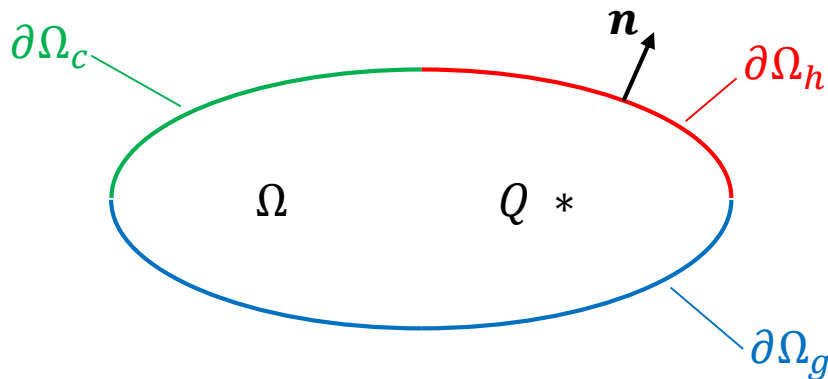


10 FE formulation heat flow



3d body occupying the domain Ω

\mathbf{n} unit outward normal vector to the boundary $\partial\Omega$

boundary $\partial\Omega = \partial\Omega_h \cup \partial\Omega_g \cup \partial\Omega_c$

Strong form of the transient heat equation $c\rho\dot{T} + \text{div}(\mathbf{q}) - Q = 0$

c heat capacity ρ mass density T temperature \mathbf{q} heat flux Q heat source

Boundary conditions

Dirichlet or essential BCs

$T = g$ given on $\partial\Omega_g$

Neumann or natural BCs

$q_n = \mathbf{q}^T \mathbf{n} = h$ given on $\partial\Omega_h$

Robin or convection BCs

$q_n = \alpha[T - T_\infty]$ given on $\partial\Omega_c$

α convection coefficient T_∞ temperature at infinity

Initial condition

$T(\mathbf{x}, t = 0) = T_{init}$ given in Ω

10 FE formulation heat flow

Strong form → **weak form** for FE formulation, we assume here no convection BCs!

Multiply by an arbitrary weight function $v(\mathbf{x})$ and apply Green-Gauss theorem

$$\int_{\Omega} v c \rho \dot{T} dV - \int_{\Omega} [\nabla v]^T \mathbf{q} dV - \int_{\Omega} v Q dV + \int_{\partial\Omega} v \mathbf{q}^T \mathbf{n} dV = 0$$

Approximate weight function by $v(\mathbf{x}) = \mathbf{c}^T \mathbf{N}^T(\mathbf{x})$ and $[\nabla v]^T = \mathbf{c}^T \mathbf{B}^T(\mathbf{x})$

$$\mathbf{c}^T \left[\int_{\Omega} \mathbf{N}^T c \rho \dot{T} dV - \int_{\Omega} \mathbf{B}^T \mathbf{q} dV - \int_{\Omega} \mathbf{N}^T Q dV + \int_{\partial\Omega} \mathbf{N}^T \mathbf{q}^T \mathbf{n} dA \right] = 0$$

Since \mathbf{c} can be chosen arbitrary the expression in the brackets $[...] = \mathbf{0}$

$$\left[\int_{\Omega} \mathbf{N}^T c \rho \dot{T} dV - \int_{\Omega} \mathbf{B}^T \mathbf{q} dV - \int_{\Omega} \mathbf{N}^T Q dV + \int_{\partial\Omega} \mathbf{N}^T \mathbf{q}^T \mathbf{n} dA \right] = \mathbf{0}$$

10 FE formulation heat flow

Strong form → **weak form** for FE formulation, we assume here no convection BCs!

$$\left[\int_{\Omega} \mathbf{N}^T c \rho \dot{T} dV - \int_{\Omega} \mathbf{B}^T \mathbf{q} dV - \int_{\Omega} \mathbf{N}^T Q dV + \int_{\partial\Omega} \mathbf{N}^T \mathbf{q}^T \mathbf{n} dA \right] = \mathbf{0}$$

