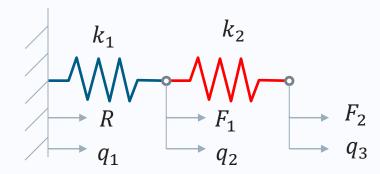


# Part 2a: Elastic Rods in Tension and Compression

FEEG3001/SESM6047 FEA in Solid Mechanics Prof A S Dickinson

From 11th October 2024

### Reminder from last time:



- We search for solutions for  $q_i$  in problems like this:
  - We express the system's Elastic Strain Energy  $\it U$  and its Potential Energy  $\it V$
  - by finding  $U_i$  for each element and V based on all applied loads.
  - We could apply PMTPE and calculate  $\frac{\partial \Pi(q_i)}{\partial q_i} = 0$  to obtain  $\{F\} = [K]\{q\}$
  - but because we can express  $U_i$  in quadratic form, we can take a shortcut where  $U_i = 1/2\{q\}^T[K_i]\{q\}$
  - and assemble our elemental  $[K_i]$  into a global [K]
  - and by expressing V in matrix form too, we find  $\{F\}$  and can therefore solve  $\{F\} = [K]\{q\}$
  - by inverting the stiffness matrix, allowing  $\{q\} = [K]^{-1}\{F\}$

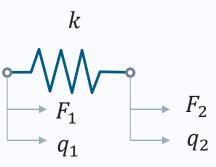
#### This week:

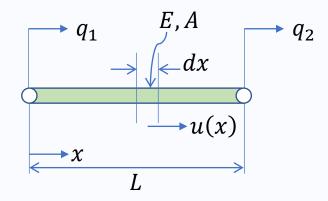
#### Similarities:

- PMTPE gives us our governing equation of equilibrium (a set of linear quadratic equations)
- How we fetch and assemble our element stiffness matrices into an assembled stiffness matrix

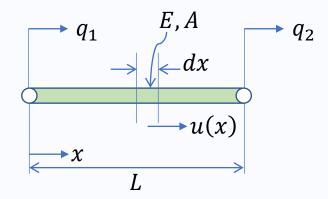
#### Differences:

- Unlike for springs, we have some approximation, for 'elastic rods'





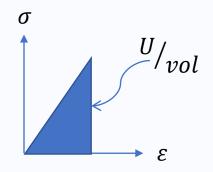
- Can deform only in tension or compression
- Unlike our spring, properties are described by E, A (cross section) and L (length), and the axial displacement of the endpoints  $q_1$  and  $q_2$  in coordinate system x
- We look at a small portion dx, and its deformation u(x) depends on where we take the slice
- Every point on the cross-section slice displaces the same amount
- So x is the 'label' (which cross section) and u(x) is the displacement.



• PMTPE requires an expression for the strain energy U in terms of the displacement field in the rod u.

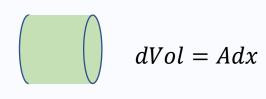
$$U = \frac{1}{2} \int_0^L EA \left(\frac{du}{dx}\right)^2 dx$$

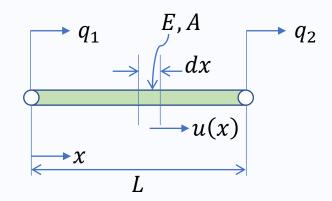
- Where does this come from? You should remember...
- Elastic strain energy per unit volume is the area under the stress-strain curve, for simplified 1D (uniaxial) elasticity:



$$U = \frac{1}{2} \int \sigma_x \varepsilon_x dVol$$

$$U = \frac{1}{2} \int \sigma_x \varepsilon_x Adx$$





• and because in 1D elasticity (highly simplified case) we can say that  $\sigma_x = E \varepsilon_x$ ,

$$U = \frac{1}{2} \int E \varepsilon_x \times \varepsilon_x \times A dx$$

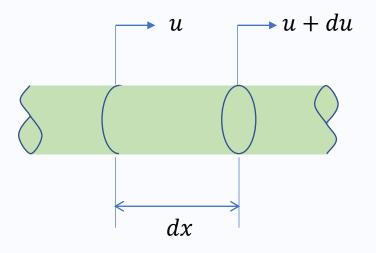
$$U = \frac{1}{2} \int E A \varepsilon_x^2 dx$$

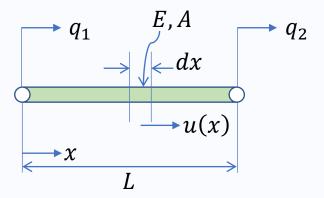
• What is  $\varepsilon_x$  given the displacement field?

$$\varepsilon_{x} = \frac{(u + du) - u}{dx} = \frac{du}{dx}$$

Hence

$$U = \frac{1}{2} \int EAu'^2 dx$$





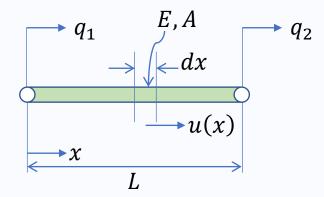
- **Aim:** we want to express the displacement throughout the element u(x) from the nodal displacements,  $q_1$  and  $q_2$ , (currently unknown).
- Now our first approximation: let's say u(x) is a linear function of x. i.e.:

$$u(x) = a + bx$$

- (We can make this linear function assumption for 1D elasticity, but we couldn't make that assumption for more complex cases).
- This is inconvenient though, as a and b have no meaning. A more convenient version, more general:

$$u(x) = N_1(x)q_1 + N_2(x)q_2$$

•  $q_1$  and  $q_2$  are unknowns,  $N_1$  and  $N_2$  are prescribed **shape functions** or **interpolation functions** for approximation: a common approach taken in all finite elements.



$$u(x) = N_1(x)q_1 + N_2(x)q_2$$

• The trick is we choose for these shape functions  $N_1$  and  $N_2$  some linear functions,  $g_1$  and  $g_2$ . (We could choose other functions...)

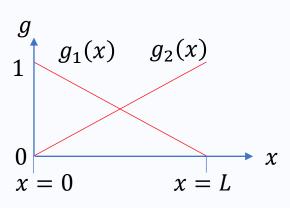
$$u(x) = g_1(x)q_1 + g_2(x)q_2$$

• We can define these by drawing our domain 0 to *L*:

$$g_1(x) = 1 - \frac{x}{L}$$

$$g_2(x) = \frac{x}{L}$$

$$u(x) = \left(1 - \frac{x}{L}\right)q_1 + \left(\frac{x}{L}\right)q_2$$



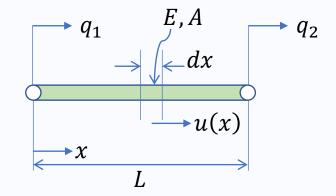
$$u(x) = \left(1 - \frac{x}{L}\right)q_1 + \left(\frac{x}{L}\right)q_2$$

- But what are these qs?
- Why didn't we just use u(x) = a + bx?
- Work out u(x) at each end:

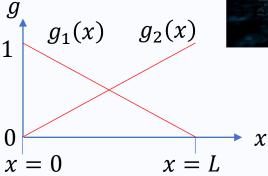
$$u(0) = \left(1 - \frac{0}{L}\right)q_1 + \left(\frac{0}{L}\right)q_2 = q_1$$

$$u(L) = \left(1 - \frac{L}{L}\right)q_1 + \left(\frac{L}{L}\right)q_2 = q_2$$

• The qs, which aren't functions of x, describe the end displacements! This is why we choose shape functions in that form.







and assuming we have solved to find the q unknowns, we can use

$$u(x) = \left(1 - \frac{x}{L}\right)q_1 + \left(\frac{x}{L}\right)q_2$$

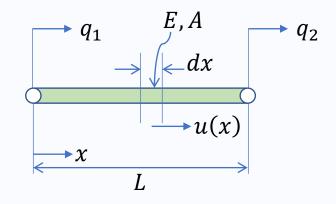
• to give us values for strain:

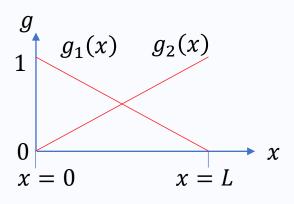
$$\varepsilon(x) = \frac{du}{dx} = \left(-\frac{1}{L}\right)q_1 + \left(\frac{1}{L}\right)q_2$$

• and in this simple 1D case, for stress:

$$\sigma(x) = E\varepsilon(x) = \left(-\frac{E}{L}\right)q_1 + \left(\frac{E}{L}\right)q_2$$

(check dimensional analysis if you like!)





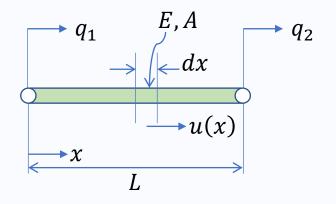
 or in matrix form (as more complex element types will require):

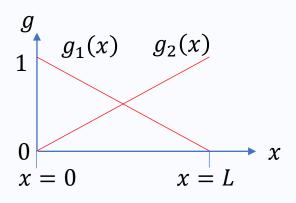
$$\varepsilon(x) = \frac{du}{dx} = \left(-\frac{1}{L}\right)q_1 + \left(\frac{1}{L}\right)q_2$$

$$\varepsilon(x) = \left[-\frac{1}{L} \quad \frac{1}{L}\right] \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}$$



$$\sigma(x) = E\varepsilon(x) = \left(-\frac{E}{L}\right)q_1 + \left(\frac{E}{L}\right)q_2$$
$$\sigma(x) = [E]\left[-\frac{1}{L} \quad \frac{1}{L}\right] \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}$$

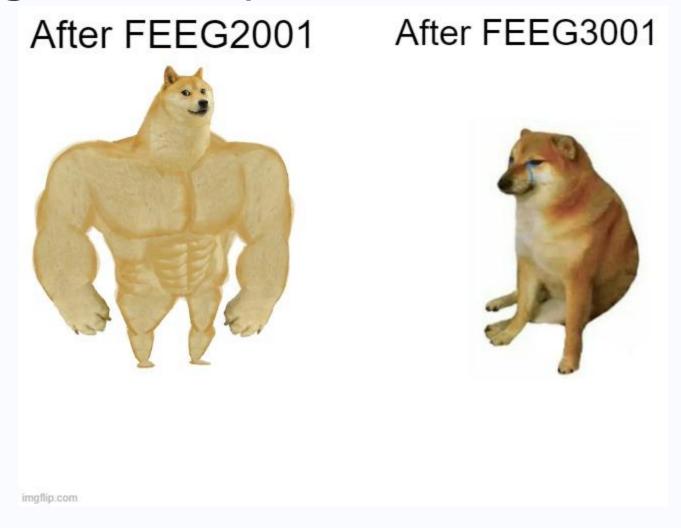




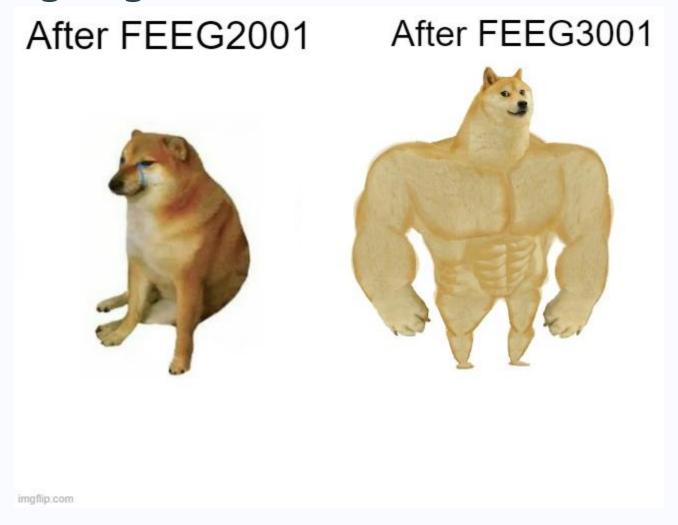
## Summary:

- We have written the internal field of displacement u(x) in terms of two chosen functions that happen here to be linear (Shape Functions)  $g_1(x)$  and  $g_2(x)$ .
- $g_1(x)$  has a property where its value is 1 at left and 0 at right, and
- $g_2(x)$  has a value 0 at left and 1 at right.
- Because they are linear functions, what about their combination? It must also be a linear function, if q values do not depend on x (they are single, nodal values).
- Next class will continue towards how we find values for q!

## How you might feel today:



## Where we are going:







# Part 2b: Elastic Rods in Tension and Compression

FEEG3001/SESM6047 FEA in Solid Mechanics Dr A S Dickinson

From 15th October 2024

- Remember we had a rod of length L: with a cross section at x where the displacement is u(x) and the properties are given by E and A.
- And we defined shape functions which vary linearly along the length and have the properties:

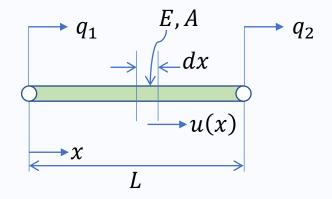
$$g_1(x) = 1 - \frac{x}{L}$$

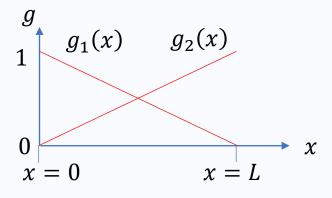
$$g_2(x) = \frac{x}{L}$$

And because we can combine them linearly:

$$u(x) = \left(1 - \frac{x}{L}\right)q_1 + \left(\frac{x}{L}\right)q_2$$

• so to estimate u(x) anywhere we are allowed to use a linear interpolation or 'approximation'.





$$u(x)=\left(1-\frac{x}{L}\right)q_1+\left(\frac{x}{L}\right)q_2$$
, or more generally: 
$$u(x)=g_1(x)q_1+g_2(x)q_2$$

