

UNIK4660 - Assignment 1

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1 Problem 1

$$u = -y$$

$$v = x$$

$$w = zt$$

We are given the velocity field above, in the right handed Cartesian coordinate system. Using this field, we are to show that, at $t = 0$, the stream lines are circles centered on the origin.

This can be done by using the dot products algebraic definition,

$$\bar{V} \bullet \bar{x} = \bar{V}_1 \bar{x}_1 + \bar{V}_2 \bar{x}_2 + \bar{V}_3 \bar{x}_3$$

where \bar{V} is the velocity vector and \bar{x} is the position vector from the origin.

As u and v are independent of time and $w = 0$ when $t = 0$, we can ignore the third dimension in this problem.

$$\bar{V} \bullet \bar{x} = \bar{V}_1 \bar{x}_1 + \bar{V}_2 \bar{x}_2$$

$$\bar{V} \bullet \bar{x} = -yx + xy$$

$$\bar{V} \bullet \bar{x} = 0$$

This means that the velocity vector is always orthogonal to the position, which would be the angular velocity.

?Further proff/explanation required?

2 Problem 2

We are to show that the streamlines are circles with center at the z-axis going through the point $(r, 0, z_0)$. We are given the definition of r as $r^2 = x^2 + y^2$,

and given $\tan\theta = \frac{y}{x}$. The streamlines can then be expressed by:

$$\begin{aligned}\frac{dx}{d\theta} &= -y \\ \frac{dy}{d\theta} &= x \\ \frac{dz}{d\theta} &= zt\end{aligned}$$

As this problem is considered at $t = 0$, the z dimension is always 0 and can be ignored. This then becomes a two dimensional problem and we can use the definition of r to find the position.

$$x = \sqrt{r^2 - y^2} \qquad y = \sqrt{r^2 - x^2}$$

This can then be inserted into the streamline equations, and since we are only concerned with the point $(r, 0, z_0)$, we get:

$$\begin{aligned}\frac{dx}{d\theta} &= -\sqrt{r^2 - x^2} &= -\sqrt{r^2 - r^2} = 0 \\ \frac{dy}{d\theta} &= \sqrt{r^2 - y^2} &= -\sqrt{r^2 - 0^2} = r \\ \frac{dz}{d\theta} &= 0\end{aligned}$$

?What was the reason this was proof enough!

No streamline that starts at $z > 0$ can cross $z = 0$ due to the ?

3 Problem 3

We are to find the formula for the streamline, SL , going through the point $P = (r, 0, w_0)$, when the time $t > 0$.

Using the result from problem 3, we see that the x and y 'elements' are independent of time and so the only change is that the z -dimension is no longer always 0. We can write the formula as:

$$SL(x, y, z) = -\sqrt{r^2 - x^2} + \sqrt{r^2 - y^2} + zt$$

and for the point P this becomes:

$$SL(r, 0, w_0) = -\sqrt{r^2 - r^2} + \sqrt{r^2 - 0^2} + w_0t$$

where r is the radius of the circle around on which z lies.

The streamlines going through $(r, 0, 0)$ are the same as the ones found in problem 3, since when $z = 0$ the z 'element' becomes independent of time.

4 Problem 4

We are to find the path line, PL for the velocity field, \bar{u} , given in problem 1, that goes through the point $P = (r, 0, z_0)$ at $t = 0$.

The path lines are defined to be the trajectories of an individual particle in a fluid. Where the suffix P indicates a single particle and not the point P .

$$\begin{cases} \frac{d\bar{x}}{dt} &= \bar{u}_P(\bar{x}_P, t) \\ \bar{x}_P(t_0) &= \bar{x}_{P0} \end{cases}$$

As $t = 0 = t_0$, the path line is simply the position of the particle at $t = 0$. To prove this we need to derive it from the velocity field at the specified point.

$$\begin{aligned} \frac{dx}{dt} &= \bar{u}_1 = -y \\ dx &= -ydt \\ \int_0^x dx &= \int_0^t -ydt \\ x &= x_0 - yt \end{aligned}$$

$$\begin{aligned} \frac{dy}{dt} &= \bar{u}_2 = x \\ dy &= xdt \\ \int_0^y dy &= \int_0^t xdt \\ y &= xt + y_0 \end{aligned}$$

$$\begin{aligned} \frac{dz}{dt} &= \bar{u}_3 = zt \\ \frac{dz}{z} &= tdt \\ \int_0^y \frac{1}{z} dz &= \int_0^t tdt \\ \ln\left(\frac{z}{z_0}\right) &= \frac{1}{2}t^2 \\ z &= z_0 e^{\frac{1}{2}t^2} \end{aligned}$$

The path line can then be written as:

$$PL(x, y, z, t) = (x_0 - yt)\hat{\mathbf{i}} + (xt + y_0)\hat{\mathbf{j}} + (z_0 e^{\frac{1}{2}t^2})\hat{\mathbf{k}}$$

The path line on point P at $t = 0$ becomes:

$$PL(r, 0, z_0, 0) = (x_0)\hat{\mathbf{i}} + (y_0)\hat{\mathbf{j}} + (z_0)\hat{\mathbf{k}}$$

5 Problem 5

We are to find the expression for the streak lines emerging from the point $(r_0, 0, z_0)$.

Streak lines are the curve connecting all particles that pass through a given point, p , from time $t_0 \leq s \leq t$, and is defined as $x = x(\xi[p, s], t)$.