Approximate time-dependent temperature field by $T(\mathbf{x}, t) = \mathbf{N}(\mathbf{x}) \mathbf{a}(t)$
and introduce Fourier's law for heat conduction $\mathbf{q} = -\mathbf{D} \nabla T$

$$\int_{\Omega} \mathbf{N}^T c \rho \mathbf{N} dV \dot{\mathbf{a}} + \int_{\Omega} \mathbf{B}^T \mathbf{D} \mathbf{B} dV \mathbf{a} - \int_{\Omega} \mathbf{N}^T Q dV + \int_{\partial\Omega} \mathbf{N}^T q_n dA = \mathbf{0}$$

Rearrange $\mathbf{C} \dot{\mathbf{a}} + \mathbf{K} \mathbf{a} = \mathbf{f}_l + \mathbf{f}_b$ with

Heat capacity matrix $\mathbf{C} = \int_{\Omega} \mathbf{N}^T c \rho \mathbf{N} dV$

,Stiffness' or conductivity matrix $\mathbf{K} = \int_{\Omega} \mathbf{B}^T \mathbf{D} \mathbf{B} dV$

Load vector $\mathbf{f}_l = \int_{\Omega} \mathbf{N}^T Q dV$

Boundary vector $\mathbf{f}_b = - \int_{\partial\Omega} \mathbf{N}^T q_n dA$

10 FE formulation heat flow

Semi-discrete equations $\mathbf{C}\dot{\mathbf{a}} + \mathbf{K}\mathbf{a} = \mathbf{f}_l + \mathbf{f}_b = \mathbf{f}$ still continuous in time

Now, discretization in time, here the θ -method.

$\theta \in [0,1]$ is a numerical parameter of this method.

Subdivide time interval $t \in [0, T_{end}]$ into $N + 1$ subintervals

$$[0, T_{end}] = \bigcup_{n=0}^N [t_n, t_{n+1}] \quad \text{and} \quad \Delta t_n = t_{n+1} - t_n$$

Approximate time derivative
$$\dot{\mathbf{a}} = \frac{\mathbf{a}(t_{n+1}) - \mathbf{a}(t_n)}{t_{n+1} - t_n} = \frac{\mathbf{a}_{n+1} - \mathbf{a}_n}{\Delta t_n}$$

Approximate temperature
$$\mathbf{a} = \theta \mathbf{a}_{n+1} + [1 - \theta] \mathbf{a}_n$$

Approximate right hand side
$$\mathbf{f} = \theta \mathbf{f}_{n+1} + [1 - \theta] \mathbf{f}_n$$

$$\frac{1}{\Delta t_n} \mathbf{C}[\mathbf{a}_{n+1} - \mathbf{a}_n] + \mathbf{K}[\theta \mathbf{a}_{n+1} + [1 - \theta] \mathbf{a}_n] = \theta \mathbf{f}_{n+1} + [1 - \theta] \mathbf{f}_n$$

\mathbf{a}_{n+1} unknown node temperatures at time t_{n+1}

10 FE formulation heat flow

$$\frac{1}{\Delta t_n} \mathbf{C}[\mathbf{a}_{n+1} - \mathbf{a}_n] + \mathbf{K}[\theta \mathbf{a}_{n+1} + [1 - \theta] \mathbf{a}_n] = \theta \mathbf{f}_{n+1} + [1 - \theta] \mathbf{f}_n$$

\mathbf{a}_{n+1} unknown node temperatures at time t_{n+1}

\mathbf{a}_n known node temperature from the previous time step t_n

initial conditions

At $t = 0$, i.e. $n = 0$, $\mathbf{a}_0 = \mathbf{a}(t = 0)$ are given values!

You need to know the temperature field in the whole domain Ω at $t = 0$

- $\theta = 1$. This is known as an implicit scheme and is often chosen due to its stability properties
- $\theta = 1/2$. Midpoint rule or Crank-Nicholson scheme. Often used due to its accuracy properties.
- $\theta = 0$. Forward scheme or explicit scheme. If additional assumption regarding the \mathbf{C} matrix is made, this method enables a very efficient scheme on large scale clusters to be obtained.