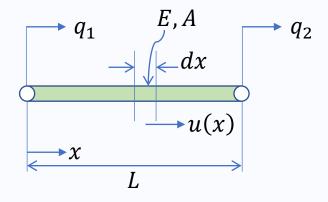
- Sometimes this will be exact, but often it is an approximation. We will see many functions of this form.
- We want to find  $q_1$  and  $q_2$  which are unknowns but which do not depend on x. Remember we had:

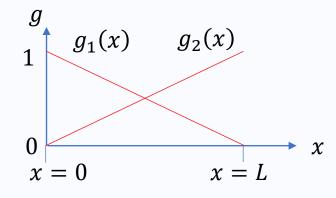
$$U = \frac{1}{2} \int EAu'^2 dx$$

$$u'(x) = g_1'(x)q_1 + g_2'(x)q_2 \text{ so}$$

$$u'(x) = \left(-\frac{1}{L}\right)q_1 + \left(\frac{1}{L}\right)q_2$$

Now we substitute into *U*:





$$U = \frac{1}{2} \int EA \left[ \left( -\frac{1}{L} \right) q_1 + \left( \frac{1}{L} \right) q_2 \right]^2 dx$$

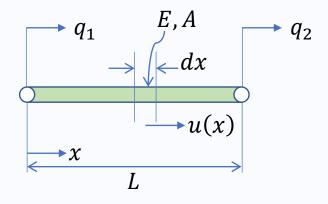
and in this case we are integrating a constant in x so

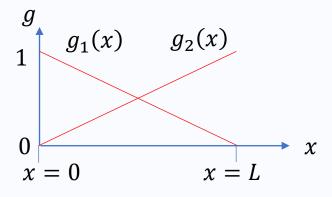
$$U = \frac{1}{2} \left( \frac{EA}{L^2} \right) (-q_1 + q_2)^2 \int_0^L 1 dx$$

$$U = \frac{1}{2} \left( \frac{EA}{L} \right) (-q_1 + q_2)^2$$
 so expanding,

$$U = \frac{1}{2} \left( \frac{EA}{L} \right) (q_1^2 - 2q_1 q_2 + q_2^2)$$

• Can we reorganise this as we did for springs, using quadratic forms?





$$U = \frac{1}{2} \left( \frac{EA}{L} \right) (q_1^2 - 2q_1 q_2 + q_2^2)$$

In quadratic form:

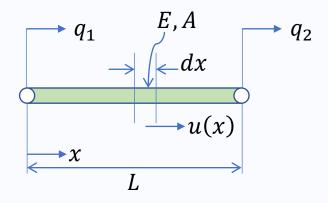
$$U = \frac{1}{2} \left( \frac{EA}{L} \right) \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}^T \begin{bmatrix} \\ \\ q_2 \end{Bmatrix}$$

$$U = \frac{1}{2} \left( \frac{EA}{L} \right) \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}^T \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}$$

• or alternatively:

$$U = \frac{1}{2} \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}^T \begin{bmatrix} \frac{EA}{L} & -\frac{EA}{L} \\ -\frac{EA}{L} & \frac{EA}{L} \end{bmatrix} \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}$$

Can you spot our new element stiffness matrix?



• Finally, and as before, we need to find the values of our generalised coordinates  $q_1$  and  $q_2$ , for a given set of loads and BCs, to solve the problem:

$$U = \frac{1}{2} \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}^T \begin{bmatrix} \frac{EA}{L} & -\frac{EA}{L} \\ -\frac{EA}{L} & \frac{EA}{L} \end{bmatrix} \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}$$

$$[K] = \begin{bmatrix} \frac{EA}{L} & -\frac{EA}{L} \\ -\frac{EA}{L} & \frac{EA}{L} \end{bmatrix}$$



# Part 2c: Solving Elastic Rods in Generalised Coordinates

FEEG3001/SESM6047 FEA in Solid Mechanics Dr A S Dickinson

From 15th October 2024

 Now again, like with the springs, we analysed a typical rod, which we can use over again, without needing to re-calculate the behaviour each time. So for any, general rod we have:

$$U = \frac{1}{2} \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}^T \begin{bmatrix} \frac{EA}{L} & -\frac{EA}{L} \\ -\frac{EA}{L} & \frac{EA}{L} \end{bmatrix} \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix}$$

- containing our element stiffness matrix, a 2x2 matrix.
- Like last time, we can assemble this into the *assembled stiffness matrix* for a *structure*.

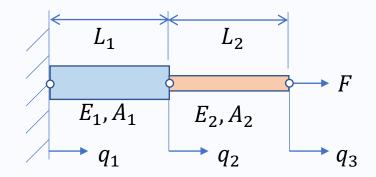
- Take a simple example of a rod with changing properties along its length
- Considering the elastic strain energy:

$$U = U_1 + U_2$$

$$U_{1} = \frac{1}{2} \begin{Bmatrix} q_{1} \\ q_{2} \end{Bmatrix}^{T} \begin{bmatrix} \frac{E_{1}A_{1}}{L_{1}} & -\frac{E_{1}A_{1}}{L_{1}} \\ -\frac{E_{1}A_{1}}{L_{1}} & \frac{E_{1}A_{1}}{L_{1}} \end{bmatrix} \begin{Bmatrix} q_{1} \\ q_{2} \end{Bmatrix}$$

$$U_{2} = \frac{1}{2} \begin{Bmatrix} q_{2} \end{Bmatrix}^{T} \begin{bmatrix} \frac{E_{2}A_{2}}{L_{2}} & -\frac{E_{2}A_{2}}{L_{2}} \\ -\frac{E_{2}A_{2}}{L_{2}} & \frac{E_{2}A_{2}}{L_{2}} \end{bmatrix} \begin{Bmatrix} q_{2} \\ q_{3} \end{Bmatrix}$$

$$U = \frac{1}{2} \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \end{Bmatrix}_{1 \times 3}^{T} \begin{bmatrix} ? & ? & ? \\ ? & ? & ? \\ ? & ? & ? \end{bmatrix}_{3 \times 3} \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \end{Bmatrix}_{3 \times 1}$$



