

# Computational Physics

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# Why do we need computers to do physics?

# An Illuminating Example - 2-Body Problem

$$\vec{F}_{ij} = \frac{Gm_i m_j}{r_{ij}^2} \hat{r} \quad \sum_j \vec{F}_{ij} = m_i \ddot{\vec{r}}_i \quad (1)$$

- 2 masses
- 1 unique force
- 2 equations of motion in 3 dimensions
- 6 coupled 2nd-order ordinary differential equations

# An Illuminating Example - 3-Body Problem

$$\vec{F}_{ij} = \frac{Gm_i m_j}{r_{ij}^2} \hat{r} \quad \sum_j \vec{F}_{ij} = m_i \ddot{\vec{r}}_i \quad (2)$$

- 3 masses
- 3 unique forces
- 3 equations of motion in 3 dimensions
- 9 coupled 2nd-order ordinary differential equations

# An Illuminating Example - “All”-Body Problem

$$\vec{F}_{ij} = \frac{Gm_i m_j}{r_{ij}^2} \hat{r} \quad \sum_j \vec{F}_{ij} = m_i \ddot{\vec{r}}_i \quad (3)$$

- 1.24 trillion masses
- $7.688 \times 10^{23}$  unique forces
- 1.24 trillion equations of motion  
in 3 dimensions
- 3.72 trillion coupled 2nd-order  
ordinary differential equations

Cosmological N-Body Simulation[22]

# Where do we need computers to do physics?

# Problems in Computational Physics

## Pillars of Physics

- Classical Mechanics → N-Body Simulations & Fluid Dynamics

## Computational Applications

# Problems in Computational Physics

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- Classical Mechanics      → N-Body Simulations & Fluid Dynamics
- Electromagnetism      → Fringing Fields & Antenna Radiation

## Computational Applications

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- Quantum Mechanics → Electronic Structure & Ultra-Cold Gases
- Relativity → Mercury's Perihelion & Black Hole Ray Tracing

## Computational Applications

# Classical Mechanics

## Remark

Given the position, velocity, and forces acting on an object, its motion is completely determined.

- The *equation of motion*:  $\sum_i \mathbf{F}_i = m\ddot{\mathbf{r}}$
- 2nd-order
- Ordinary
- 3 dimensions

# Classical Mechanics - 3-Body Equations of Motion

$$\frac{Gm_2}{((x_1 - x_2)^2 + (y_1 - y_2)^2 + (z_1 - z_2)^2)^{3/2}} \begin{bmatrix} x_2 - x_1 \\ y_2 - y_1 \\ z_2 - z_1 \end{bmatrix} + \frac{Gm_3}{((x_1 - x_3)^2 + (y_1 - y_3)^2 + (z_1 - z_3)^2)^{3/2}} \begin{bmatrix} x_3 - x_1 \\ y_3 - y_1 \\ z_3 - z_1 \end{bmatrix} = \begin{bmatrix} \ddot{x}_1 \\ \ddot{y}_1 \\ \ddot{z}_1 \end{bmatrix} \quad (4)$$

$$\frac{Gm_1}{((x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2)^{3/2}} \begin{bmatrix} x_1 - x_2 \\ y_1 - y_2 \\ z_1 - z_2 \end{bmatrix} + \frac{Gm_3}{((x_2 - x_3)^2 + (y_2 - y_3)^2 + (z_2 - z_3)^2)^{3/2}} \begin{bmatrix} x_3 - x_2 \\ y_3 - y_2 \\ z_3 - z_2 \end{bmatrix} = \begin{bmatrix} \ddot{x}_2 \\ \ddot{y}_2 \\ \ddot{z}_2 \end{bmatrix} \quad (5)$$

$$\frac{Gm_1}{((x_3 - x_1)^2 + (y_3 - y_1)^2 + (z_3 - z_1)^2)^{3/2}} \begin{bmatrix} x_1 - x_3 \\ y_1 - y_3 \\ z_1 - z_3 \end{bmatrix} + \frac{Gm_2}{((x_3 - x_2)^2 + (y_3 - y_2)^2 + (z_3 - z_2)^2)^{3/2}} \begin{bmatrix} x_2 - x_3 \\ y_2 - y_3 \\ z_2 - z_3 \end{bmatrix} = \begin{bmatrix} \ddot{x}_3 \\ \ddot{y}_3 \\ \ddot{z}_3 \end{bmatrix} \quad (6)$$

