

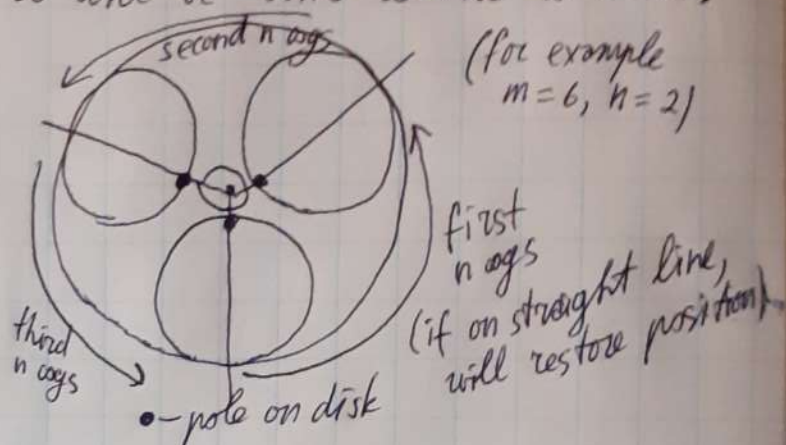
Problem 4.1 Hypotrochoids

a) Closed line: because if we take $\text{LCM}(m,n)$ - the amount of cogs when both rotations/revolutions result in same position, after this distance the center of disk will be at the same position, and pole position on it will be same relative to center.

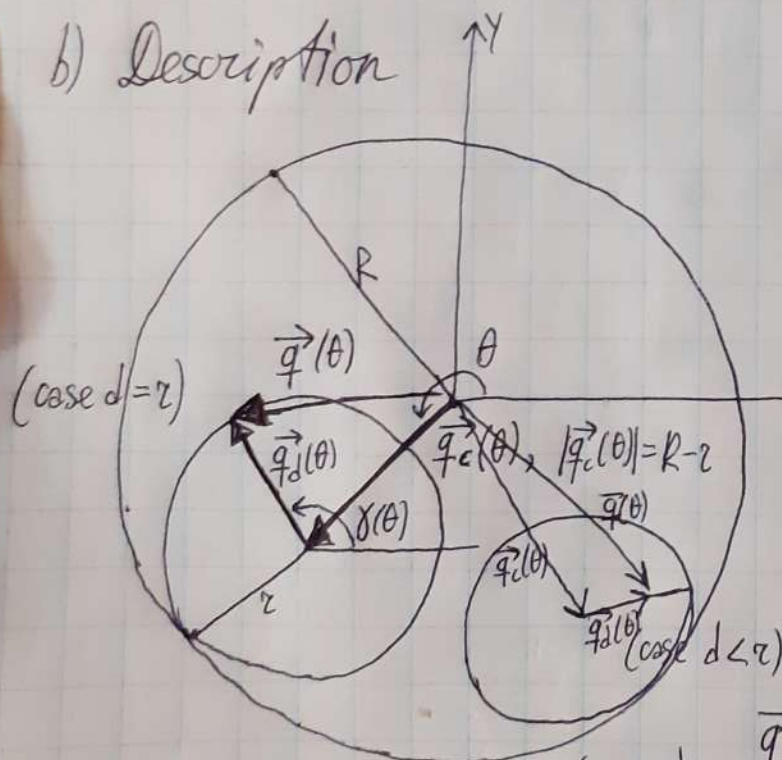
$$N_{\text{rev}} = \frac{\text{LCM}(m,n)}{m}$$

$$\text{Symmetry} = \frac{\text{LCM}(m,n)}{n} \text{-fold}$$

for example



b) Description



m, n cogs measure circumference,
 $\frac{m}{n} = \frac{2\pi R}{2\pi r} = \frac{R}{r} \rightarrow r = \frac{n}{m} R.$

d is distance from point $\vec{q}_c(\theta)$ to $\vec{q}(\theta)$.

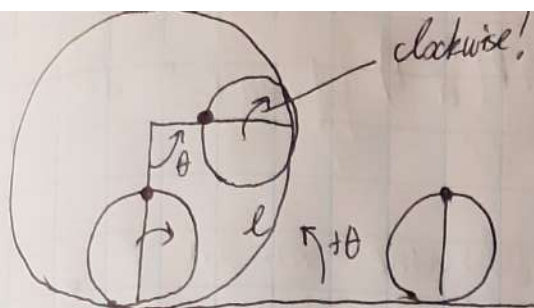
$$\rho = \frac{d}{r},$$

$$\text{So } d = \rho r = \rho \frac{n}{m} R.$$

$$\vec{q}(\theta) = \vec{q}_c(\theta) + \vec{q}_d(\theta),$$

$$\text{here } \vec{q}_c(\theta) = (R-r) \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix} = R \left(1 - \frac{n}{m}\right) \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix}$$

$\vec{q}_d(\theta) = d \begin{pmatrix} \cos \gamma \\ \sin \gamma \end{pmatrix}$, γ is rotation relative to disk, must find its relation to θ .



disk moved $l = \theta R$,

$$\text{total rot.} = \frac{l}{2\pi r} = \frac{\theta R}{2\pi r}, \text{ angle is}$$

$$\frac{\theta R}{2\pi r} \cdot (-2\pi) = -\frac{\theta R}{r}$$

Since not straight line, but "curved"

for θ , total is $-\frac{\theta R}{r} + \theta = \theta(1 - \frac{R}{r}) = \gamma$.

