

# Comparing numerical results with experimental data

February 24, 2020

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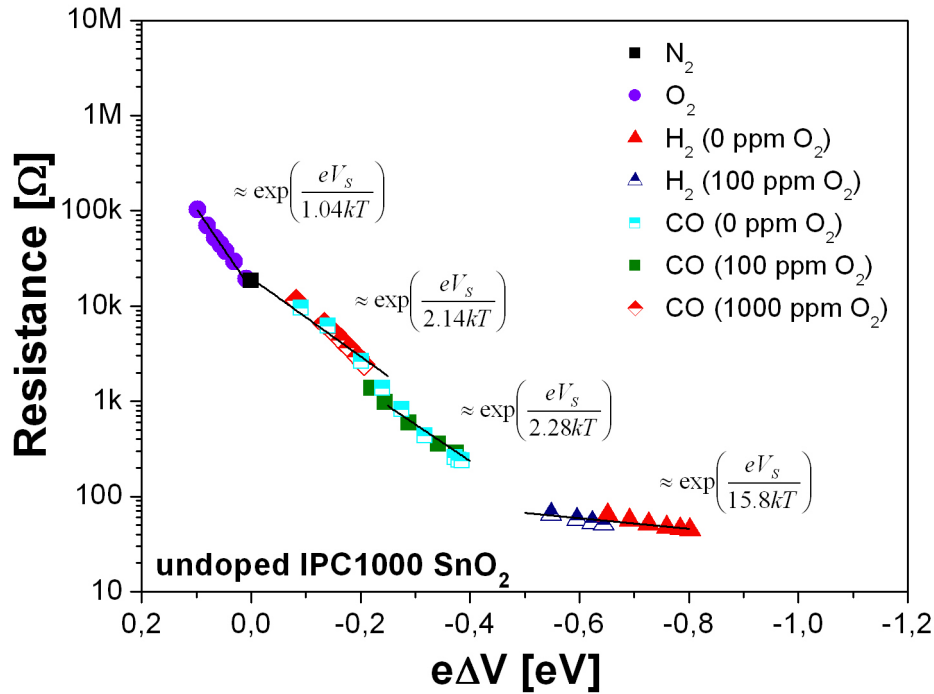
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## 1 Abstract

Experimental data from simultaneous work function and resistance measurements will be compared with the results from the numerical calculations. Results from an  $\text{SnO}_2$  gas sensor measured at  $300^\circ\text{C}$  will be used to demonstrate, how numerical data can be used to gain more insights about the measured material. The chosen dataset the the graphical representation originates from the Phd. thesis of Julia Rebholz: [Reb16].

The data was generated by exposing the sensor to various gas compositions of  $\text{H}_2$ ,  $\text{CO}$ ,  $\text{O}_2$  and  $\text{N}_2$ . The surface potential changes  $\Delta V$  resulting from the different gas atmosphere have been obtained with the Kelvin probe technique. Simultaneously the corresponding resistance was measured. The data point point at  $0e\Delta V$  corresponds to the situation in nitrogen.

These experimental data points will be compared with the results obtained from the numerical model.



This figure shows the dependency of the resistance on the band bending changes for the undoped  $\text{SnO}_2$  nanopowder based gas sensor.  $e\Delta = 0$  is denoted as the situation in  $\text{N}_2$

## 2 Fitting experimental data to numerical results

### 2.1 Importing experimental and numerical data

The data is saved in an excel file, which will be loaded by using the tools provided by pandas.

```
[3]: #Setting up the env.
from part2 import *
import pandas as pd
%pylab inline

#importing the data
calc_dF = pd.read_hdf('numerical_sol.h5','raw')
dF_1000 = pd.read_excel('Kelvin_Data.xlsx', sheet_name='ipc1000').
    ↳sort_values(by='dV')

#instead of unsing the row number
#each row has the value of dV as index
dF_1000.index = dF_1000['dV']
```