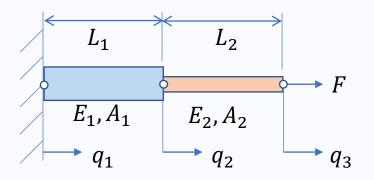
- Take a simple example of a rod with changing properties along its length
- Considering the elastic strain energy:

$$U = U_1 + U_2$$

$$U_{1} = \frac{1}{2} \begin{Bmatrix} q_{1} \\ q_{2} \end{Bmatrix}^{T} \begin{bmatrix} \frac{E_{1}A_{1}}{L_{1}} & -\frac{E_{1}A_{1}}{L_{1}} \\ -\frac{E_{1}A_{1}}{L_{1}} & \frac{E_{1}A_{1}}{L_{1}} \end{bmatrix} \begin{Bmatrix} q_{1} \\ q_{2} \end{Bmatrix}$$

$$U_{2} = \frac{1}{2} \begin{Bmatrix} q_{2} \\ q_{3} \end{Bmatrix}^{T} \begin{bmatrix} \frac{E_{2}A_{2}}{L_{2}} & -\frac{E_{2}A_{2}}{L_{2}} \\ -\frac{E_{2}A_{2}}{L_{2}} & \frac{E_{2}A_{2}}{L_{2}} \end{bmatrix} \begin{Bmatrix} q_{2} \\ q_{3} \end{Bmatrix}$$

$$U = \frac{1}{2} \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \end{Bmatrix}^T_{1 \times 3} \begin{bmatrix} \frac{E_1 A_1}{L_1} & -\frac{E_1 A_1}{L_1} & 0 \\ -\frac{E_1 A_1}{L_1} & \frac{E_1 A_1}{L_1} + \frac{E_2 A_2}{L_2} & -\frac{E_2 A_2}{L_2} \\ 0 & -\frac{E_2 A_2}{L_2} & \frac{E_2 A_2}{L_2} \end{bmatrix}_{3 \times 3} \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \end{Bmatrix}_{3 \times 1}$$



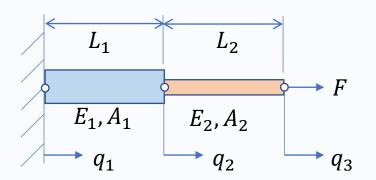
- Considering the work done by the external forces:
- F is the only force which is doing work, so

$$V = -Fq_3$$



$$V = -\{ \}_{1 \times 3} \{q\}_{3 \times 1}$$

$$V = -\{R \quad 0 \quad F\} \begin{cases} q_1 \\ q_2 \\ q_3 \end{cases}$$



 So we have been able to represent the whole system in the matrix form:

$$U = \frac{1}{2} \{q\}^T [K] \{q\}$$

and

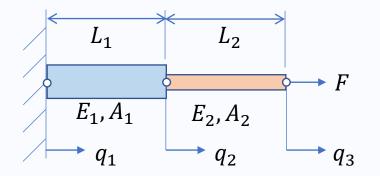
$$V = -\{F\}^T \{q\}$$

and since:

$$\Pi = U + V$$
 so
$$\Pi = \frac{1}{2} \{q\}^T [K] \{q\} - \{F\}^T \{q\}$$

 We want to apply PMTPE, where in i notation Equilibrium says

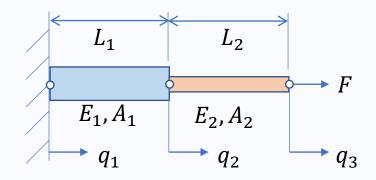
$$\delta\Pi(q_i) = 0 \Rightarrow \frac{\partial\Pi}{\partial q_i} = 0, i = 1, 2, ...$$



· And we know that where we have this quadratic form,

$$\Pi = \frac{1}{2} \{q\}^T [K] \{q\} - \{F\}^T \{q\}$$

• that the partial derivative of the total potential energy leads to  $[K]{g} = {F}$ 

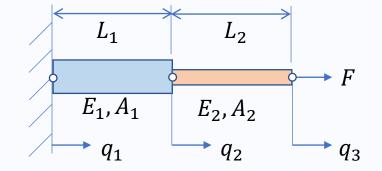


• (i.e. we don't need to derive it all over again, we can just write out our equation of equilibrium):

$$\begin{bmatrix} \frac{E_1 A_1}{L_1} & -\frac{E_1 A_1}{L_1} & 0\\ -\frac{E_1 A_1}{L_1} & \frac{E_1 A_1}{L_1} + \frac{E_2 A_2}{L_2} & -\frac{E_2 A_2}{L_2} \\ 0 & -\frac{E_2 A_2}{L_2} & \frac{E_2 A_2}{L_2} \end{bmatrix} \begin{pmatrix} q_1\\ q_2\\ q_3 \end{pmatrix} = \begin{pmatrix} R\\ 0\\ F \end{pmatrix}$$

And recall we cannot solve this - our stiffness matrix [K] is singular and cannot be inverted (its determinant is zero), so alone it cannot be solved. So what do we do?

• We can apply our boundary condition because  $q_1 = 0$ :



$$\begin{bmatrix} \frac{E_1A_1}{I_1} & \frac{E_1A_1}{L_1} & 0 \\ -\frac{E_1A_1}{L_1} & \frac{E_1A_1}{L_1} + \frac{E_2A_2}{L_2} & -\frac{E_2A_2}{L_2} \\ -\frac{E_2A_2}{L_2} & \frac{E_2A_2}{L_2} \end{bmatrix} \begin{pmatrix} \frac{q_1}{q_2} & \frac{R}{q_2} \\ q_3 \end{pmatrix} = \begin{pmatrix} R \\ 0 \\ F \end{pmatrix} \text{ leaving}$$

$$\begin{bmatrix} \frac{E_1 A_1}{L_1} + \frac{E_2 A_2}{L_2} & -\frac{E_2 A_2}{L_2} \\ -\frac{E_2 A_2}{L_2} & \frac{E_2 A_2}{L_2} \end{bmatrix} \begin{Bmatrix} q_2 \\ q_3 \end{Bmatrix} = \begin{Bmatrix} 0 \\ F \end{Bmatrix}$$

leaving a 2x2 system (two equations, two unknowns) which we can solve

## Why do you solve matrix equations like that?

#### Question:

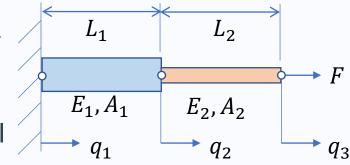
– Why do we solve our equations in a strange way, by crossing out rows and columns? Why is it OK to ignore those rows and columns?

#### Answer:

- We are not actually ignoring them. We are taking advantage of a simplification possible due to a 4<sup>th</sup> equation (in a 3x3 system) which comes from our boundary condition (BC), i.e.  $q_1 = 0$
- If you set non-zero displacement BCs, you would need another solution method...!



• To calculate the reaction force *R*, we could use intuition – and here it is easy enough we could do it by inspection, but with a complex system we don't have that luxury.



 Instead, we use the rows we crossed out from the general equation of equilibrium:

$$\begin{bmatrix} \frac{E_1 A_1}{L_1} & -\frac{E_1 A_1}{L_1} & 0\\ -\frac{E_1 A_1}{L_1} & \frac{E_1 A_1}{L_1} + \frac{E_2 A_2}{L_2} & -\frac{E_2 A_2}{L_2} \\ 0 & -\frac{E_2 A_2}{L_2} & \frac{E_2 A_2}{L_2} \end{bmatrix} \begin{pmatrix} q_1\\ q_2\\ q_3 \end{pmatrix} = \begin{pmatrix} R\\ 0\\ F \end{pmatrix}$$

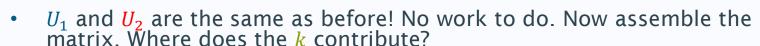
$$R = \frac{E_1 A_1}{L_1} q_1 - \frac{E_1 A_1}{L_1} q_2 + 0 q_3 = -\frac{E_1 A_1}{L_1} q_2$$

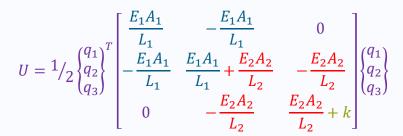
 That is, you can calculate your reaction solution by using the boundary condition rows.

