

Basics of Heat Transfer

Modes of heat transfer

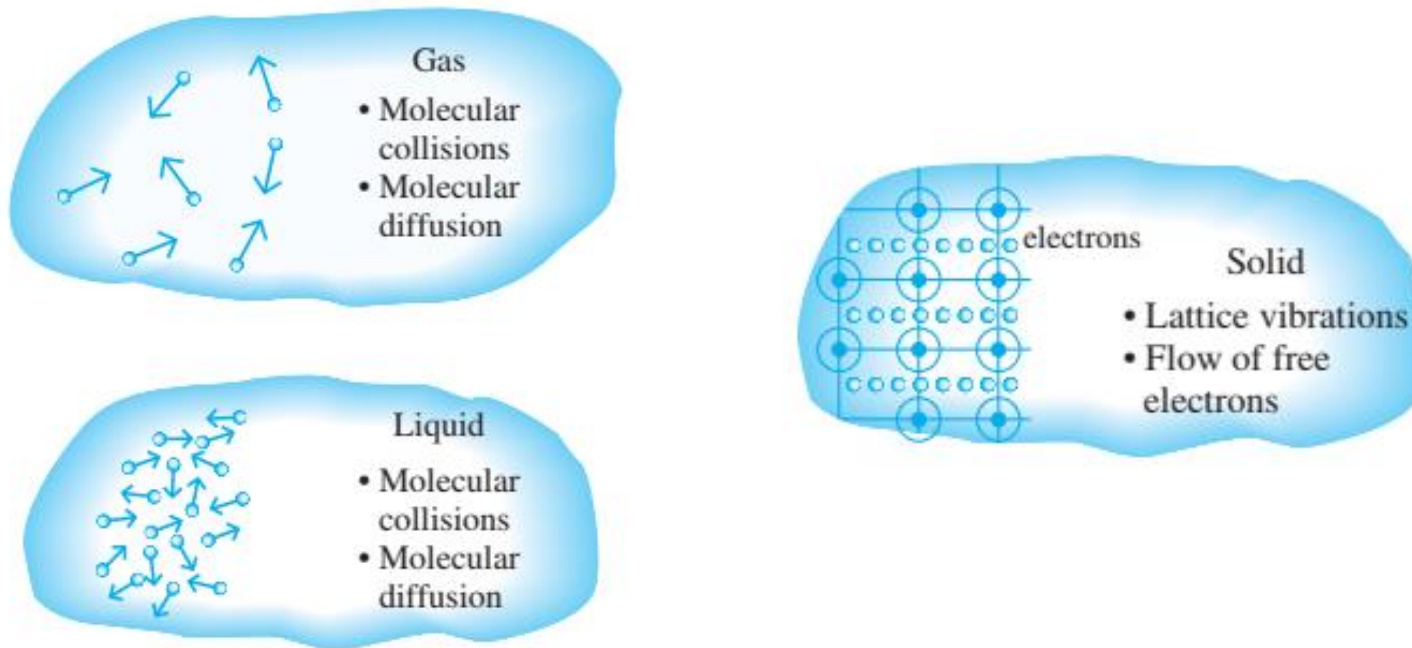
- we defined *heat* as the form of energy that can be transferred from one system to another as a result of temperature difference.
- Thermodynamics is concerned with the *amount* of heat transfer as a system undergoes a process from one equilibrium state to another.
- The science that deals with the determination of the *rates* of such energy transfers is the *heat transfer*.

Heat can be transferred in three different modes: *conduction*, *convection*, and *radiation*.

CONDUCTION

- **Conduction** is the transfer of energy from the more energetic particles of a substance to the adjacent less energetic ones as a result of interactions between the particles.

mechanisms of heat conduction in different phases of a substance :



Experiments have shown that :

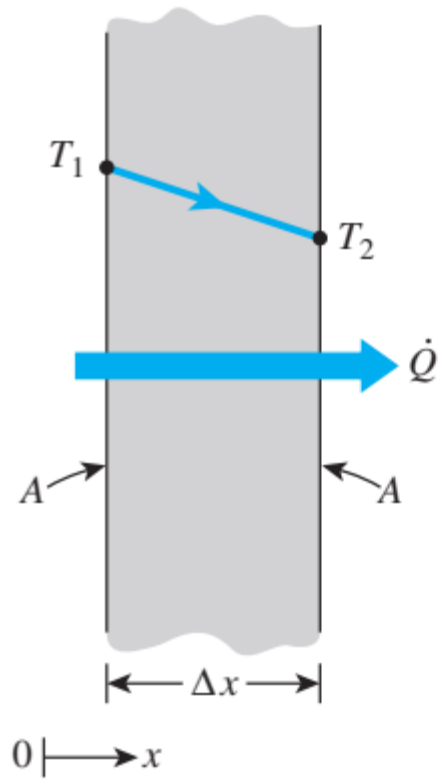


FIGURE 1–23

Heat conduction through a large plane wall of thickness Δx and area A .

Rate of heat conduction $\propto \frac{(\text{Area})(\text{Temperature difference})}{\text{Thickness}}$

$$\begin{aligned}\dot{Q}_{\text{cond}} &= kA \frac{T_1 - T_2}{\Delta x} \\ &= -kA \frac{\Delta T}{\Delta x}\end{aligned}$$

where the constant of proportionality k is the **thermal conductivity** of the material, which is a *measure of the ability of a material to conduct heat*

