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Quantum HW #1
Due 1/16/19

WELL DONE!
10/10

1.) Using the dual correspondence definition:

$$\langle \alpha | c \rangle \xleftrightarrow{DC} \langle \alpha | c^* \rangle$$

Consider $|\alpha\rangle\langle\beta|$ acting on $|\psi\rangle$, and take the dual correspondence \Rightarrow

$$(|\alpha\rangle\langle\beta|\underbrace{\langle\psi|}_{c})^\dagger = (\langle\alpha|c\rangle)^\dagger \xleftrightarrow{DC} \langle\alpha|c^*\rangle$$

By the associative axiom of multiplication, we could also have written

$$(\underbrace{(|\alpha\rangle\langle\beta|)\cdot|\psi\rangle}_{\hat{X}})^\dagger = \hat{X}^\dagger |\psi\rangle \xleftrightarrow{DC} \langle\psi|\hat{X}^\dagger$$

Using: $c^* = \langle\beta|\psi\rangle^* = \langle\psi|\beta\rangle$, we know that

$$\langle\alpha|c^*\rangle = \langle\psi|\beta\rangle\langle\alpha| = \langle\psi|\hat{X}^\dagger$$

Thus, again by associativity, $\hat{X}^\dagger = (|\alpha\rangle\langle\beta|)^\dagger = |\beta\rangle\langle\alpha|$
 $\langle\psi|\cdot(|\beta\rangle\langle\alpha|)$

Q.E.D.



2.) Starting with RHS.

$$\sum_{b', b''} \langle c' | b' \rangle \langle b' | a' \rangle \langle a' | b'' \rangle \langle b'' | c' \rangle \quad (i)$$

If $[\hat{A}, \hat{B}] = 0$, and if the system is non-degenerate,
then \hat{A} and \hat{B} have a common eigenbasis. ✓

This implies: $\sum_{b''} \langle a' | b'' \rangle = \sum_{b''} \delta_{b', b''} \langle a' | b'' \rangle \}$

This is an odd way of writing the relation, but it will be convenient \Rightarrow

So (i) = $\sum_{b', b''} \langle c' | b' \rangle \langle b' | a' \rangle \delta_{b', b''} \langle a' | b'' \rangle \langle b'' | c' \rangle$

The δ collapses the sum over b'' , leaving us with

$$(i) = \sum_{b'} \underbrace{\langle c' | b' \rangle \langle b' | a' \rangle \langle a' | b' \rangle}_{\text{These } = 1 \text{ b/c}} \langle b' | c' \rangle = \text{L.H.S.}$$

$\{ |a'^{(i)}\rangle \} = \{ |b'^{(i)}\rangle \}$ Q.E.D.

But I leave them in this form for clarity.

(The proof holds for $[B, C] = 0$ as well. Simply replace

$\langle b'' | c' \rangle$ with $\delta_{b', b''} \langle b'' | c' \rangle \dots$)

jk!

3.)

For $|S_z; +\rangle$,

$$\langle S_x \rangle = \langle S_z; + | \hat{S}_x | S_z; + \rangle$$

Expressed in the S_z basis, this gives

$$\langle S_x \rangle = (1 \ 0) \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{\hbar}{2} (1 \ 0) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = 0$$

Which is what one would expect since measurement of S_x yields $+\frac{\hbar}{2}$ and $-\frac{\hbar}{2}$ with equal probability.

Now,

$$\langle S_x^2 \rangle = \langle S_z; + | \hat{S}_x \hat{S}_x | S_z; + \rangle \Rightarrow \text{again, choose } S_z \text{ basis}$$

$$= (1 \ 0) \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$= \frac{\hbar^2}{4} (1 \ 0) \underbrace{\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}}_{\text{II}} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{\hbar^2}{4}$$

$$\text{Thus, } \langle S_x^2 \rangle - \langle S_x \rangle^2$$

$$= \frac{\hbar^2}{4} - (0)^2 = \frac{\hbar^2}{4} \quad \checkmark$$

Yes!

$$4.) H = H_1 |1\rangle\langle 1| + H_2 |2\rangle\langle 2| + H_n (|1\rangle\langle 2| + |2\rangle\langle 1|)$$

Expressing H as a matrix in the $1, 2$ basis \Rightarrow

$$H \doteq \begin{pmatrix} H_1 & H_n \\ H_n & H_2 \end{pmatrix}$$

To Find eigenvalues, $H|\psi\rangle = \lambda|\psi\rangle$, require $\begin{vmatrix} H_1 - \lambda & H_n \\ H_n & H_2 - \lambda \end{vmatrix} = 0$
for non-trivial solutions

The characteristic equation gives

$$\lambda^2 - (H_1 + H_2)\lambda + (H_1 H_2 - H_n^2) = 0$$

Plugging this in to the quadratic formula yields 2 eigenvalues

$$\lambda_{+,-} = \frac{1}{2} \left(H_1 + H_2 \pm \sqrt{H_1^2 + H_2^2 - 2H_1 H_2 + 4H_n^2} \right)$$

To find the corresponding eigenvectors, Solve the two systems

$$H|\Psi_{+,-}\rangle = \lambda_{+,-}|\Psi_{+,-}\rangle \Rightarrow \begin{pmatrix} H_1 & H_n \\ H_n & H_2 \end{pmatrix} \begin{pmatrix} \alpha_{+,-} \\ \beta_{+,-} \end{pmatrix} = \lambda_{+,-} \begin{pmatrix} \alpha_{+,-} \\ \beta_{+,-} \end{pmatrix}$$