# Classical Mechanics - Computational Model

- Finite Differences
- Approximates the derivative
- Oldest
- Most straightforward

Central Finite Difference Coefficients

$$\begin{bmatrix} 1 & 1 & \dots & 1 & 1 \\ -p & -p+1 & \dots & p-1 & p \\ (-p)^2 & (-p+1)^2 & \dots & (p-1)^2 & p^2 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ (-p)^{2p} & (-p+1)^{2p} & \dots & (p-1)^{2p} & p^{2p} \end{bmatrix} \begin{bmatrix} a-p \\ a-p+1 \\ \vdots \\ \vdots \\ a_p \end{bmatrix} = \begin{bmatrix} 0 \\ \vdots \\ m! \\ \vdots \\ 0 \end{bmatrix} \quad (7)$$

Ex. The 2-Body Problem

$$\frac{Gm_1m_2}{r_{12}^2} \hat{r}_{12} = m_1 \ddot{\mathbf{r}}_1 \implies \frac{Gm_1m_2}{r_{12,n}^2} \hat{r}_{12} = m_1 \frac{\mathbf{r}_1^{n-1} - 2\mathbf{r}_1^n + \mathbf{r}_1^{n+1}}{h^2} \quad (8)$$

# Classical Mechanics - Computational Model

## Velocity Verlet

- Velocity Verlet Integration
- Useful for equations of motion
- Conserves energy
- The right choice for long time scales

$$n = 0, 1, 2, 3, \dots$$

$$\mathbf{r}_1^{n+1} = \mathbf{r}_1^n + \dot{\mathbf{r}}_1^n \Delta t + \frac{1}{2} \ddot{\mathbf{r}}_1 \Delta t^2$$

$$\dot{\mathbf{r}}_1^{n+1} = \dot{\mathbf{r}}_1^n + \frac{\ddot{\mathbf{r}}_1^n + \ddot{\mathbf{r}}_1^{n+1}}{2}$$

# Classical Mechanics - Computational Model

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## Algorithm 1: N-Body Gravitational Simulation

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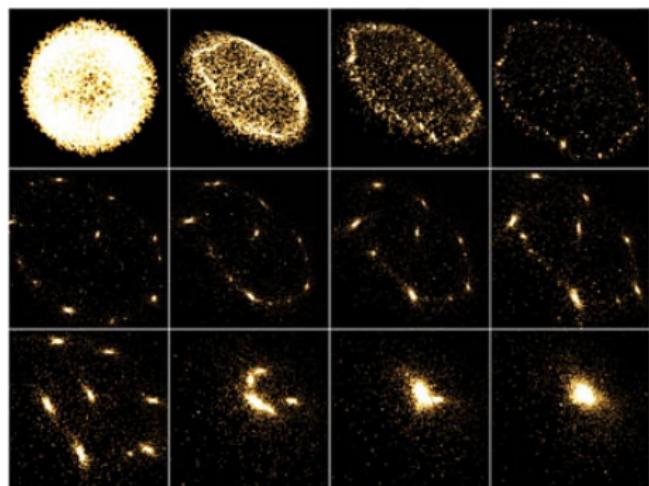
**Input:** Initial Positions, Initial Velocities, Masses

**Output:** The trajectory of the objects

- 1 Discretize the equations of motion
  - 2 Use intial values  $\mathbf{r}^0$  and  $\dot{\mathbf{r}}^0$  to get  $\mathbf{r}^1$  through somthing like Euler integration
  - 3 **while**  $t^n < t^f$  **do**
  - 4     **for**  $i \leq N$  **do**
  - 5         Use discretized equation of motion for object  $i$  to get acceleration  $\ddot{\mathbf{r}}_i^n$
  - 6         Use Velocity Verlet to get new velocity  $\dot{\mathbf{r}}_i^{n+1}$  and position  $\mathbf{r}_i^{n+1}$
  - 7     **end**
  - 8 **end**
-

# Classical Mechanics - Computational Model

- Forces can be calculated in parallel
- N-Body simulations can benefit from CUDA implementations[31]
- “50 times as fast as a highly tuned serial implementation (Elsen et al. 2006)”
- “250 times faster than our portable C implementation”



“Frames from an Interactive 3D Rendering of a 16,384-Body System...”