$$\dot{Q}_{\text{cond}} = -kA \frac{dT}{dx}$$

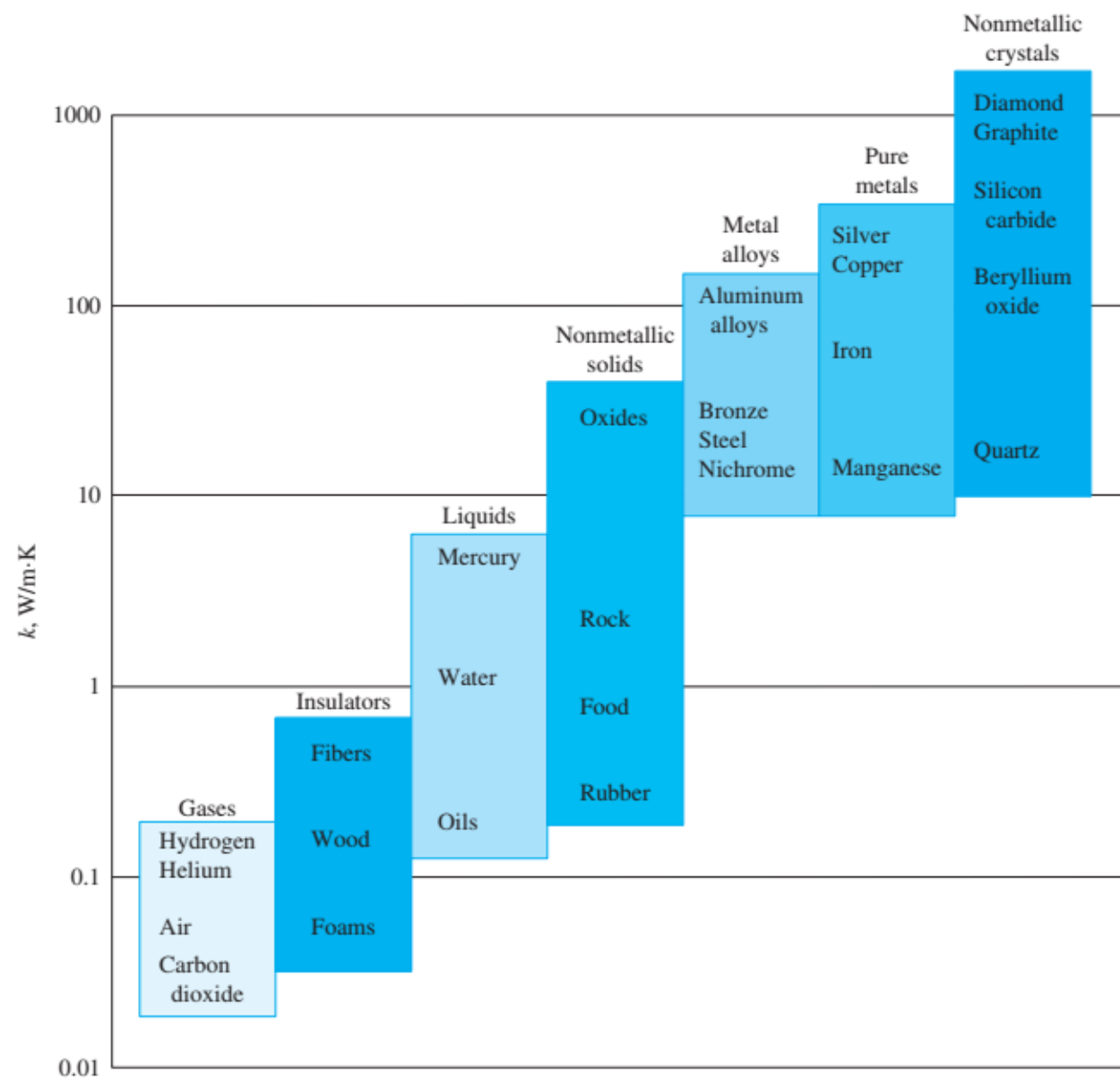
In the limiting case of $\Delta x \rightarrow 0$, the equation above reduces to the differential form:

$$\dot{Q}_{\text{cond}} = -kA \frac{dT}{dx} \quad \text{Fourier's law of heat conduction}$$

Here dT/dx is the **temperature gradient**, which is the slope of the temperature curve on a T - x diagram

Thermal Conductivity K

- the **thermal conductivity** of a material can be defined as *the rate of heat transfer through a unit thickness of the material per unit area per unit temperature difference*.
- The thermal conductivity of a material is a measure of the ability of the material to conduct heat.
- A high value of k indicates that the material is a good heat conductor, and a low value indicates that the material is a poor heat conductor or *insulator*.



Thermal Diffusivity

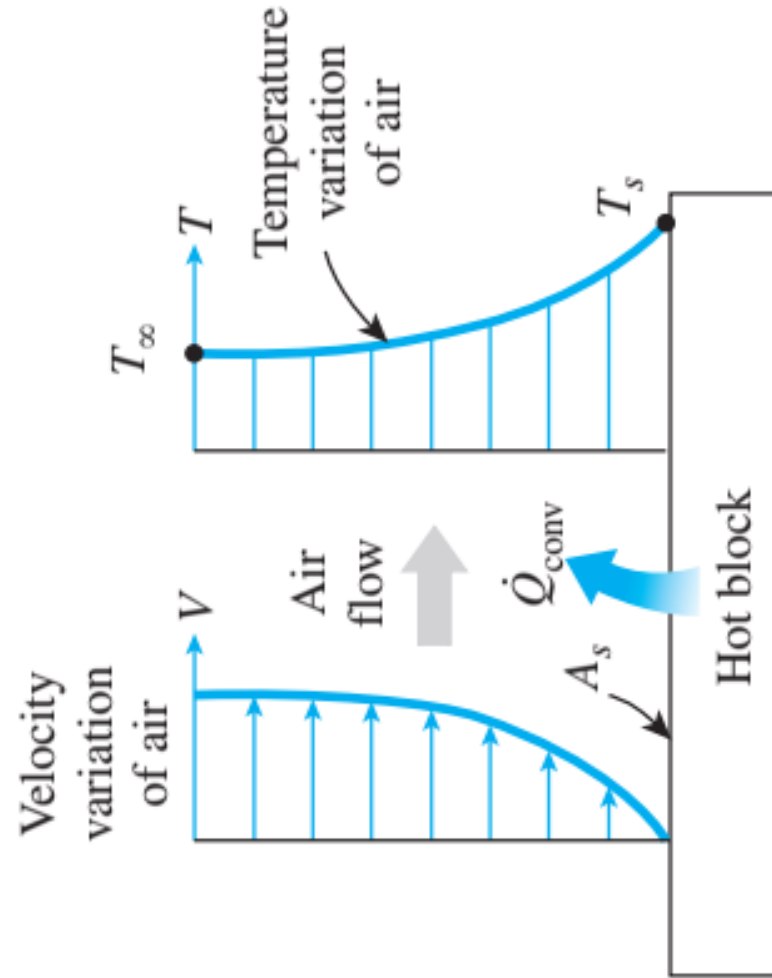
The material property that appears in the transient heat conduction analysis is the **thermal diffusivity**, which represents how fast heat diffuses through a material and is defined as

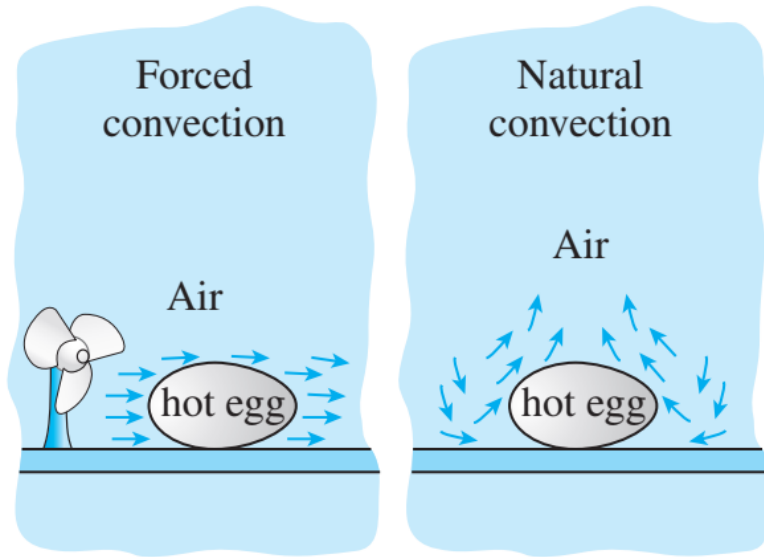
$$\alpha = \frac{\text{Heat conduction}}{\text{Heat storage}} = \frac{k}{\rho c_p}$$

The thermal conductivity k represents how well a material conducts heat, and the heat capacity ρc_p represents how much energy a material stores per unit volume.