## Finally, a new problem:

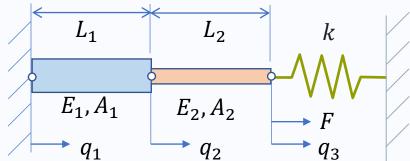
• To show you how quickly this can be done, for a less intuitive problem, add a grounded spring (i.e. there is no  $q_4$ ):

$$U = U_1 + U_2 + U_{spring}$$
$$U_{spring} = \frac{1}{2} kq_3^2$$



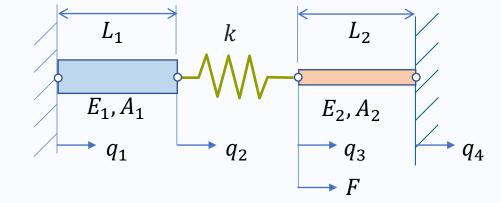


- This spring's stiffness contribution is 'lumped', or lies on the diagonal.
- Note this would NOT affect the force vector, as the spring support is considered 'yielding' and therefore provides no reaction. The ground reaction at the spring would be found using the spring constitutive equation,  $R_{spring} = kq_3$



## To try at home:

- and what if the spring was not grounded?
- Write down the equation of equilibrium in matrix form, with:
  - the degrees of freedom as shown
  - just one element for each of the two elastic rods in tension/compression, and the spring





# Part 2d: Elastic Rods with Distributed Loading

FEEG3001/SESM6047 FEA in Solid Mechanics Dr A S Dickinson

From 18th October 2024

## Recap

- Last time we looked at how to derive and then assemble stiffness matrices for elastic rods in tension and compression
- This was quite simple because we assumed a linear interpolation function (or 'shape function') for how the elements deform
- This is the simplest interpolation we can get away with
- You could try quadratic or cubic... here linear is adequate, and the stiffness matrix becomes messy if we make it more complicated.
- In the (near) future, we will meet situations where linear interpolation is not sufficient...

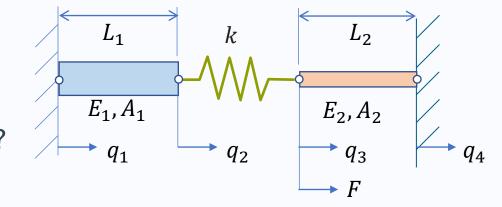
## But why...?

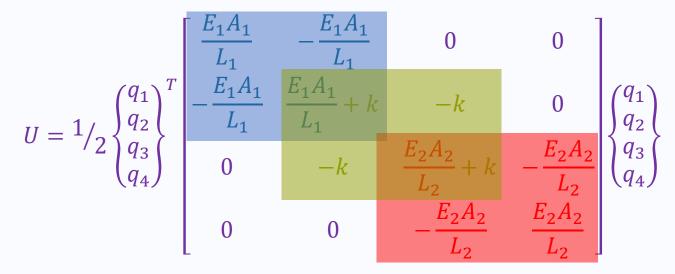




#### Solution:

- and what if the spring was not grounded?
- Short-Hand:
- Where does the k contribute to the stiffness matrix?





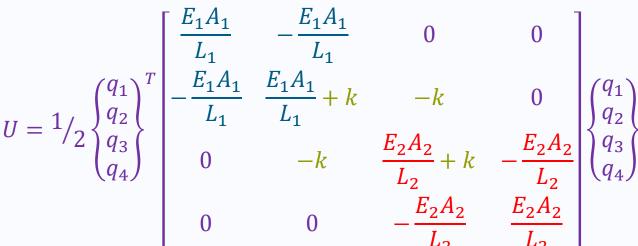
Why?

#### Solution:

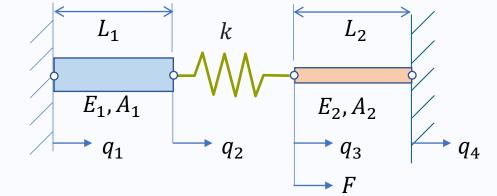
Long-hand:

$$U = U_1 + U_2 + U_{spring}$$

$$U = U_1 + U_2 + \frac{1}{2}k(q_3 - q_2)^2$$



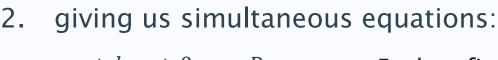




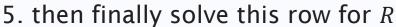
## Recap of our solution sequence in FEM:

- If our problem is as sketched on the right, we can formulate it without the fixed wall on the left, with a variable for that node's displacement, and an unknown reaction force.
- 1. We can assemble a stiffness matrix to solve it:

$$\begin{bmatrix} a & b & 0 \\ c & d & e \\ 0 & f & g \end{bmatrix} \begin{pmatrix} q_1 \\ q_2 \\ q_3 \end{pmatrix} = \begin{pmatrix} R \\ F \\ P \end{pmatrix}$$

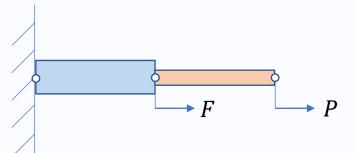


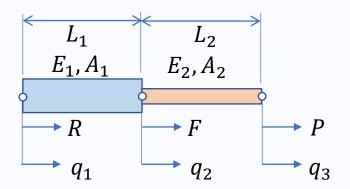
$$aq_1 + bq_2 + 0q_3 = R$$
 $cq_1 + dq_2 + eq_3 = F$ 
 $0q_1 + fq_2 + gq_3 = P$ 



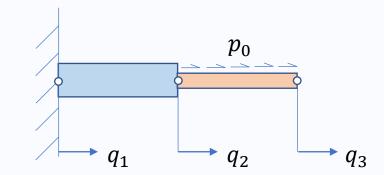
3. Apply BCs: i.e. say 
$$q_1 = 0$$

4. Solve these for  $q_2$  and  $q_3$ 





#### Today a new problem:



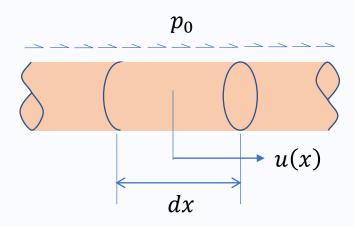
- Three generalised coordinates, variables needed to describe the configuration of the mechanical system
- The properties of the rods don't matter yet.
- How will the Stiffness Matrix be affected?
- As before:

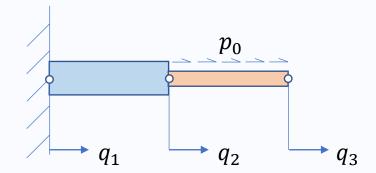
$$U = \frac{1}{2} \{q\}^T [K] \{q\}$$

- but how to deal with a distributed force/unit length,  $p_0$ ?
- Where do the forces come from? Unlike the previous version, the forces do not act at a node here...

