

Assignment 10

OPTI 570 Quantum Mechanics

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Problem I

- a) The probabilities for x and y directions are:

$$P_x = |\langle \phi_x | \phi_\mu \rangle|^2 = \cos^2 \mu, \quad \text{and} \quad P_y = |\langle \phi_y | \phi_\mu \rangle|^2 = \sin^2 \mu.$$

For circular polarization we do the same:

$$\begin{aligned} P_{\sigma_+} &= |\langle \sigma_+ | \phi_\mu \rangle|^2 = \left| -\frac{1}{\sqrt{2}} [(\langle \phi_x | - i \langle \phi_y |)(\cos \mu | \phi_x \rangle + \sin \mu | \phi_y \rangle)] \right|^2 = \frac{1}{2} (\sin^2 \mu + \cos^2 \mu) = \frac{1}{2} \\ P_{\sigma_-} &= 1 - P_{\sigma_+} = \frac{1}{2}. \end{aligned}$$

- b) The state produced by the source is:

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|\phi_x^L\rangle |\phi_y^R\rangle - |\phi_x^R\rangle |\phi_y^L\rangle).$$

The joint measurement is performed through the following projector which is the tensorproduct of the measurement for each photon:

$$P_{jk} = P_j^L \otimes P_k^R = |\phi_j^L\rangle |\phi_k^R\rangle \langle \phi_j^L | \langle \phi_k^R |, \quad j, k \in \{x, y\}.$$

We are let j, k to vary in either direction so that we will have four types of measurements:

$$\begin{aligned} P_{xy} &= \langle \psi | P_{xy} | \psi \rangle = |\langle \phi_x^L \phi_y^R | \psi \rangle|^2 = \frac{1}{2} \\ P_{yx} &= \langle \psi | P_{yx} | \psi \rangle = |\langle \phi_y^L \phi_x^R | \psi \rangle|^2 = \frac{1}{2} \\ P_{xx} &= \langle \psi | P_{xx} | \psi \rangle = |\langle \phi_x^L \phi_x^R | \psi \rangle|^2 = 0 \\ P_{yy} &= \langle \psi | P_{yy} | \psi \rangle = |\langle \phi_y^L \phi_y^R | \psi \rangle|^2 = 0 \end{aligned}$$

Only the non-zero probabilities are possible state the system may be left after the measurement, which are then:

$$|\psi\rangle \xrightarrow{P_{xy}} \frac{P_{xy} |\psi\rangle}{\sqrt{\langle \psi | P_{xy} | \psi \rangle}} = |\phi_x^L \phi_y^R\rangle, \quad \text{and} \quad |\psi\rangle \xrightarrow{P_{yx}} \frac{P_{yx} |\psi\rangle}{\sqrt{\langle \psi | P_{yx} | \psi \rangle}} = |\phi_y^L \phi_x^R\rangle.$$

On the other hand, measurement of a single photon is kind of a marginal measurement. The probabilities can be constructed from the above probabilities:

$$\begin{aligned} P_x^L &= P_{xy} + P_{xx} = \frac{1}{2} + 0 = \frac{1}{2} \\ P_x^R &= P_{xx} + P_{yx} = 0 + \frac{1}{2} = \frac{1}{2} \\ P_y^L &= P_{yx} + P_{yy} = \frac{1}{2} + 0 = \frac{1}{2} \\ P_y^R &= P_{xy} + P_{yy} = \frac{1}{2} + 0 = \frac{1}{2} \end{aligned}$$

c) We first change of basis, by passing from $\phi_{x,y}$ to $\phi_{+,-}$. The transformation matrix is:

$$M = \frac{-1}{\sqrt{2}} \begin{bmatrix} 1 & i \\ 1 & -i \end{bmatrix} \longrightarrow M^\dagger = \frac{-1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ -i & i \end{bmatrix}$$

. By extracting the equation from M^\dagger and substitute them in $\phi_{x,y}$ we find the following equivalent state:

$$|\psi\rangle \propto \frac{1}{\sqrt{2}} (|\phi_+^L \phi_-^R\rangle - |\phi_-^L \phi_+^R\rangle).$$

The measurements are analogous to the previous part, but now in the circular polarized basis, with the exactly same measurement projector, but with j, k changed:

$$P_{jk} = P_j^L \otimes P_k^R = |\phi_j^L\rangle |\phi_k^R\rangle \langle \phi_j^L| \langle \phi_k^R|, \quad j, k \in \{-, +\}.$$

Therefore,

$$\begin{aligned} P_{++} &= |\langle \phi_+^L \phi_+^R | \psi \rangle|^2 = 0 \\ P_{+-} &= |\langle \phi_+^L \phi_-^R | \psi \rangle|^2 = \frac{1}{2} \\ P_{-+} &= |\langle \phi_-^L \phi_+^R | \psi \rangle|^2 = \frac{1}{2} \\ P_{--} &= |\langle \phi_-^L \phi_-^R | \psi \rangle|^2 = 0. \end{aligned}$$

We see the polarization for either basis is exactly the same, but the states after the measurement will be different.

