

From smox-grains to resistance

November 28, 2019

Contents

1	Review	1
2	Load the results	1
3	From charge distribution to resistance	3
3.1	The “numerical” grain	4
3.2	Precalc the numerical grains for all conditions	6
4	Relaxation	7
5	Bibliography section	11

1 Review

In the last notebook the semiconductor part of the SMOX grains was addressed. This included the numerical calculation the charge carrier density as a function of the conduction band bending. Additionally the Poisson equation for spherical grains was solved. The grain results are saved to file and can now be here used again. Additionally a Python module was created. In this module all the parts we need to recycle from the previous notebook are merged together. By merging the relevant classes and functions into one python file, the command: `from part2 import *` will execute all the commands in this file and add them to the main namespace. By this the classes `material` and `grain` will again be available for further evaluations.

```
[1]: %pylab inline
```

Populating the interactive namespace from numpy and matplotlib

```
[2]: from part2 import *
```

2 Load the results

```
[3]: calc_dF = pd.read_hdf('results.h5', 'raw')
      calc_dF.index = range(len(calc_dF))
```

```

[21]: def create_grain_from_data(dF):
    if type(dF)==pd.Series:
        dF = pd.DataFrame([dF])

    if len(dF['temp'].unique())==1:
        T_C = dF['temp'].unique()[0]
    else:
        raise Exception('Multiple paramters for one grain are invalid.')

    if len(dF['ND'].unique())==1:
        ND = dF['ND'].unique()[0]
    else:
        raise Exception('Multiple paramters for one grain are invalid.')

    if len(dF['mass_eff'].unique())==1:
        mass_e_eff_factor = dF['mass_eff'].unique()[0]/CONST.MASS_E
    else:
        raise Exception('Multiple paramters for one grain are invalid.')

    if len(dF['R'].unique())==1:
        grainsize_radius = dF['R'].unique()[0]
    else:
        raise Exception('Multiple paramters for one grain are invalid.')

    EDCF_eV = calc_EDCF_by_temp(T_C, ND, mass_e_eff_factor)

    material = Material(T_C,DIFF_EF_EC_evolt=EDCF_eV)
    grain = Grain(grainsize_radius=grainsize_radius,material=material)

    return grain

```

```

[5]: fig, axes= subplots(2,2,figsize = (16,9))
for ax_i, (R, calc_dF_grainsize) in enumerate(calc_dF.groupby('R')):
    axe = fig.axes[ax_i]

    grain = create_grain_from_data(calc_dF_grainsize)

    axe.axhline(-grain.material.J_to_kT(grain.material.Diff_EF_EC),
                linestyle='--',color='r', label='Fermi Level')
    axe.set_ylim(-10,10)

    for vinit, ser_temp in calc_dF_grainsize.iterrows():

        r = ser_temp['r']
        v = ser_temp['v']

```

```

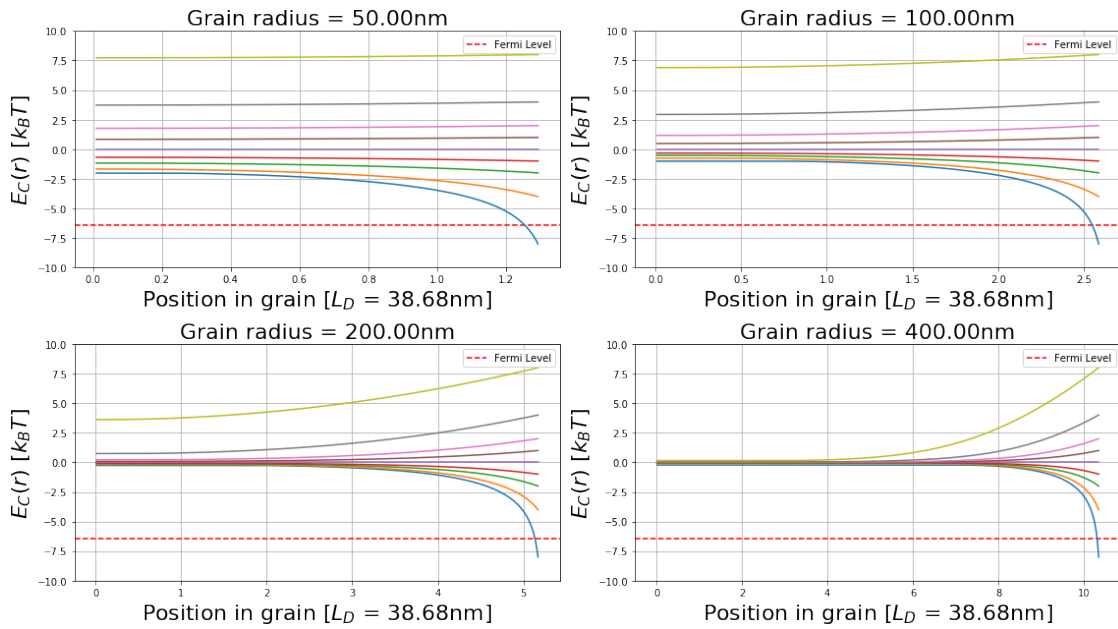
vdot = ser_temp['v_dot']

axe.set_title(f'Grain radius = {grain.R*1e9:.2f}nm', fontsize=22)

axe.plot(r,v, '-', label = "")
axe.set_ylabel('$E_C(r)$ [$k_BT$]', fontsize =22)
axe.set_xlabel(f'Position in grain [$L_D$ = {grain.material.LD*1e9:.2f}nm]',
               fontsize =22)

axe.legend()
axe.grid(b=True)
fig.tight_layout()