So $\vec{q}(\theta) = R(1 - \frac{n}{m}) \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix} + d \begin{pmatrix} \cos \theta(1 - \frac{R}{r}) \\ \sin \theta(1 - \frac{R}{r}) \end{pmatrix} = R(1 - \frac{n}{m}) \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix} + d \begin{pmatrix} \cos \theta(1 - \frac{m}{n}) \\ \sin \theta(1 - \frac{m}{n}) \end{pmatrix}$

If disk center 2 equation is same with negative θ .

General form: $\vec{q}(\theta) = R(1 - \frac{n}{m}) \begin{pmatrix} \cos(\theta_0 + \theta) \\ \sin(\theta_0 + \theta) \end{pmatrix} + \rho \frac{n}{m} R \begin{pmatrix} \cos \langle (1 - \frac{m}{n})(\theta + \theta_0) \rangle \\ \sin \langle (1 - \frac{m}{n})(\theta + \theta_0) \rangle \end{pmatrix}$

R, n, m, ρ - parameters.

θ_0, θ_0 - initial conditions (0 below).

c) Pyth. program in .pdf \rightarrow

d) length can be evaluated for given θ . Length of one closed line means $\theta \in [0, N_{\text{rev}}] = [0, \frac{\text{LCM}(m, n)}{m}]$.

$$L = \int_{t_{\text{start}}}^{t_{\text{end}}} |\dot{\vec{q}}(t)| dt, \quad \vec{q}(\theta) = R(1 - \frac{n}{m}) \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix} + \rho \frac{n}{m} R \begin{pmatrix} \cos \theta(1 - \frac{m}{n}) \\ \sin \theta(1 - \frac{m}{n}) \end{pmatrix}$$

transform to θ

$$\dot{\vec{q}}(\theta) = R(1 - \frac{n}{m}) \begin{pmatrix} -\sin \theta \cdot \dot{\theta} \\ \cos \theta \cdot \dot{\theta} \end{pmatrix} + \rho \frac{n}{m} R \begin{pmatrix} -\sin \theta(1 - \frac{m}{n}) \cdot \dot{\theta} \\ (\cos \theta(1 - \frac{m}{n})) \cdot \dot{\theta} \end{pmatrix} \cdot \frac{1}{(1 - \frac{m}{n})}$$

$$= R(1 - \frac{n}{m}) \begin{pmatrix} -\sin \theta \cdot \dot{\theta} \\ \cos \theta \cdot \dot{\theta} \end{pmatrix} + \rho R \left(\frac{n}{m} - 1 \right) \begin{pmatrix} -\sin \langle \theta(1 - \frac{m}{n}) \rangle \cdot \dot{\theta} \\ \cos \langle \theta(1 - \frac{m}{n}) \rangle \cdot \dot{\theta} \end{pmatrix} =$$

$$= R \dot{\theta} \left(\frac{n}{m} - 1 \right) \begin{pmatrix} \sin \theta - \rho \sin \langle \theta(1 - \frac{m}{n}) \rangle \\ -\cos \theta + \rho \cos \langle \theta(1 - \frac{m}{n}) \rangle \end{pmatrix}, \quad |\dot{\vec{q}}(\theta)| = R |\dot{\theta}| \left(\frac{m}{n} - 1 \right) \sqrt{\frac{\sin^2 \theta + \rho^2 \sin^2 \langle \dots \rangle + \dots}{\dots}} =$$

$$= R \dot{\theta} \left(\frac{m}{n} - 1 \right) \sqrt{1 + \rho^2 - 2\rho \cos(\theta - \theta(1 - \frac{m}{n}))} = R \dot{\theta} \left(\frac{m}{n} - 1 \right) \sqrt{1 + \rho^2 - 2\rho \left(\theta \frac{m}{n} \right)}$$

