SESA3029 Aerothermodynamics



Lecture 2.7 Under/over-expanded nozzles and supersonic wind tunnels



Environment is at 'back' pressure p_h

Reservoir conditions p₀, T₀

 p_{b1} – low subsonic throughout (effectively incompressible) p_{b2} –subsonic throughout

 p_{b2} –subsonic throughout (compressible)

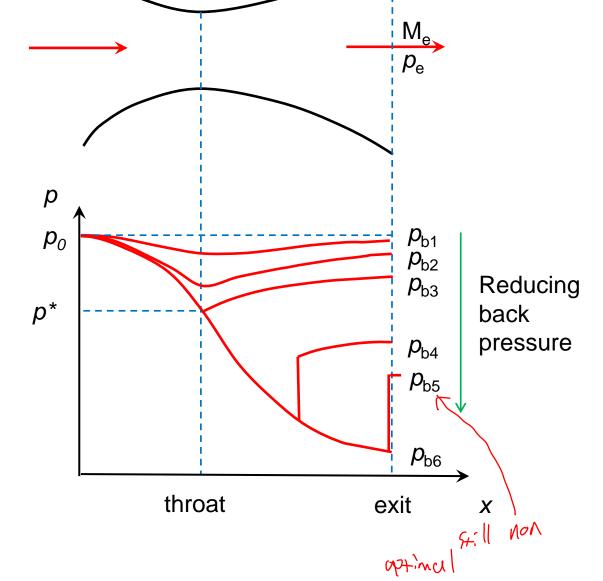
 p_{b3} —sonic at throat, subsonic in diverging section

 p_{b4} –shock in diverging section

 p_{b5} –shock at exit

 p_{b6} –design condition

(supersonic exit)



Example

For a Laval nozzle with a design exit Mach number of $M_e=2.2$ find p_{b6} , p_{b3} and p_{b5} in terms of p_0

	Isentropic-flow table ($\gamma = 1.4$):					
	M	p/p_0	ρ/ρ_0	T/T_0	v (deg.)	A/A*
2	2.2000	0.0935	0.1841	0.5081	31.7325	2.0050
- 2	2.2200	0.0906	0.1800	0.5036	32.2494	2.0409
2	2.2400	0.0878	0.1760	0.4991	32.7629	2.0777
2	2.2600	0.0851	0.1721	0.4947	33.2730	2.1153
2	2.2800	0.0825	0.1683	0.4903	33.7796	2.1538

Design condition $M_e=2.2$, $A_e/A^*=2.005$ and $p_{b6}/p_0=0.0935$

For same area ratio there is a subsonic solution

	Isentropic-flow table ($\gamma = 1.4$):							
M	p/p ₀	ρ/ρ_0	T/T_0	ν (deg.)	A/A*			
0.3000	0.9395	0.9564	0.9823	n/a	2.0351			
0.3200	0.9315	0.9506	0.9799	n/a	1.9219			
0.3400	0.9231	0.9445	0.9774	n/a	1.8229			
0.3600	0.9143	0.9380	0.9747	n/a	1.7358			
0.3800	0.9052	0.9313	0.9719	n/a	1.6587			

Linear interpolation for $A_e/A^*=2.005$ gives $M_e=0.305$ and $p_{b3}/p_0=0.937$

To get p_{b5} we add a normal shock at the exit i.e. from M=2.2

NST for M=2.2 gives p_{b5}/p_{b6} =5.480

$$p_{b5} = p_{b6} \times p_{b5}/p_{b6} = 0.0935p_0 \times 5.480 = 0.512p_0$$

Flow regimes

- $p_b > p_{b3}$ Subsonic flow
- $p_{b5} < p_b < p_{b3}$ Shock wave in diverging section
- $p_b < p_{b5}$ Supersonic flow at exit
- $p_{b6} < p_b < p_{b5}$ Flow needs to recompress after exit to get to ambient conditions ('over-expanded' case)
- $p_b < p_{b6}$ Flow needs to expand further to get to ambient conditions ('under-expanded' case)

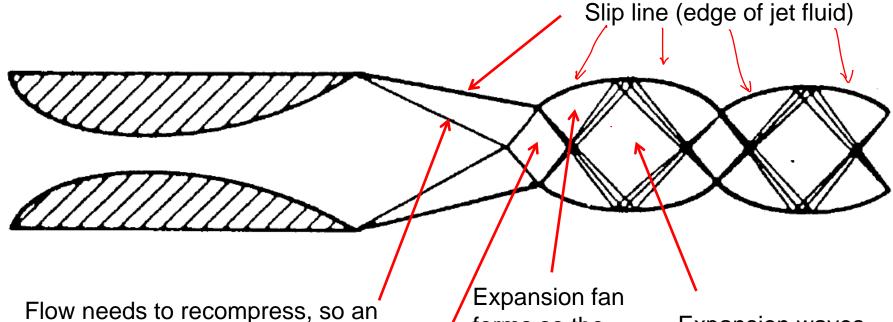
How does the adjustment in pressure occur?