```



This graph shows how a surface potential is shielded by the remaining ionized donors. In the case of on deletion layer ($E_{C_{Surface}} > 0$), the total number of charges shielding the surface potential is rather small compared to the amount of charges in an accumulation layer ($E_{C_{Surface}} < 0$). The result of such an asymmetry is visible in the graph. The width of the accumulation layer is by far smaller than the width of the depleted are.

3 From charge distribution to resistance

With the previous tools and calculation it is now possible to assign each point inside the grain a certain charge density. From this charge density a specific resistivity can be assigned to this area. First the ratio of $\frac{n(r)}{n_b}$ inside the grain will be represented for the different starting conditions regarding $E_{C_{Surface}}$. To evaluate the $n(r)$ at arbitrary points r inside the grain, one additional step is needed. To calculate values between the points, where already solution exist, the additional

value can be retrieved by interpolating between the neighbors. This process is generally called interpolation. Again, SciPy and Python offer here also a easy to use and robust solution. `from scipy import interpolate` adds the interpolate module into the kernel. The `interp1d` function of this module is described ([here](#)) as follows: >Interpolate a 1-D function. >>x and y are arrays of values used to approximate some function f: $y = f(x)$. This class returns a function whose call method uses interpolation to find the value of new points.

This is how the function will be used for our needs. First we select one of the initial parameters for E_{init} and $E_{dot_{init}}$ to recalculate the solution of the Poisson equation with these correct start parameters. In a second step we use `interp1d` to create a function which uses an interpolation algorithm to find the right value for any position inside the grain.

```
[114]: from scipy import interpolate
import ipywidgets as widgets
from ipywidgets import interact, interactive, fixed, interact_manual

def get_interpolated_n_v(ser, grain):

    v = ser['v']
    r = ser['r']
    n = ser['n']

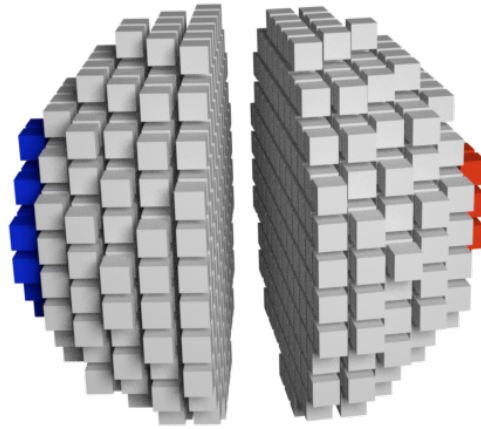
    r[0] = 0

    n_int = interpolate.interp1d(r*grain.material.LD, n, kind='linear')
    v_int = interpolate.interp1d(r*grain.material.LD, v, kind='linear')

    return n_int, v_int
```