$L = \int_0^{\theta_{\text{max}} = \frac{\text{LCM}(m, n)}{m}, 2\pi} R \left(\frac{m}{n} - 1 \right) \sqrt{1 + \rho^2 - 2\rho \left(\theta \frac{m}{n} \right)} d\theta$. Check for $\rho=1$ (point on edge of disk):

```

# usage:
# launch in command terminal
# $ python hypotrochoid.py <m> <n> <d/r>

# specifying sizes of large and small circles, and ratio where pole on small circle is

import sys
import numpy as np
import matplotlib.pyplot as plt
import math

R=1

m=int(sys.argv[1])
n=int(sys.argv[2])
ro=float(sys.argv[3])

def LCM(m,n):
    return m*n/GCD(m,n)

def GCD(m,n):
    if n==0:
        return m
    else:
        return GCD(n, m%n)

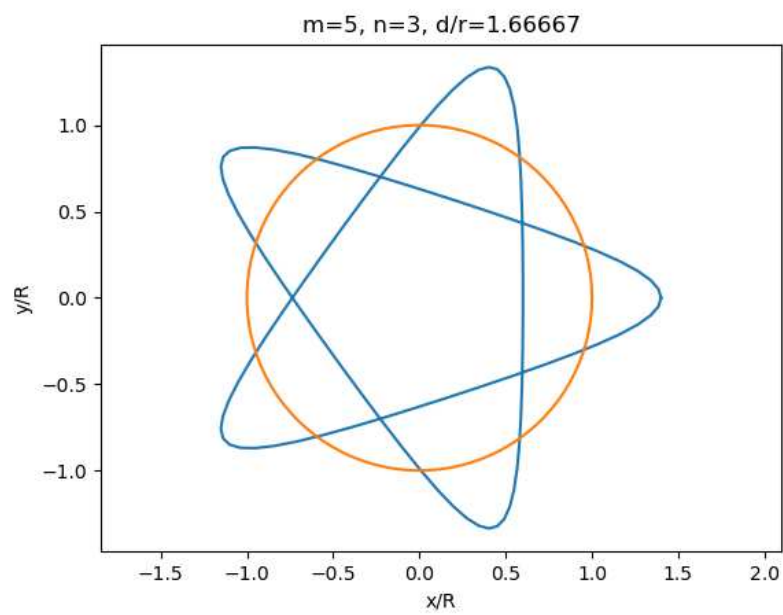
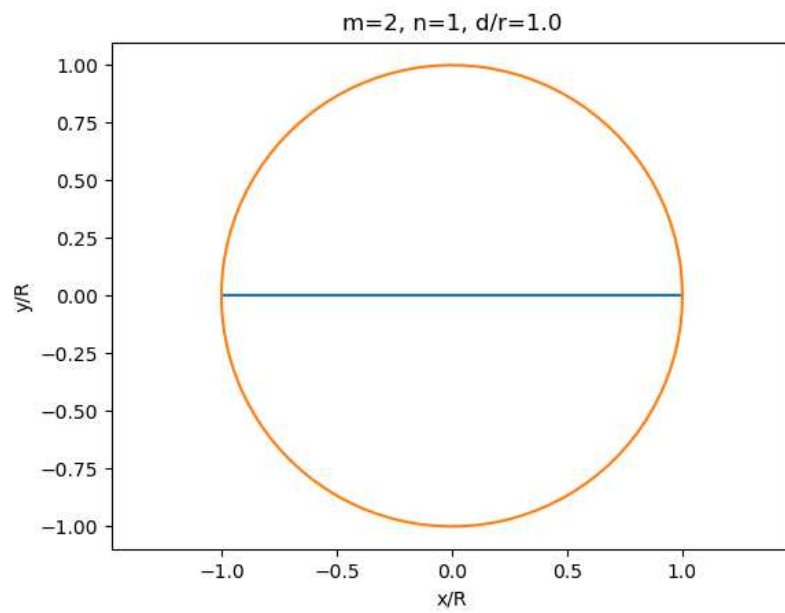
theta = np.linspace(0, LCM(m,n)/m * 2*math.pi, 100)
x = R*(1-n/m)*np.cos(theta) + ro*(n/m)*R*np.cos((1-m/n)*theta)
y = R*(1-n/m)*np.sin(theta) + ro*(n/m)*R*np.sin((1-m/n)*theta)

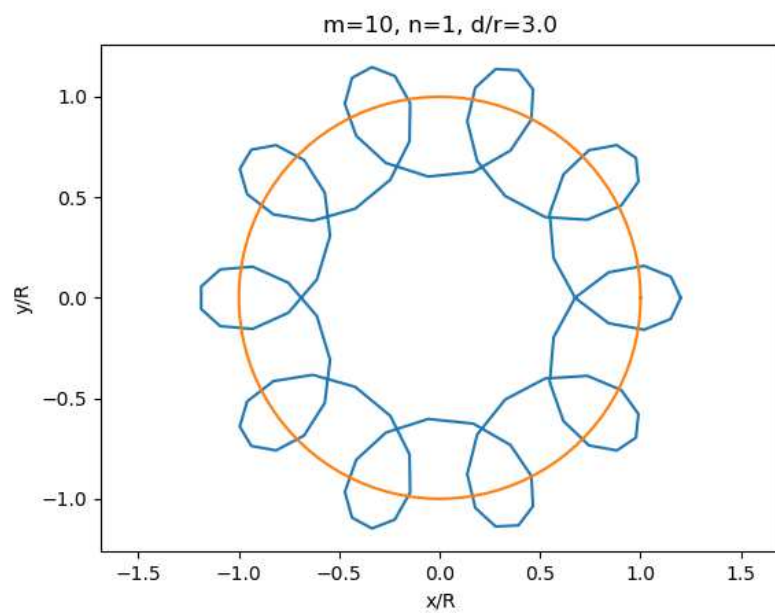
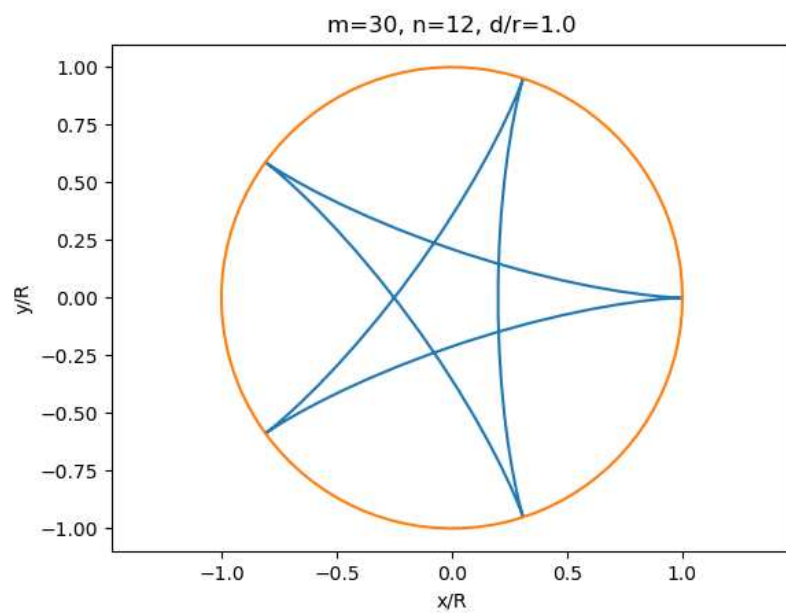
plt.plot(x, y)

# external circle to compare
theta = np.linspace(0, 2*math.pi, 100)
x = R*np.cos(theta)
y = R*np.sin(theta)
plt.plot(x, y)
plt.axis('equal')
plt.ylabel('y/R')
plt.xlabel('x/R')
plt.title("m=" + str(m) + ", n=" + str(n) + ", d/r=" + str(ro))

plt.show()