Populating the interactive namespace from numpy and matplotlib

## 2.2 Representing the raw data

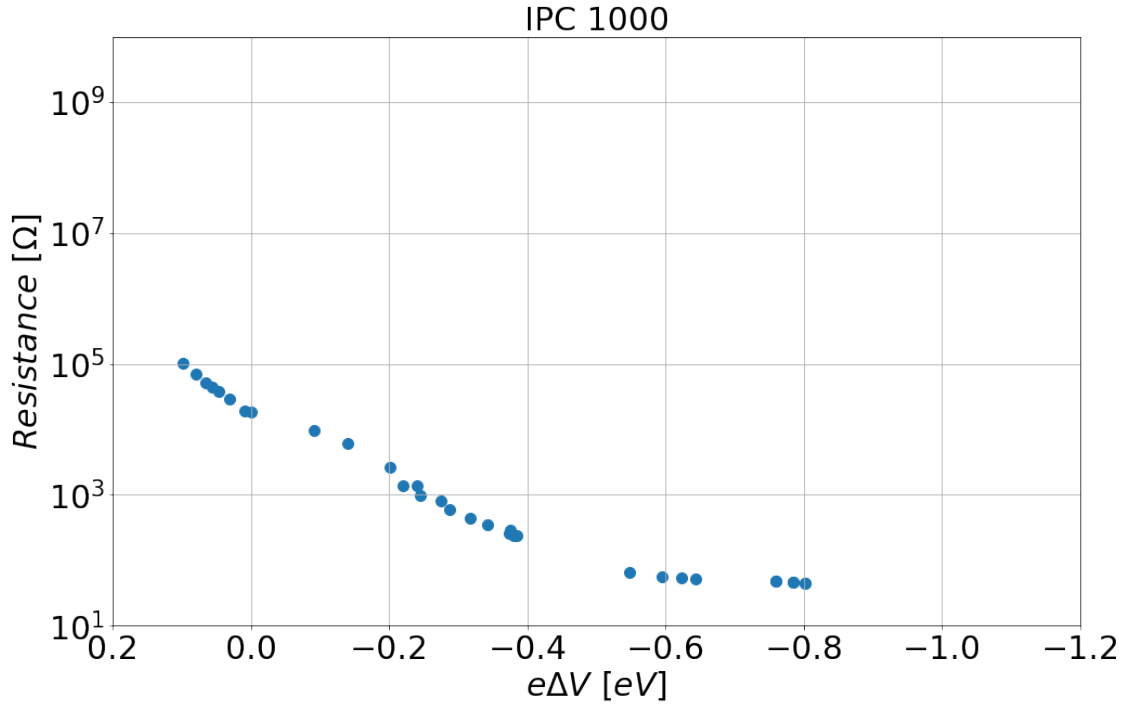
```
[14]: def format_axe(axe, ylabel = None, set_ylim=False):
    labelsz = 30
    if set_ylim:
        axe.set_ylim((1e-4,1e3))
    axe.set_yscale('log')
    axe.set_xlim((0.2,-1.2))
    if ylabel:
        axe.set_ylabel(ylabel, fontsize = labelsz)
    else:
        axe.set_ylabel(r'$\frac{R_{V_S}}{R_{(V_S=0)}}$', fontsize = labelsz)
    axe.set_xlabel('$e\Delta V$ [eV]', fontsize = labelsz)
    axe.tick_params(axis='both', which='both', labelsz=labelsz)
    axe.grid()

fig, axe = subplots(figsize=(16,10))

sens, dF = 'IPC 1000',dF_1000

v_exp = dF['dV']

res_exp = dF['res']
axe.set_title(sens, fontsize = 30)
axe.scatter(v_exp,res_exp, s=100)
format_axe(axe,ylabel='$Resistance$ [$\Omega$]')
axe.set_ylim(res_exp.min()/2,res_exp.max()*2);
axe.set_ylim(10,10e9);
```



This figure represents the experimental data as shown before plotted in the Python environment. The dependency of the resistance on the band bending changes is shown for the undoped  $\text{SnO}_2$  nanopowder based gas sensor.  $q\Delta = 0$  is denoted as the situation in  $\text{N}_2$ .