# Classical Mechanics - Numerical Analysis

- Discretization Error
- Time step  $\Delta t$
- Verlet error =  $\mathcal{O}(\Delta t^2)$
- $\Delta t \rightarrow \Delta t/2 \implies \text{error} \rightarrow \text{error}/4$
- Verify results with energy conservation and energy drift

# Electromagnetism

# Electromagnetism - Mathematical Model

In general, the dynamics of the electric and magnetic fields are described by Maxwell's Equations

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \frac{\rho}{\epsilon_0} & \nabla \cdot \mathbf{B} &= 0 \\ \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} & \nabla \times \mathbf{B} &= \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}.\end{aligned}\tag{9}$$

# Electromagnetism - Mathematical Model

Expanded

$$\begin{aligned}\partial_x E_x(t, \mathbf{r}) + \partial_y E_y(t, \mathbf{r}) + \partial_z E_z(t, \mathbf{r}) &= \frac{\rho}{\epsilon_0} & \partial_x B_x(t, \mathbf{r}) + \partial_y B_y(t, \mathbf{r}) + \partial_z B_z(t, \mathbf{r}) &= 0 \\ \begin{bmatrix} \partial_y E_z - \partial_z E_y \\ \partial_z E_x - \partial_x E_z \\ \partial_x E_y - \partial_y E_x \end{bmatrix} &= -\partial_t \begin{bmatrix} B_x \\ B_y \\ B_z \end{bmatrix} & \begin{bmatrix} \partial_y B_z - \partial_z B_y \\ \partial_z B_x - \partial_x B_z \\ \partial_x B_y - \partial_y B_x \end{bmatrix} &= \mu_0 \rho \begin{bmatrix} v_x \\ v_y \\ v_z \end{bmatrix} + \mu_0 \epsilon_0 \partial_t \begin{bmatrix} E_x \\ E_y \\ E_z \end{bmatrix}\end{aligned}\tag{10}$$

as a coupled system of eight *partial* differential equations

# Electromagnetism - Computational Model

- Finite Differences → Finite-Difference time-domain (Yee's Method)
- Multiple spatial dimensions ⇒ discretized spatial mesh
- Also in time ⇒ multiple runs over mesh computing gradients

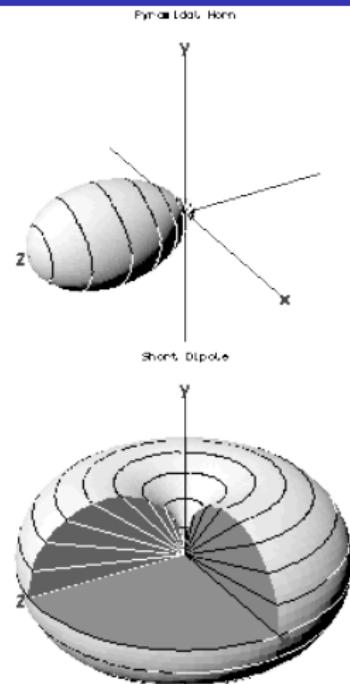


Figure: Antenna radiation pattern

# Electromagnetism - Computational Model

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## Algorithm 2: Electromagnetism Evolution

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**Input:** Initial State of the electromagnetic field

**Output:** Final state of the electromagnetic field

```
1 Discretize space into mesh
2 Discretize time into temporal chunks
3 for  $t_n < t_f$  do
4     foreach cell in mesh do
5         Calculate spatial derivatives
6         Combine derivatives via Maxwell's Equations
7         Get the electric and magnetic field
8     end
9     Calculate the time derivative of the whole mesh
10 end
```

## CFL Condition

Information in the simulation cannot travel faster than it would in reality.

- Rigorously,

$$\frac{\partial w}{\partial t} = u \frac{\partial w}{\partial x} \implies C = \frac{u \Delta t}{\Delta x} \leq C_{\max} \quad (11)$$

- For explicit (time-marching) algorithms,  $C_{\max} = 1$ .
- $C_{\max}$  is called the *Courant Number*.
- If the CFL condition is not met, the results are unphysical.

