

IV Probability postulate

$$A \leftrightarrow \hat{A}, \quad |\psi\rangle \in \mathcal{E}$$

$$P(a_m) = \langle \psi | \hat{P}_{a_m} | \psi \rangle$$

Cases: (i) discrete, non-degenerate

$$\hat{P}_{a_m} = |a_m\rangle\langle a_m|$$

$$P(a_m) = \langle \psi | a_m \rangle \langle a_m | \psi \rangle =$$

$$= \underbrace{|\langle a_m | \psi \rangle|^2}_{\text{expansion coefficients of state } |\psi\rangle \text{ in the eigenbasis.}}$$

$$|\psi\rangle = \sum_n \hat{P}_{a_n} |\psi\rangle = \sum_n |a_n\rangle \langle a_n | \psi \rangle =$$

$$= \sum_n \underbrace{\langle a_n | \psi \rangle}_{\text{expansion coefficients.} \Rightarrow \text{probability amplitudes}} |a_n\rangle$$

(ii) discrete, degenerate:  $\hat{A} |a_m^i\rangle = \underline{a_m} |a_m^i\rangle$

$$\hat{P}(a_m) = \sum_{i=1}^{g_m} |a_m^i\rangle\langle a_m^i|$$

$$\Rightarrow \underline{P}(a_m) = \sum_{i=1}^{g_m} |\langle a_m^i | \psi \rangle|^2$$

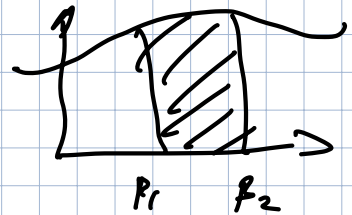
(iii) Continuous spectrum  $\hat{B}|\beta\rangle = \beta|\beta\rangle$

$$\hat{P}_{\beta_1 \leq \beta < \beta_2} = \int_{\beta_1}^{\beta_2} |\beta\rangle\langle\beta| d\beta$$

$$P(\beta_1 \leq \beta < \beta_2) = \langle \psi | \hat{P} | \psi \rangle = \int_{\beta_1}^{\beta_2} \langle \psi | \beta \rangle \langle \beta | \psi \rangle d\beta =$$

$$P(\beta_1 \leq \beta < \beta_2) = \int_{\beta_1}^{\beta_2} |\langle \beta | \psi \rangle|^2 d\beta$$

probability density



Normalization:  $\int_{\text{all } \beta} dP = 1$

$$dP = |\langle \beta | \psi \rangle|^2 d\beta$$

$dP$  - prob of meas. between  $\beta$  and  $\beta+d\beta$

Q:  $\psi(x)$  meas  $x_0$ ?

prob of meas particle between  $x_0$  and  $x_0+dx$  for small  $dx$ ?

$$|\psi(x_0)|^2 dx$$

(iv) continuous, but degenerate spectrum.

ex: 3D Space coordinates in  $\Sigma$  basis?

$$|x, y, z\rangle, \{ |x\rangle, |y\rangle, |z\rangle \}$$

Formally: -tensor product:  $\Sigma = \Sigma_x \otimes \Sigma_y \otimes \Sigma_z$

$$\Rightarrow \{ |x\rangle \otimes |y\rangle \otimes |z\rangle \}$$

infinity.

closure relation:

$$1 = \int_{\text{all } x, y, z} |x, y, z\rangle \langle x, y, z| dx dy dz.$$

$$\hat{P}_{x_1 \leq x < x_2} = \int_{x_1}^{x_2} dx \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dz \underbrace{|x, y, z\rangle \langle x, y, z|}_{1(x) 1(y) 1(z)}$$

Colloppse:

$$\hat{A} |a_m\rangle = a_m |a_m\rangle \quad \text{state } |\psi\rangle \in \mathcal{E}$$

$$\boxed{\text{I}} \quad P(a_m) = \hat{P}_{a_m} |\psi\rangle$$

$$\boxed{\text{II}} \quad \text{re-normalize}$$

Consequences of postulates

1. Constant global phase factors have no physical significance.

$$\begin{aligned} \text{ex: } |\psi\rangle &= c_1 e^{i\theta_1} |u_1\rangle + c_2 e^{i\theta_2} |u_2\rangle = \\ &= \underbrace{e^{i\theta_1}}_{\substack{\text{global phase} \\ \text{no physical}}} \left[ c_1 |u_1\rangle + c_2 \underbrace{e^{i(\theta_2 - \theta_1)}}_{\substack{\text{relative phase} \\ \text{factor}}} |u_2\rangle \right] \end{aligned}$$

