

$$\vec{A} = \begin{pmatrix} A_1 \\ A_2 \\ A_3 \end{pmatrix} = A_1 \hat{e}_1 + A_2 \hat{e}_2 + A_3 \hat{e}_3 = \sum_{i=1}^3 A_i \hat{e}_i$$

Implicitly using an orthonormal basis

$$\hat{e}_i \cdot \hat{e}_j = \delta_{ij} = \begin{cases} 1, & i=j \\ 0, & i \neq j \end{cases}$$

Dot product and cross product (See also Tutorial 24 Nov)

$$\vec{A} \cdot \vec{B} = A_1 B_1 + A_2 B_2 + A_3 B_3, \quad \vec{A} \times \vec{B} = \begin{pmatrix} A_2 B_3 - B_2 A_3 \\ A_3 B_1 - B_3 A_1 \\ A_1 B_2 - B_1 A_2 \end{pmatrix}$$

You can define "cross product" in arbitrary dim (exterior product), but it will be a vector only in 3 dim

Einstein summation convention: we sum over a repeated index

$$\vec{A} = A_i \hat{e}_i$$

$$\vec{A} \cdot \vec{B} = (A_i \hat{e}_i) \cdot (B_j \hat{e}_j) = A_i B_j \hat{e}_i \cdot \hat{e}_j = A_i B_j \delta_{ij} = A_i B_i$$

Don't do this!

$$\vec{x} \cdot \vec{y} = (x_i \hat{e}_i) \cdot (y_i \hat{e}_i)$$

Don't know what are included in which summation

dummy indices
(μνδσλ)

$$\text{Length } \sqrt{\vec{A} \cdot \vec{A}} = \sqrt{A_i A_i} = \sqrt{A_1^2 + A_2^2 + A_3^2}$$

$$\hat{e}_i \times \hat{e}_j = \hat{e}_k \text{ where } i, j, k \text{ are cyclic}$$

Antisymmetry

$$\hat{e}_i \times \hat{e}_j = -\hat{e}_j \times \hat{e}_i \text{ (In particular } \hat{e}_i \times \hat{e}_i = 0)$$

More generally, we can compactly define the cross product as

$$\vec{A} \times \vec{B} = \epsilon_{ijk} A_i B_j \hat{e}_k$$

Levi-Civita symbol

$$\epsilon_{ijk} = \begin{cases} 1 & (i, j, k) \text{ is cyclic} \\ -1 & (i, j, k) \text{ is anticyclic} \\ 0 & \text{otherwise} \end{cases}$$

The kth component

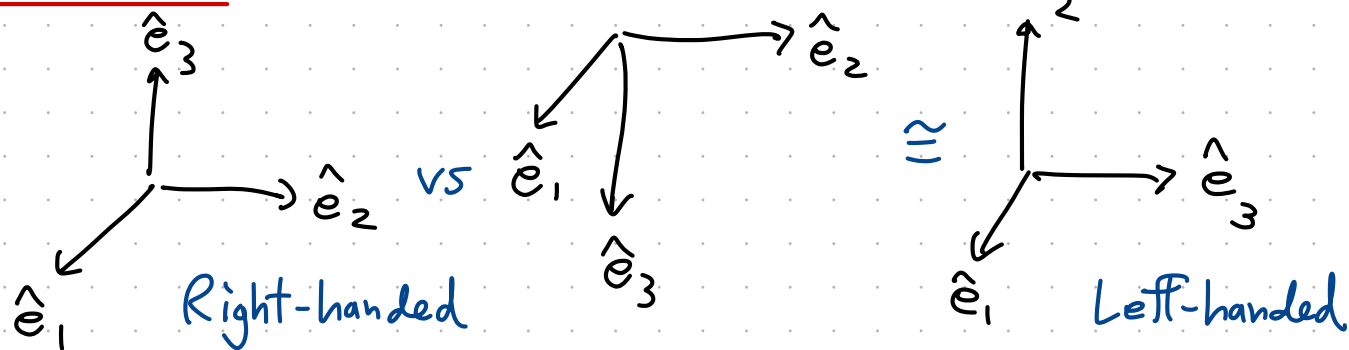
$$(\vec{A} \times \vec{B})_k = \epsilon_{ijk} A_i B_j$$