# Quantum Mechanics

# Quantum Mechanics - Mathematical Model

The Schrödinger equation is the quantum equation of motion

$$H\psi = E\psi \quad (12)$$

and for atoms and molecules

$$H = \underbrace{-\sum_i \frac{\hbar^2}{2m_e} \nabla_i^2}_{e \text{ motion}} - \underbrace{\sum_k \frac{\hbar^2}{2m_e} \nabla_k^2}_{Z \text{ motion}} - \underbrace{\sum_{i,k} \frac{e^2 Z_k}{r_{ik}}}_{e-Z \text{ interaction}} + \underbrace{\sum_{i < k} \frac{e^2}{r_{ik}}}_{e-e \text{ interaction}} + \underbrace{\sum_{k < l} \frac{e^2 Z_k Z_l}{r_{kl}}}_{Z-Z \text{ interaction}} \quad (13)$$

# Quantum Mechanics - Computational Model

- Molecular Mechanics → classical n-body simulation
- ab-initio Self-Consistent Field theory → iteratively solve the Schrödinger equation

# Relativity

# Relativity - Mathematical Model

## The Einstein Field Equations

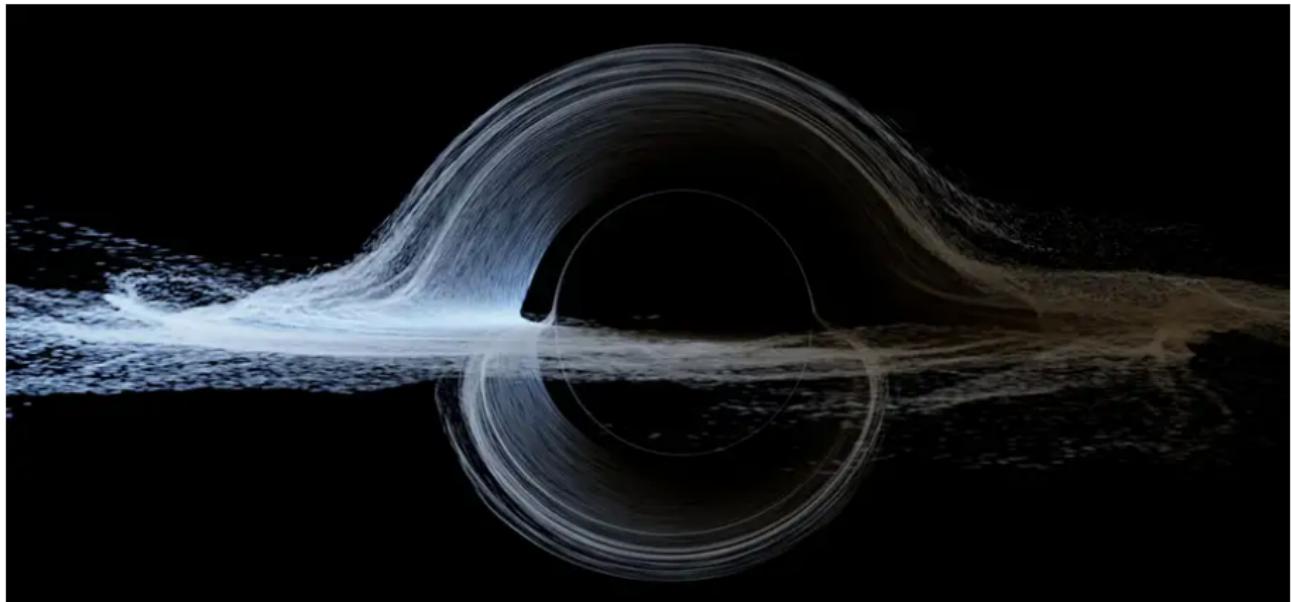
$$\underbrace{G_{\mu\nu}}_{\text{Curvature of Spacetime}} + \underbrace{\Lambda g_{\mu\nu}}_{\text{Expanding Spacetime}} = \kappa \underbrace{T_{\mu\nu}}_{\text{Mass and Energy}} \quad (14)$$

Actually 10 highly nonlinear, coupled partial differential equations in 4 dimensions.

# Relativity - Computational Model

- OpenRelativity visualization[26]
- The physical scenarios considered, such as binary black hole mergers[35] and ray tracing near a black hole[23], yield forms of the EFEs that must be numerically solved with high-performance computing[27, 3, 14].

# Relativity - Results



# Conclusion

- All physics can use computers
- Even the simplest models need computation
- More complex mathematical models require high-performance computing
- Numerical analysis is of the utmost important when developing computational models and results

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