10 FE formulation heat flow

Global FE formulation

Integration has to be performed over the complete domain Ω or boundary $\partial\Omega$ like global matrices $\mathbf{K} = \int_{\Omega} \mathbf{B}^T \mathbf{D} \mathbf{B} dV$ or vectors $\mathbf{f}_b = - \int_{\partial\Omega} \mathbf{N}^T q_n dA$ etc.

Element FE formulation

Practical implementation: split integrals over the complete domain Ω into a sum of integrals over each individual element e . Then **assemble in a loop** the **element matrices** and vectors to **global matrices** and vectors.

One introduces element matrices and vectors like

$$\mathbf{K}^e = \int_{\Omega^e} \mathbf{B}^{eT} \mathbf{D} \mathbf{B}^e dV \quad \text{and} \quad \mathbf{f}_l^e = - \int_{\partial\Omega^e} \mathbf{N}^{eT} q_n dA \quad \text{and} \quad \mathbf{f}_b^e = - \int_{\Omega^e} \mathbf{N}^{eT} Q dA$$

e.g. for a triangular 2d element

$$\mathbf{N}^e = [N_1^e \quad N_2^e \quad N_3^e] \quad \mathbf{B}^e = \nabla \mathbf{N}^e = \begin{bmatrix} \frac{\partial N^e}{\partial x} \\ \frac{\partial N^e}{\partial y} \end{bmatrix}$$

10 FE formulation heat flow

Element FE formulation

$$\mathbf{K}^e = \int_{\Omega^e} \mathbf{B}^{eT} \mathbf{D} \mathbf{B}^e dV \quad \text{and} \quad \mathbf{f}_b^e = - \int_{\partial\Omega^e} \mathbf{N}^{eT} q_n dA \quad \text{and} \quad \mathbf{f}_b^e = - \int_{\Omega^e} \mathbf{N}^{eT} Q dA$$

$\dim \mathbf{K}^e = [3 \times 3]$ for a 3-node triangular element

Expanded element matrices and vectors

Define expanded element stiffness matrix by

$$\mathbf{K}^{ee} = \int_{\Omega^e} \mathbf{B}^T \mathbf{D} \mathbf{B} dV$$

Attention: e.g. 2d mesh of 3-node triangular elements with total node number n
 we use the global \mathbf{B} matrix with $\dim \mathbf{B} = [2 \times n]$ but **integrate only over one element** Ω^e
 Thus, $\dim \mathbf{K}^{ee} = [n \times n]$ but only $[3 \times 3] = 9$ entries are non-zero.

We do this for all n_{el} elements in the mesh and get the global stiffness matrix

$$\mathbf{K} = \sum_{e=1}^{n_{el}} \mathbf{K}^{ee} \quad \text{analog} \quad \mathbf{f}_l = \sum_{e=1}^{n_{el}} \mathbf{f}_l^{ee} \quad \text{etc.}$$