We derive an expression for the spatial coordinates from the velocity field given in the same way as we did in problem 4.

$$\begin{aligned} \frac{dx}{dt} &= u \\ x(x, y, z, t) &= (x_0 - yt)\hat{\mathbf{i}} + (y_0 + xt)\hat{\mathbf{j}} + (z_0 e^{\frac{1}{2}t^2})\hat{\mathbf{k}} \end{aligned}$$

We can replace the initial coordinates, x_0 , y_0 and z_0 , with the material coordinates ξ_i , since ξ is equal to the initial spatial coordinates for each particle in the flow.

$$x(x, y, z, t) = (\xi_1 - yt)\hat{\mathbf{i}} + (\xi_2 + xt)\hat{\mathbf{j}} + (\xi_3 e^{\frac{1}{2}t^2})\hat{\mathbf{k}}$$

Reshuffling this gives us:

$$\xi_i = (x + yt)\hat{\mathbf{i}} + (y - xt)\hat{\mathbf{j}} + (ze^{-\frac{1}{2}t^2})\hat{\mathbf{k}} \quad (5.1)$$

The expression for the particle at position $p(x, y, z)$ and time s is then:

$$\xi_i = (x + ys)\hat{\mathbf{i}} + (y - xs)\hat{\mathbf{j}} + (ze^{-\frac{1}{2}s^2})\hat{\mathbf{k}} \quad (5.2)$$

The streak line expression is then found by inserting (5.1) into (5.2) and using the coordinates for the emerging point, $(r_0, 0, z_0)$:

$$(x + yt)\hat{\mathbf{i}} + (y - xt)\hat{\mathbf{j}} + (ze^{-\frac{1}{2}t^2})\hat{\mathbf{k}} = (x + ys)\hat{\mathbf{i}} + (y - xs)\hat{\mathbf{j}} + (ze^{-\frac{1}{2}s^2})\hat{\mathbf{k}}$$

$$x_1 = r_0 \quad x_2 = r_0 t - r_0 s \quad x_3 = z_0 e^{\frac{1}{2}t^2 - \frac{1}{2}s^2}$$

$$Streakline = r_0\hat{\mathbf{i}} + (r_0 t - r_0 s)\hat{\mathbf{j}} + (z_0 e^{\frac{1}{2}t^2 - \frac{1}{2}s^2})\hat{\mathbf{k}}$$

6 Problem 6

Using the velocity field given in problem 1, we are to calculate the vorticity field. The vorticity field is defined as $\nabla \times \bar{v}$. This gives us:

$$\begin{aligned}\nabla \times \bar{v} &= \begin{vmatrix} \bar{i} & \bar{j} & \bar{k} \\ \frac{d}{dx} & \frac{d}{dy} & \frac{d}{dz} \\ -y & x & zt \end{vmatrix} \\ &= \left(\frac{dzt}{dy} - \frac{dx}{dz}\right)\bar{i} + \left(\frac{dy}{dz} - \frac{dzt}{dx}\right)\bar{j} + \left(\frac{dx}{dx} + \frac{dy}{dy}\right)\bar{k} \\ \text{Vorticity} &= 0\bar{i} + 0\bar{j} + 2\bar{k}\end{aligned}$$

7 Problem 7

Since

$$\frac{\partial u_i}{\partial t_j} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial y_i} \right) + \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial y_i} \right) = s_{ij} + \frac{1}{2} \omega_{ij}$$

and

$$\omega_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial y_i} \right)$$

we can find the strain rate tensor by simply expanding the tensor formula and solving for it using our two vectors given by the velocity field

$$\vec{x} = \langle x, y, z \rangle$$

$$\vec{u} = \langle -y, x, zt \rangle$$

$$s_{ij} = \begin{bmatrix} \frac{\partial -y}{\partial x} & \frac{1}{2} \left(\frac{\partial x}{\partial x} + \frac{\partial -y}{\partial y} \right) & \frac{1}{2} \left(\frac{\partial zt}{\partial x} + \frac{\partial -y}{\partial z} \right) \\ \frac{1}{2} \left(\frac{\partial -y}{\partial y} + \frac{\partial x}{\partial x} \right) & \frac{\partial x}{\partial y} & \frac{1}{2} \left(\frac{\partial zt}{\partial y} + \frac{\partial x}{\partial z} \right) \\ \frac{1}{2} \left(\frac{\partial -y}{\partial z} + \frac{\partial zt}{\partial x} \right) & \frac{1}{2} \left(\frac{\partial x}{\partial z} + \frac{\partial zt}{\partial y} \right) & \frac{\partial zt}{\partial z} \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & t \end{bmatrix}$$

The rotation tensor can similarly be found using the same method.

$$\omega_{ij} = \begin{bmatrix} \frac{\partial -y}{\partial x} - \frac{\partial -y}{\partial x} & \frac{1}{2} \left(\frac{\partial x}{\partial x} - \frac{\partial -y}{\partial y} \right) & \frac{1}{2} \left(\frac{\partial zt}{\partial x} - \frac{\partial -y}{\partial z} \right) \\ \frac{1}{2} \left(\frac{\partial -y}{\partial y} - \frac{\partial x}{\partial x} \right) & \frac{\partial x}{\partial y} - \frac{\partial x}{\partial y} & \frac{1}{2} \left(\frac{\partial zt}{\partial y} - \frac{\partial x}{\partial z} \right) \\ \frac{1}{2} \left(\frac{\partial -y}{\partial z} - \frac{\partial zt}{\partial x} \right) & \frac{1}{2} \left(\frac{\partial x}{\partial z} - \frac{\partial zt}{\partial y} \right) & \frac{\partial zt}{\partial z} - \frac{\partial zt}{\partial z} \end{bmatrix} = \begin{bmatrix} 0 & 2 & 0 \\ -2 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

8 Problem 8

We are to calculate the relative dilatation following a particle path.