CONVECTION

- **Convection** is the mode of energy transfer between a solid surface and the adjacent liquid or gas that is in motion,





- Convection is called **forced convection** if the fluid is forced to flow over the surface by external means such as a fan, pump, or the wind.
- Convection is called **natural** (or **free**) **convection** if the fluid motion is caused by buoyancy forces that are induced by density differences due to the variation of temperature in the fluid.

The rate of *convection heat transfer* is

$$\dot{Q}_{\text{conv}} = hA_s (T_s - T_{\infty}) \quad \text{Newton's law of cooling}$$

where, h is the *convection heat transfer coefficient*

A_s is the surface area through which convection heat transfer takes place,

T_s is the surface temperature, and

T_{∞} is the temperature of the fluid sufficiently far from the surface.

RADIATION

Radiation is the energy emitted by matter in the form of *electromagnetic waves* (or *photons*) as a result of the changes in the electronic configurations of the atoms or molecules

In heat transfer studies we are interested in *thermal radiation*, which is the form of radiation emitted by bodies because of their temperature. All bodies at a temperature above absolute zero emit thermal radiation

The maximum rate of radiation that can be emitted from a surface at a thermodynamic temperature T_s (in K) is given by the **Stefan–Boltzmann law** as

$$\dot{Q}_{\text{emit, max}} = \sigma A_s T_s^4$$

where $\sigma = 5.670 \times 10^{-8} \text{ W/m}^2\cdot\text{K}^4$ is the *Stefan–Boltzmann constant*.

The idealized surface that emits radiation at this maximum rate is called a **blackbody**, and the radiation emitted by a blackbody is called **blackbody radiation**

The radiation emitted by all real surfaces is less than the radiation emitted by a blackbody at the same temperature, and is expressed as

$$\dot{Q}_{\text{emit}} = \epsilon \sigma A_s T_s^4$$

where ϵ is the **emissivity** of the surface. The property emissivity, whose value is in the range 0 to 1, is a measure of how closely a surface approximates a blackbody for which $\epsilon=1$

When a surface of emissivity ϵ and surface area A_s at a *thermodynamic temperature* T_s is *completely enclosed* by a much larger (or black) surface at thermodynamic temperature T_{surr} separated by a gas (such as air) that does not intervene with radiation, the net rate of radiation heat transfer between these two surfaces is given by (Fig. 1–40)

$$\dot{Q}_{\text{rad}} = \epsilon \sigma A_s (T_s^4 - T_{\text{surr}}^4)$$

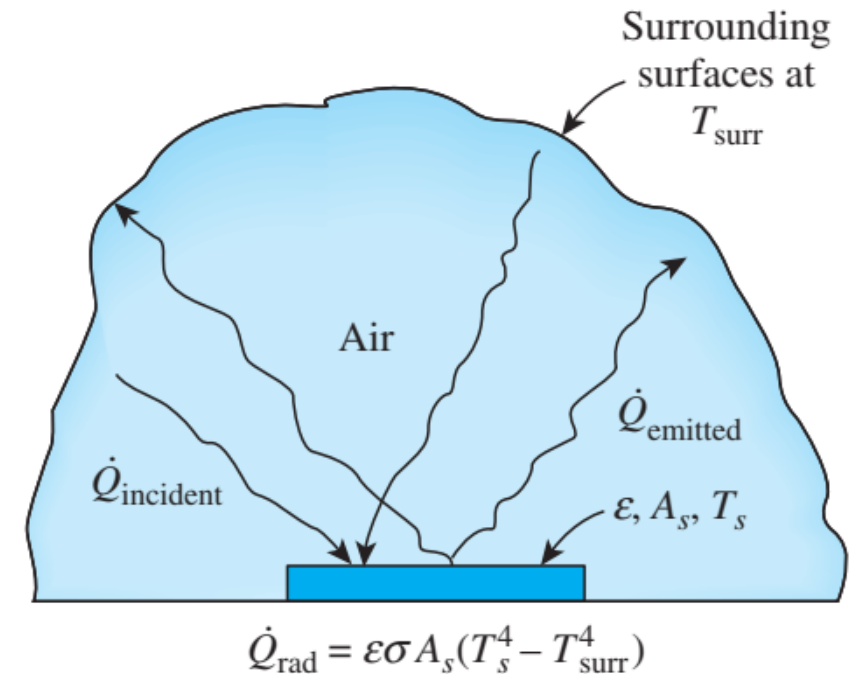
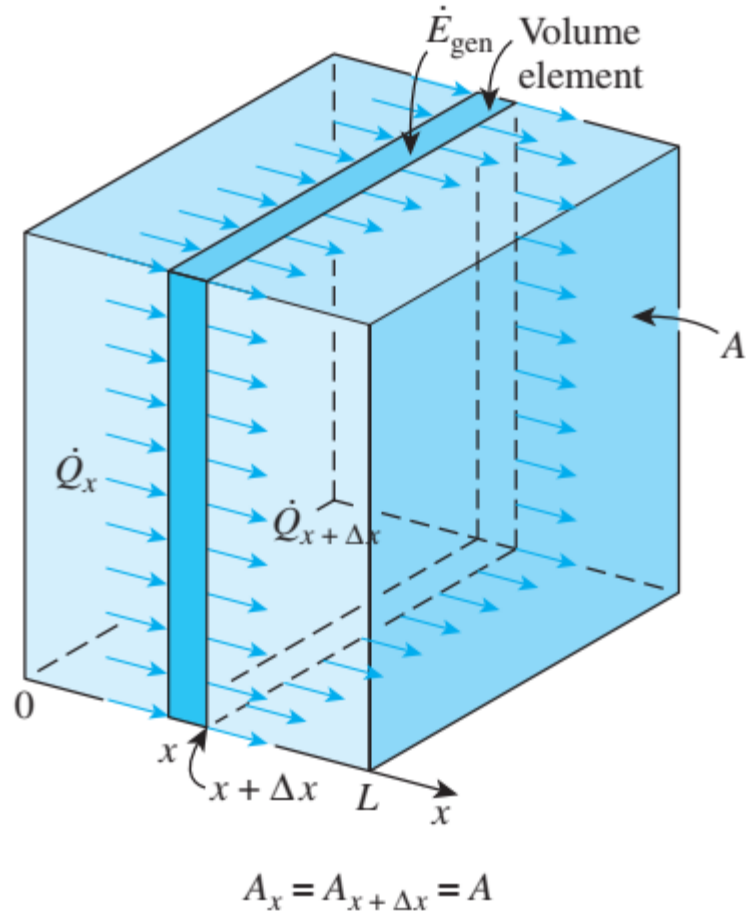


FIGURE 1–40

Radiation heat transfer between a surface and the surfaces surrounding it.