3.1 The “numerical” grain

With the gain in information from the last notebook, we can calculate the number of free charge carriers



```
[114]: def pos_to_r(xi,yi,grain, d):
        dx = 2*grain.R/d.shape[0]
        cx = d.shape[0]//2
        cy = d.shape[1]//2
        ri = ((xi-cx)**2+(yi-cy)**2)**0.5
        return(ri*dx)

def create_numerical_grain_matrix( grain, ser,size_n=100,):
    n_int, v_int = get_interpolated_n_v(ser, grain)
    nx = ny = size_n

    d_v = np.zeros((2*nx+1,2*ny+1))
    d_cond = np.zeros((2*nx+1,2*ny+1))
    d_mask = np.zeros((2*nx+1,2*ny+1))

    for xi in range(d_cond.shape[0]):
        for yi in range(d_cond.shape[1]):
            r = float(pos_to_r(xi,yi,grain, d_cond))
            try:
                condu = n_int(r)/grain.material.nb
                d_cond[xi, yi] = condu
                d_mask[xi,yi] = 1

            except ValueError:
                #outside the grain
```

```

        d_cond[xi, yi] = 0
        d_mask[xi,yi] = 0
    return d_v, d_cond, d_mask

```

3.2 Precalc the numerical grains for all conditions

```

[119]: for i, ser in calc_dF.iterrows():
        print(f'Initalized {i+1} of {len(calc_dF)}.', end='\r')
        grain = create_grain_from_data(ser)
        d_v, d_cond,d_mask = create_numerical_grain_matrix( grain, ser,size_n=100)
        d_cond_plot = d_cond.copy()
        d_cond_plot[np.where(d_mask==0)]=None
        calc_dF.loc[ser.name, 'd_cond'] = [d_cond_plot]

```

Initalized 36 of 36.

```

[125]: %matplotlib inline

```

This means in the accumulation layer there are a the specific position 5.39 times more free charge carriers available as in the center. This results directly in a lower resistivity of the same factor. Instead of just calculating just individual values, the graphical representation is in this case more helpful.

```

[126]: def plot_grain_states(calc_dF_grainsize, vmax=None, vmin=None):
        fig, axes = subplots(3,3, figsize = (16,9))

        grain = create_grain_from_data(calc_dF_grainsize)

        for ax_i, (vinit, ser) in enumerate(calc_dF_grainsize.iterrows()):
            axe = fig.axes[ax_i]

            Einit_kT = ser['Einit_kT']

            axe.set_title(r'$E_{C_{Surface}}$'+f'{Einit_kT}kT')
            axe.set_ylabel('x [nm]')
            axe.set_xlabel('y [nm]')

            d_cond_plot = calc_dF.loc[ser.name, 'd_cond']
            im = axe.imshow(np.log(d_cond_plot), interpolation='bicubic',
                            extent=(-grain.R*1e9, grain.R*1e9, -grain.R*1e9, grain.
→R*1e9),
                            vmax=vmax, vmin=vmin, cmap='hot')
            fig.tight_layout()

def plot_currents(GrainRadius):

```

```

R = GrainRadius/1e9
calc_dF_grainsize = calc_dF.groupby('R').get_group(R)
max_n = np.log(calc_dF_grainsize['d_cond'].apply(lambda r:np.nanmax(r))).
→max()
min_n = np.log(calc_dF_grainsize['d_cond'].apply(lambda r:np.nanmin(r))).
→min()
print(max_n, min_n)

plot_grain_states(calc_dF_grainsize, vmax = max_n, vmin = min_n)

grainsizes = list(calc_dF['R'].unique())
interact(plot_currents, GrainRadius=np.array(grainsizes)*1e9, text='Select a_
→grainsize:');

```

```

interactive(children=(Dropdown(description='GrainRadius', options=(50.0, 100.0, 200.0, 400.0), v

```

4 Relaxation

```

[71]: calc_dF.to_hdf('size_test_delete_please.h5', 'raw')

```

```

/usr/lib/python3.8/site-packages/pandas/core/generic.py:2530:

```

```

PerformanceWarning:

```

```

your performance may suffer as PyTables will pickle object types that it cannot
map directly to c-types [inferred_type->mixed,key->block1_values] [items->['n',
'r', 'v', 'v_dot', 'd_cond']]

```

```

pytables.to_hdf(path_or_buf, key, self, **kwargs)

```

```

[139]: from scipy import signal

```

```

def initaliz_d_v(d_v):
    d_v[:,0] = -1000
    d_v[:,-1] = +1000
    d_v = d_v*d_mask
    return d_v

def solve_relaxation(d_v, d_cond, d_mask):

    for i in range(10000):

        conv = [[0,1,0],[1,0,1],[0,1,0]]
        numerator = signal.convolve2d(d_v*d_cond, conv, boundary='fill',
                                     mode='same', fillvalue=0)
        denominator = signal.convolve2d(d_cond, conv, boundary='fill',