Fixed index Dummy

$$\text{Examples } \epsilon_{312} = 1, \epsilon_{321} = -1$$

$$\epsilon_{112} = 0$$

Handedness



The difference is encapsulated in the triple product $\hat{e}_i \cdot (\hat{e}_j \times \hat{e}_k)$

$$\vec{A} \cdot (\vec{B} \times \vec{C}) = \epsilon_{ijk} A_i B_j C_k = \det \begin{pmatrix} A_1 & B_1 & C_1 \\ A_2 & B_2 & C_2 \\ A_3 & B_3 & C_3 \end{pmatrix}$$

Proof (optional)

Determinant
in N dimension

$$\det M = \sum_{\sigma \in S_N} \left(\text{sgn } \sigma \prod_{i=1}^N M_{i, \sigma(i)} \right)$$

Permutation group
of N objects

$$|S_N| = N!$$

$$3 \text{ dim} \Rightarrow \text{sgn } \sigma = \begin{cases} 1 & \text{cyclic} \\ -1 & \text{anticyclic} \end{cases}$$

$$|S_3| = 3! = 6$$

and specialize the notation to $\begin{pmatrix} A_1 & A_2 & A_3 \\ B_1 & B_2 & B_3 \\ C_1 & C_2 & C_3 \end{pmatrix}$ (I can write this because $\det M^T = \det M$)

$$\sigma = (i, j, k) \text{ means } \begin{matrix} 1 & 2 & 3 \\ \downarrow & \downarrow & \downarrow \\ i & j & k \end{matrix}$$

σ	$\prod M_{i, \sigma(i)} \Rightarrow$ New notation	$\text{sgn } \sigma$
$(1, 2, 3)$	$M_{11} M_{22} M_{33} \Rightarrow A_1 B_2 C_3$	$+1$
$(1, 3, 2)$	$M_{11} M_{23} M_{32} \Rightarrow A_1 B_3 C_2$	-1
	and so on :	

$$\epsilon_{ijk} A_i B_j C_k \parallel \vec{A} \cdot (\vec{B} \times \vec{C})$$

Linear transformations

$$\text{Matrix } M = \begin{pmatrix} M_{11} & M_{12} & M_{13} \\ M_{21} & M_{22} & M_{23} \\ M_{31} & M_{32} & M_{33} \end{pmatrix}$$

$$\begin{aligned} & \text{(Matrix multiplication)} \\ & (MN)_{ij} = M_{ik} N_{kj} \end{aligned}$$

So if M maps the set of coordinate axes to a new one, $\det M$ determines whether M changes the handedness of the axes.

$$\hat{e}'_1 = M \hat{e}_1$$

$$\hat{e}'_2 = M \hat{e}_2$$

$$\hat{e}'_3 = M \hat{e}_3$$

$$\hat{e}'_1 \cdot (\hat{e}'_2 \times \hat{e}'_3) = \underbrace{\vec{m}_1 \cdot (\vec{m}_2 \times \vec{m}_3)}_{\text{(columns)}} = \det M$$

Transformations that preserve the dot product (lengths, angles)

$$\begin{aligned} \vec{A} \cdot \vec{B} & \stackrel{?}{=} (O\vec{A}) \cdot (O\vec{B}) = (O\vec{A})_i (O\vec{B})_i = O_{ij} A_j O_{ik} B_k \\ & = A_j B_k O_{ij} O_{ik} = A_j B_k \underbrace{O_{ji} O_{ik}}_{\text{Want } \delta_{jk}} \end{aligned}$$

$$\text{Want } \delta_{jk} \Leftrightarrow \boxed{O^T O = \mathbb{1}}$$

Can be proven using only \hat{e}_i

$$\begin{aligned} \delta_{ij} & \stackrel{?}{=} \hat{e}_i \cdot \hat{e}_j = (O \hat{e}_i) \cdot (O \hat{e}_j) \\ & = (O \hat{e}_i)_k (O \hat{e}_j)_k \\ & = O_{ki} O_{kj} = O_{ik}^T O_{kj} \end{aligned}$$

$$\Leftrightarrow O^T O = \mathbb{1}$$

Multiplication with a vector

$$\text{ith entry} \rightarrow \begin{array}{|c|} \hline \text{green dot} \\ \hline \end{array} = \begin{array}{|c|c|} \hline \text{green row} & \text{green column} \\ \hline \end{array}$$

$$\text{Let } \vec{B} = M\vec{A}, \quad B_i = M_{ij} A_j$$

$$\text{If } \vec{C} = M\vec{e}_i$$

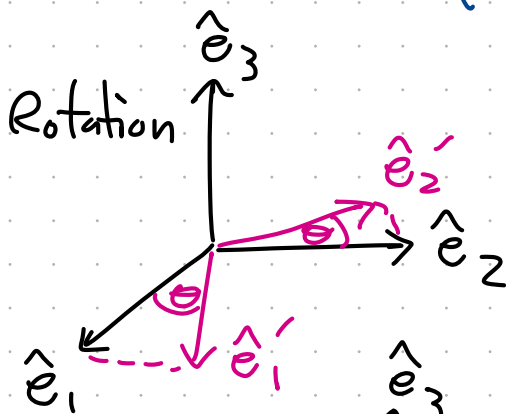
$$C_j = M_{ji} \Leftrightarrow \vec{C} = \vec{m}_i, \text{ the } i\text{th column of } M.$$