```

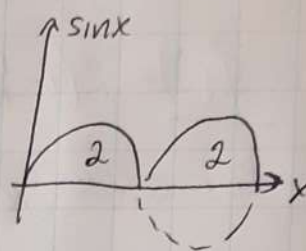




$$L_{p=1} = R \left(\frac{m}{n} - 1 \right) \int_0^{\theta_{\max}} \sqrt{2(1 - \cos(\frac{\theta m}{n})} d\theta = \sqrt{2} R \left(\frac{m}{n} - 1 \right) \int_0^{\theta_{\max}} \sqrt{2 \sin^2(\frac{\theta m}{2n})} d\theta =$$

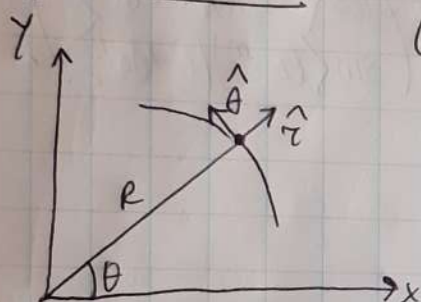
$$= 2R \left(\frac{m}{n} - 1 \right) \int_0^{\theta_{\max}} |\sin(\frac{\theta m}{2n})| d\theta. \text{ Assume } m=2n$$

$$\theta_{\max} = \frac{L(N(2n, n))}{2n} \approx 2\pi = 2\pi, \quad L = 2R \int_0^{2\pi} |\sin \theta| d\theta = 8R.$$



6.1

Problem 4.2 Circular track ($R = \text{const}$)



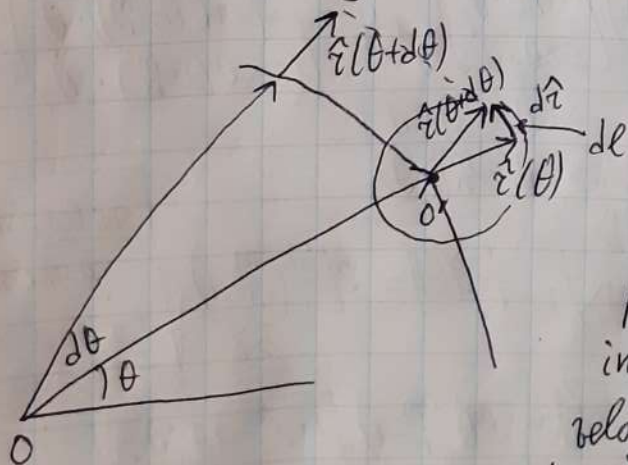
$$\hat{r} = \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix} \quad \hat{\theta} = \begin{pmatrix} -\sin \theta \\ \cos \theta \end{pmatrix}$$

(Since $\hat{r}, \hat{\theta}$ make ONB, $\hat{r} \hat{\theta} = -\sin \cos + \sin \cos = 0$,
 $\hat{r} \hat{r} = \cos^2 + \sin^2 = 1$
 $\hat{\theta} \hat{\theta} = \sin^2 + \cos^2 = 1$
 $\hat{r}, \hat{\theta}$ orthogonal vectors

$$a) \quad \frac{\partial}{\partial \theta} \hat{r} = \frac{d}{d\theta} \hat{r} = \begin{pmatrix} -\sin \theta \\ \cos \theta \end{pmatrix} = \hat{\theta}$$

$$\frac{\partial^2}{\partial \theta^2} \hat{r} = \frac{d}{d\theta} (\hat{\theta}) = \begin{pmatrix} -\cos \theta \\ -\sin \theta \end{pmatrix} = -\hat{r}.$$

Velocity is directed as $\hat{\theta}$, so tangentially, acceleration as $-\hat{r}$ to center.
 Geometry:



In the limit (\hat{r}) becomes \perp ,
 and since it's on unit circle
 $0' \left| \frac{d\hat{r}}{d\theta} \right| \approx \left| \frac{d\ell}{d\theta} \right| = |\hat{r}(\theta)| = 1.$

For $\hat{\theta}$ similar reasoning but instead of this plot - plot velocity vectors in their own "space", and $(\hat{\theta}) = \hat{r}$ will be \perp to $\hat{\theta}$, so $\uparrow \uparrow$ to original $\hat{r}, \hat{\theta}$.
 (can also say - each different. rotates $\frac{\pi}{2}$)

6.2

7.1

b) $\vec{q}(t) = R \hat{e}(\theta(t))$, $\dot{\vec{q}}, \ddot{\vec{q}}?$

Polaric way $\vec{q}(\theta(t)) = \frac{d\vec{q}}{d\theta} \cdot \frac{d\theta}{dt} = \frac{d}{d\theta}(R \cdot \hat{e}(\theta)) \cdot \dot{\theta} = R \hat{\theta} \dot{\theta} =$
 $\vec{q}'(\theta(t)) = (R \hat{\theta}) = R(\ddot{\theta} \hat{e} + \dot{\theta}(\hat{e}'(t))) = R(\ddot{\theta} \hat{e} + \dot{\theta} \frac{d\hat{e}}{d\theta} \cdot \frac{d\theta}{dt}) =$
 $= R(\ddot{\theta} \hat{e} + \dot{\theta}(-\hat{e}) \dot{\theta}) = R(\ddot{\theta} \hat{e} - (\dot{\theta})^2 \hat{e}) = R \ddot{\theta} \hat{e} - R(\dot{\theta})^2 \hat{e}.$

Cartesian way $\vec{q}(\theta(t)) = \begin{pmatrix} R \cos \theta(t) \\ R \sin \theta(t) \end{pmatrix}$, $\vec{q}'(\theta(t)) = \begin{pmatrix} -R \sin \theta(t) \cdot \dot{\theta} \\ R \cos \theta(t) \cdot \dot{\theta} \end{pmatrix} = R \dot{\theta} \hat{\theta}$ exactly
 $\vec{q}''(\theta(t)) = R \begin{pmatrix} -\cos \theta(t) \cdot (\dot{\theta})^2 - \sin \theta(t) \ddot{\theta} \\ -\sin \theta(t) \cdot (\dot{\theta})^2 + \cos \theta(t) \ddot{\theta} \end{pmatrix} = \underbrace{-R(\dot{\theta})^2 \hat{e}}_{\text{split}} + \underbrace{R \ddot{\theta} \hat{\theta}}_{\text{same}}.$

c) $\theta(t) = \omega t$, $\dot{\theta} = \omega$, $\ddot{\theta} = 0$.