- $p_0$   $q_1$   $q_2$   $q_3$
- Where do the forces come from? Unlike the previous version, the forces do not act at a node here...
- What is the work done by distributed force  $p_0$ ?
- What is the work done on a small slice dx?

$$V = -W = -\int_{x=0}^{L_E} p_0 dx \times u(x)$$
 the force the in terms of the unknown element's  $x$ 



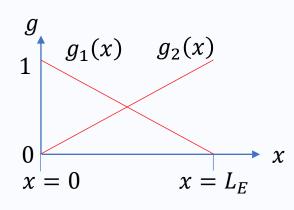


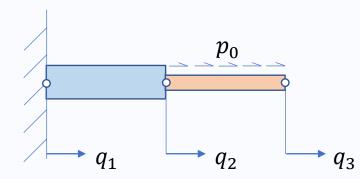
• What is the work done on a small slice u(x)?

$$V = -W = -\int_{x=0}^{L_E} p_0 dx \times u(x)$$

• Recall u(x) is unknown, but we agreed before to approximate it as a linear function, locally within the element. So we can interpolate as:

$$u(x) = \left(1 - \frac{x}{L_E}\right)q_2 + \left(\frac{x}{L_E}\right)q_3$$





So substituting:

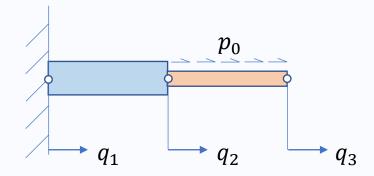
$$V = -W = -\int_{x=0}^{L_E} p_0 \left[ \left( 1 - \frac{x}{L_E} \right) q_2 + \left( \frac{x}{L_E} \right) q_3 \right] dx$$

$$V = -p_0 \left[ \left( x - \frac{x^2}{2L_E} \right) q_2 + \left( \frac{x^2}{2L_E} \right) q_3 \right]_0^{L_E}$$

$$V = -p_0 \left[ \left( L_E - \frac{L_E}{2} \right) q_2 + \left( \frac{L_E}{2} \right) q_3 \right]$$

$$V = -p_0 \left[ \left( \frac{L_E}{2} \right) q_2 + \left( \frac{L_E}{2} \right) q_3 \right]$$

so we see an intuitive (but not general!)
 phenomenon: the total force can be split half and
 half between the degrees of freedom.



$$V = -\left(p_0 \frac{L_E}{2}\right) q_2 - \left(p_0 \frac{L_E}{2}\right) q_3$$

and recall that we expected V to have a form something like:

$$V = -\{F\}^T\{q\}$$

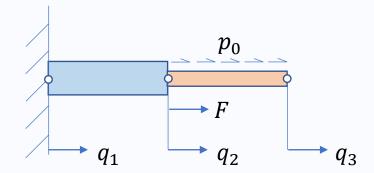
and in that case we known PMTPE will give us a solution of the form:

$$[K]{q} = {F}$$

 So without needing to solve every step, we can say that our generalised equation of equilibrium for the system is:

$$\begin{bmatrix} q_1 \\ q_2 \\ q_3 \end{bmatrix} = \begin{bmatrix} R \\ p_0 \frac{L_E}{2} \\ p_0 \frac{L_E}{2} \end{bmatrix}$$

#### Another example (to show it's easy!):



• Suppose we add a further nodal force *F*:

$$V = -\left(p_0 \frac{L_E}{2} + F\right) q_2 - \left(p_0 \frac{L_E}{2}\right) q_3$$

SO

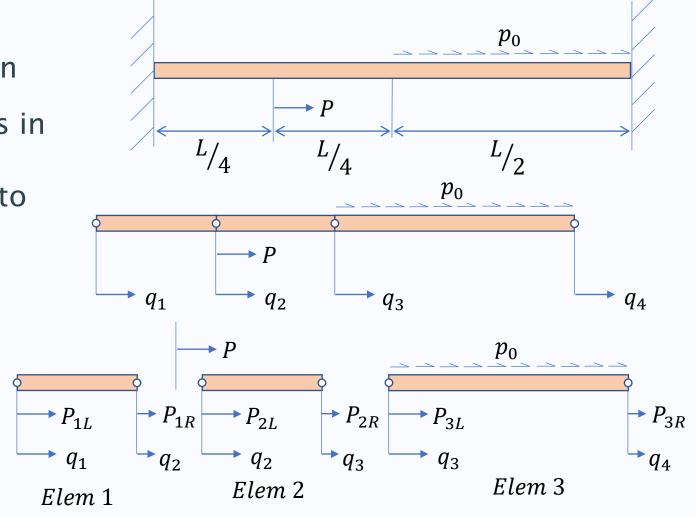
$$\begin{bmatrix} q_1 \\ q_2 \\ q_3 \end{bmatrix} = \begin{bmatrix} R \\ p_0 \frac{L_E}{2} + F \\ p_0 \frac{L_E}{2} \end{bmatrix}$$

#### **Practice Questions**

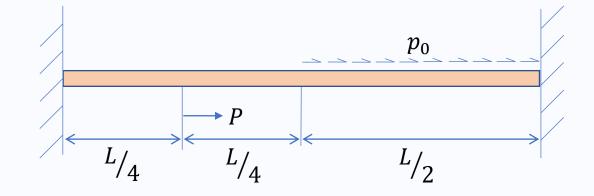
- A typical sequence:
- 1. Draw a Diagram
- 2. Idealise the Elements, calculating equivalent nodal forces
- 3. Assemble the Global Stiffness Matrix and Equilibrium Equation
- 4. Invert the Stiffness Matrix and solve for Nodal Displacements
- 5. Solve for Displacement across elements
- 6. Solve for Strain across elements
- 7. Solve for Stress across elements

- The choice of elements might be defined by the loads we wish to apply, and our required resolution of deformation, strain and stress results, instead of by the changes in rod cross section.
- How many elements do we need to model this scenario?
- At least 3 elements
- (not all working is shown!)
- Nodal forces:

$$P_{1L} = R_L$$
  $P_{1R} + P_{2L} = P$   $P_{2R} + P_{3L} = \frac{p_{0L}}{4}$   $P_{3R} = R_R + \frac{p_{0L}}{4}$ 



Find the generalised equation of equilibrium



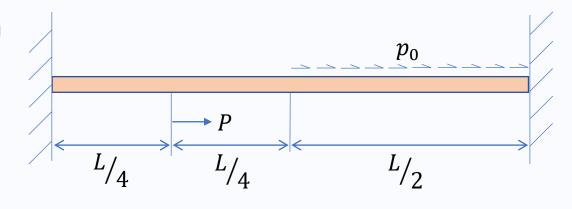
 Apply boundary conditions to obtain the reduced equation of equilibrium

$$\frac{EA}{L} \begin{bmatrix} ? & ? \\ ? & ? \end{bmatrix} \begin{Bmatrix} q_2 \\ q_3 \end{Bmatrix} = \begin{Bmatrix} ? \\ ? \end{Bmatrix}$$

which can be solved, as  $|K^{reduced}| \neq 0$ .