Problem II

In the deuterium atom we have that $I = 1$ and $S = 1/2$.

a) In the state $1s$, we have $n = 1$ and $l = 0$. The quantum number J is:

$$J \in \{|L - S|, |L - S| + 1, \dots, L + S - 1, L + S\} = \left\{ \frac{1}{2} \right\}.$$

For this unique value, we have

$$J = \frac{1}{2} : \quad F \in \{|J - I|, |J - I| + 1, \dots, J + I - 1, J + I\} = \left\{ \frac{1}{2}, \frac{3}{2} \right\}.$$

b) For the state $2p$, we have $n = 2$ and $l = 1$. The quantum number J is:

$$J \in \left\{ \left| 1 - \frac{1}{2} \right|, 1 + \frac{1}{2} \right\} = \left\{ \frac{1}{2}, \frac{3}{2} \right\}.$$

For each J , we have F :

$$\begin{aligned} J = \frac{1}{2} : \quad F &\in \left\{ \left| \frac{1}{2} - 1 \right|, \frac{1}{2} + 1 \right\} = \left\{ \frac{1}{2}, \frac{3}{2} \right\} \\ J = \frac{3}{2} : \quad F &\in \left\{ \left| \frac{3}{2} - 1 \right|, \left| \frac{3}{2} - 1 \right| + 1, \frac{3}{2} + 1 \right\} = \left\{ \frac{1}{2}, \frac{3}{2}, \frac{5}{2} \right\}. \end{aligned}$$

Problem III

We are asked to compute the matrix element of A , but the initial and final states are in the TAM basis, which is not suitable due to the action of A only on $|L, M_L\rangle$; in the TAM basis those are encoded. We need therefore to convert from TAM basis to TP basis each $|\phi\rangle$ and $|\phi'\rangle$. We need to do the following for each state:

$$|F, M_F\rangle \xrightarrow{CG} |J, M_J\rangle |I, M_I\rangle \xrightarrow{CG} |L, M_L\rangle |S, M_S\rangle |I, M_I\rangle.$$

We do that for each state then,

a) We first have that

$$5^2 S_{1/2} \implies n = 5, S = \frac{1}{2}, L = 0, J = 1/2.$$

We also omit those quantities as are the same during all the computations. We then need to express $|F = 1, M_F = 1\rangle$, and we select $J_1 = I = 3/2, J_2 = J = 1/2$ for convention. To decouple this ket in J and I , we use the coefficient from the table to say that:

$$3/2 \times 1/2 \quad |1, 1\rangle_F = \sqrt{\frac{3}{4}} |1/2, -1/2\rangle_J |3/2, 3/2\rangle_I - \sqrt{\frac{1}{4}} |1/2, 1/2\rangle_J |3/2, 1/2\rangle_I.$$

If we look at the following decomposition in the table, with $J_1 = S = 1/2, J_2 = L = 0$, we won't get anything. However, we can use the conservation of angular momentum $M_J = M_S + M_L$ to select the states that corresponds. For $L = 0$ we only have $M_L = 0$. For $S = 1/2$, we have $M_S = \pm 1/2$ and only the negative one satisfies the conservation of AM. Therefore,

$$M_J = M_S + M_L \quad |1/2, -1/2\rangle_J = |0, 0\rangle_L |1/2, -1/2\rangle_S.$$

We do the same for the remaining ket:

$$M_J = M_S + M_L \quad |1/2, 1/2\rangle_J = |0, 0\rangle_L |1/2, 1/2\rangle_S.$$

Thus, putting all together:

$$|\phi\rangle = |1, 1\rangle_F = \sqrt{\frac{3}{4}} |0, 0\rangle_L |1/2, -1/2\rangle_S |3/2, 3/2\rangle_I - \sqrt{\frac{1}{4}} |0, 0\rangle_L |0, 1/2\rangle_S |3/2, 1/2\rangle_I.$$

We do this process for $|\phi'\rangle$, which is completely analogous to this one.

$$3/2 \times 1/2 \quad |1, 1\rangle_{F'} = \sqrt{\frac{3}{4}} |1/2, -1/2\rangle_{J'} |3/2, 3/2\rangle_{I'} - \sqrt{\frac{1}{4}} |1/2, 1/2\rangle_{J'} |3/2, 1/2\rangle_{I'}$$