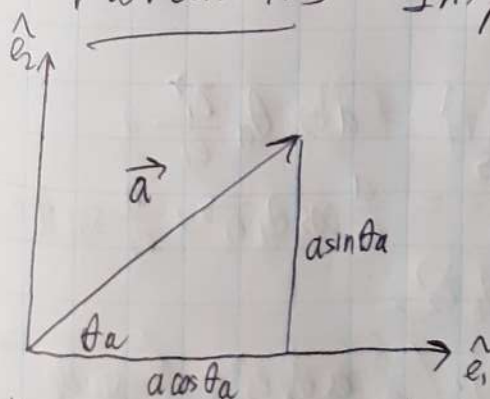
7.2

$\vec{q}'' = \underbrace{R \cdot 0 \cdot \hat{e}}_0 - \underbrace{R(\dot{\theta})^2 \hat{e}}_{\omega^2 R \hat{e}} = -\omega^2 R \hat{e}$

$\vec{q}' \cdot \vec{q}'' = R \dot{\theta} \hat{\theta} \cdot (-\omega^2 R \hat{e}) = -\omega^2 R \dot{\theta} (\hat{\theta} \cdot \hat{e}) = 0 \rightarrow \text{orthogonal}$
 Abs. value of velocity.

$\frac{d(\vec{v}^2)}{dt} = 2\vec{v} \cdot \frac{d\vec{v}}{dt} = 2\vec{q}' \cdot \vec{q}'' = 0$, $(\vec{v})^2 = v^2 \text{ const}$, $\boxed{v \text{ const.}}$
 $|\vec{q}'| = \omega^2 R$ also const,
 $|\vec{v}| = |R \dot{\theta} \hat{\theta}| = R|\omega| = \sqrt{\omega^2 R^2} =$
 $\boxed{|\vec{q}'| = \frac{v^2}{R} \text{ all const.}}$ $\sqrt{|\vec{q}'| R}$

Problem 4.3 In. product



a) Let $\vec{b} = b \cos \theta_b \hat{e}_1 + b \sin \theta_b \hat{e}_2$,

$\vec{a} \cdot \vec{b} = (a \cos \theta_a \hat{e}_1 + a \sin \theta_a \hat{e}_2) \cdot (b \cos \theta_b \hat{e}_1 + b \sin \theta_b \hat{e}_2) =$
 $= a b \cos \theta_a \cos \theta_b + a b \sin \theta_a \sin \theta_b =$

$= a b \cos(\theta_a - \theta_b)$, [formula for cos we covered when Euler identity & isomorphism on \mathbb{C}]

b) $\vec{a} = \begin{pmatrix} a \cos \theta_a \\ a \sin \theta_a \end{pmatrix} = \begin{pmatrix} a_1 \\ a_2 \end{pmatrix}$

$\vec{b} = \begin{pmatrix} b \cos \theta_b \\ b \sin \theta_b \end{pmatrix} = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}$

then $\vec{a} \cdot \vec{b} = \underline{a_1 b_1 + a_2 b_2}$.

Since does not depend, it is enough to calculate $\sum_{i=1}^n a_i b_i$ in any orthon. basis, and then $\sum a_i^{(e)} b_i^{(e)} = \sum a_i^{(n)} b_i^{(n)}$ etc.

c) Let $\theta_{ai} = \angle(\vec{a}, \hat{e}_i)$, $\theta_{bi} = \angle(\vec{b}, \hat{e}_i)$

$$\vec{a} \cdot \vec{b} = \sum_{i=1}^D (a \cos \theta_{ai} \cdot \hat{e}_i) \cdot \sum_{i=1}^D (b \cos \theta_{bi} \hat{e}_i) = \sum_{i,j=1}^D a b \cos \theta_{ai} \cos \theta_{bj} \delta_{ij}$$

$$= \sum_{i=1}^D a b \cos \theta_{ai} \cos \theta_{bi}$$

Denoting $a \cos \theta_{ai} = a_i$,
 $b \cos \theta_{bi} = b_i$,

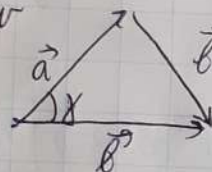
$$\vec{a} \cdot \vec{b} = \begin{pmatrix} a_1 \\ \vdots \\ a_D \end{pmatrix} \cdot \begin{pmatrix} b_1 \\ \vdots \\ b_D \end{pmatrix} = \sum_{i=1}^D a_i b_i$$

8.1

d) Special bases. (show on \mathbb{R}^2 example)