ONE-DIMENSIONAL HEAT CONDUCTION EQUATION



Consider a thin element of thickness Δx in a large plane wall, as shown.

Assume the density of the wall is ρ , the specific heat is c , and the area of the wall normal to the direction of heat transfer is A .

An *energy balance* on this thin element during a small time interval Δt can be expressed as

$$\left(\begin{array}{c} \text{Rate of heat} \\ \text{conduction} \\ \text{at } x \end{array} \right) - \left(\begin{array}{c} \text{Rate of heat} \\ \text{conduction} \\ \text{at } x + \Delta x \end{array} \right) + \left(\begin{array}{c} \text{Rate of heat} \\ \text{generation} \\ \text{inside the} \\ \text{element} \end{array} \right) = \left(\begin{array}{c} \text{Rate of change} \\ \text{of the energy} \\ \text{content of the} \\ \text{element} \end{array} \right)$$

$$\dot{Q}_x - \dot{Q}_{x+\Delta x} + \dot{E}_{\text{gen, element}} = \frac{\Delta E_{\text{element}}}{\Delta t}$$

Heat Conduction Equation in a Large Plane Wall

the change in the energy content of the element

$$\begin{aligned}\Delta E_{\text{element}} &= E_{t + \Delta t} - E_t \\ &= mc(T_{t + \Delta t} - T_t) \\ &= \rho c A \Delta x (T_{t + \Delta t} - T_t)\end{aligned}$$

the rate of heat generation within the element

$$\begin{aligned}\dot{E}_{\text{gen, element}} &= \dot{e}_{\text{gen}} V_{\text{element}} \\ &= \dot{e}_{\text{gen}} A \Delta x\end{aligned}$$

On substituting,

$$\dot{Q}_x - \dot{Q}_{x + \Delta x} + \dot{e}_{\text{gen}} A \Delta x = \rho c A \Delta x \frac{T_{t + \Delta t} - T_t}{\Delta t}$$

Dividing by $A\Delta x$ gives

$$-\frac{1}{A} \frac{\dot{Q}_{x+\Delta x} - \dot{Q}_x}{\Delta x} + \dot{e}_{\text{gen}} = \rho c \frac{T_{t+\Delta t} - T_t}{\Delta t}$$

Taking the limit as $\Delta x \rightarrow 0$ and $\Delta t \rightarrow 0$,

$$\begin{aligned} \lim_{\Delta x \rightarrow 0} \frac{\dot{Q}_{x+\Delta x} - \dot{Q}_x}{\Delta x} &= \frac{\partial \dot{Q}}{\partial x} \\ &= \frac{\partial}{\partial x} \left(-kA \frac{\partial T}{\partial x} \right) \quad \text{(From Fourier's law of heat conduction)} \end{aligned}$$

On substituting,

$$\frac{1}{A} \frac{\partial}{\partial x} \left(kA \frac{\partial T}{\partial x} \right) + \dot{e}_{\text{gen}} = \rho c \frac{\partial T}{\partial t}$$

The *one-dimensional transient* heat conduction equation in a plane wall becomes,

$$\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \dot{e}_{\text{gen}} = \rho c \frac{\partial T}{\partial t}$$

For constant thermal conductivity k ,

$$\frac{\partial^2 T}{\partial x^2} + \frac{\dot{e}_{\text{gen}}}{k} = \frac{1}{\alpha} \frac{\partial T}{\partial t}$$

Further simplifications:

- | | |
|--|--|
| (1) <i>Steady-state:</i>
($\partial/\partial t = 0$) | $\frac{d^2 T}{dx^2} + \frac{\dot{e}_{\text{gen}}}{k} = 0$ |
| (2) <i>Transient, no heat generation:</i>
($\dot{e}_{\text{gen}} = 0$) | $\frac{\partial^2 T}{\partial x^2} = \frac{1}{\alpha} \frac{\partial T}{\partial t}$ |
| (3) <i>Steady-state, no heat generation:</i>
($\partial/\partial t = 0$ and $\dot{e}_{\text{gen}} = 0$) | $\frac{d^2 T}{dx^2} = 0$ |

Consider a large plane wall of thickness $L = 0.2 \text{ m}$, thermal conductivity $k = 1.2 \text{ W/m}\cdot\text{K}$, and surface area $A = 15 \text{ m}^2$. The two sides of the wall are maintained at constant temperatures of $T_1 = 120^\circ\text{C}$ and $T_2 = 50^\circ\text{C}$, respectively. Determine (i) the variation of temperature within the wall and the value of temperature at $x = 0.1 \text{ m}$ and (ii) the rate of heat conduction through the wall under steady conditions.

