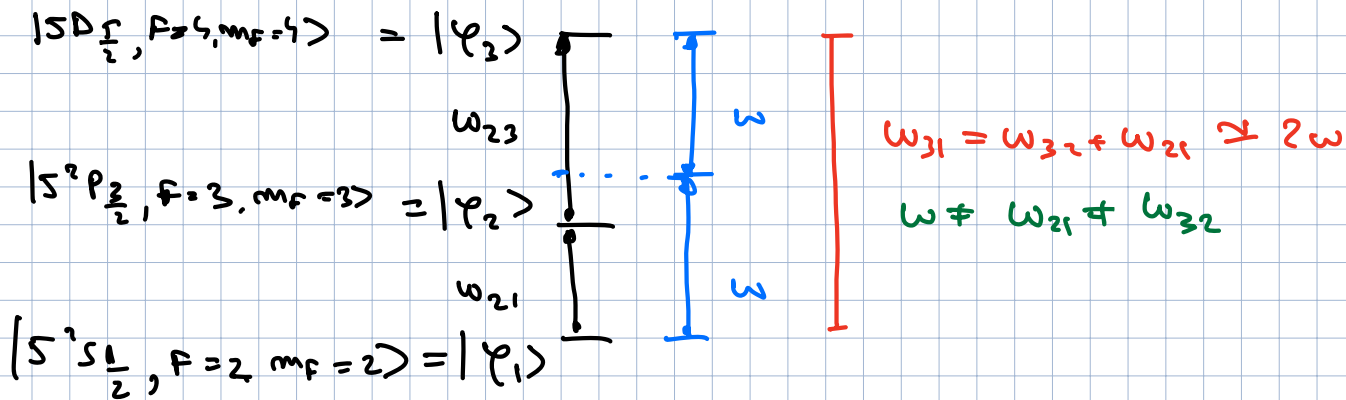


# OPTI 570 LECTURE Tu DEC 9

## Last time

$8^2$  Rb atoms that are interacting w/  $\sigma^+$  light



atom starts in  $|\psi_1\rangle$

TDPT - prob. of transition from  $|\psi_1\rangle$  to  $|\psi_3\rangle$

$$P_3^{(2)}(t) = \left| \frac{\Omega_{21} \Omega_{32}}{4 \Delta_{21}} \right|^2 t^2 \text{sinc}^2\left(\frac{\delta \cdot t}{2}\right)$$

$\delta = 2\omega - \omega_{31}$  "two-photon detuning"

$\delta = 0 \Leftrightarrow \omega = \frac{\omega_{31}}{2}$  maximizes  $P_3^{(2)}(t)$

$$P_3^{(2)}(t) = \left| \frac{\Omega_{21} \Omega_{32}}{\Delta_{21}} \right|^2 \left(\frac{t}{2}\right)^2 \quad \Delta_{21} = \omega - \omega_{21}$$

Q: What is  $\lambda^2 P_2^{(2)} = ?$

$$\lambda^2 b_2^{(2)}(t) = \frac{1}{i\hbar} \int_0^t dt' e^{i\omega_{32}t'} \underbrace{\omega_{22}}_0(t') \left[ \lambda \underbrace{b_2^{(1)}}(t') \right]$$

$$\lambda^2 b_2^{(2)}(t) = 0$$

$$\Rightarrow P_2^{(2)}(t) = P_2^{(1)}(t) = \left| \frac{\Omega_{21}}{\Delta_{21}} \right|^2 \text{sinc}^2\left(\frac{\Delta_{21}t}{2}\right)$$

$$\frac{P_3^{(2)}}{P_2^{(2)}} = \left| \frac{\Omega_{21} \cdot \Omega_{32}}{\Delta_{21}} \right|^2 \left( \frac{t}{2} \right)^2 \left| \frac{\Delta_{21}}{\Omega_{21}} \right|^2 \frac{1}{\sin^2(\Delta_{21} t/2)} =$$

$$= \left| \frac{P_{32} \epsilon_0}{\hbar \Delta_{21}} \right|^2 \cdot \frac{1}{\sin^2\left(\frac{\Delta_{21} t}{2}\right)} \frac{\sin^2(\Delta_{21} t/2)}{\Delta_{21}^2 \cdot t^2/4} \cdot \frac{t^2}{4} =$$

$$= \frac{t^2}{4} \cdot \sin^2(\Delta_{21} t/2)$$

$P_3^{(2)}, P_2^{(2)} \ll 1$  necessary

but: ratio  $\gg 1$  possible.