See Chapter 2 exploration p. 35

I rely that on \mathbb{R}^D always $\vec{a} \cdot \vec{b} = |\vec{a}| |\vec{b}| \cos \angle(\vec{a}, \vec{b})$, from cosine Law



$$|\vec{b} - \vec{a}|^2 = (\vec{b} - \vec{a})^2 = b^2 + a^2 - 2\vec{a} \cdot \vec{b} = a^2 + b^2 - 2ab \cos x$$

but also

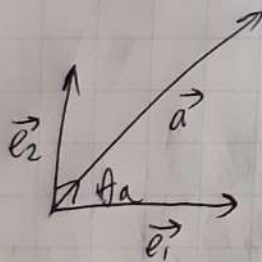
$$a^2 + b^2 - 2ab \cos x \quad \swarrow \quad \vec{a} \cdot \vec{b} = ab \cos x$$

Orthogonal, not orthonormal basis!

$$\exists \vec{e}_1, \vec{e}_2, \vec{e}_i \cdot \vec{e}_j = \delta_{ij}$$

$$\vec{e}_1 = e_1 \cdot \hat{e}_1, \quad \vec{e}_2 = e_2 \cdot \hat{e}_2$$

our basis



$$\vec{a} = a \cos \theta_a \cdot \hat{e}_1 + a \sin \theta_a \hat{e}_2 = a \cos \theta_a \frac{\vec{e}_1}{e_1} + a \sin \theta_a \frac{\vec{e}_2}{e_2} = \frac{a}{e_1} \cos \theta_a \vec{e}_1 + \frac{a}{e_2} \sin \theta_a \vec{e}_2$$

$$\vec{b} = \frac{b}{e_1} \cos \theta_b \vec{e}_1 + \frac{b}{e_2} \sin \theta_b \vec{e}_2$$

Then $\vec{a} \cdot \vec{b} = \left(\frac{a}{e_1} \cos \theta_a \vec{e}_1 + \frac{a}{e_2} \sin \theta_a \vec{e}_2 \right) \cdot \left(\frac{b}{e_1} \cos \theta_b \vec{e}_1 + \frac{b}{e_2} \sin \theta_b \vec{e}_2 \right)$

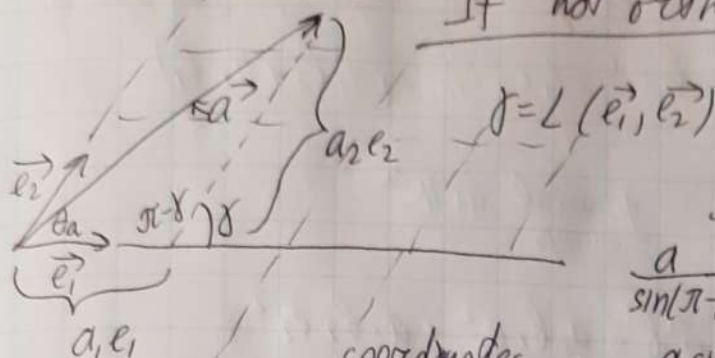
$$= \frac{a}{e_1} \cos \theta_a \cdot \frac{b}{e_1} \cos \theta_b \cdot e_1^2 + \frac{a}{e_2} \sin \theta_a \cdot \frac{b}{e_2} \sin \theta_b \cdot e_2^2 = ab \cos(\theta_a - \theta_b)$$

same as in any

basis! Denote $a_1 = \frac{a \cos \theta_a}{e_1}$, $a_2 = \frac{a \sin \theta_a}{e_2}$, then

$$\vec{a} = a_1 e_1 + a_2 e_2$$

If not orthogonal b.



Sine law:

$$\frac{a}{\sin(\pi - \gamma)} = \frac{a_2 e_2}{\sin \theta_a} = \frac{a_1 e_1}{\sin(\gamma - \theta_a)} \rightarrow$$

coordinates $a_1 = \frac{a \sin(\gamma - \theta_a)}{(\sin \gamma) e_1}$, $a_2 = \frac{a \sin \theta_a}{(\sin \gamma) e_2}$

Special cases: $\gamma = \frac{\pi}{2}$ (orthogonal),

$$a_1 = \frac{a \cos \theta_a}{e_1}, a_2 = \frac{a \sin \theta_a}{e_2} \checkmark$$

$\gamma = \frac{\pi}{2} \wedge e_1 = e_2 = 1$ (orthonormal)

$$a_1 = a \cos \theta_a, a_2 = a \sin \theta_a \checkmark$$

General formula for inner product:

$$\vec{a} \cdot \vec{b} = (a_1 \vec{e}_1 + a_2 \vec{e}_2) \cdot (b_1 \vec{e}_1 + b_2 \vec{e}_2) = a_1 b_1 e_1^2 + a_2 b_2 e_2^2 + e_1 e_2 \cos \gamma$$

$$= (\text{now check its value is same}) = \frac{a \sin(\gamma - \theta_a)}{e_1 \sin \gamma} \cdot \frac{b \sin(\gamma - \theta_b)}{e_1 \sin \gamma} e_1^2 + \frac{a \sin \theta_a}{e_2 \sin \gamma} \cdot \frac{b \sin \theta_b}{e_2 \sin \gamma} e_2^2 + e_1 e_2 \cos \gamma \left(\frac{a \sin(\gamma - \theta_a)}{e_1 \sin \gamma} \cdot \frac{b \sin \theta_b}{e_2 \sin \gamma} + \frac{a \sin \theta_a}{e_2 \sin \gamma} \cdot \frac{b \sin(\gamma - \theta_b)}{e_1 \sin \gamma} \right) = ab \left[\frac{\sin(\gamma - \theta_a) \sin(\gamma - \theta_b)}{\sin^2 \gamma} + \frac{\sin \theta_a \sin \theta_b}{\sin^2 \gamma} + \frac{\cos \gamma \sin(\gamma - \theta_a) \sin \theta_b}{\sin^2 \gamma} + \frac{\cos \gamma \sin \theta_a \sin(\gamma - \theta_b)}{\sin^2 \gamma} \right] = \text{must be } \cos(\theta_a - \theta_b)!$$

