

## Subsystem stabilizer codes

Start again with two sets of  $n$  hyperbolic pairs in  $\mathcal{P}_n$ :

$$\begin{array}{ll}
 Z_1 \longleftrightarrow X_1 & \bar{Z}_1 \longleftrightarrow \bar{X}_1 \\
 Z_2 \longleftrightarrow X_2 & \bar{Z}_2 \longleftrightarrow \bar{X}_2 \\
 \vdots & \vdots \\
 Z_k \longleftrightarrow X_k & \bar{Z}_k \longleftrightarrow \bar{X}_k \\
 \left. \begin{array}{l} \text{n-k stabilizers} \\ \text{for generators} \\ Q = \{ | \phi \rangle \langle 0 |, | \phi \rangle \langle e^k | \} \end{array} \right\} & \left. \begin{array}{l} \text{n-k stabilizers} \\ \text{for some other Qs} \\ g_1 \longleftrightarrow h_1 \\ \vdots \\ g_{n-k} \longleftrightarrow h_{n-k} \end{array} \right\} \\
 Z_{k+1} \longleftrightarrow X_{k+1} & \\
 \vdots & \\
 Z_n \longleftrightarrow X_n & 
 \end{array}$$

There is a unitary  $U$  s.t.  $U \underbrace{\left( \begin{smallmatrix} \text{operator in} \\ \text{first set} \end{smallmatrix} \right)}_{Q} U^\dagger = \text{operator in 2nd set}$

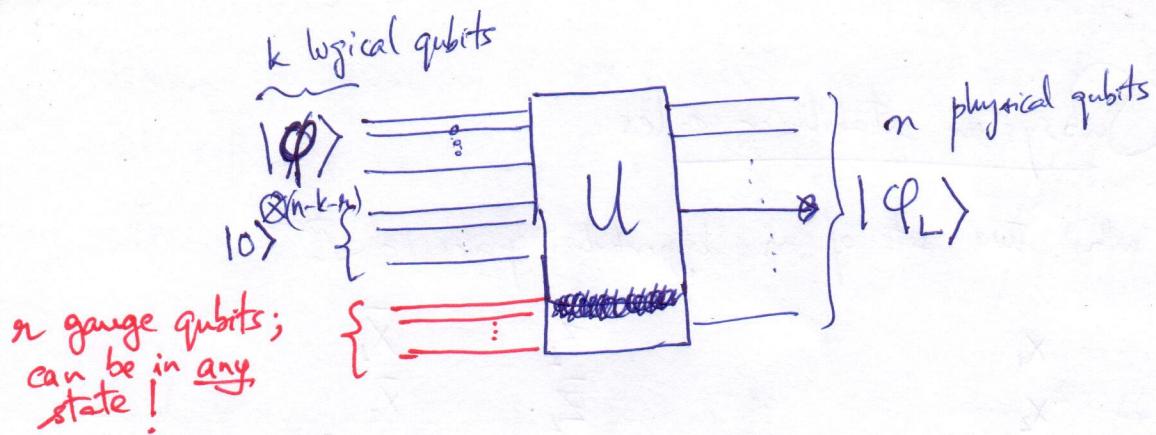
$$\begin{array}{ccc}
 | \phi \rangle \langle 0 |^{(n-k)} & \xrightarrow{U} & | \phi \rangle \langle 0 | \\
 \hat{Q} & & \hat{Q}_S
 \end{array}$$

Then, for an arbitrary  $E \in \mathcal{P}_n$ , we can write it as  $(\cancel{c} \mod \langle \omega \rangle)$

$$E = g(E) h(E) L(E)$$

We convert some of the stabilizers on either side into gauge operators.