The definition of an orthogonal transformation (Equivalently)
 $O O^T = \mathbb{1}$

What is the det of an orthogonal transform.?

$$1 = \det \mathbb{1} = \det(O^T O) = \det(O^T) \det(O) \\ = \det(O) \det(O) = (\det O)^2$$

$$\Rightarrow \det O = \begin{cases} +1 & \text{Preserves the handedness} \\ -1 & \text{Switches " " " "} \end{cases}$$

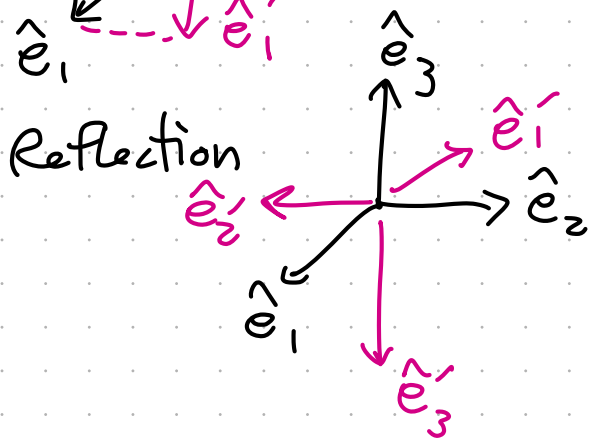


$$\hat{e}_1' = \cos \theta \hat{e}_1 + \sin \theta \hat{e}_2$$

$$\hat{e}_2' = -\sin \theta \hat{e}_1 + \cos \theta \hat{e}_2$$

$$O = \begin{pmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\det O = 1$$



$$O = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad \det O = -1$$

Let us agree that $\det(\vec{A} \vec{B} \vec{C})$ means the det of a 3×3 matrix with columns \vec{A} , \vec{B} , and \vec{C} .

Under the reflection

$$\left. \begin{aligned} \vec{A} &\mapsto \vec{A}' = -\vec{A} \\ \vec{B} &\mapsto \vec{B}' = -\vec{B} \\ \vec{C} &\mapsto \vec{C}' = -\vec{C} \end{aligned} \right\} \textcircled{1} \det(\vec{A}' \vec{B}' \vec{C}') = \det[(-\vec{A})(-\vec{B})(-\vec{C})] = -\det(\vec{A} \vec{B} \vec{C})$$

\Rightarrow Pseudoscalars change signs under reflection.

$$\textcircled{2} \text{ Let } \vec{A} \times \vec{B} = \vec{D}, \text{ then } \vec{A}' \times \vec{B}' = (-\vec{A}) \times (-\vec{B}) = \vec{D}$$

Pseudovectors remain unchanged under reflections

(A product of two vectors act like a scalar in the same way the dot product is a scalar)

Examples

Scalars	$\vec{A} \cdot \vec{B}$
Pseudo scalars	$\vec{A} \cdot (\vec{B} \times \vec{C})$
Vectors	$\vec{A}, \vec{A} \times (\vec{B} \times \vec{C})$
Pseudo vectors	$\vec{A} \times \vec{B}$

$\vec{F}_{\text{magnetic}} = q \underbrace{\vec{v}}_{\text{Vector}} \times \underbrace{\vec{B}}_{\text{Pseudo vector}}$ because \vec{B} is a cross product
 $\vec{B} = \vec{\nabla} \times \vec{A}$ (Curl of a vector potential)
 or $d\vec{B}(\vec{r}) = \frac{I d\vec{\ell} \times \vec{r} - \vec{r}_0}{|\vec{r} - \vec{r}_0|^3}$ from Biot-Savart

It is worth emphasizing that these notions are *only defined with respect to a given class of transformations*

(Here, they are orthogonal transformations in 3 dimensions)

As an example, unitary transformations are privileged in QM
 $U^\dagger U = U U^\dagger = \mathbb{1}$ where $U^\dagger = (U^T)^*$ ← complex conjugation

Tensors

Now we will look at matrices (and more generally multi-index objects) from the point of view of their transformation properties.