$$= \frac{ab}{\sin^2 \gamma} \left[(\sin \gamma \cos \theta_a - \cos \gamma \sin \theta_a)(\sin \gamma \cos \theta_b - \cos \gamma \sin \theta_b) + \sin \theta_a \sin \theta_b + \cos \gamma (\sin \gamma \cos \theta_a - \cos \gamma \sin \theta_a) \sin \theta_b + \cos \gamma \sin \theta_a (\sin \gamma \cos \theta_b - \cos \gamma \sin \theta_b) \right] =$$

$$= \frac{ab}{\sin^2 \gamma} \left[\sin^2 \gamma \cos \theta_a \cos \theta_b - \sin \gamma \cos \gamma \cos \theta_a \sin \theta_b - \sin \gamma \cos \gamma \sin \theta_a \cos \theta_b + \cos^2 \gamma \sin \theta_a \sin \theta_b + \cos^2 \gamma \sin \theta_a \sin \theta_b + \sin \theta_a \sin \theta_b + \cos \gamma \sin \gamma \cos \theta_a \sin \theta_b - \cos^2 \gamma \sin \theta_a \sin \theta_b + \cos \gamma \sin \gamma \sin \theta_a \cos \theta_b - \cos^2 \gamma \sin \theta_a \sin \theta_b \right] = \frac{ab}{\sin^2 \gamma} \left[\sin^2 \gamma (\cos \theta_a \cos \theta_b + \sin \theta_a \sin \theta_b) + (1 - \cos^2 \gamma) \sin \theta_a \sin \theta_b \right] =$$

$= ab \cos(\theta_a - \theta_b) \rightarrow$ still some number!
(calculation are terrible, do not know how to simplify)

In coordinates formula is not scalar product but (see in sheet):

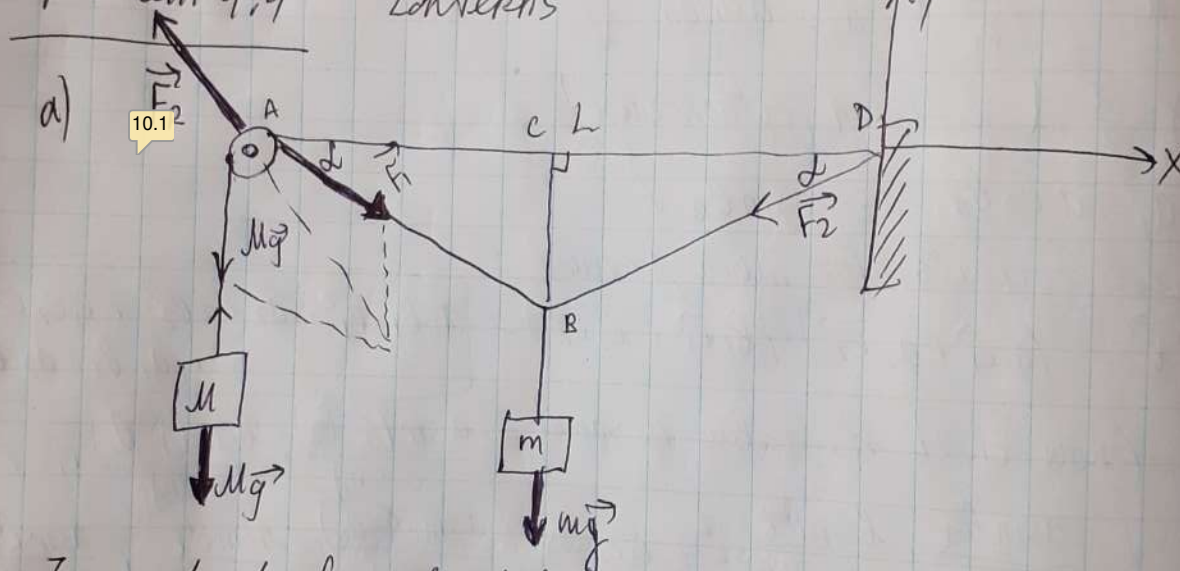
$$\vec{a} \cdot \vec{b} = ab \cos(\theta_a - \theta_b) = \underbrace{a_1 b_1 e_1^2 + a_2 b_2 e_2^2}_{\text{always in } \mathbb{R}^D} + \underbrace{e_1 e_2 \cos \gamma}_{\text{general always in } \mathbb{R}^2} (a_1 b_2 + a_2 b_1)$$

Special case:

$$\gamma = \frac{\pi}{2} \text{ (orthogonal)}; \vec{a} \cdot \vec{b} = a_1 b_1 e_1^2 + a_2 b_2 e_2^2$$

$$\gamma = \frac{\pi}{2}, e_1 = e_2 = 1; a_1 b_1 + a_2 b_2 \text{ (orthonormal)}$$

Problem 4.4 Lanterns



Important for all tasks:

b) ~~Lantern~~ Roll is stationary, so

$$\vec{F}_R + \vec{Mg} + \vec{F}_1 = \vec{0}, \quad \vec{F}_R = -(\vec{Mg} + \vec{F}_1)$$

transmitted from M

c) Since m is on the middle, $F_1 = F_2$ reason 1 $\rightarrow mg + \vec{F}_1 + \vec{F}_2 = \vec{0}$
 triangles ABC and DBC are equal,
 $\angle A = \angle B = \angle C$
reason 2 $mg_x + F_{1x} + F_{2x} = 0$
 $0 + F_1 \cos \alpha - F_2 \cos \alpha = 0 \rightarrow F_1 = F_2$
 from symmetry, if we

look at the picture from other side, it looks same. There is no way to distinguish F_1 and F_2 .