### 2.3 From $R_{V_s}$ to $\Delta R_{V_s}$

In the experimental dataset the value at  $0qV_s$  represent the data points measured under nitrogen. Therefore  $\Delta R_{V_s} = \frac{R_{V_s}}{R_0}$  is calculated by :

- First derive the resistance under nitrogen  $R_0$
- Second divide all resistance values by this value

```
[16]: from scipy.optimize import curve_fit
      from scipy.interpolate import interp1d

      fig, axe = subplots(1, figsize=(16,10))

      #get the value of the flatband (if needed)
      #by interpolation
      interp_res = interp1d(v_exp,res_exp)
      res_flatband = interp_res(0)

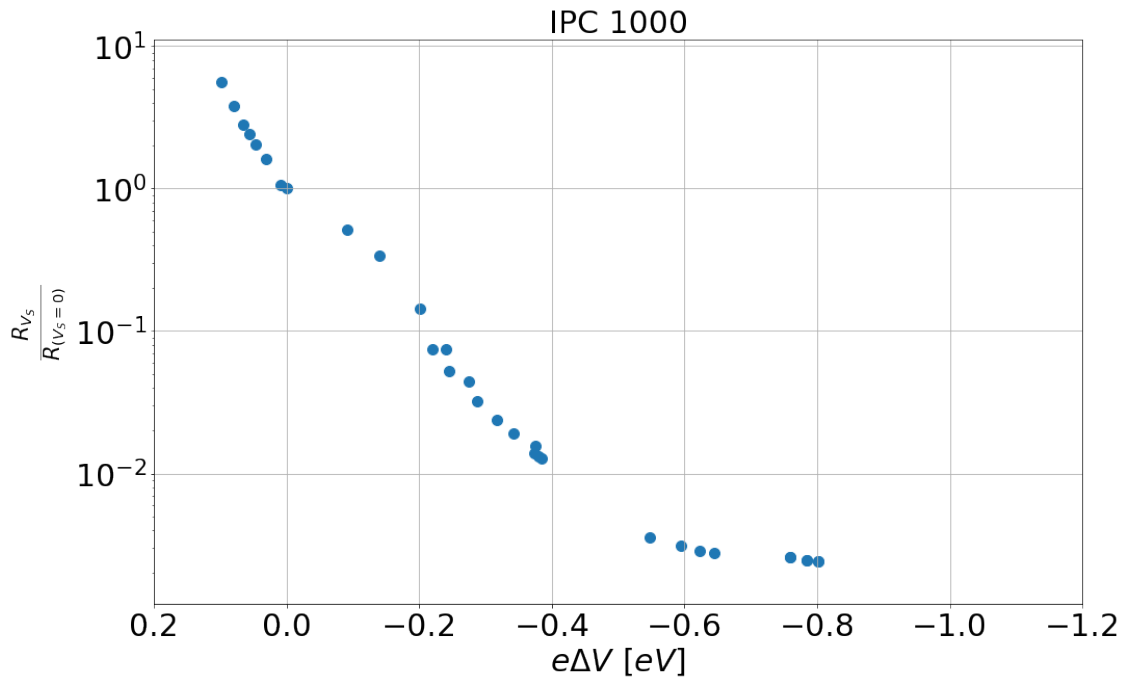
      #calculate the rel. res change
      rel_res_exp = dF['res']/res_flatband
```

```

#represent it
format_axe(axe)
axe.scatter(dF['dV'],rel_res_exp, s=100)
axe.set_ylim(rel_res_exp.min()/2,
             rel_res_exp.max()*2);
axe.set_title(sens, fontsize = 30)

```

```
[16]: Text(0.5, 1.0, 'IPC 1000')
```



The resistances have been normalized to the resistance at  $q\Delta = 0$

## 2.4 Interpolating the numerical values

In the previous section, the numerical solution for multiple start parameters have been calculated. Nevertheless most probably the calculated dataset will not hold exactly the same values gathered from the experiment. To obtain the numerical value for a specific experimental value of  $qV_s$  a interpolation between of the existing numerical values will be used again. For all different numerical grains  $\Delta R_{V_s}$  will be calculated for all the experimental values of  $qV_s$ . Once this is done, the different between the numerical model and the exp. data can be calculated and evaluated.

```

[17]: #The dataframe to hold the different
      #of the exp. values to the numerical ones
      #Will be used to find the best fitting num. solution

