

Solution to Problem Sheet 1

Noor E Mustafa Ferdous

Solve for problem no. 1

Given

$$\begin{aligned}\eta_{\mu\nu}x^\mu x^\nu &= \eta_{\mu\nu}x'^\mu x'^\nu \\ &= \eta_{\mu\nu}\Lambda^\mu_\sigma x^\sigma \Lambda^\nu_\tau x^\tau \\ &= \eta_{\mu\nu}\Lambda^\mu_\sigma \Lambda^\nu_\tau x^\sigma x^\tau \\ &= \eta_{\sigma\tau}\Lambda^\sigma_\mu \Lambda^\tau_\nu x^\mu x^\nu \quad (\sigma \rightarrow \mu, \tau \rightarrow \nu)\end{aligned}$$

Therefore,

$$\eta_{\mu\nu} = \eta_{\sigma\tau}\Lambda^\sigma_\mu \Lambda^\tau_\nu \quad (1)$$

Again, given

$$\Lambda^\mu_\nu = \delta^\mu_\nu + \omega^\mu_\nu \quad (2)$$

putting this into eqn(1)

$$\begin{aligned}\eta_{\mu\nu} &= \eta_{\sigma\tau}(\delta^\sigma_\mu + \omega^\sigma_\mu)(\delta^\tau_\nu + \omega^\tau_\nu) \\ &= \eta_{\sigma\tau}(\delta^\sigma_\mu \delta^\tau_\nu + \delta^\sigma_\mu \omega^\tau_\nu + \delta^\tau_\nu \omega^\sigma_\mu) \quad [\text{ignoring higher order } \omega] \\ &= \eta_{\mu\nu}\eta_{\mu\tau}\omega^\tau_\nu + \eta_{\sigma\nu}\omega^\sigma_\mu \\ &= \eta_{\mu\nu} + \omega_{\mu\nu} + \omega_{\nu\mu} \\ &= \eta_{\mu\nu} \quad [\omega_{\mu\nu} = -\omega_{\nu\mu}]\end{aligned}$$

Therefore infinitesimal transformation around identity of the form of eqn(2) is a Lorentz transformation.

Solve for problem no. 2

For real scalar field ϕ the Euler-Langrange eqn will be

$$\frac{\partial \mathcal{L}}{\partial \phi} - \partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} = 0 \quad (3)$$

putting

$$\mathcal{L} = \partial_\mu \phi^* \partial^\mu \phi - m^2 \phi^* \phi - \frac{\lambda^2}{2} (\phi^* \phi)^2 \quad (4)$$

in eqn (3), we get,

$$\begin{aligned} -m^2 \phi^* - \partial_\mu \partial^\mu \phi^* &= 0 \\ (\partial_\mu \partial^\mu + m^2 + \lambda^2 \phi^* \phi) \phi^* &= 0 \end{aligned}$$

And for complex scalar field ϕ^* the Euler-Langrange eqn will be

$$\frac{\partial \mathcal{L}}{\partial \phi^*} - \partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^*)} = 0 \quad (5)$$

putting eqn (4) in eqn (5) will be

$$\begin{aligned} -m^2 \phi - \partial_\mu \partial^\mu \phi &= 0 \\ (\partial_\mu \partial^\mu + m^2 + \lambda^2 \phi^* \phi) \phi &= 0 \end{aligned}$$

For deriving Noether theorem of given Langrangian. The infinitesimal transformation of the fields are

$$\begin{aligned} \phi(x) &\rightarrow \phi'(x) = e^{i\alpha} \phi(x) \\ \phi^\dagger(x) &\rightarrow \phi'^\dagger(x) = \phi^\dagger(x) e^{-i\alpha} \end{aligned}$$

where θ is a real constant parameter of a transformation (a global transformation). Such a transformation is not a space-time symmetry transformation since the space-time coordinates are not changed by this transformation, such a transformation is known as an internal symmetry transformation, Infinitesimally, the transformation takes the form of

$$\begin{aligned}
\delta\phi(x) &= \phi'(x) - \phi(x) = i\alpha\phi \\
\delta\phi^\dagger(x) &= \phi'^\dagger(x) - \phi^\dagger(x) = -i\alpha\phi^\dagger \\
x^\mu &= x'^\mu
\end{aligned}$$

so the least action principle will be

$$\begin{aligned}
&\delta S = 0 \\
\Rightarrow \int d^4x' \mathcal{L}(\phi'(x'), \partial_\mu^\dagger \phi'(x'), \phi'^\dagger(x'), \partial_\mu^\dagger \phi'^\dagger(x')) - \int d^4x \mathcal{L}(\phi(x), \partial_\mu \phi(x), \phi^\dagger(x), \partial_\mu \phi^\dagger(x)) &= 0 \\
\Rightarrow \int d^4x \mathcal{L}(\phi'(x), \partial_\mu \phi'(x), \phi'^\dagger(x), \partial_\mu^\dagger \phi'^\dagger(x)) - \int d^4x \mathcal{L}(\phi(x), \partial_\mu \phi(x), \phi^\dagger(x), \partial_\mu \phi^\dagger(x)) &= 0 \\
\Rightarrow \mathcal{L}(\phi'(x), \partial_\mu \phi'(x), \phi'^\dagger(x), \partial_\mu^\dagger \phi'^\dagger(x)) - \mathcal{L}(\phi(x), \partial_\mu \phi(x), \phi^\dagger(x), \partial_\mu \phi^\dagger(x)) &= K^\mu
\end{aligned} \tag{6}$$