Further, $F_1 = F_2 = F$, act on same rope and must equal Mg (mass of rope/wire not given, can neglect and it acts only as transmitter)

Then:

c.1) $2F \sin \alpha = 2Mg \sin \alpha = mg$, (equilib. for lantern)
 $2M \sin \alpha = m$, $\alpha = \arcsin\left(\frac{m}{2M}\right)$.

c.2) Equilib. for roll: $F_{ry} = Mg + \underbrace{Mg \cos\left(\frac{\pi}{2} - \alpha\right)}_{F_1} = Mg + Mg \sin \alpha = Mg(1 + \sin \alpha)$.

$F_{rx} = -F_1 \cos \alpha = -Mg \cos \alpha$.

Then $|\vec{F}_r| = \sqrt{F_{rx}^2 + F_{ry}^2} = Mg \sqrt{\cos^2 \alpha + (1 + \sin \alpha)^2} = Mg \sqrt{2 + 2 \sin \alpha} =$

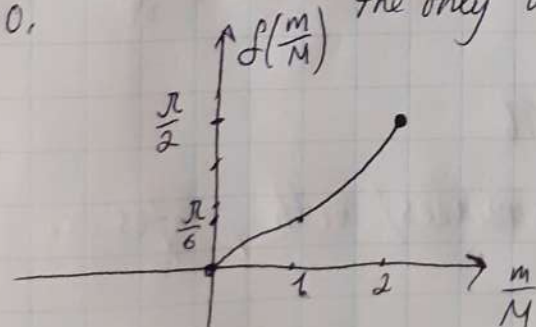
(special cases: $\alpha = \frac{\pi}{2}$ $\uparrow F_r$ $F_{ry} = 2Mg$, $F_{rx} = 0$)
 $\alpha = 0$, $F_{ry} = Mg$, $F_{rx} = -Mg$ \checkmark

Actual values:

$\alpha = \arcsin\left(\frac{15 \text{ kg}}{2 \cdot 80 \text{ kg}}\right) = \arcsin\left(\frac{3 \text{ kg}}{32 \text{ kg}}\right) \approx \arcsin \frac{1}{10}$, almost horizontal

$|\vec{F}_r| = \sqrt{2} \cdot 80 \text{ kg} \cdot 10 \frac{\text{N}}{\text{kg}} \cdot \sqrt{1 + \frac{15}{160}} \approx 800\sqrt{2} \approx \text{weight of big man.}$

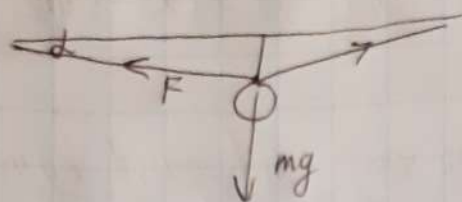
d) $\alpha = \arcsin\left(\frac{m}{2M}\right)$ — it should be dimensionless, and this is the only way with given data, up to f , and $\frac{1}{2}$, and 1.



If $m > 2M$, it should break.

e) $F_{\text{max}} = Mg = 14 \text{ kN}$, $m_{\text{max}} = 2M_{\text{max}} = 2 \frac{F_{\text{max}}}{g} = \frac{2 \cdot 14 \cdot 10^3 \text{ N}}{g} = 28 \cdot 10^2 \text{ kg} \approx 3 \text{ tonne.}$
 α becomes $\frac{\pi}{2}$.

f)

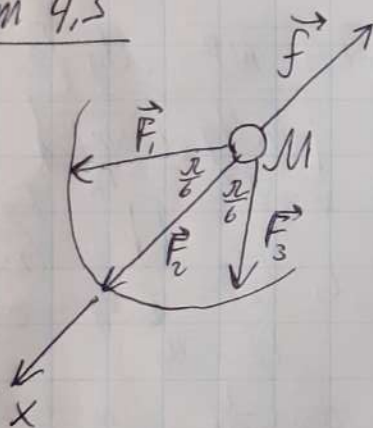


$mg = 2F \sin \theta \approx 2\theta F$ for small angles,
 then $F = \frac{mg}{2\theta}$ becomes $> F_{max}$
 described above, and ropes break.

With the roll we can adjust L , remedying what was described above.

I promise I will try to do experiment for d, but very unhandy I did not find suitable tools yet.

Self Test Problem 4.5



Max. force will be exactly at required angle.

a) $F + 2F \cdot \frac{\sqrt{3}}{2} = (\sqrt{3}+1)F$ total force,
 ≈ 3 times more

b) $(\sqrt{3}+1)F = \mu Mg$,
 $\mu = \frac{(\sqrt{3}+1)F}{Mg}$

$1 \text{ cwt} = 8 \text{ st} = 8 \cdot 14 \text{ lb} \approx 56 \text{ kg}$

Actual values

$\mu \approx \frac{3 \cdot 300 \cdot \frac{1}{2} \cdot 70 \text{ N}}{1120 \text{ kg} \cdot 10 \frac{\text{km}}{\text{st}}} \approx \frac{450}{1120} = \frac{45}{112} \approx \frac{3}{7}$

Bonus problem — problem 4.6

To discuss on seminars/with professor.

Index of comments

6.1 4.2: 11.5/12

6.2 3.5/4

7.1 4/4

7.2 4/4

8.1 Good try, perhaps think about it in these lines :
Let the first two basis vectors, e^1 and e^2 , span this plane. Then the coordinates of all other vectors are zero, and the same expressions are found as in the 2D case.
5/5 +1 bonus :)
Also nice plots in the earlier exercises :D

10.1 4.4a: 4/4

10.2 4.4b: 3/3

11.1 4.4c: 4/4

11.2 4.4d: 3/3