

OPTI 570 FINAL EXAM SOLUTIONS

Problem 1

a. Need to keep at least 3×3 , since we are going to x^4 :

$$\hat{X}^2 = \frac{\sigma^2}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & \sqrt{2} \\ 0 & \sqrt{2} & 0 \\ \dots & \dots & \dots \end{pmatrix} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & \sqrt{2} \\ 0 & \sqrt{2} & 0 \\ \dots & \dots & \dots \end{pmatrix} = \frac{\sigma^2}{2} \begin{pmatrix} 1 & 0 & \sqrt{2} \\ 0 & 3 & 0 \\ \sqrt{2} & 0 & 2 \\ \dots & \dots & \dots \end{pmatrix}$$

$$\hat{X}^4 = \frac{\sigma^4}{4} \begin{pmatrix} 1 & 0 & \sqrt{2} \\ 0 & 3 & 0 \\ \sqrt{2} & 0 & 2 \\ \dots & \dots & \dots \end{pmatrix} \begin{pmatrix} 1 & 0 & \sqrt{2} \\ 0 & 3 & 0 \\ \sqrt{2} & 0 & 2 \\ \dots & \dots & \dots \end{pmatrix} = \frac{\sigma^4}{4} \begin{pmatrix} 3 & 0 & 3\sqrt{2} \\ 0 & 9 & 0 \\ 3\sqrt{2} & 0 & 6 \\ \dots & \dots & \dots \end{pmatrix}$$

notice these already need more states

Alternative:

$$\hat{X} = \frac{\sigma}{\sqrt{2}} (a^\dagger + a) \quad \hat{X}^2 = \frac{\sigma^2}{2} (a^{\dagger 2} + a^2 + a^\dagger a + a a^\dagger)$$

keep only terms that bring $|0\rangle$ back to $|0\rangle$

$$\hat{X}^4 = \frac{\sigma^4}{4} (a^2 a^{\dagger 2} + a^\dagger a a a^\dagger + \dots)$$

$$\begin{aligned} \langle 0 | \hat{X}^4 | 0 \rangle &= \langle 0 | \frac{\sigma^4}{4} (a^2 a^{\dagger 2} + a^\dagger a a a^\dagger) | 0 \rangle = \frac{\sigma^4}{4} (\sqrt{1} \cdot \sqrt{2} \cdot \sqrt{2} \cdot \sqrt{1} + \sqrt{1} \cdot \sqrt{1} \cdot \sqrt{1} \cdot \sqrt{1}) \\ &= \frac{\sigma^4}{4} \cdot 3 \end{aligned}$$

b. $E_0^{(1)} = \langle 0,0 | W | 0,0 \rangle = \lambda \hbar \omega \cdot \cancel{\frac{4}{\sigma^4}} \cdot \cancel{\frac{\sigma^4}{4}} \cdot 3 = 3 \lambda \hbar \omega$

$$E_0^{(2)} = \lambda^2 \sum_{p_x \neq 0} \frac{|\langle p_x, 0 | \hbar \omega \left(\frac{\sqrt{2} x}{\sigma} \right)^4 | 0,0 \rangle|^2}{\hbar \omega - \hbar \omega (p_x + 1)}$$

Summing only non-zero matrix elements

$$\begin{aligned} E_0^{(2)} &= \lambda^2 \hbar^2 \omega^2 \cdot \left(\underbrace{\frac{72}{-2\hbar\omega}}_{p_x=2} + \underbrace{\frac{24}{-4\hbar\omega}}_{p_x=4} \right) = \lambda^2 \hbar \omega (-36 - 6) = \\ &= -42 \lambda^2 \hbar \omega \end{aligned}$$

$$E_0 = \hbar\omega + 3\lambda\hbar\omega - 12\lambda^2\hbar\omega + \dots$$

$$\underline{c.} \quad |\psi_0\rangle \approx |0,0\rangle + \lambda \sum_{p_x \neq 0} \frac{\langle p_x, 0 | \hbar\omega \left(\frac{\sqrt{2}x}{\sigma}\right)^4 | 0,0\rangle}{\hbar\omega(-p_x)} |p_x, 0\rangle$$

$$\approx |0,0\rangle + \lambda \cdot \left(\underbrace{\frac{6\sqrt{2}}{-2}}_{p_x=2} |2,0\rangle + \underbrace{\frac{2\sqrt{6}}{-4}}_{p_x=4} |4,0\rangle \right)$$

$$|\psi_0\rangle \approx |0,0\rangle - 3\sqrt{2}\lambda |2,0\rangle - \frac{\sqrt{6}}{2}\lambda |4,0\rangle$$