Find the generalised equation of equilibrium

$$\frac{EA}{L} \begin{bmatrix} 4 & -4 & 0 & 0 \\ -4 & 4+4 & -4 & 0 \\ 0 & -4 & 4+2 & -2 \\ 0 & 0 & -2 & 2 \end{bmatrix} \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{Bmatrix} = \begin{Bmatrix} R_L \\ P \\ \frac{p_0 L}{4} \\ \frac{p_0 L}{4} + R_R \end{Bmatrix}$$

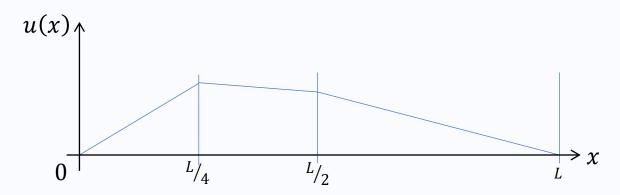


Apply boundary conditions to obtain the reduced equation of equilibrium

$$\frac{EA}{L} \begin{bmatrix} 8 & -4 \\ -4 & 6 \end{bmatrix} \begin{Bmatrix} q_2 \\ q_3 \end{Bmatrix} = \begin{Bmatrix} P \\ p_0 L \\ 4 \end{Bmatrix}$$

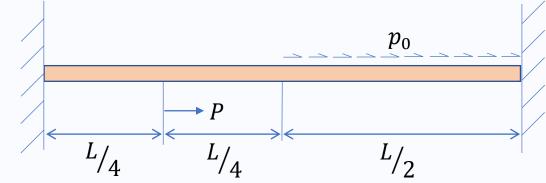
which can be solved, as  $|K^{reduced}| \neq 0$ .

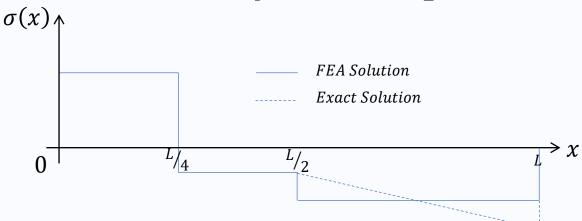
 Finally you could calculate and plot the displacement field and the stress field in the structure.



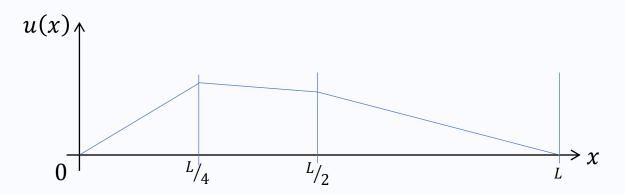


How could correspondence be improved?



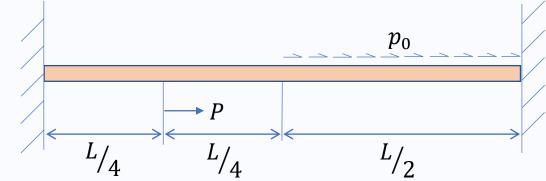


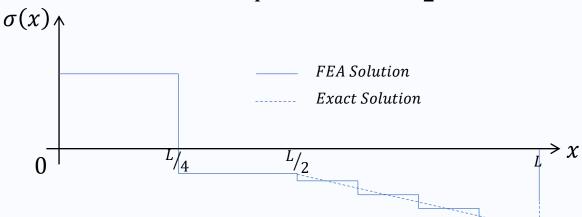
 Finally you could calculate and plot the displacement field and the stress field in the structure.





More subdivisions to improve agreement?





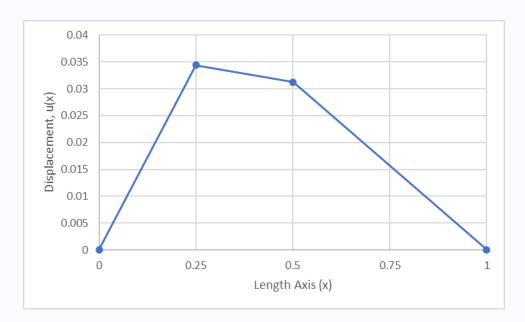


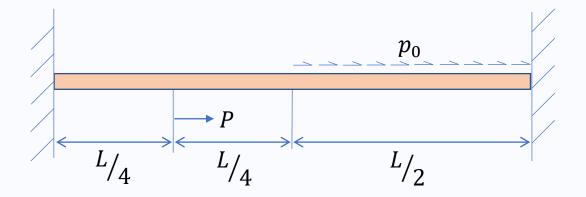
# Part 2e: Elastic Rod Example Questions

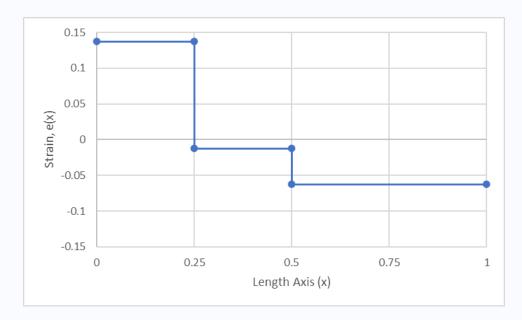
FEEG3001/SESM6047 FEA in Solid Mechanics Dr A S Dickinson

From 25th October 2024

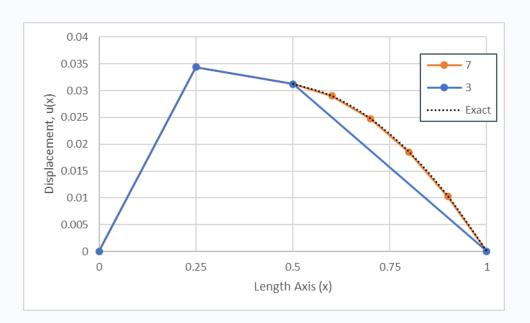
- Where do the lines on this graph come from?
- What do u(x),  $q_1$ ,  $q_2$  etc. mean?
- *x* refers to the undeformed length