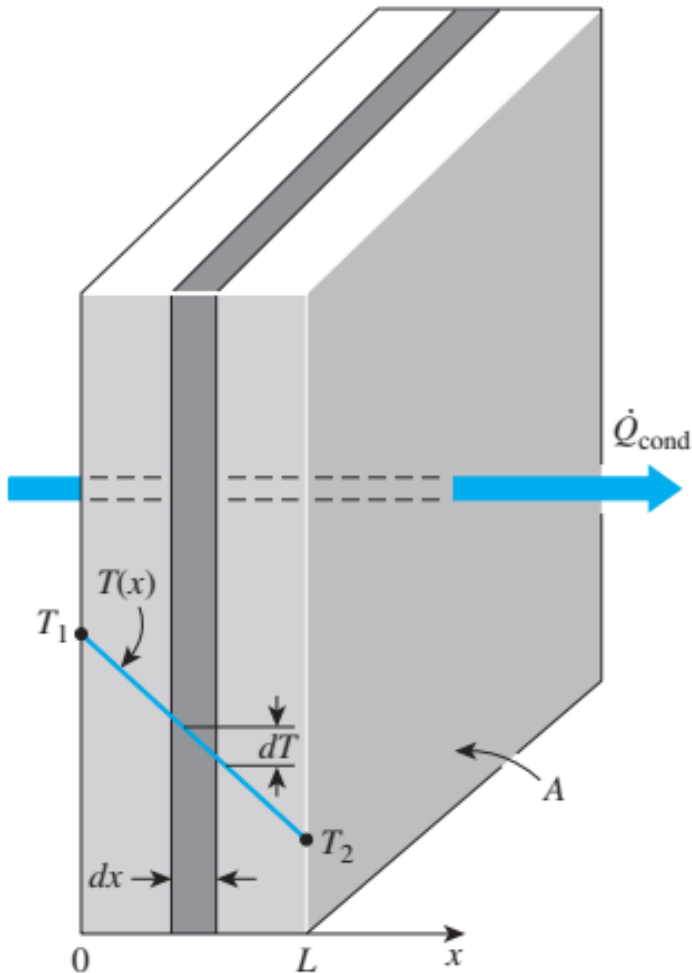
ONE DIMENSIONAL HEAT CONDUCTION EQUATION

$$\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \dot{e}_{\text{gen}} = \rho c \frac{\partial T}{\partial t}$$

THREE DIMENSIONAL HEAT CONDUCTION EQUATION

$$\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left(k \frac{\partial T}{\partial z} \right) + \dot{e}_{\text{gen}} = \rho c \frac{\partial T}{\partial t}$$

STEADY HEAT CONDUCTION IN PLANE WALLS



$$\left(\begin{array}{c} \text{Rate of} \\ \text{heat transfer} \\ \text{into the wall} \end{array} \right) - \left(\begin{array}{c} \text{Rate of} \\ \text{heat transfer} \\ \text{out of the wall} \end{array} \right) = \left(\begin{array}{c} \text{Rate of change} \\ \text{of the energy} \\ \text{of the wall} \end{array} \right)$$

$$\dot{Q}_{\text{in}} - \dot{Q}_{\text{out}} = \frac{dE_{\text{wall}}}{dt} \quad \text{0}$$

$$\dot{Q}_{\text{cond, wall}} = \text{constant}$$

Fourier's law of heat conduction for the wall can be expressed as

$$\dot{Q}_{\text{cond, wall}} = -kA \frac{dT}{dx}$$

On re-arranging the variables and integrating between 1 and 2,

$$\int_{x=0}^L \dot{Q}_{\text{cond, wall}} dx = - \int_{T=T_1}^{T_2} kA dT$$

$$\dot{Q}_{\text{cond, wall}} = kA \frac{T_1 - T_2}{L}$$

the rate of heat conduction through a plane wall is proportional to the average thermal conductivity, the wall area, and the temperature difference, but is inversely proportional to the wall thickness

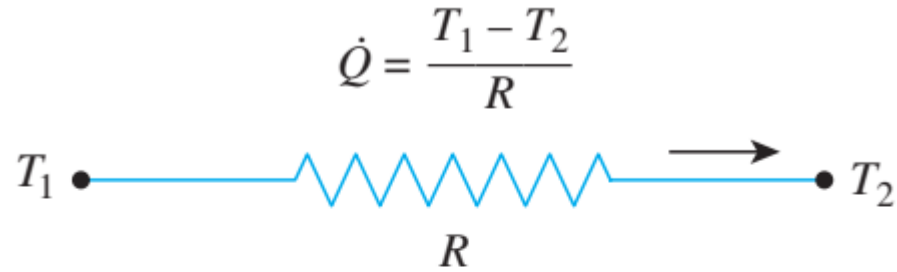
Thermal Resistance Concept

heat conduction through a plane wall can be rearranged as

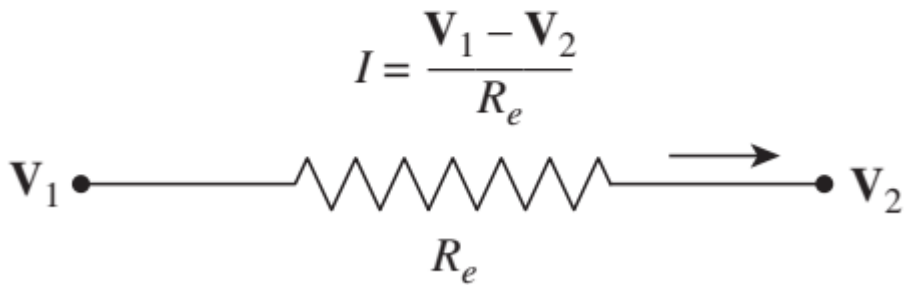
$$\dot{Q}_{\text{cond, wall}} = \frac{T_1 - T_2}{R_{\text{wall}}}$$

where, $R_{\text{wall}} = \frac{L}{kA}$ is the *thermal resistance* of the wall against heat conduction or simply the **conduction resistance** of the wall

Analogy between thermal and electrical resistance concepts



(a) Heat flow



(b) Electric current flow

Thermal resistance R can be conduction resistance, convective resistance or radiation resistance

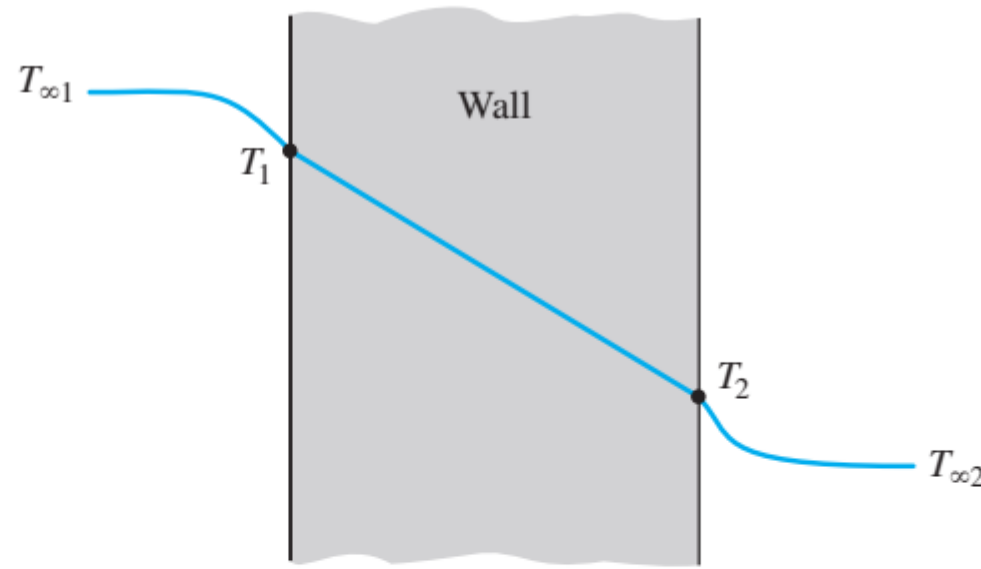
$$\dot{Q}_{\text{cond, wall}} = \frac{T_1 - T_2}{R_{\text{wall}}} \quad R_{\text{wall}} = \frac{L}{kA}$$

$$\dot{Q}_{\text{conv}} = \frac{T_s - T_{\infty}}{R_{\text{conv}}} \quad R_{\text{conv}} = \frac{1}{hA_s}$$

$$\dot{Q}_{\text{rad}} = \frac{T_s - T_{\text{surr}}}{R_{\text{rad}}} \quad R_{\text{rad}} = \frac{1}{h_{\text{rad}} A_s}$$

$$\text{Where, } h_{\text{rad}} = \frac{\dot{Q}_{\text{rad}}}{A_s(T_s - T_{\text{surr}})} = \varepsilon\sigma(T_s^2 + T_{\text{surr}}^2)(T_s + T_{\text{surr}})$$