```

```

num_data_at_exp_pos_dF = pd.DataFrame(index = v_exp)

#group the num. data by its paramters (T, R and ND)
data_by_grain = calc_dF.groupby(['temp', 'R', 'ND'])

for (T, R, ND), calc_dF_grain in data_by_grain:

    num_data_at_exp_pos_dF[(T, R, ND)] = None

    grain = create_grain_from_data(calc_dF_grain)

    flat_band_data = calc_dF_grain[calc_dF_grain['Einit_kT']==0].iloc[0]

    rel_res_num = calc_dF_grain['rel_res_change']

    #express the surace potential in eV
    #to be comparable with the exp. data
    v_num = calc_dF_grain['Einit_kT']*CONST.J_to_eV(grain.material.kT)

    #use interpolation to get the values for the positions
    #of the experiment data points
    interp_rs_num = interp1d(v_num, rel_res_num, bounds_error=False)
    interp_v_num = interp1d(rel_res_num, v_num, bounds_error=False)

    #caculate the numerical value of rel. res at the position
    # of V from the experiment
    res_num_at_exp_pos = interp_rs_num(v_exp)

    #save those values in the new DataFrame
    num_data_at_exp_pos_dF.loc[:, (T, R, ND)] = res_num_at_exp_pos

```

## 2.5 Calculating the fit error

num\_data\_at\_exp\_pos\_dF contains now the values of  $\Delta R_{V_S}$  at the positions  $qV_S$ . From these values the relative error needs to be calculated. The following formula is used to derive the error:

$$\epsilon_{V_S} = \left( \frac{R_{numerical}(qV_S) - R_{experiment}(qV_S)}{R_{experiment}(qV_S)} \right)^2 \quad (1)$$

The sum of all  $\epsilon_{V_S}$  is the total error of the fit. The numerical model with the lowest value of  $\sum \epsilon_{V_S}$  is the model which fits best to the experimental data. The average grain diameter of the material “IPC1000” is known to be in average radius of 55nm (diameter of 110nm) (Nanoparticle engineering for gas sensor optimization: Improved sol-gel fabricated nanocrystalline SnO<sub>2</sub> thick film gas sensor for NO<sub>2</sub> detection by calcination, catalytic metal introduction and grinding treatments, 1999). The dataset we created in the previous section includes models for grains with radii of

50nm and 100nm. Therefore we can narrow the fit algorithm down, to take only models with a radius of 50nm and 100nm in account

```
[25]: abs_error = num_data_at_exp_pos_dF.subtract(rel_res_exp, axis='index')
rel_error = abs_error.divide(rel_res_exp, axis='index')
rel_error_square = rel_error**2
sum_of_squares = rel_error_square.sum()

valid_index = [i for i in sum_of_squares.index if i[1] in [50e-9,100e-9]]

sum_of_squares_grainsize = sum_of_squares.loc[valid_index].sort_values()

grain_min_error_tuple = sum_of_squares_grainsize.idxmin()
display(pd.DataFrame({'error':sum_of_squares_grainsize}))
```

	error
(300.0, 5e-08, 1e+22)	6.822288
(300.0, 1e-07, 1e+22)	7.665906
(300.0, 1e-07, 1e+21)	14.358742
(300.0, 5e-08, 1e+21)	18.771491
(300.0, 5e-08, 1e+23)	175.204658
(300.0, 1e-07, 1e+23)	1256.744498
(300.0, 5e-08, 1e+24)	9905.878196
(300.0, 1e-07, 1e+24)	39709.781227

