UNIK4660 - Assignment 1

Wilhelm Karlsen and Magnus Eden

March 31, 2016

1 Problem 1

$$u = -y$$
$$v = x$$

w = zt

We are given the velocity field above, in the right handed Cartesian coordinate system. Using this field, we are to show that, at t=0, the stream lines are circles centered on the origin.

This can be done by using the dot products algebraic definition,

$$\bar{V} \bullet \bar{x} = \bar{V}_1 \bar{x}_1 + \bar{V}_2 \bar{x}_2 + \bar{V}_3 \bar{x}_3$$

where \bar{V} is the velocity vector and \bar{x} is the position vector from the origin.

As u and v are independent of time and w = 0 when t = 0, we can ignore the third dimension in this problem.

$$\bar{V} \bullet \bar{x} = \bar{V}_1 \bar{x}_1 + \bar{V}_2 \bar{x}_2$$

$$\bar{V} \bullet \bar{x} = -yx + xy$$

$$\bar{V} \bullet \bar{x} = 0$$

This means that the velocity vector is always orthogonal to the position, which would be the angular velocity.

?Further proff/explanation required?

2 Problem 2

We are to show that the streamlines are circles with center at the z-axis going through the point $(r, 0, z_0)$. We are given the definition of r as $r^2 = x^2 + y^2$,

and given $tan\theta = \frac{y}{x}$. The streamlines can then be expressed by:

$$\frac{dx}{d\theta} = -y$$
$$\frac{dy}{d\theta} = x$$
$$\frac{dz}{d\theta} = zt$$

As this problem is considered at t = 0, the z dimension is always 0 and can be ignored. This then becomes a two dimensional problem and we can use the definition of r to find the position.

$$x = \sqrt{r^2 - y^2} \qquad \qquad y = \sqrt{r^2 - x^2}$$

This can then be inserted into the streamline equations, and since we are only concerned with the point $(r, 0, z_0)$, we get:

$$\frac{dx}{d\theta} = -\sqrt{r^2 - x^2} \qquad = -\sqrt{r^2 - r^2} = 0$$

$$\frac{dy}{d\theta} = \sqrt{r^2 - y^2} \qquad = -\sqrt{r^2 - 0^2} = r$$

$$\frac{dz}{d\theta} = 0$$

?What was the reason this was proof enough!

No streamline that starts at z > 0 can cross z = 0 due to the ?

3 Problem 3

We are to find the formula for the streamline, SL, going through the point $P = (r, 0, w_0)$, when the time t > 0.

Using the result from problem 3, we see that the x and y 'elements' are independent of time and so the only change is that the z-dimension is no longer always 0. We can write the formula as:

$$SL(x, y, z) = -\sqrt{r^2 - x^2} + \sqrt{r^2 - y^2} + zt$$

and for the point P this becomes:

$$SL(r, 0, w_0) = -\sqrt{r^2 - r^2} + \sqrt{r^2 - 0^2} + w_0 t$$

where r is the radius of the circle around on which z lies.

The streamlines going through (r,0,0) are the same as the ones found in problem 3, since when z=0 the z 'element' becomes independent of time.

4 Problem 4

We are to find the path line, PL for the velocity field, \bar{u} , given in problem 1, that goes through the point $P = (r, 0, z_0)$ at t = 0.

The path lines are defined to be the trajectories of an individual particle in a fluid. Where the suffix P is indicates a single particle and not the point P.

$$\begin{cases} \frac{d\bar{x}}{dt} &= \bar{u}_P(\bar{x}_P, t) \\ \bar{x}_P(t_0) &= \bar{x}_{P0} \end{cases}$$

As $t = 0 = t_0$, the path line is simply the position of the particle at t = 0. To prove this we need to derive it from the velocity field at the specified point.