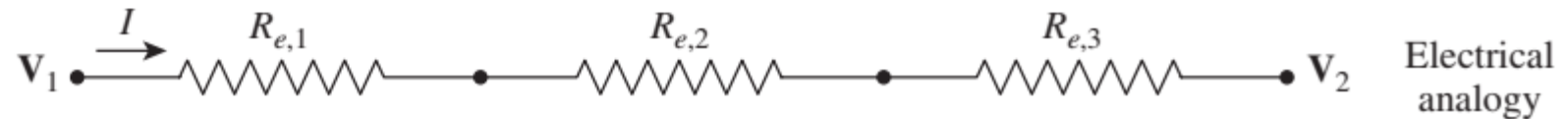
Thermal Resistance Network



$$\dot{Q} = \frac{T_{\infty 1} - T_{\infty 2}}{R_{\text{conv},1} + R_{\text{wall}} + R_{\text{conv},2}}$$



$$I = \frac{V_1 - V_2}{R_{e,1} + R_{e,2} + R_{e,3}}$$



$$\dot{Q} = \frac{T_{\infty 1} - T_{\infty 2}}{R_{\text{total}}} \quad \text{where, } R_{\text{total}} = R_{\text{conv},1} + R_{\text{wall}} + R_{\text{conv},2} = \frac{1}{h_1 A} + \frac{L}{kA} + \frac{1}{h_2 A}$$

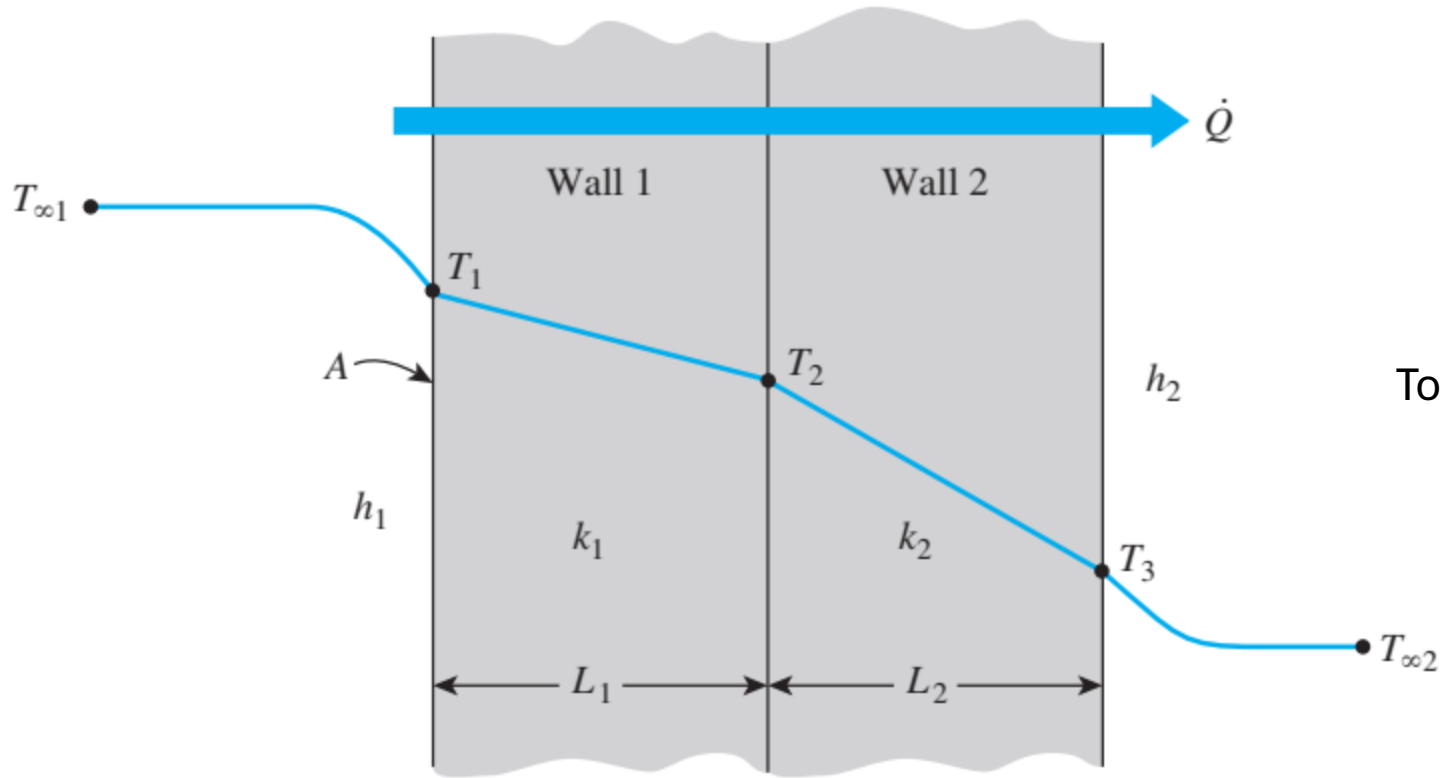
$$\left(\begin{array}{c} \text{Rate of} \\ \text{heat convection} \\ \text{into the wall} \end{array} \right) = \left(\begin{array}{c} \text{Rate of} \\ \text{heat conduction} \\ \text{through the wall} \end{array} \right) = \left(\begin{array}{c} \text{Rate of} \\ \text{heat convection} \\ \text{from the wall} \end{array} \right)$$

$$\dot{Q} = h_1 A(T_{\infty 1} - T_1) = kA \frac{T_1 - T_2}{L} = h_2 A(T_2 - T_{\infty 2})$$

which can be rearranged as

$$\begin{aligned} \dot{Q} &= \frac{T_{\infty 1} - T_1}{1/h_1 A} = \frac{T_1 - T_2}{L/kA} = \frac{T_2 - T_{\infty 2}}{1/h_2 A} \\ &= \frac{T_{\infty 1} - T_1}{R_{\text{conv}, 1}} = \frac{T_1 - T_2}{R_{\text{wall}}} = \frac{T_2 - T_{\infty 2}}{R_{\text{conv}, 2}} \end{aligned}$$

Once Q is known, this can be used to determine the intermediate temperatures T_1 or T_2 .