## 2.6 Representation of the fit

Finally the best fit results can be represented graphically.

```
[24]: fig, axe = subplots(figsize = (16,10))
#for grain_tuple in num_data_at_exp_pos_dF.keys():
for grain_tuple in sum_of_squares.index:
    if grain_tuple == grain_min_error_tuple:
        linestyle = '*-'
        linewidth = 5
        alpha = 0.5
        label = 'Best Fit'
    elif grain_tuple in sum_of_squares_grainsize.index[0:2]:
        linestyle = '-o'
        linewidth = 5
        alpha = 0.3
        label = 'Second best fit'
    else:
        linestyle = '-.'
        linewidth = 1
        alpha = 0.3
        label = 'Other solution'
```

```

    ax.plot(num_data_at_exp_pos_dF.index,
            num_data_at_exp_pos_dF[grain_tuple],
            linestyle, linewidth=linewidth, alpha = alpha,
            label =label)

    last_x  = num_data_at_exp_pos_dF.index[0]
    last_y  = num_data_at_exp_pos_dF.iloc[0][grain_tuple]

    if grain_tuple in sum_of_squares_grainsize.index[0:2]:
        ax.text(last_x-0.05,last_y,
                f'''Radius:{grain_tuple[1]*1e9:.0f}nm\n$N_D$:{grain_tuple[2]:.2} 1/
    ↪m3''' )
format_axe(ax)

ax.scatter(rel_res_exp.index,
          rel_res_exp,
          s=100,
          label = 'Exp. data'
          )

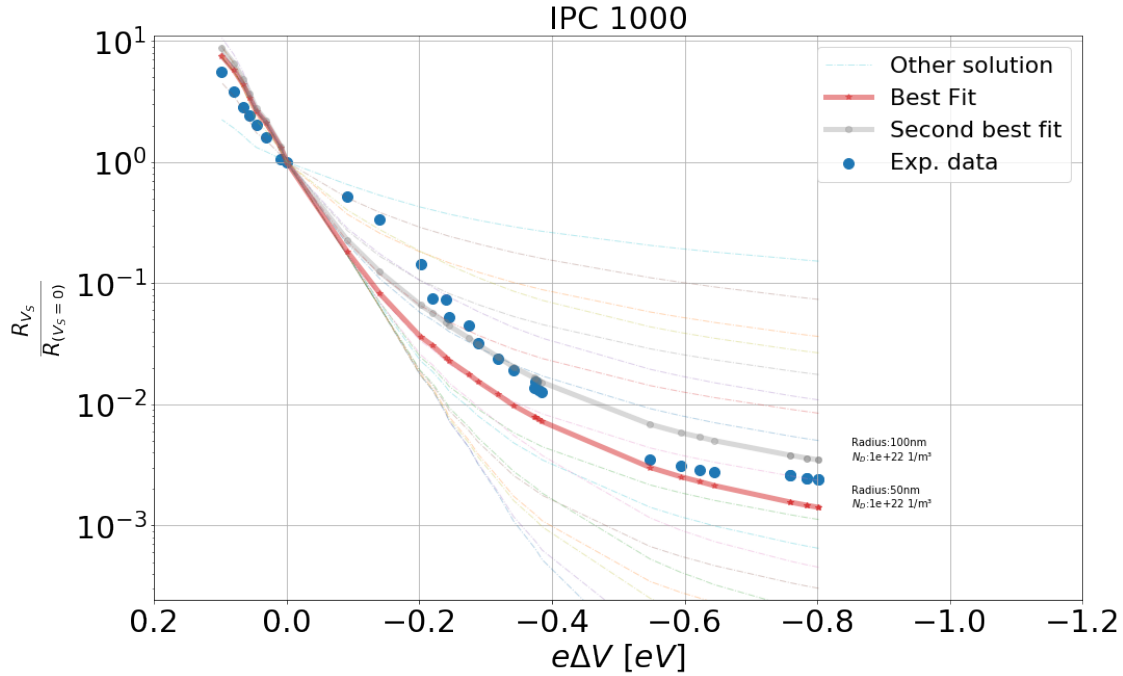
ax.set_ylim(rel_res_exp.min()/10,
            rel_res_exp.max()*2);

l = {h[1]:h[0] for h in zip(*ax.get_legend_handles_labels())}.keys()
h = {h[1]:h[0] for h in zip(*ax.get_legend_handles_labels())}.values()
ax.legend(h,l,loc=1, fontsize = 22)
ax.set_title(sens, fontsize = 30)

```

[24]: Text(0.5, 1.0, 'IPC 1000')





Comparison of the numerical results with the experimental data points

### 3 Conclusion

We can see, that both fits do not fit perfectly to the experimental data. One obvious reason is the coarse screening of the grain size and  $N_D$ . A second iteration of creating models with radii between 50nm and 100nm might find a better fit. Also a finer screening of  $N_D$  will turn out to be helpful for a better result. On the other side the fitting shows, that the experimental data fits will to a grain with approximately 50nm radius and a defect concentration of around  $N_D = 1 * 10^{22} 1/m$  at 300°C.

### 4 Bibliography section

#### References

- [Reb16] REBHOLZ, Julia M.: *Influence of Conduction Mechanism Changes and Related Effects on the Sensing Performance of Metal Oxide Based Gas Sensors*. Shaker Verlag, 2016. – 127 S. – ISBN 978-3-8440-4832-2

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