In this case we have $J_1 = L = 1$ and $J_2 = S = 1/2$, so that we use the $1 \times 1/2$ table.

$$\begin{aligned} 1 \times 1/2 \rfloor \quad |1/2, -1/2\rangle_{J'} &= \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{2}{3}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'} \\ |1/2, 1/2\rangle_{J'} &= \sqrt{\frac{2}{3}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'} \end{aligned}$$

Putting all together:

$$\begin{aligned} |\phi'\rangle &= |1, 1\rangle_{F'} = \sqrt{\frac{3}{4}} \left[\sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{2}{3}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'} \right] |3/2, 3/2\rangle_{I'} \\ &\quad - \sqrt{\frac{1}{4}} \left[\sqrt{\frac{2}{3}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'} \right] |3/2, 1/2\rangle_{I'} \\ &= \sqrt{\frac{1}{4}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'}|3/2, 3/2\rangle_{I'} - \sqrt{\frac{1}{2}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'}|3/2, 3/2\rangle_{I'} \\ &\quad - \sqrt{\frac{1}{6}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'}|3/2, 1/2\rangle_{I'} + \sqrt{\frac{1}{12}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'}|3/2, 1/2\rangle_{I'}. \end{aligned}$$

Finally, we compute the matrix element considering that only $L' - L = \pm 1$ and $M'_L - M_L = 0$ for z polarization. Also, recall that A only acts on the L -basis:

$$\langle\phi'|A|\phi\rangle = \left[\sqrt{\frac{3}{4}}\sqrt{\frac{1}{4}} - \frac{1}{4} + \sqrt{\frac{1}{12}}\sqrt{\frac{3}{4}} - \sqrt{\frac{1}{12}}\sqrt{\frac{1}{4}} \right] A_{1000} = 0.2889 A_{1000}.$$

Therefore,

$$T = |\langle\phi'|A|\phi\rangle|^2 = 0.0834|A_{1000}|^2.$$

b) For $|\phi\rangle$, $n = 5$, $S = 1/2$, $L = 0$, $J = 1/2$:

$$3/2 \times 1/2 \rfloor \quad |1, 0\rangle_F = \sqrt{\frac{1}{2}}|1/2, -1/2\rangle_J|3/2, 1/2\rangle_I - \sqrt{\frac{1}{2}}|1/2, 1/2\rangle_J|3/2, -1/2\rangle_I.$$

$$\begin{aligned} M_J &= M_L + M_S \rfloor \quad |1/2, -1/2\rangle_J = |0, 0\rangle_L|1/2, -1/2\rangle_S \\ M_J &= M_L + M_S \rfloor \quad |1/2, 1/2\rangle_J = |0, 0\rangle_L|1/2, 1/2\rangle_S \end{aligned}$$

$$|\phi\rangle = |1, 0\rangle_F = \sqrt{\frac{1}{2}}|0, 0\rangle_L|1/2, -1/2\rangle_S|3/2, 1/2\rangle_I - \sqrt{\frac{1}{2}}|0, 0\rangle_L|1/2, 1/2\rangle_S|3/2, -1/2\rangle_I.$$

For $|\phi'\rangle$, $n = 5$, $S = 1/2$, $L = 1$, $J = 1/2$:

$$3/2 \times 1/2 \rfloor \quad |1, 0\rangle_{F'} = \sqrt{\frac{1}{2}}|1/2, -1/2\rangle_{J'}|3/2, 1/2\rangle_{I'} - \sqrt{\frac{1}{2}}|1/2, 1/2\rangle_{J'}|3/2, -1/2\rangle_{I'}.$$

$$\begin{aligned} 1 \times 1/2 \rfloor \quad |1/2, -1/2\rangle_{J'} &= \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{2}{3}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'} \\ 1 \times 1/2 \rfloor \quad |1/2, 1/2\rangle_{J'} &= \sqrt{\frac{2}{3}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'} \end{aligned}$$

$$|\phi'\rangle = |1, 0\rangle_{F'} = \sqrt{\frac{1}{2}} \left[\sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{2}{3}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'} \right] |3/2, 1/2\rangle_{I'} \\ - \sqrt{\frac{1}{2}} \left[\sqrt{\frac{2}{3}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'} \right] |3/2, -1/2\rangle_{I'}$$