After some messy algebra, we find (see attached)

$$\lambda_+ : |\Psi_+\rangle \doteq \begin{pmatrix} \frac{-2 - H_{12}}{H_{11} - H_{12} - \gamma} \\ 1 \end{pmatrix} \quad \text{where } \gamma = \sqrt{H_1^2 + H_2^2 - 2H_1 H_2 + 4H_n^2}$$

$$\lambda_- : |\Psi_-\rangle \doteq \begin{pmatrix} \frac{-2 \cdot H_{12}}{H_{11} - H_{12} + \gamma} \\ 1 \end{pmatrix} \quad \checkmark$$

4.) Messy Algebra

$$\lambda_+ : \begin{pmatrix} H_{11} & H_{12} \\ H_{21} & H_{22} \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \frac{1}{2} (H_{11} + H_{22} + \gamma) \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$

where $\gamma = \sqrt{H_{11}^2 + H_{22}^2 - 2H_{11}H_{22} + 4H_{12}^2}$

$$\Rightarrow \begin{pmatrix} H_{11} - \frac{1}{2}(H_{11} + H_{22} + \gamma) & H_{12} \\ H_{12} & H_{22} - \frac{1}{2}(H_{11} + H_{22} + \gamma) \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = 0$$

$$\text{eq1. } \frac{1}{2}(H_{11} - H_{22} - \gamma)\alpha + H_{12}\beta = 0 \Rightarrow \beta = \frac{-\frac{1}{2}(H_{11} - H_{22} - \gamma)\alpha}{H_{12}}$$

$$\text{eq2. } H_{12}\alpha + \frac{1}{2}(H_{22} - H_{11} - \gamma)\beta = 0 \Rightarrow \text{Plug in for } \beta$$

$$H_{12}\alpha + \frac{-\alpha}{4H_{12}} (H_{11}H_{22} - H_{12}^2 - H_{12}\gamma - H_{11}\gamma + H_{11}H_{22} + H_{12}\gamma - H_{11}\gamma + H_{12}\gamma + \gamma^2) = 0$$

$$\alpha [H_{12} - \frac{1}{4H_{12}} (2H_{11}H_{22} - H_{12}^2 - H_{11}^2 + \gamma^2)] = 0$$

$$\alpha [H_{12} - \frac{1}{4H_{12}} (2H_{11}H_{22} - H_{12}^2 - H_{11}^2 + H_{11}^2 + H_{12}^2 - 2H_{12}H_{22} + 4H_{11}^2)] = 0$$

$$\alpha [H_{12} - H_{11}] = 0 \Rightarrow \text{Redundancy in solution.}$$

$$\text{choose } \alpha = \frac{-2H_{12}}{H_{11} - H_{12} - \gamma}, \text{ Then } \beta = 1$$

$$\text{For } \lambda_-, \text{ it is simple to see. } \alpha = \frac{-2H_{12}}{H_{11} - H_{12} + \gamma}, \beta = 1$$

5.)

We can construct a transformation matrix which connects the natural S_z basis to the natural S_x basis as follows

$$U_{z \rightarrow x} = \sum_k |x^{(k)}\rangle \langle z^{(k)}|$$

In the S_z basis, the matrix elements of $U_{z \rightarrow x}$ are represented as $\langle z^{(k)} | U_{z \rightarrow x} | z^{(l)} \rangle$, which, by our construction of $U_{z \rightarrow x}$ is equivalent to $\langle z^{(k)} | x^{(l)} \rangle$

$$\text{So, } U_{z \rightarrow x} \doteq \begin{pmatrix} \langle +z | +x \rangle & \langle +z | -x \rangle \\ \langle -z | +x \rangle & \langle -z | -x \rangle \end{pmatrix}$$

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} (1 0) \begin{pmatrix} 1 \\ 1 \end{pmatrix} & (1 0) \begin{pmatrix} 1 \\ -1 \end{pmatrix} \\ (0 1) \begin{pmatrix} 1 \\ 1 \end{pmatrix} & (0 1) \begin{pmatrix} 1 \\ -1 \end{pmatrix} \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

Check unitarity:

$$U^\dagger U = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 2 & 0 \\ 0 & 2 \end{pmatrix} = 11 \quad \checkmark$$

✓

6.)

Initially, when α is measured, the eigenvalue a_1 is obtained. Therefore, we know the initial state to be $|\phi_1\rangle$ since $A|\phi_1\rangle = a_1|\phi_1\rangle$.

The state corresponding to a_2 is

Now B is measured, collapsing the state into either $|x_1\rangle$ or $|x_2\rangle$ with probabilities given by:

$$P(|x_1\rangle) = |\langle x_1 | \phi_1 \rangle|^2 = \left| \langle x_1 | \left(\frac{2|x_1\rangle + 3|x_2\rangle}{\sqrt{13}} \right) \right|^2 = \frac{4}{13}$$

$$P(|x_2\rangle) = |\langle x_2 | \phi_1 \rangle|^2 = \left| \langle x_2 | \left(\frac{2|x_1\rangle + 3|x_2\rangle}{\sqrt{13}} \right) \right|^2 = \frac{9}{13}$$

Since $|x\rangle$ eigenstates are orthonormal.

If α is measured again, the probability of observing a_1 a second time is the probability of the state being in either $|x_1\rangle$ or $|x_2\rangle$ times the probability of that state collapsing into $|\phi_1\rangle$:

$$P(a_1 \text{ again}) = |\langle \phi_1 | x_1 \rangle|^2 \cdot \frac{4}{13} + |\langle \phi_1 | x_2 \rangle|^2 \cdot \frac{9}{13}$$

$$= \left(\frac{4}{13} \right)^2 + \left(\frac{9}{13} \right)^2 = \frac{97}{169} \underset{\underline{\hspace{10em}}}{\simeq} 57\%$$

Ans!