10 FE formulation heat flow

Example expanded element matrices and vectors

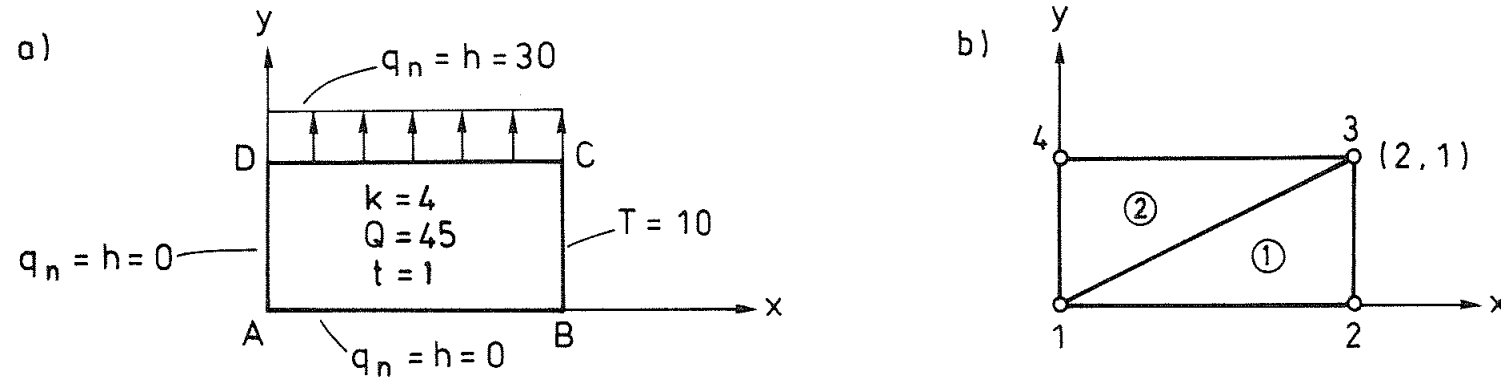


Figure 10.4 (a) Problem reformulation; (b) finite element mesh

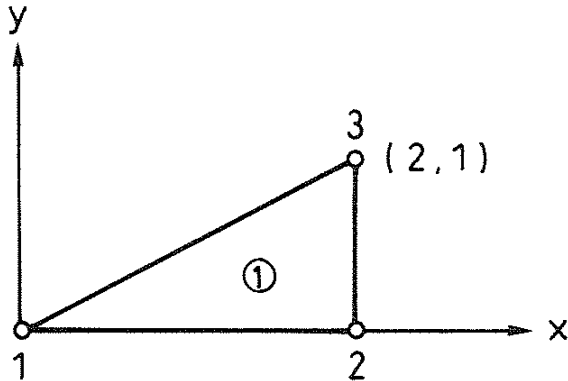
$$\mathbf{K}^e = k \mathbf{B}^{eT} \mathbf{B}^e t A_\alpha$$

$$\mathbf{B}^e = \frac{1}{2A_\alpha} \begin{bmatrix} y_j - y_k & y_k - y_i & y_i - y_j \\ x_k - x_j & x_i - x_k & x_j - x_i \end{bmatrix}$$

$$\mathbf{f}_1^e = Qt \int_{A_\alpha} \mathbf{N}^{eT} dA$$

k thermal conductivity
 t thickness of the plate
 A_α area of the element α

10 FE formulation heat flow



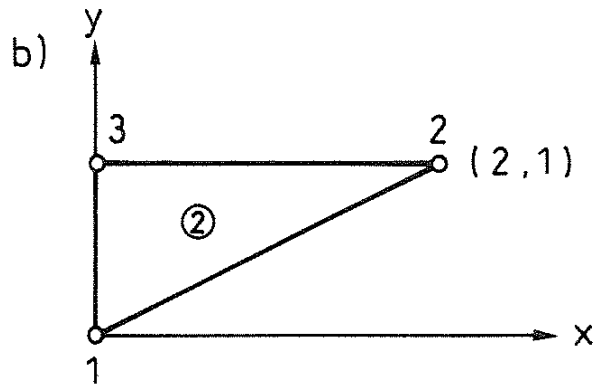
$$x_i = y_i = 0; \quad x_j = 2, y_j = 0; \quad x_k = 2, y_k = 1$$

$$\mathbf{B}^e = \frac{1}{2} \begin{bmatrix} -1 & 1 & 0 \\ 0 & -2 & 2 \end{bmatrix}$$

$$\mathbf{K}^e = \begin{bmatrix} 1 & -1 & 0 \\ -1 & 5 & -4 \\ 0 & -4 & 4 \end{bmatrix}; \quad \mathbf{f}_1^e = 15 \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

$$\mathbf{K}_1^{ee} = \begin{bmatrix} 1 & -1 & 0 & 0 \\ -1 & 5 & -4 & 0 \\ 0 & -4 & 4 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}; \quad \mathbf{f}_1 = \mathbf{f}_{11}^{ee} = 15 \begin{bmatrix} 1 \\ 1 \\ 1 \\ 0 \end{bmatrix}$$