Over-expanded nozzle/jet

Slip line (edge of jet fluid) Tvery high pressure (nver expanded) Dresture Flow needs to recompress, so an oblique shock forms, with turning angle so that pressure is continuous across the slip line

Shocks cross. Now the pressure in this region is too high

Over-expanded nozzle/jet



oblique shock forms, with turning angle so that pressure is continuous across the slip line

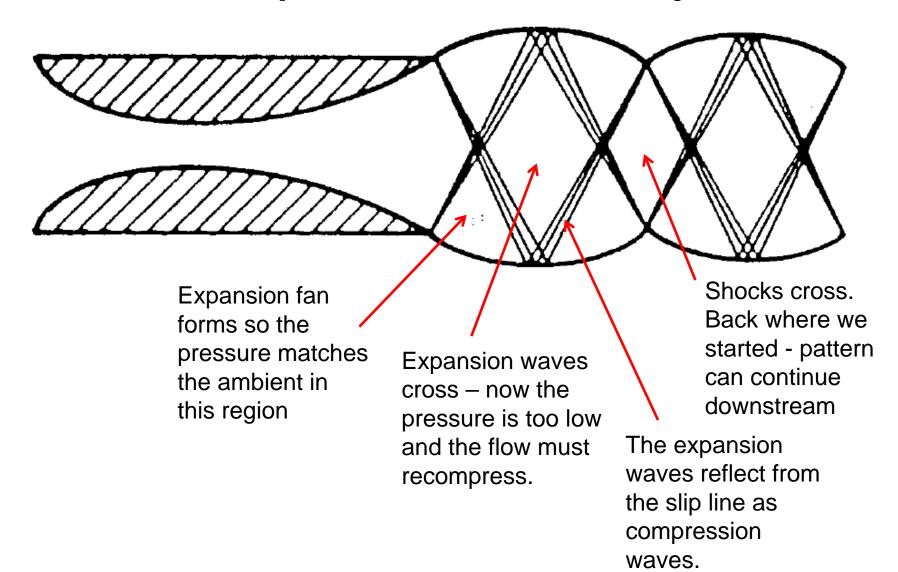
forms so the pressure matches cross - now the the ambient in this region

Expansion waves pressure is too low and the flow must recompress.

Shocks cross. Now the pressure in this region is too high

Back where we started. The expansion waves reflect from the slip line as compression waves.

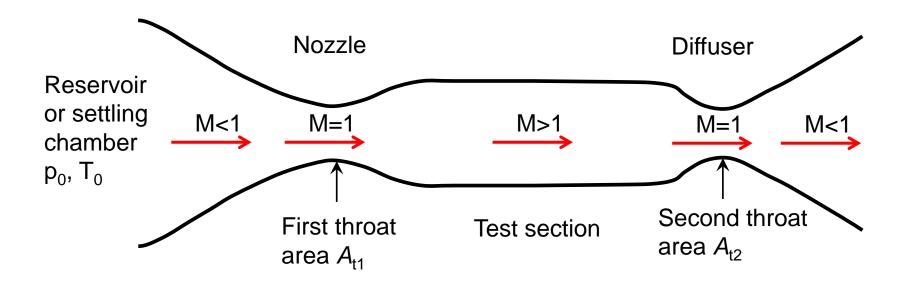
Under-expanded nozzle/jet



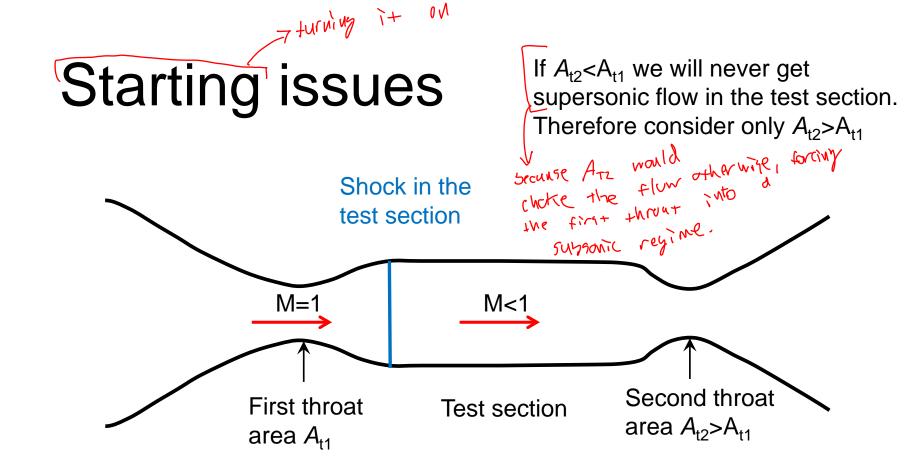
Notes

- After the first cycle over- and under-expanded jets follow the same mechanism.
- Slip line follows a barrel shape
- Shock waves follow a diamond pattern
- Turning angle may be too great to get a regular shock crossing.
 - Instead we can get a Mach disk.