$$\frac{dx}{dt} = \bar{u_1} = -y$$
$$dx = -ydt$$
$$\int_0^x dx = \int_0^t -ydt$$
$$x = x_0 - yt$$

$$\frac{dy}{dt} = \bar{u}_2 = x$$
$$dy = xdt$$
$$\int_0^y dy = \int_0^t xdt$$
$$y = xt + y_0$$

$$\frac{dz}{dt} = \bar{u_3} = zt$$

$$\frac{dz}{z} = tdt$$

$$\int_0^y \frac{1}{z} dz = \int_0^t tdt$$

$$ln(\frac{z}{z_0}) = \frac{1}{2}t^2$$

$$z = z_0 e^{\frac{1}{2}t^2}$$

The path line can then be written as:

$$PL(x, y, z, t) = (x_0 - yt)\hat{\mathbf{i}} + (xt + y_0)\hat{\mathbf{j}} + (z_0e^{\frac{1}{2}t^2})\hat{\mathbf{k}}$$

The path line on point P at As t = 0 becomes:

$$PL(r, 0, z_0, 0) = (x_0)\hat{\mathbf{i}} + (y_0)\hat{\mathbf{j}} + (z_0)\hat{\mathbf{k}}$$

5 Problem 5

We are to find the expression for the streak lines emerging from the point $(r_0, 0, z_0)$.

Streak lines are the curve connecting all particles that pass through a given point, p, from time $t_0 \le s \le t$, and is defined as $x = x(\xi[p, s], t)$.

We derive an expression for the spatial coordinates from the velocity field given in the same way as we did in problem 4.

$$\frac{dx}{dt} = u$$

$$x(x, y, z, t) = (x_0 - yt)\hat{\mathbf{i}} + (y_0 + xt)\hat{\mathbf{j}} + (z_0 e^{\frac{1}{2}t^2})\hat{\mathbf{k}}$$

We can replace the initial coordinates, x_0 , y_0 and z_0 , with the material coordinates ξ_i , since ξ is equal to the initial spatial coordinates for each particle in the flow.

$$x(x, y, z, t) = (\xi_1 - yt)\mathbf{\hat{i}} + (\xi_2 + xt)\mathbf{\hat{j}} + (\xi_3 e^{\frac{1}{2}t^2})\mathbf{\hat{k}}$$

Reshuffling this gives us:

$$\xi_i = (x+yt)\hat{\mathbf{i}} + (y-xt)\hat{\mathbf{j}} + (ze^{-\frac{1}{2}t^2})\hat{\mathbf{k}}$$
 (5.1)

The expression for the particle at position p(x, y, z) and time s is then:

$$\xi_i = (x+ys)\hat{\mathbf{i}} + (y-xs)\hat{\mathbf{j}} + (z^{-\frac{1}{2}s^2})\hat{\mathbf{k}}$$
 (5.2)

The streak line expression is then found by inserting (5.1) into (5.2) and using the coordinates for the emerging point, $(r_o, 0, z_0)$:

$$(x+yt)\hat{\mathbf{i}} + (y-xt)\hat{\mathbf{j}} + (ze^{-\frac{1}{2}t^2})\hat{\mathbf{k}} = (x+ys)\hat{\mathbf{i}} + (y-xs)\hat{\mathbf{j}} + (z^{-\frac{1}{2}s^2})\hat{\mathbf{k}}$$

$$x_1 = r_0$$
 $x_2 = r_0 t - r_0 s$ $x_3 = z_0 e^{\frac{1}{2}t^2 - \frac{1}{2}s^2}$

$$Streakline = r_0 \hat{\mathbf{i}} + (r_0 t - r_0 s) \hat{\mathbf{j}} + (z_0 e^{\frac{1}{2}t^2 - \frac{1}{2}s^2}) \hat{\mathbf{k}}$$

6 Problem 6

Using the velocity field given in problem 1, we are to calculate the vorticity field. The vorticity field is defined as $\nabla \times \bar{v}$. This gives us:

$$\nabla \times \bar{v} = \begin{vmatrix} \bar{i} & \bar{j} & \bar{k} \\ \frac{d}{dx} & \frac{d}{dy} & \frac{d}{dz} \\ -y & x & zt \end{vmatrix}$$

$$= (\frac{dzt}{dy} - \frac{dx}{dz})\bar{i} + (\frac{dy}{dz} - \frac{dzt}{dx})\bar{j} + (\frac{dx}{dx} + \frac{dy}{dy})\bar{k}$$

$$Vorticity = 0\bar{i} + 0\bar{j} + 2\bar{k}$$