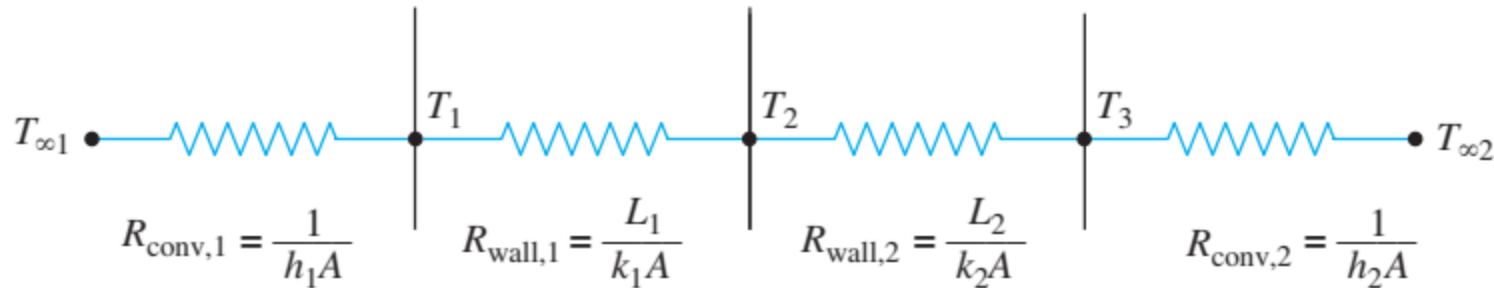


Heat transfer rate

$$\dot{Q} = \frac{T_{\infty 1} - T_{\infty 2}}{R_{\text{total}}}$$

Total thermal resistance

$$\begin{aligned} R_{\text{total}} &= R_{\text{conv},1} + R_{\text{wall},1} + R_{\text{wall},2} + R_{\text{conv},2} \\ &= \frac{1}{h_1 A} + \frac{L_1}{k_1 A} + \frac{L_2}{k_2 A} + \frac{1}{h_2 A} \end{aligned}$$

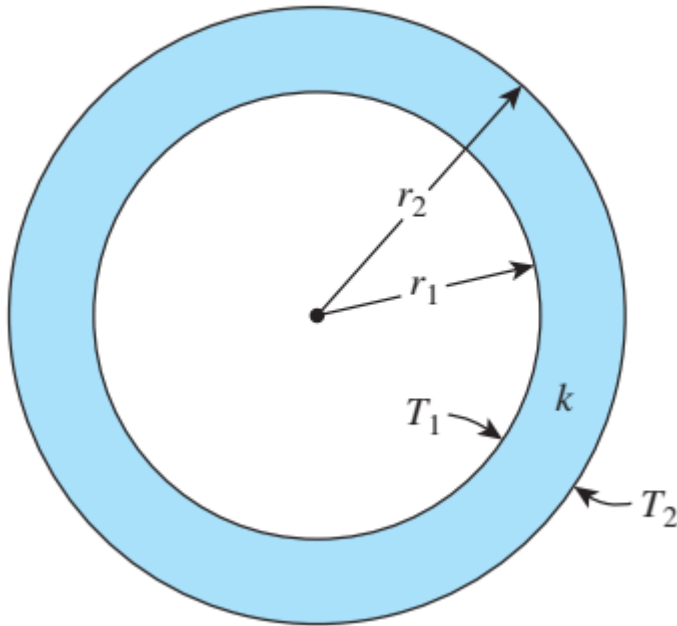


Consider a 0.8-m-high and 1.5-m-wide glass window with a thickness of 8 mm and a thermal conductivity of $k = 0.78 \text{ W/m}\cdot\text{K}$. Determine the steady rate of heat transfer through this glass window and the temperature of its inner surface for a day during which the room is maintained at 20°C while the temperature of the outdoors is 210°C . Take the heat transfer coefficients on the inner and outer surfaces of the window to be $h_1 = 10 \text{ W/m}^2\cdot\text{K}$ and $h_2 = 40 \text{ W/m}^2\cdot\text{K}$.

HEAT CONDUCTION IN CYLINDERS

Rate of heat conduction through a cylindrical pipe

$$\dot{Q}_{\text{cond, cyl}} = -kA \frac{dT}{dr}$$

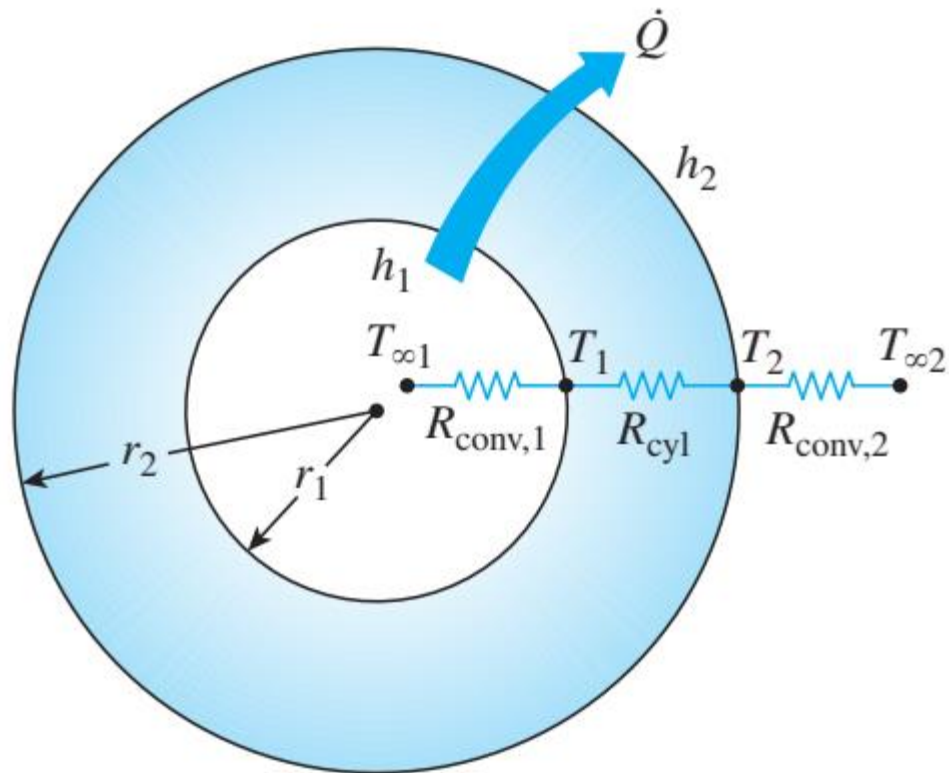


On rearranging and integrating between inner radius and outer radius

$$\int_{r=r_1}^{r_2} \frac{\dot{Q}_{\text{cond, cyl}}}{A} dr = - \int_{T=T_1}^{T_2} k dT$$

