cdot \frac{\sqrt{3}}{2} \right] = \frac{\sqrt{3}}{6} A_{1000} \longrightarrow T = \frac{1}{12} |A_{1000}|^2.$$

The process is exactly the same for the others, we just show the results.

b) s

$$\begin{aligned} |\phi\rangle &= \frac{\sqrt{2}}{2} |0, 0; \frac{1}{2}, \frac{1}{2}; \frac{3}{2}, -\frac{1}{2}\rangle - \frac{\sqrt{2}}{2} |0, 0; \frac{1}{2}, -\frac{1}{2}; \frac{3}{2}, \frac{1}{2}\rangle \\ |\phi'\rangle &= -\frac{\sqrt{6}}{6} |1, 0; \frac{1}{2}, -\frac{1}{2}\rangle - \frac{\sqrt{6}}{6} |1, 0; -\frac{1}{2}, \frac{1}{2}\rangle \\ T &= 0. \end{aligned}$$

c) asgas

$$\begin{aligned}
|\phi\rangle &= \frac{1}{2}|0, 0; \frac{1}{2}; \frac{1}{2}\rangle - \frac{\sqrt{3}}{2}|0, 0; -\frac{1}{2}; \frac{3}{2}\rangle \\
|\phi'\rangle &= -\frac{1}{2}|1, 0; \frac{1}{2}; \frac{1}{2}\rangle + \frac{\sqrt{3}}{6}|1, 0; -\frac{1}{2}; \frac{3}{2}\rangle \\
T &= \frac{1}{4}|A_{1000}|^2.
\end{aligned}$$

d) asgasg

$$\begin{aligned}
|\phi\rangle &= \frac{\sqrt{2}}{2}|0, 0; \frac{1}{2}; -\frac{1}{2}\rangle - \frac{\sqrt{2}}{2}|0, 0; -\frac{1}{2}; \frac{1}{2}\rangle \\
|\phi'\rangle &= -\frac{\sqrt{6}}{6}|1, 0; \frac{1}{2}; -\frac{1}{2}\rangle + \frac{\sqrt{6}}{6}|1, 0; -\frac{1}{2}; \frac{1}{2}\rangle \\
T &= \frac{1}{3}|A_{1000}|^2.
\end{aligned}$$