10 FE formulation heat flow



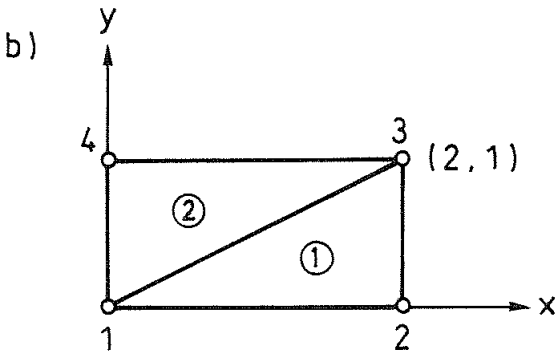
$$x_i = y_i = 0; \quad x_j = 2, y_j = 1; \quad x_k = 0, y_k = 1$$

$$\mathbf{B}^e = \frac{1}{2} \begin{bmatrix} 0 & 1 & -1 \\ -2 & 0 & 2 \end{bmatrix}$$

$$\mathbf{K}^e = \begin{bmatrix} 4 & 0 & -4 \\ 0 & 1 & -1 \\ -4 & -1 & 5 \end{bmatrix}; \quad \mathbf{f}_1^e = 15 \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

$$\mathbf{K}_2^{ee} = \begin{bmatrix} 4 & 0 & 0 & -4 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ -4 & 0 & -1 & 5 \end{bmatrix}; \quad \mathbf{f}_{12}^{ee} = 15 \begin{bmatrix} 1 \\ 0 \\ 1 \\ 1 \end{bmatrix}$$

10 FE formulation heat flow

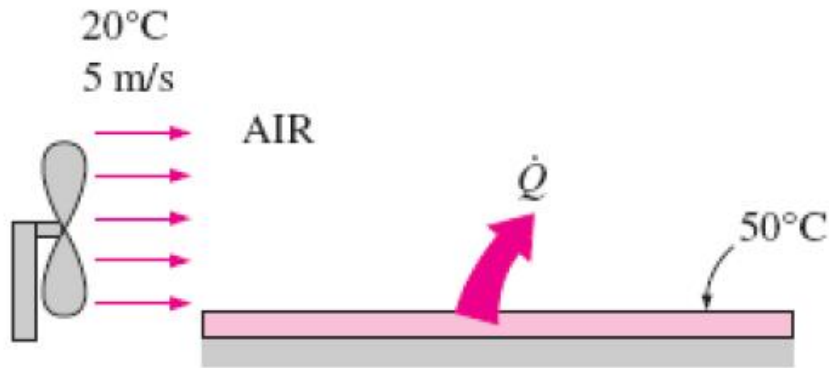


$$\mathbf{K}_1^{ee} = \begin{bmatrix} 1 & -1 & 0 & 0 \\ -1 & 5 & -4 & 0 \\ 0 & -4 & 4 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}; \quad \mathbf{f}_1 = \mathbf{f}_{11}^{ee} = 15 \begin{bmatrix} 1 \\ 1 \\ 1 \\ 0 \end{bmatrix}$$

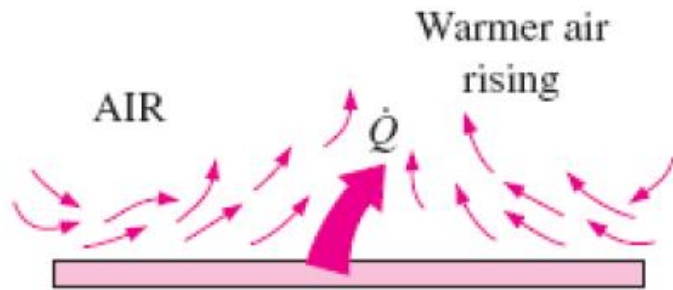
$$\mathbf{K}_2^{ee} = \begin{bmatrix} 4 & 0 & 0 & -4 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ -4 & 0 & -1 & 5 \end{bmatrix}; \quad \mathbf{f}_{12}^{ee} = 15 \begin{bmatrix} 1 \\ 0 \\ 1 \\ 1 \end{bmatrix}$$

$$\begin{aligned} \mathbf{K} &= \sum_{e=1}^{n_{el}} \mathbf{K}^{ee} \\ \mathbf{f}_l &= \sum_{e=1}^{n_{el}} \mathbf{f}_l^{ee} \end{aligned} \quad \Rightarrow \quad \mathbf{K} = \begin{bmatrix} 5 & -1 & 0 & -4 \\ -1 & 5 & -4 & 0 \\ 0 & -4 & 5 & -1 \\ -4 & 0 & -1 & 5 \end{bmatrix}; \quad \mathbf{f}_l = 15 \begin{bmatrix} 2 \\ 1 \\ 2 \\ 1 \end{bmatrix}$$