```

```

mode='same', fillvalue=0)

d_v_new = (numerator/denominator)*d_mask
d_v_new = np.nan_to_num(d_v_new,0)

d_v_prev = d_v.copy()
d_v = d_v_new.copy()
d_v = inicializ_d_v(d_v)
return d_v, d_cond, d_mask

def plot_num_grain(d_v, d_cond, d_mask):
    fig, axes =subplots(1,3)
    d_v_plot = d_v.copy()
    d_v_plot[np.where(d_mask==0)]=None

    axes[0].imshow(d_mask)
    axes[1].imshow(d_cond)
    axes[2].imshow(d_v_plot,interpolation= 'nearest')

def plot_voltage_1d(d_v):
    fig, axe = subplots()
    center = d_v[d_v.shape[0]//2,:]
    axe.plot(center)

def calc_current_center(d_v, d_cond, d_mask):
    center_pos = d_v.shape[0]//2
    center_current = (d_v[:,center_pos+1]-d_v[:,center_pos-1])*d_cond[:,
→,center_pos]

    r = np.array([float(pos_to_r(xi,center_pos,grain, d_v)) for xi in_
→range(len(center_current))])

    center_current_tot = np.sum(center_current*r*2*pi)
    return center_current_tot, center_current, r

def plot_center_current(r, center_current):
    fig, axe = subplots()
    axe.plot(r*1e9,center_current)
    return

#vinit, current = calc_current(d_v, d_cond, d_mask)

```

```
[ ]: calc_dF['current'] = None
```

```
[248]: def inicializ_d_v(d_v):
        d_v[:,0] = -1000
        d_v[:,-1] = +1000
        d_v = d_v*d_mask

```



```

return d_v

#for vinit, ser in calc_dF.iterrows():
for i, (ind,ser) in enumerate(calc_dF.iterrows()):
    break
    print(f'Initalized {i+1} of {len(calc_dF)}.', end='\r')
    vinit = ser.name

    grain = create_grain_from_data(ser)

    d_v, d_cond, d_mask = create_numerical_grain_matrix( grain, ser,size_n=50,)
    d_v = initaliz_d_v(d_v)

    d_v, d_cond, d_mask = solve_relaxation(d_v, d_cond, d_mask)

    center_current_tot, center_current, r = calc_current_center(d_v, d_cond,
→d_mask)

    #plot_num_grain(d_v, d_cond, d_mask)
    #plot_voltage_1d(d_v)
    #plot_center_current(r, center_current)

    calc_dF.loc[vinit, 'current'] = center_current_tot

```

[150]: `calc_dF.to_hdf('res.h5', 'raw')`

```

/usr/lib/python3.8/site-packages/pandas/core/generic.py:2530:
PerformanceWarning:
your performance may suffer as PyTables will pickle object types that it cannot
map directly to c-types [inferred_type->mixed,key->block1_values] [items->['n',
'r', 'v', 'v_dot', 'd_cond', 'current']]

pytables.to_hdf(path_or_buf, key, self, **kwargs)

```

[257]:

```

def overlay_res_vs_band(GrainRadius):
    print(GrainRadius)
    fig, axe = subplots(figsize = (16,9))

    for ax_i, (R, calc_dF_grainsize) in enumerate(calc_dF.groupby('R')):
        if R*1e9==GrainRadius:
            linewidth = 5
        else:
            linewidth = 1

```

```

        flat_band = calc_dF_grainsize[calc_dF_grainsize['Einit_kT']==0].
        →iloc[0]['current']
        res = flat_band/calc_dF_grainsize['current']
        v = calc_dF_grainsize['Einit_kT']
        axe.plot(v, res, 'o-', label = f'R={R*1e9}mn', linewidth=linewidth)
        axe.set_yscale('log')
        axe.grid(b=True)

    fig.tight_layout()
    axe.legend()
    display(fig)

grainsizes = list(calc_dF['R'].unique())
interact(overlay_res_vs_band, GrainRadius=np.array(grainsizes)*1e9,
        text='Select a grainsize:');

```

```

interactive(children=(Dropdown(description='GrainRadius', options=(50.0, 100.0, 200.0, 400.0), v