$$\begin{aligned} \text{ex: } |\psi\rangle &= \frac{e^{i\bar{a}/5}}{\sqrt{2}} |u_1\rangle + \frac{e^{-i\bar{a}/5}}{\sqrt{2}} |u_2\rangle = \\ &= \cancel{e^{i\bar{a}/5}} \cdot \left( \frac{1}{\sqrt{2}} |u_1\rangle - \frac{1}{\sqrt{2}} |u_2\rangle \right) \end{aligned}$$

$$\psi'(x) = e^{i\theta} \psi(x)$$

$$|\psi'(x)|^2 = \psi'^*(x) \psi'(x) = e^{-i\theta} \psi^*(x) e^{i\theta} \psi(x) =$$

$$\boxed{|\psi'(x)|^2 = |\psi(x)|^2}$$

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = \hat{H}(t) |\psi(t)\rangle \Leftrightarrow i\hbar \frac{d}{dt} (e^{i\theta} |\psi(t)\rangle) = \hat{H}(t) e^{i\theta} |\psi(t)\rangle$$

2. Measurement of a large # of measurements:

weighted sum of the possible outcomes:

eigenvalues of  $\hat{A}$ :  $a_m$

$$\begin{aligned} \underbrace{\langle \hat{A} \rangle}_{\text{expectation value}}_{\psi} &= \sum_{\text{all } m} a_m P(a_m) = \\ &= \sum_m a_m \langle \psi | \hat{P}_{a_m} | \psi \rangle = \\ &= \sum_m a_m \langle \psi | a_m \times a_m | \psi \rangle = \\ &= \sum \langle \psi | \underbrace{a_m}_{\text{eigenvalue}} \underbrace{a_m}_{\text{eigenvalue}} | \psi \rangle = \\ &= \sum \langle \psi | \hat{A} | a_m \times a_m | \psi \rangle = \\ &= \langle \psi | \hat{A} | \underbrace{\sum_m a_m \times a_m}_{=1} | \psi \rangle = \\ \boxed{\langle \hat{A} \rangle} &= \langle \psi | \hat{A} | \psi \rangle \end{aligned}$$

expectation value  $\neq$  expected value

3. Expected spread of measurement results

$$\begin{aligned}\hat{\Delta A} &= \sqrt{\langle \hat{A}^2 \rangle - \langle A \rangle^2} = \\ &= \sqrt{\langle \Psi | \hat{A}^2 | \Psi \rangle - (\langle \Psi | A | \Psi \rangle)^2}\end{aligned}$$

Ex:  $(\hat{\Delta A})(\hat{\Delta B}) \geq \frac{1}{2} |\langle [\hat{A}, \hat{B}] \rangle|$

ex:  $\hat{x}, \hat{p}$   $\hat{\Delta x} \hat{\Delta p} \geq \frac{\hbar}{2}$

Q:  $\hat{H} |u_n\rangle = \underline{E_n} |u_n\rangle$

$\Delta H$  for  $|u_n\rangle = ?$   $\Delta H = 0$  because  $E_n$

$\hat{\Delta A} = 0$  for any  $|u_n\rangle$  always.

language:

• Compatible observable  $\Rightarrow$  CSCO

$\hat{A} \wedge \hat{B} \Leftrightarrow [\hat{A}, \hat{B}] = 0 \Leftrightarrow$  share set of eigenvectors

ex: meas  $x$  and then  $p$   $[\hat{x}, \hat{p}] = i\hbar \Rightarrow$  loss information

Cons. measurement w/ compatible observables - no loss of information.

• Bohr frequencies : differences b/w eigenvectors of  $\hat{H}$

$\gamma_{nm} = (E_n - E_m) / \hbar$  cycle freq.  $\text{Hz.}$   
 $\omega_{nm} = (E_n - E_m) / \hbar$  ang. freq.  $\frac{\text{rad}}{\text{s}}$

ANO:  $\omega = 2\pi \gamma = \frac{\text{rad}}{\text{s}}$

## Time - evolution operator

$$|\psi(t_0)\rangle \rightarrow |\psi(t)\rangle$$

$$|\psi(t)\rangle = \hat{U}(t, t_0) |\psi(t_0)\rangle$$

$$\hat{U}(t, t_0) = \text{(i) } \hat{H} \text{ is time-independent.}$$

$$\hat{U}(t, t_0) = e^{-i\hat{H}(t-t_0)/\hbar}$$

$$\text{(ii) } \hat{H} \text{ is time-dependent but } [\hat{H}(t), \hat{H}(t')] = 0$$

then

$$\hat{U}(t, t_0) = e^{-i \int_{t_0}^t dt' \hat{H}(t')/\hbar}$$

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