$$I = \frac{1}{2} c \epsilon_0 |\epsilon_0|^2$$

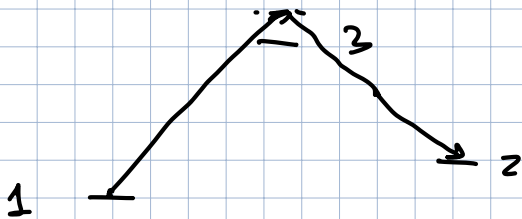
$$P_2^{(2)} \propto I$$

$$P_3^{(2)} \propto I^2 \rightarrow \text{non-linear in intensity}$$

two-photon process

Large I: two-photon process  $(|1\rangle \rightarrow |2\rangle)$  could be stronger than single photon absorption  $(|1\rangle \rightarrow |2\rangle)$

Lamb level structure



Stimulated Raman Transitions

Pulse perturbations

- separate  $\hat{W}(t)$  into time independent + time dependent.

$$\lambda \hat{W}(t) = \lambda \hat{W}_i f(t)$$

$|\psi\rangle_i$ , let  $t_0 = -\infty$ , evaluate  $P_{i \rightarrow f}^{(1)}(t = \infty)$

$$\lambda b_f^{(1)}(\infty) = \frac{1}{i\hbar} \int_{-\infty}^{\infty} e^{i\omega_{fi}t'} (\lambda \hat{W}_{fi}) f(t') dt'$$

$$= \frac{\lambda \hat{w}_{ji}}{i\hbar} \underbrace{\int_{-\infty}^{+\infty} e^{-i(-\omega_{ji})t'} f(t') dt'} =$$

$$= \frac{\lambda \hat{w}_{ji}}{i\hbar} \text{FT} \{ f(t) \} \big|_{\omega = -\omega_{ji}}$$

$$P_{i \rightarrow j}^{(1)}(\infty) = \frac{|\lambda \hat{w}_{ji}|^2}{\hbar^2} \left| \text{FT} \{ f(t) \} \big|_{\omega = -\omega_{ji}} \right|^2$$

Overview:

Goal in QM:

given  $|\Psi(t_0)\rangle$

Find state  $|\Psi(t)\rangle$  - dynamics

Schrodinger eq. - time evolution.

$\hat{H}$  - Hamiltonian -

$$|\Psi(t)\rangle = \sum_n \langle u_n | \Psi(t_0) \rangle e^{-iE_n(t-t_0)/\hbar} |u_n\rangle$$

Learning QM

I. symbols:  $|\rangle, \langle|, \langle|, \langle|, \langle|, \langle|, \frac{\partial}{\partial t}, \int \dots$

representations & transformations

II. Postulates of QM

Physical Mean

Math symbol.

1. physical states

$$|\Psi\rangle \in \Sigma$$

2. physical observables

operators (Hermitian)

3. States evolve over time

SE

4. Collapse: after meas. state of the system - eigenstate

after  
eigenstate of operators

- J. measure = quantity, - result  
 G. probability of each measurement

eigenstate of the operator  
 projector operator

## II. Examples & techniques

$H(x, p)$  1D  $\rightarrow x, p$  representations.

Specific examples :  $H_{\text{quo}}$ ,  $H_{\text{img}}$ ,  $H_{\text{free}}$ ,  $H_{\text{delta}}$ .

$H(\vec{R}, \vec{P})$  2D, 3D

$H(|\vec{R}|, \vec{P})$  - central potential problem

ex. hydrogen

$\rightarrow a^+, a$

$H(\vec{S}, \vec{L}, \dots)$  AM

spin -

$H = H_0 + W$

approx. methods

$H(|\vec{R}|, \vec{P}^2, \vec{S}, \vec{L}, \vec{I}, \vec{F} \dots)$  hydrogen

$H(t)$