```

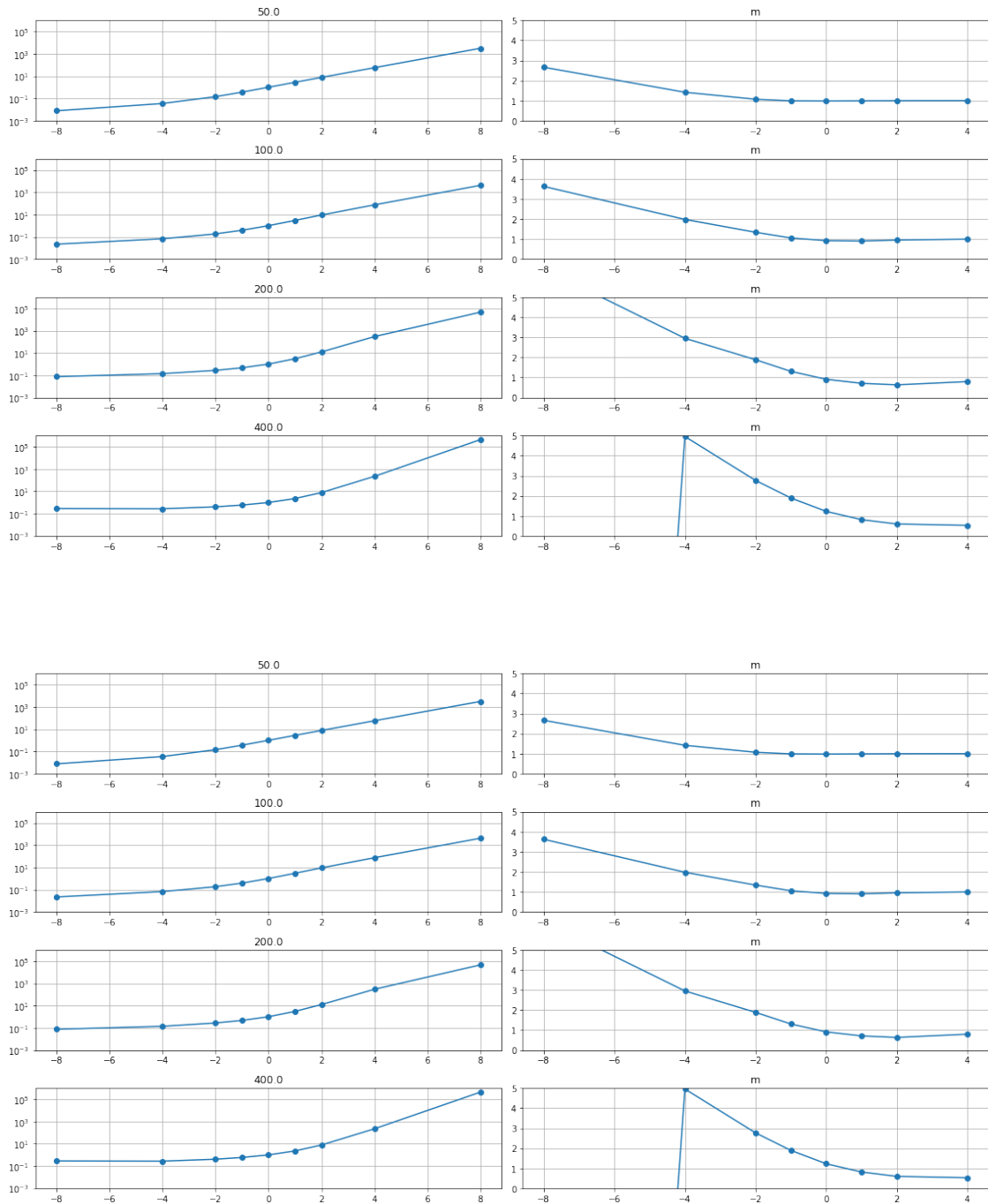
```

[256]: fig, axess = subplots(4,2,figsize = (16,9))

for ax_i, (R, calc_dF_grainsize) in enumerate(calc_dF.groupby('R')):
    axes = axess[ax_i,:]
    axes[0].set_title(R*1e9)
    flat_band = calc_dF_grainsize[calc_dF_grainsize['Einit_kT']==0].
    →iloc[0]['current']
    res = flat_band/calc_dF_grainsize['current']
    v = calc_dF_grainsize['Einit_kT']
    axes[0].plot(v, res, 'o-', label = f'R={R*1e9}mn')
    axes[0].set_yscale('log')
    lnr = np.log([float(r) for r in res.values])
    axes[1].plot(v[0:-1], 1/(np.diff(lnr)/np.diff(v)), 'o-')
    axes[1].set_ylim(0,5)
    axes[0].set_ylim(0.001,1e6)
    axes[0].grid(b=True)
    axes[1].grid(b=True)
    axes[1].set_title('m')

    axe.legend()
fig.tight_layout()
display(fig)

```



5 Bibliography section

References