$$\begin{array}{ccccc}
 Z_1 & X_1 & & \bar{Z}_1 & \bar{X}_1 \\
 \vdots & \vdots & & \vdots & \vdots \\
 Z_k & X_k & & \bar{Z}_k & \bar{X}_k \\
 \left. \begin{array}{l} \text{n-k-n} \\ \text{stabilizer} \\ \text{generators} \end{array} \right\} & \left. \begin{array}{l} X_{k+1} \\ \vdots \\ X_{n-n} \end{array} \right\} & \xleftarrow{\text{Unitary } U} & \left. \begin{array}{l} \text{n-k-n} \\ \text{stabilizer} \\ \text{generators} \end{array} \right\} & \left. \begin{array}{l} \bar{h}_1 \\ \vdots \\ \bar{h}_{n-k-n} \end{array} \right\} \\
 \left. \begin{array}{l} 2r \\ \text{gauge} \\ \text{operators} \end{array} \right\} & \left. \begin{array}{l} Z_{n-r+1} \\ \vdots \\ Z_n \end{array} \right\} & & \left. \begin{array}{l} Z_{n-r+1} \\ \vdots \\ Z_n \end{array} \right\} & \left. \begin{array}{l} \bar{h}_{n-k+1} \\ \vdots \\ \bar{h}_{n-k} \end{array} \right\} \\
 & & & & \left. \begin{array}{l} h_{n-k+1} \\ \vdots \\ h_{n-k} \end{array} \right\} \\
 & & & & \left. \begin{array}{l} G_1^z \\ \vdots \\ G_n^z \end{array} \right\} \quad \left. \begin{array}{l} \bar{G}_1^z \\ \vdots \\ \bar{G}_n^z \end{array} \right\} \quad \left. \begin{array}{l} h_{n-k+1} \\ \vdots \\ h_{n-k} \end{array} \right\} \quad \left. \begin{array}{l} G_1^x \\ \vdots \\ G_n^x \end{array} \right\} \quad \left. \begin{array}{l} \bar{G}_1^x \\ \vdots \\ \bar{G}_n^x \end{array} \right\} \\
 & & & & \left. \begin{array}{l} \text{2r} \\ \text{gauge} \\ \text{operators} \end{array} \right\} \quad \left. \begin{array}{l} \text{2r} \\ \text{gauge} \\ \text{operators} \end{array} \right\} \quad \left. \begin{array}{l} \text{2r} \\ \text{gauge} \\ \text{operators} \end{array} \right\} \quad \left. \begin{array}{l} \text{2r} \\ \text{gauge} \\ \text{operators} \end{array} \right\} \quad \left. \begin{array}{l} \text{2r} \\ \text{gauge} \\ \text{operators} \end{array} \right\}
 \end{array}$$



The ~~sub~~ space on the left (input end) of the encoder  $U$  is of the form  $\underbrace{A}_{|φ\rangle} \otimes \underbrace{B}_{|0\rangle \otimes(n-k-n)}$   $n_1$ -qubit space.

The logical qubits come from subsystem  $A$ ; hence the term "subsystem codes".

Now, any Pauli error  $E$  can be decomposed as

$$E = \underbrace{g(E)}_{\text{from the gauge operators}} \underbrace{h(E)}_{\text{from the logical operators}} \underbrace{L(E)}_{\substack{\text{here come from} \\ g_1 \dots g_{n-k-n} \\ h_1 \dots h_{n-k-n}}} \underbrace{g(E)}_{\substack{\text{from the gauge operators} \\ G_1^x \dots G_{n-k-n}^x \\ G_1^z \dots G_{n-k-n}^z}}$$

The syndrome  $s = (s_1, \dots, s_{n-k-n})$  of  $E$  is obtained by measuring the stabilizer generators  $g_1 \dots g_{n-k-n}$ . This, as before, determines  $h(E) = \prod_{i=1}^{n-k-n} h_i^{s_i}$ .

Again, as before,  $g(E)$  need not be recovered exactly, as this part of  $E$  has no effect on the encoded state.

But now, even  $L(E)$  need not be recovered exactly, as this too has no effect on the logical qubits encoded!

To see why:

Applying, say,  $G_i^x$  to  $|φ_L\rangle$ , an encoding of the  $k$  logical qubits:

For any  $1 \leq i \leq n$

$$\begin{aligned} G_i^x |φ_L\rangle &= U \otimes_{n+1-i} U^\dagger \cdot U (|φ\rangle |0\rangle \otimes_{n-k-n} |φ\rangle) \\ &= U (|φ\rangle |0\rangle \otimes_{n-k-n} X_{n+1-i} |φ\rangle) \end{aligned}$$

which is also an encoding of  $|φ\rangle$ .

Removing some of the stabilizers implies that less information about the error  $E$  is available to the decoder.

However, this loss is compensated by an increase in the code's degeneracy,

meaning ~~there~~ there are more errors that do not affect the <sup>encoded</sup> qubits.

Also, the min. dist. is unaffected as the logical operators do not change.