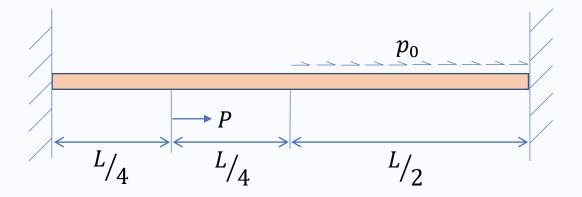


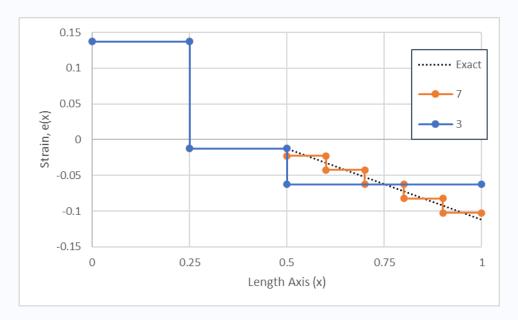




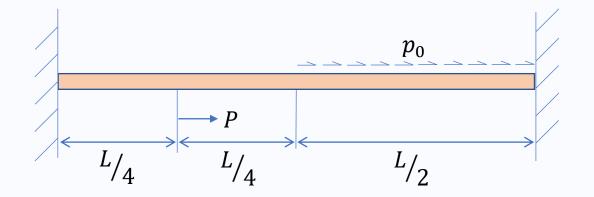
- Where do the lines on this graph come from?
- What do u(x),  $q_1$ ,  $q_2$  etc. mean?
- *x* refers to the undeformed length







How did we get from here:

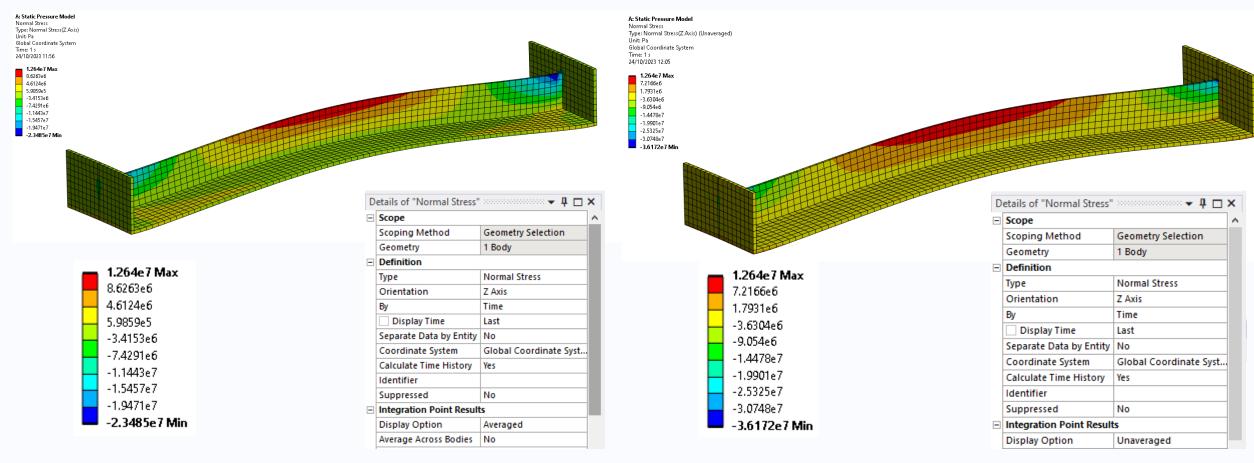


• to here:

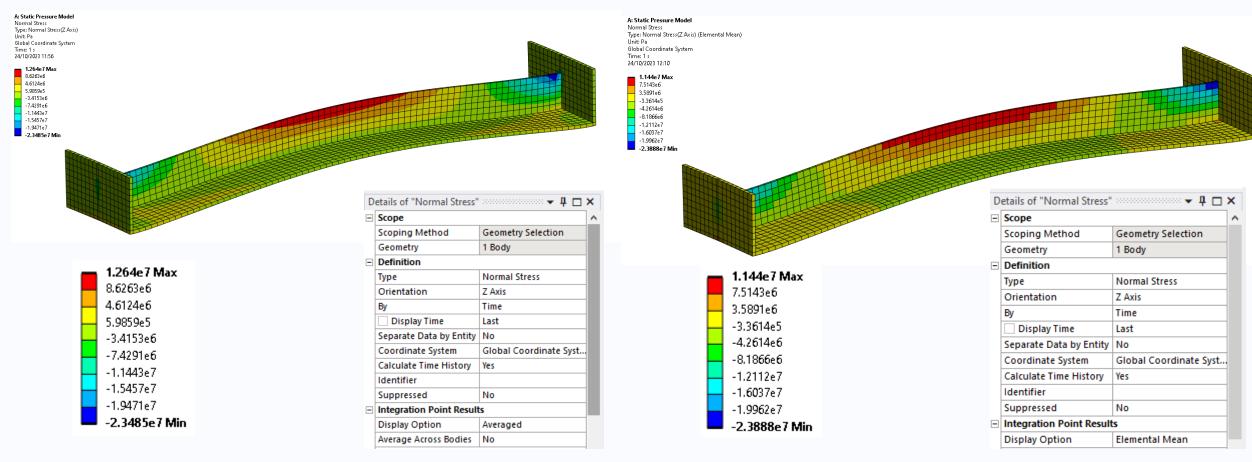
$$\frac{EA}{L} \begin{bmatrix} 4 & -4 & 0 & 0 \\ -4 & 4+4 & -4 & 0 \\ 0 & -4 & 4+2 & -2 \\ 0 & 0 & -2 & 2 \end{bmatrix} \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{Bmatrix} = \begin{Bmatrix} \frac{R_L}{P} \\ \frac{p_0 L}{4} \\ \frac{p_0 L}{4} + R_R \end{Bmatrix}$$

• ?

What did he mean by averaged and unaveraged, or nodal and elemental results?



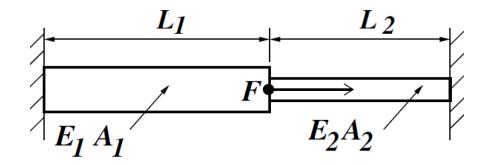
What did he mean by averaged and unaveraged, or nodal and elemental results?



#### Past Exam Questions:

- 2017/18 Question 3
- 2018/19 Question 5
- Q3. Use two finite elements to model the stepped structure in tension and compression as shown in Figure Q3. Calculate reactions at the two fixed ends as well as displacement at the point of application of the external load F, which is the junction of step change in the cross-section.

[12]



**Q5.** (a) A rod, of total length 2L, in tension and compression, is fixed at the left end and sprung at the right end, as shown in Figure Q5. Determine the displacement at the centre of the rod and at the point of connectivity of the rod with the spring. A concentrated force F acts at the centre of the rod, whereas a distributed force  $p_0$  per unit length acts on the right half of the rod, in the directions shown in Figure Q5. The axial stiffness of each half is EA, as shown.

[17]

**(b)** Determine the reaction at the fixed end of the rod.

[3]

(c) Determine the value of the force F for which the spring does not store any strain energy. Comment on this value, concerning the relationship between F and  $p_0$ , for the condition of zero energy stored in the spring.

[3]