Supersonic wind tunnel: ideal configuration



An efficient diffuser section is used to get the M>1 test section for a lower stagnation pressure compared to an open jet configuration



If A_{t2} is too small we have a problem starting this configuration – the second throat must be large enough to accommodate the mass flow with a shock in position as shown

It is also possible that the first get's choked, to provent that we increase was flow rate.

Example

For an M=3 wind tunnel, find A_{t2}/A_{t1} sufficient to swallow a shock sitting in the test section.

Isentropic-flow table ($\gamma = 1.4$):							
<i>M</i>	p/p_0	ho/ ho 0	T/T_0	ν (deg.)	A/A*		
3.0000	0.0272	0.0762	0.3571	49.7573	4.2346		

For M=3 in test section we know (IFT) $A_{test}/A_{t1}=4.23$

Normal-shock table ($\gamma = 1.4$):

$M_{n,1}$	M_{n2}	p_2/p_1	$ ho_2/ ho_1$	T_2/T_1	p_{02}/p_{01}	p_{02}/p_{1}
3.0000	0.4752	10.3333	3.8571	2.6790	0.3283	12.0610

A normal shock in the test section would have (NST) M_2 =0.475 and p_{02}/p_{01} =0.3283

Recall the result for mass flow rate

$$\dot{m} = \frac{p_0}{\sqrt{RT_0}} A^* \sqrt{\gamma} \left(\frac{2}{\gamma + 1}\right)^{\frac{\gamma + 1}{2(\gamma - 1)}}$$

$$\uparrow \qquad \uparrow \qquad \uparrow$$
Stagnation Throat Gas conditions area properties

If p₀ drops we need a bigger A to let the mass flow pass through the second throat at the required rate

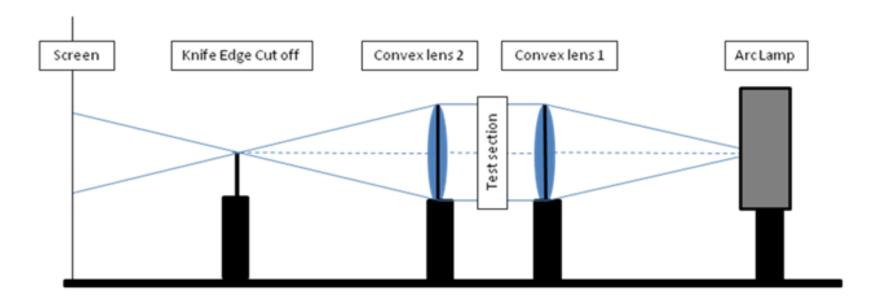
In this case we need $A_{t2}/A_{t1} > p_{01}/p_{02} = 1/0.3283 = 3.046$

For slightly higher A_{t2} the shock gets 'swallowed' by the second throat, leaving supersonic flow in the test section

After we have established supersonic flow we would ideally reduce A₁₂ to make the diffuser more efficient

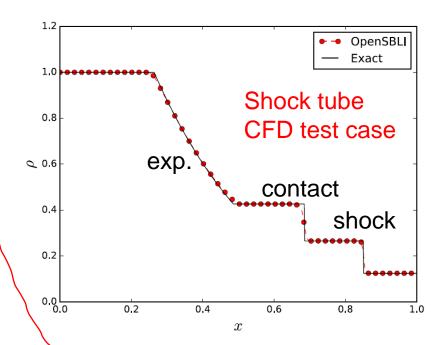
Flow visualisation

- Light is refracted by density variations
 - Shadowgraph: shine a parallel beam of light though test section 70 Mayi onto screen/camera
 - Image intensity proportional to second derivative of density
 - Schlieren: use lenses to focus light beam onto a knife edge that oreterred L cuts refracted light
 - Image intensity proportional to density gradient.



Ludwieg tube

- Shock tube = long straight tube with diaphragm separating high pressure gas from low pressure (or vacuum)
- Bursting the diaphragm leads to a shock wave and contact discontinuity moving from left to right and an expansion fan moving right to left
- Add a constitution nozzle to accelerate the flow behind the contact discontinuity to supersonic = relatively inexpensive supersonic/hypersonic wind tunnel
- Short duration (of order 0.1s): can run until the expansion wave reflects back from the end of the pressurised tube





TU Braunschweig

(Add image of pressure with time

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