Outer product $\vec{A} \vec{B}^T = \begin{pmatrix} A_1 \\ A_2 \\ A_3 \end{pmatrix} (B_1, B_2, B_3) = \begin{pmatrix} A_1 B_1 & A_1 B_2 & A_1 B_3 \\ A_2 B_1 & A_2 B_2 & A_2 B_3 \\ A_3 B_1 & A_3 B_2 & A_3 B_3 \end{pmatrix}$

(Also called dyadics and can be written as $\vec{A} \otimes \vec{B}^T$ or $|A\rangle\langle B|$ for QM aficionados) or $T_{ij} = A_i B_j$ Example of a rank-2 tensor

Recall $A'_i = O_{ij} A_j$

$$T'_{ij} = O_{ik} O_{jl} A_k B_l = O_{ik} O_{jl} T_{kl} \quad \leftarrow \begin{array}{l} \text{want to have} \\ j \text{ as the last} \\ \text{index} \end{array}$$

$$= O_{ik} T_{kl} (O^T)_{lj} = (O T O^T)_{ij}$$

$$\Leftrightarrow T' = O T O^T = O T O^{-1}$$

So it transforms under O like an ordinary matrix.
But the transformation is **reducible**.

vector \mapsto vector' mixes all components = irreducible
scalar \mapsto scalar' ——— " ——— (trivially)

You will show in the homework a special case of the fact that any rank-2 tensor can be decomposed into 3 irreducible parts:

The scalar part: $T_{ij}^{(0)} = \frac{\text{Tr } T}{3} \delta_{ij}$

The vector part: $T_{ij}^{(1)} = \frac{T_{ij} - T_{ji}}{2}$

The rest: $T_{ij}^{(2)} = \frac{T_{ij} + T_{ji}}{2} - \frac{\text{Tr } T}{3} \delta_{ij}$

where the trace $\text{Tr } T = T_{ii} = T_{11} + T_{22} + T_{33}$

More generally, the components of a rank- n tensor transform as

$$T'_{i_1 i_2 \dots i_n} = \bigcirc_{i_1 j_1} \bigcirc_{i_2 j_2} \dots \bigcirc_{i_n j_n} T_{j_1 j_2 \dots j_n}$$

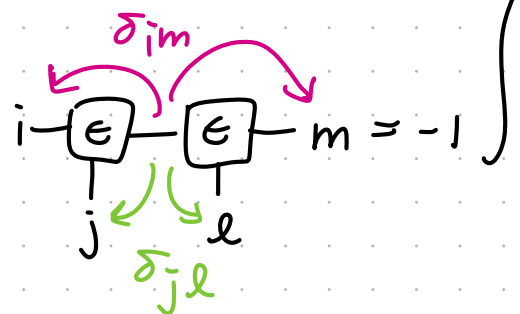
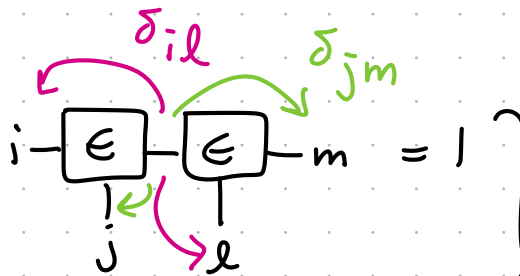
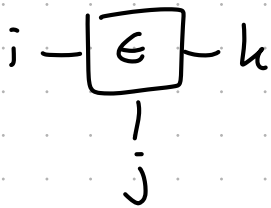
A contraction by setting a pair of indices to be the same turns a rank n tensor to a rank $(n-2)$ tensor

Example: Rank-2 $T_{ij} \mapsto T_{ii} = \text{Tr } T$ scalar (rank-0)

Rank-3 $T_{ijk} = A_i B_j C_k \mapsto T_{iik} = A_i B_i C_k$
 $T \mapsto (\vec{A} \cdot \vec{B}) \vec{C}$ vector

How to remember the product-of-Levi-Civita identity (Personal note)

Levi-Civita



$$\begin{aligned} &\Rightarrow \epsilon_{ijk} \epsilon_{klm} \\ &= \delta_{il} \delta_{jm} - \delta_{im} \delta_{jl} \end{aligned}$$

Can we use this graphical notation to contract any two indices

Ultimately, it's a contraction of two special 3-tensors. This kind of diagrammatical notation is used frequently in many-body quantum physics (and to some extent in QFT/HEP)