d. degenerate states $|1,0\rangle$ and $|0,1\rangle$ with energy $2\hbar\omega$

subspace matrix is:
$$\begin{pmatrix} \langle 01 | W | 01 \rangle & \langle 01 | W | 10 \rangle \\ \langle 10 | W | 01 \rangle & \langle 10 | W | 10 \rangle \end{pmatrix} = \begin{pmatrix} \langle 01 | W | 0 \rangle \langle 1 | 1 \rangle & \langle 01 | W | 1 \rangle \langle 1 | 0 \rangle \\ \langle 11 | W | 0 \rangle \langle 0 | 1 \rangle & \langle 11 | W | 1 \rangle \langle 0 | 0 \rangle \end{pmatrix} =$$

$$= \lambda\hbar\omega \begin{pmatrix} 3 & 0 \\ 0 & 15 \end{pmatrix}. \text{ No mixing, so already diagonal } \Rightarrow$$

$$E_+ = 2\hbar\omega + 15\lambda\hbar\omega$$

$$E_- = 2\hbar\omega + 3\lambda\hbar\omega$$

Problem 2

$$W(t) = i\frac{\hbar\Omega}{2} [a^2 e^{2i\omega t} - (a^\dagger)^2 e^{-2i\omega t}]$$

a. we need matrix elements that are non-zero

$$\langle 2 | a^{\dagger 2} | 0 \rangle = \sqrt{2} \text{ is the only non-zero}$$

$$\Rightarrow \lambda b_2^{(1)}(t) = \frac{1}{i\hbar} \int_0^t dt' e^{i\omega_{20}t'} W_{20}(t') =$$

$$= \frac{1}{i\hbar} \cdot i\frac{\hbar\Omega}{2} \int_0^t dt' e^{i \cdot 2\omega t'} \cdot (-\sqrt{2}) \cdot e^{-2i\omega t'}$$

$$= -\frac{\Omega}{\sqrt{2}} t \Rightarrow p_2^{(1)}(t) = \frac{\Omega t^2}{2}$$

Second order, we can couple to 2 or 4

$$\langle 4 | a^{+2} | 2 \rangle = \sqrt{3} \cdot \sqrt{4}$$

$$\begin{aligned} \lambda^2 b_4^{(2)}(t) &= \frac{1}{i\hbar} \int_0^t dt' e^{i\omega_4 t'} W_{42}(t') \cdot \lambda b_2^{(1)}(t) = \\ &= \frac{1}{i\hbar} \int_0^t dt' e^{i2\omega t'} \cdot \cancel{\sqrt{3}} \frac{i\hbar \Omega}{\cancel{\sqrt{2}}} (-e^{-i2\omega t'}) \left(-\frac{\Omega}{\sqrt{2}} \cdot t\right) = \\ &= \Omega^2 \cdot \sqrt{\frac{3}{2}} \cdot \frac{t^2}{2} = \\ &= \frac{\sqrt{6}}{4} (\Omega t)^2 \end{aligned} \quad P_4^{(2)}(t) = \frac{3}{8} (\Omega t)^4$$

b. $\max [P_{0 \rightarrow 2}(\tau)] = 0.005 = \frac{\Omega^2 \tau_{\max}^2}{2}$

$$\tau_{\max} = \frac{\Omega}{10}$$

$$\gamma = \sigma e^{-\frac{1}{10}} = \sigma \left(1 - \frac{1}{10}\right) \sim \frac{9}{10} \sigma. \text{ A little bit of squeezing.}$$

Problem 3

$$W(t) = \frac{\lambda}{2} m \omega_0^2 x^2 \exp(-|t|/\tau)$$

a. $\lambda b_1^{(1)}(0) = \frac{1}{i\hbar} \int_{-\infty}^0 dt' e^{i\omega_0 t'} e^{-|t'|/\tau} \underbrace{\langle 1 | \frac{\lambda}{2} m \omega_0^2 x^2 | 0 \rangle}_{=0}$

$$= 0$$

$\Rightarrow P_1 = 0$ to all orders since W only couples states 2 away from each other

b. $\lambda b_2^{(1)}(0) = \frac{1}{i\hbar} \int_{-\infty}^0 dt' e^{i2\omega_0 t'} e^{t'/\tau} \langle 2 | \frac{\lambda}{2} m \omega_0^2 x^2 | 0 \rangle =$

$$= \frac{\lambda m \omega_0^2}{i\hbar} \int_{-\infty}^0 e^{i2\omega_0 t'} e^{t'/\tau} \cdot \frac{\sigma^2}{\sqrt{2}} dt' =$$

