

Homework Assignment 5a

20 points

DUE: TBA

1. Noisy measurement (10 points).

- (a) Suppose you make a von Neumann measurement of the spin component of a spin-1/2 particle along one of the three Cartesian axes, tossing a fair three-sided die to determine which of the three measurements to make. Determine the Kraus operators, quantum operations, and POVM elements for this six-outcome measurement.
- (b) Suppose now that all you know after the measurement of part (a) is whether the result is +1 or -1, but not which axis was used. Determine the Kraus operators, quantum operations, and POVM elements for this two-outcome measurement. This is a model of a very noisy measurement, the noise coming from not knowing how the measuring apparatus is oriented.
- (c) Find at least one other pair of quantum operations that corresponds to the POVM of part (b).

2. From Neumark extension to a measurement model (10 points).

The goal of this problem is to convert the Neumark extension of a rank-one POVM in **P3** of **HW4** to a measurement model. Recall that in the original setting, we considered an N -element, rank-one POVM for a qubit system, and denoted the POVM elements by $E_\alpha = |\bar{\psi}_\alpha\rangle\langle\bar{\psi}_\alpha|$, $\alpha = 1, \dots, N$, where the vectors $|\bar{\psi}_\alpha\rangle = \sqrt{\mu_\alpha}|\psi_\alpha\rangle$ are subnormalized, with the norm squared $\mu_\alpha = \langle\bar{\psi}_\alpha|\bar{\psi}_\alpha\rangle$. The Neumark-extended vectors, $|\hat{\psi}_\alpha\rangle$, live in a Hilbert space of dimension N , which is the *direct sum* of the qubit Hilbert space and a Hilbert space of $N - 2$ additional dimensions. The Neumark-extended vectors satisfy $P|\hat{\psi}_\alpha\rangle = |\bar{\psi}_\alpha\rangle$, where $P = \sum_{\alpha=1}^N E_\alpha$ is the projector onto the qubit Hilbert space.

The immediate issue we run into when we want to convert the Neumark extension to a measurement model is that, in the latter, the total Hilbert space is the *tensor product* of the qubit Hilbert space with the ancillary Hilbert space. So now it does not make sense to talk about a projector P onto the qubit Hilbert space anymore.

We begin by choosing the ancillae to consist of $n - 1$ additional qubits, where n is the smallest integer such that $2^n \geq N$. We then add $2^n - N$ zero POVM elements (that have zero probability of occurring) so that we formally have a 2^n -outcome, rank-one POVM. The index α now takes on 2^n values, so we replace it by a bitstring $\mathbf{a} = a_1 \dots a_n$. The idea is to choose a two-dimensional subspace embedded in the n -qubit tensor product space to be our system Hilbert space. Consider the Hilbert space spanned by $|0\rangle \otimes |0\rangle^{\otimes(n-1)}$ and $|1\rangle \otimes |0\rangle^{\otimes(n-1)}$. The POVM elements take the form $E_{\mathbf{a}} = |\bar{\psi}_{\mathbf{a}}\rangle\langle\bar{\psi}_{\mathbf{a}}| \otimes (|0\rangle\langle 0|)^{\otimes(n-1)}$. The projector onto the system Hilbert space is

$$P = \sum_{\mathbf{a}} E_{\mathbf{a}} = \sum_{\mathbf{a}} |\bar{\psi}_{\mathbf{a}}\rangle\langle\bar{\psi}_{\mathbf{a}}| \otimes (|0\rangle\langle 0|)^{\otimes(n-1)} = \mathbb{1} \otimes (|0\rangle\langle 0|)^{\otimes(n-1)}$$

The Neumark-extended vector $|\hat{\psi}_{\mathbf{a}}\rangle$ live in the n -qubit tensor product space and satisfy $P|\hat{\psi}_{\mathbf{a}}\rangle = |\bar{\psi}_{\mathbf{a}}\rangle \otimes |0\rangle^{\otimes(n-1)}$.

(a) By considering the way the Neumark extension is constructed from the computational basis $|x\rangle = |x_1 \dots x_n\rangle$, draw an n -qubit quantum circuit that corresponds to a measurement of the POVM.

(b) The circuit of part (a) does not fit our definition of a measurement model because it involves a direct measurement on the original qubit, instead of just measurements on the ancilla qubits. Construct a proper measurement model by adding one more ancilla qubit to your circuit. Draw the resulting $(n + 1)$ -qubit circuit, and determine the Kraus operators for the measurement model.

(c) By controlling on the outputs of the n measurements in part (b), modify the circuit so that the Kraus operators are $A_a = \sqrt{E_a} = \sqrt{\mu_a} |\psi_a\rangle\langle\psi_a| = |\psi_a\rangle\langle\bar{\psi}_a|$.