$$\dot{Q}_{\text{cond, cyl}} = 2\pi Lk \frac{T_1 - T_2}{\ln(r_2/r_1)}$$

Thermal Resistance Network in cylinder



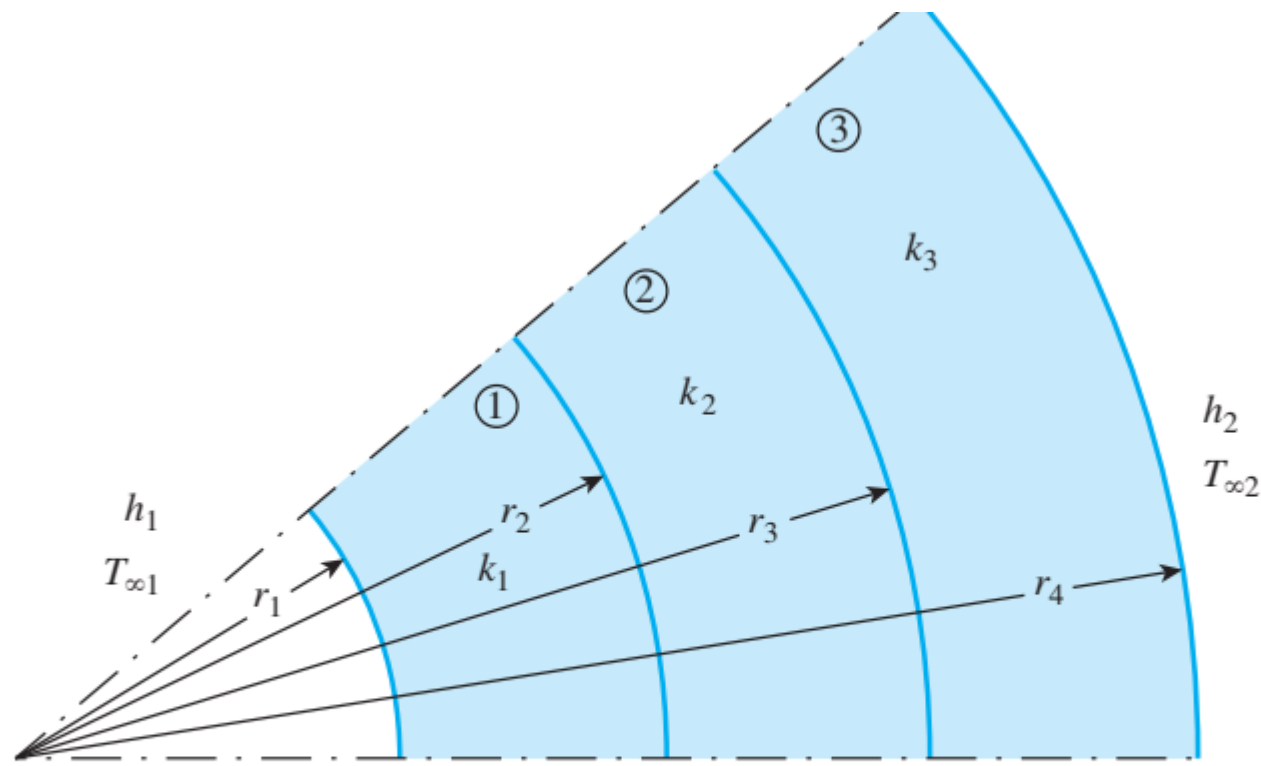
$$R_{\text{total}} = R_{\text{conv},1} + R_{\text{cyl}} + R_{\text{conv},2}$$

Heat transfer rate

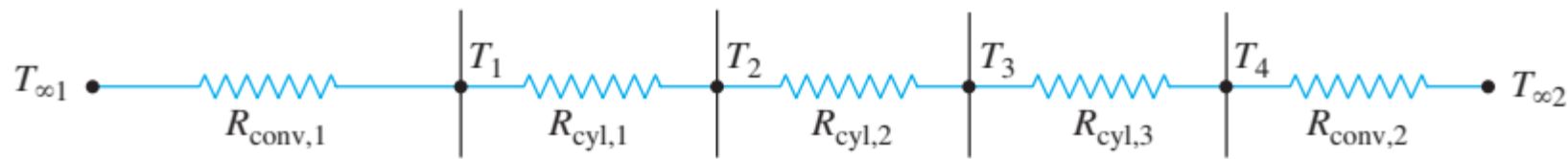
$$\dot{Q} = \frac{T_{\infty 1} - T_{\infty 2}}{R_{\text{total}}}$$

Total thermal resistance is

$$\begin{aligned} R_{\text{total}} &= R_{\text{conv},1} + R_{\text{cyl}} + R_{\text{conv},2} \\ &= \frac{1}{(2\pi r_1 L)h_1} + \frac{\ln(r_2/r_1)}{2\pi Lk} + \frac{1}{(2\pi r_2 L)h_2} \end{aligned}$$



$$\dot{Q} = \frac{T_{\infty 1} - T_{\infty 2}}{R_{\text{total}}}$$



$$\begin{aligned} R_{\text{total}} &= R_{\text{conv},1} + R_{\text{cyl},1} + R_{\text{cyl},2} + R_{\text{cyl},3} + R_{\text{conv},2} \\ &= \frac{1}{h_1 A_1} + \frac{\ln(r_2/r_1)}{2\pi L k_1} + \frac{\ln(r_3/r_2)}{2\pi L k_2} + \frac{\ln(r_4/r_3)}{2\pi L k_3} + \frac{1}{h_2 A_4} \end{aligned}$$

- 4 An insulated steel pipe carrying a hot liquid. Inner diameter of the pipe is 25 cm, wall thickness is 2 cm, thickness of insulation is 5 cm, temperature of hot liquid is 100°C , temperature of surrounding is 20°C , inside heat transfer co-efficient is $730\text{ W/m}^2\text{K}$ and outside heat transfer co-efficient is $12\text{ W/m}^2\text{K}$. Calculate the heat loss per metre length of the pipe.

Take $k_{\text{steel}} = 55\text{ W/mK}$, $k_{\text{insulating material}} = 0.22\text{ W/mK}$