$$x^2 = \frac{\sigma^2}{2} (a^{+2} + a^2) \Rightarrow \langle 2 | x^2 | 0 \rangle = \frac{\sigma^2}{2} \sqrt{2}$$

$$= \frac{\lambda m \omega_0^2 \sigma^2}{2\sqrt{2} i \hbar} \int_{-\infty}^0 e^{\left(2i\omega_0 + \frac{1}{\tau}\right)t'} dt'$$

$$= \frac{\lambda m \omega_0^2 \sigma^2}{2\sqrt{2} i \hbar} \left. \frac{e^{2i\omega_0 + \frac{1}{\tau}t'}}{2i\omega_0 + \frac{1}{\tau}} \right|_{-\infty}^0 =$$

$$\sigma = \sqrt{\frac{\hbar}{m\omega}}$$

$$= \frac{\lambda m \omega_0^2 \sigma^2}{2\sqrt{2} i \hbar} \cdot \frac{\tau}{1 + 2i\omega_0 \tau} = \frac{\lambda m \omega_0^2 \frac{\hbar}{m\omega_0}}{2\sqrt{2} i \hbar} \cdot \frac{\tau}{1 + 2i\omega_0 \tau}$$

$$P_2^{(1)}(0) = \frac{\lambda^2 \omega_0^2 \tau^2}{8 \cdot (1 + 4\omega_0^2 \tau^2)}$$

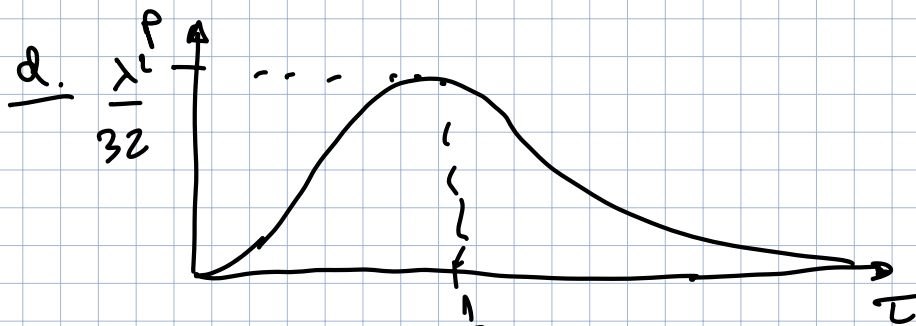
$$\underline{c.} \quad \lambda b_2^{(1)}(\infty) = \frac{1}{i\hbar} \int_{-\infty}^{+\infty} dt' e^{i\omega_0 t'} e^{-|t'|/\tau} \frac{\lambda}{2} m \omega_0^2 \cdot \frac{\sigma^2}{\sqrt{2}} =$$

$$= -i \frac{\lambda m \omega_0^2 \sigma^2}{2\sqrt{2} \hbar} \int_{-\infty}^{+\infty} e^{-|t'|/\tau} e^{2i\omega_0 t'} dt =$$

$$= -i \frac{\lambda m \omega_0^2 \frac{\hbar}{m\omega_0}}{2\sqrt{2} \cdot \hbar} \left. \frac{2\tau}{1 + \Omega^2 \tau^2} \right|_{\Omega = 2\omega_0}$$

$$= -i \frac{\lambda \omega_0 \tau}{2\sqrt{2}} \frac{2\tau}{1 + 4\omega_0^2 \tau^2} = -i \frac{\lambda \omega_0 \tau}{\sqrt{2} (1 + 4\omega_0^2 \tau^2)}$$

$$P_2^{(1)}(\infty) = \frac{\lambda^2 \omega_0^2 \tau^2}{2 (1 + 4\omega_0^2 \tau^2)^2}$$



$$\frac{d P_2^{(1)}(\infty)}{d \tau} = 0 \Rightarrow$$

$$\frac{2\tau}{(1+4\omega_0^2\tau^2)^2} - \frac{2\tau^2}{(1+4\omega_0^2\tau^2)^3} \cdot 4\omega_0^2 \cdot 2 \cdot \tau = 0$$

$$2\tau(1+4\omega_0^2\tau^2) = 8\omega_0^2\tau^3$$

$$4\omega_0^2\tau^2 = 1 \quad \omega_0^2\tau_{\max}^2 = \frac{1}{4} \quad \tau_{\max} = \frac{1}{2\omega_0}$$

$$P_{\max} = \frac{\lambda^2 \cdot \frac{1}{4}}{2\left(1+4 \cdot \frac{1}{4}\right)^2} = \frac{\lambda^2 \cdot \frac{1}{4}}{2 \cdot 4} = \frac{\lambda^2}{32}$$

