# SESA3029 Aerothermodynamics

Lecture 4.4
Prandtl-Glauert transformation

#### Recap:

Velocity potential equation

$$\left(1 - M_{\infty}^{2}\right) \frac{\partial^{2} \phi}{\partial \mathbf{x}^{2}} + \frac{\partial^{2} \phi}{\partial \mathbf{y}^{2}} = 0$$

- Reduces to Laplace's equation in incompressible flow
- Not valid near M<sub>∞</sub>=1
- Not applicable for M<sub>∞</sub>>5 (hypersonic flow)
- We will consider solution to this equation in two flow regimes
  - M<sub>∞</sub><0.8 (elliptic equation: Prandtl-Glauert transformation)</li>
  - 1.2<M<sub>∞</sub><5 (hyperbolic equation: Ackeret theory)</li>

# Prandtl-Glauert transformation

Start from  $\left(1 - M_{\infty}^{2}\right) \frac{\partial^{2} \phi}{\partial x^{2}} + \frac{\partial^{2} \phi}{\partial y^{2}} = 0$  Let  $\phi_{0} = \beta \phi$ 

$$\mathbf{X}_0 = \mathbf{X}$$

$$\mathbf{y}_0 = \beta \mathbf{y}$$

$$\left(1 - M_{\infty}^{2}\right) \frac{\partial^{2} \phi}{\partial \mathbf{x}^{2}} + \frac{\partial^{2} \phi}{\partial \mathbf{y}^{2}} = 0$$

$$\beta = \sqrt{1 - M_{\infty}^2}$$

$$\beta^{2} \frac{\partial^{2}(\phi_{0}/\beta)}{\partial \mathbf{x}_{0}^{2}} + \frac{\partial^{2}(\phi_{0}/\beta)}{\partial (\mathbf{y}_{0}/\beta)^{2}} = 0$$
i.e. a simple transform takes

$$\frac{\partial^2 \phi_0}{\partial \mathbf{x}_0^2} + \frac{\partial^2 \phi_0}{\partial \mathbf{y}_0^2} = 0$$

 $\frac{\partial^2 \phi_0}{\partial \mathbf{x}_0^2} + \frac{\partial^2 \phi_0}{\partial \mathbf{y}_0^2} = 0$  i.e. a simple transform takes us back to Laplace's equation  $\mathbf{x} = \mathbf{x} + \mathbf{y} = \mathbf{y}$ 

Since 
$$C_p = -2\frac{1}{I}$$

As in incompressible flow

and

$$u' = \frac{\partial \phi}{\partial x}$$

$$C_{p} = -\frac{2}{U_{\infty}} \frac{\partial \phi}{\partial x}$$

$$= -\frac{2}{U_{\infty}\beta} \frac{\partial \phi_0}{\partial \mathbf{x}_0}$$

$$=\frac{C_{p0}}{\sqrt{1-M_{\infty}^2}}$$

where  $C_{p0}$  is the incompressible pressure coefficient

$$C_{p} = \frac{C_{p0}}{\sqrt{1 - M_{\infty}^{2}}}$$

**Prandtl-Glauert** transformation

#### Prandtl-Glauert transformation

- Allows incompressible results to be used in the compressible flow regime
- Applies also to quantities derived from  $C_P$ , for example lift and moment coefficients
- Example: If  $C_L=0.4$  in incompressible flow, find  $C_L$  at  $M_{\infty}=0.6$

$$\frac{C_{L0}}{P} = C_L = \frac{C_{L0}}{\sqrt{1 - M_{\infty}^2}} = \frac{0.4}{\sqrt{1 - 0.6^2}} = 0.5$$

## Use of P-G to estimate $M_{crit}$

and for isentropic flow

$$\frac{\boldsymbol{p}}{\boldsymbol{p}_{\infty}} = \frac{\boldsymbol{p}_{0}/\boldsymbol{p}_{\infty}}{\boldsymbol{p}_{0}/\boldsymbol{p}} = \left(\frac{1 + \frac{\gamma - 1}{2}\boldsymbol{M}_{\infty}^{2}}{1 + \frac{\gamma - 1}{2}\boldsymbol{M}^{2}}\right)^{\frac{\gamma}{\gamma - 1}}$$

Critical Mach number is when the local M=1 and  $M_{\infty}=M_{crit}$  Also apply P-G transform

$$C_{p0} = \frac{C_{p0}}{\sqrt{1 - M_{crit}^2}} = \frac{2}{\gamma M_{crit}^2} \left( \left( \frac{1 + \frac{\gamma - 1}{2} M_{crit}^2}{1 + \frac{\gamma - 1}{2}} \right)^{\frac{\gamma}{\gamma - 1}} - 1 \right)$$

i.e. Knowing the worst case  $C_{p0}$  (the lowest  $C_p$  on the airfoil in incompressible flow), we can estimate  $M_{\rm crit}$ .

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#### Example

• If  $C_{p,min}$ =-0.4 in low speed (incompressible) flow of air, estimate  $M_{crit}$ 

$$\frac{C_{p0}}{\sqrt{1 - M_{crit}^2}} = \frac{2}{\gamma M_{crit}^2} \left( \left( \frac{1 + \frac{\gamma - 1}{2} M_{crit}^2}{1 + \frac{\gamma - 1}{2}} \right)^{\frac{\gamma}{\gamma - 1}} - 1 \right)$$

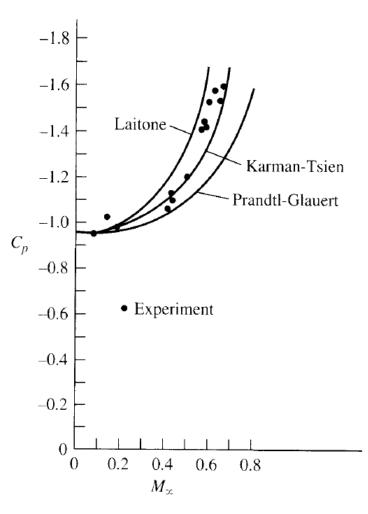
M <sub>crit</sub>	LHS	RHS
0.6	-0.50	-1.29
0.7	-0.56	-0.78
0.747	-0.60	-0.60

Critical Mach number is 0.747

### Improved correction: Karman-Tsien

$$C_{p} = \frac{C_{p,0}}{\sqrt{1 - M_{\infty}^{2}} + \frac{C_{p,0}}{2} \left[ \frac{M_{\infty}^{2}}{1 + \sqrt{1 - M_{\infty}^{2}}} \right]}$$

Improved fit to experimental data – the formula attempts to include nonlinear effects.



(Anderson, 5<sup>th</sup> Ed)