Heat flow with convection



(a) Forced convection $\alpha = 12 \dots 120 \text{ W}/(\text{m}^2\text{K})$



(b) Free convection $\alpha = 5 \dots 25 \text{ W}/(\text{m}^2\text{K})$

The convection coefficient α has to be measured for convective heat transfer!!!

$$q_n = \alpha[T - T_\infty]$$

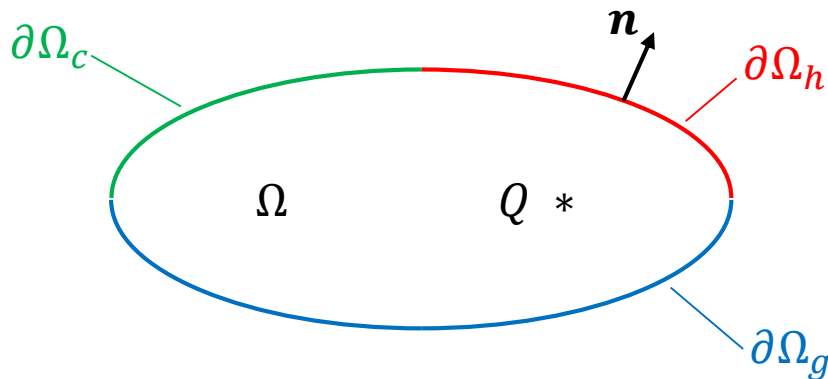


Air-cooled engine



CPU-cooler

10 FE formulation heat flow



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\mathbf{n} unit outward normal vector to the boundary $\partial\Omega$

boundary $\partial\Omega = \partial\Omega_h \cup \partial\Omega_g \cup \partial\Omega_c$

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α convection coefficient T_∞ temperature at infinity

Initial condition

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10 FE formulation heat flow

Strong form → **weak form** for FE formulation, we assume here no convection BCs!

$$\left[\int_{\Omega} \mathbf{N}^T c \rho \dot{T} dV - \int_{\Omega} \mathbf{B}^T \mathbf{q} dV - \int_{\Omega} \mathbf{N}^T Q dV + \int_{\partial\Omega} \mathbf{N}^T \mathbf{q}^T \mathbf{n} dA \right] = \mathbf{0}$$

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$$\int_{\Omega} \mathbf{N}^T c \rho \mathbf{N} dV \dot{\mathbf{a}} + \int_{\Omega} \mathbf{B}^T \mathbf{D} \mathbf{B} dV \mathbf{a} - \int_{\Omega} \mathbf{N}^T Q dV + \int_{\partial\Omega} \mathbf{N}^T q_n dA = \mathbf{0}$$

Rearrange $\mathbf{C} \dot{\mathbf{a}} + \mathbf{K} \mathbf{a} = \mathbf{f}_l + \mathbf{f}_b$ with

Heat capacity matrix $\mathbf{C} = \int_{\Omega} \mathbf{N}^T c \rho \mathbf{N} dV$

,Stiffness' or conductivity matrix $\mathbf{K} = \int_{\Omega} \mathbf{B}^T \mathbf{D} \mathbf{B} dV$

Load vector $\mathbf{f}_l = \int_{\Omega} \mathbf{N}^T Q dV$

Boundary vector $\mathbf{f}_b = - \int_{\partial\Omega} \mathbf{N}^T q_n dA$

Assignment

- Project should be solved in **groups of two**.
- Write a technical report
- Send your project as a zipped folder to FHLF20@solid.lth.se
 - Report as a pdf-file
 - Matlab files
- Hand in a printed version of your report at the Division of Solid Mechanics
Use the box close to our printer on the 5th floor, M-Huset
- Deadline 22nd May 2018, 16:00
- Max. 5 bonus points only valid for the exam end of May/June 2018
- Approval of your report until 12 June 2018 if it contains substantial mistakes
- Do **not** use subroutines which are included in CALFEM!
- You **can** use the matlab program skeletons provided for the Computer Exercises