$$\langle\phi'|A|\phi\rangle = \left[\sqrt{\frac{1}{6}}\sqrt{\frac{1}{2}} - \sqrt{\frac{1}{6}}\sqrt{\frac{1}{2}} + \sqrt{\frac{1}{6}}\sqrt{\frac{1}{2}} - \sqrt{\frac{1}{6}}\sqrt{\frac{1}{2}} \right] A_{1000} = 0 A_{1000} \implies T = |\langle\phi'|A|\phi\rangle|^2 = 0.$$

c) For $|\phi\rangle$, $n = 5$, $S = 1/2$, $L = 0$, $J = 1/2$:

$$|\phi\rangle = |1, 1\rangle_F = \sqrt{\frac{3}{4}}|0, 0\rangle_L|1/2, -1/2\rangle_S|3/2, 3/2\rangle_I - \sqrt{\frac{1}{4}}|0, 0\rangle_L|0, 1/2\rangle_S|3/2, 1/2\rangle_I.$$

For $|\phi'\rangle$, $n = 5$, $S = 1/2$, $L = 1$, $J = 1/2$:

$$|3/2 \times 1/2\rangle \quad |2, 1\rangle_{F'} = \sqrt{\frac{1}{4}}|1/2, -1/2\rangle_{J'}|3/2, 3/2\rangle_{I'} + \sqrt{\frac{3}{4}}|1/2, 1/2\rangle_{J'}|3/2, 1/2\rangle_{I'}.$$

$$|1 \times 1/2\rangle \quad |1/2, -1/2\rangle_{J'} = \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{2}{3}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'} \\ |1 \times 1/2\rangle \quad |1/2, 1/2\rangle_{J'} = \sqrt{\frac{2}{3}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'}.$$

$$|\phi'\rangle = |2, 1\rangle_{F'} = \sqrt{\frac{1}{4}} \left[\sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{2}{3}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'} \right] |3/2, 3/2\rangle_{I'} \\ + \sqrt{\frac{3}{4}} \left[\sqrt{\frac{2}{3}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'} \right] |3/2, 1/2\rangle_{I'}$$

$$\langle\phi'|A|\phi\rangle = \left[\sqrt{\frac{1}{12}}\sqrt{\frac{3}{4}} - \sqrt{\frac{1}{12}}\sqrt{\frac{1}{4}} - \sqrt{\frac{1}{3}}\sqrt{\frac{3}{4}} + \sqrt{\frac{1}{3}}\sqrt{\frac{1}{4}} \right] A_{1000} = -0.1056 A_{1000} \implies 0.0111|A_{1000}|^2.$$

d) For $|\phi\rangle$, $n = 5$, $S = 1/2$, $L = 0$, $J = 1/2$:

$$|1, 0\rangle_F = \sqrt{\frac{1}{2}}|0, 0\rangle_L|1/2, -1/2\rangle_S|3/2, 1/2\rangle_I - \sqrt{\frac{1}{2}}|0, 0\rangle_L|1/2, 1/2\rangle_S|3/2, -1/2\rangle_I.$$

For $|\phi'\rangle$, $n = 5$, $S = 1/2$, $L = 1$, $J = 1/2$:

$$|3/2 \times 1/2\rangle \quad |2, 0\rangle_{F'} = \sqrt{\frac{1}{2}}|1/2, -1/2\rangle_{J'}|3/2, 1/2\rangle_{I'} + \sqrt{\frac{1}{2}}|1/2, 1/2\rangle_{J'}|3/2, -1/2\rangle_{I'}.$$

$$1 \times 1/2] \quad |1/2, -1/2\rangle_{J'} = \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{2}{3}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'}$$

$$1 \times 1/2] \quad |1/2, 1/2\rangle_{J'} = \sqrt{\frac{2}{3}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'}.$$

$$|\phi'\rangle = |2, 0\rangle_{F'} = \sqrt{\frac{1}{2}} \left[\sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{2}{3}}|1, -1\rangle_{L'}|1/2, 1/2\rangle_{S'} \right] |3/2, 1/2\rangle_{I'}$$

$$+ \sqrt{\frac{1}{2}} \left[\sqrt{\frac{2}{3}}|1, 1\rangle_{L'}|1/2, -1/2\rangle_{S'} - \sqrt{\frac{1}{3}}|1, 0\rangle_{L'}|1/2, 1/2\rangle_{S'} \right] |3/2, -1/2\rangle_{I'}.$$

$$\langle\phi'|A|\phi\rangle = \left[\sqrt{\frac{1}{6}}\sqrt{\frac{1}{2}} - \sqrt{\frac{1}{6}}\sqrt{\frac{1}{2}} - \sqrt{\frac{1}{6}}\sqrt{\frac{1}{2}} + \sqrt{\frac{1}{6}}\sqrt{\frac{1}{2}} \right] A_{1000} = 0A_{1000} \implies T = 0.$$