For internal symmetry $K^\mu = 0$ And

$$\begin{aligned}
\delta(\partial_\mu \phi(x)) &= \partial_\mu \phi'(x) - \partial_\mu \phi(x) \\
&= \partial_\mu \delta\phi(x) \\
\delta(\partial_\mu^\dagger \phi(x)) &= \partial_\mu \phi'^\dagger(x) - \partial_\mu \phi^\dagger(x) \\
&= \partial_\mu \delta\phi^\dagger(x)
\end{aligned}$$

Therefore

$$\begin{aligned}
&\mathcal{L}(\phi'(x), \partial_\mu \phi'(x), \phi'^\dagger(x), \partial_\mu^\dagger \phi'^\dagger(x)) - \mathcal{L}(\phi(x), \partial_\mu \phi(x), \phi^\dagger(x), \partial_\mu \phi^\dagger(x)) \\
&= \mathcal{L}(\phi(x), \partial_\mu \phi(x), \phi^\dagger(x), \partial_\mu \phi^\dagger(x)) + \delta\phi(x) \frac{\partial \mathcal{L}}{\partial \phi(x)} + \delta(\partial_\mu \phi(x)) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi(x)} \\
&+ \delta\phi^\dagger(x) \frac{\partial \mathcal{L}}{\partial \phi^\dagger(x)} + \delta(\partial_\mu \phi^\dagger(x)) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^\dagger(x)} - \mathcal{L}(\phi(x), \partial_\mu \phi(x), \phi^\dagger(x), \partial_\mu \phi^\dagger(x)) \\
&= \delta\phi(x) \frac{\partial \mathcal{L}}{\partial \phi(x)} + \partial_\mu(\delta\phi(x)) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi(x)} + \delta\phi^\dagger(x) \frac{\partial \mathcal{L}}{\partial \phi^\dagger(x)} + \partial_\mu(\delta\phi^\dagger(x)) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^\dagger(x)}
\end{aligned}$$

$$= \partial_\mu \left(\delta\phi(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi(x)} + \delta\phi^\dagger(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^\dagger(x)} \right) \quad (7)$$

Comparing eqn (6) eqn (7) we get,

$$\partial_\mu \left(\delta\phi(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi(x)} + \delta\phi^\dagger(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^\dagger(x)} \right) = \partial K_\mu \quad (8)$$

$$\partial_\mu \left(\delta\phi(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi(x)} + \delta\phi^\dagger(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^\dagger(x)} - K_\mu \right) = 0 = \partial_\mu J^\mu \quad (9)$$

Which is Noether current. Because $K_\mu = 0$, we get

$$\partial_\mu \left(\delta\phi(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi(x)} + \delta\phi^\dagger(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^\dagger(x)} \right) = \partial_\mu J^\mu \quad (10)$$

$$\delta\phi(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi(x)} + \delta\phi^\dagger(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^\dagger(x)} = J^\mu \quad (11)$$

Now

$$\begin{aligned} \delta\phi(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi(x)} + \delta\phi^\dagger(x) \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^\dagger(x)} \\ = i\alpha\phi(x)\partial^\mu\phi^\dagger(x) - i\alpha\phi^\dagger(x)\partial^\mu\phi(x) \\ = i\alpha(\phi(x)\partial^\mu\phi^\dagger(x) - \phi^\dagger(x)\partial^\mu\phi(x)) \end{aligned}$$

$$J^\mu = i\alpha\phi(x)\overleftrightarrow{\partial}_\mu\phi^\dagger \quad (12)$$

Given Lagrangian is

$$\mathcal{L} = \partial_\mu\phi^*\partial^\mu\phi - m^2\phi^*\phi - \frac{\lambda^2}{2}(\phi^*\phi)^2 \quad (13)$$

where

$$\phi(x) = \frac{1}{\sqrt{2}}(\phi_1(x) + i\phi_2(x)) \quad (14)$$

$$\phi(x) = \frac{1}{\sqrt{2}}(\phi_1(x) - i\phi_2(x)) \quad (15)$$

putting these in eqn (13), we get

$$\mathcal{L} = \frac{1}{2} \partial_\mu (\phi_1 + i\phi_2) \frac{1}{2} \partial^\mu (\phi_1 - i\phi_2) - \frac{m^2}{2} (\phi_1 + i\phi_2) (\phi_1 - i\phi_2) - \frac{\lambda^2}{4} ((\phi_1 + i\phi_2) (\phi_1 - i\phi_2))^2$$

Therefore

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi_1 \partial^\mu \phi_1 - \frac{m^2}{2} \phi_1^2 - \frac{\lambda^2}{4} \phi_1^4 + \frac{1}{2} \partial_\mu \phi_2 \partial^\mu \phi_2 - \frac{m^2}{2} \phi_2^2 - \frac{\lambda^2}{4} \phi_2^4 \quad (16)$$

Solve for problem no. 3

From Wick's theory we know that

$$T(\phi_1 \phi_2 \dots \phi_n) = N(\phi_1 \phi_2 \dots \phi_n) + \text{all possible contractions} \quad (17)$$

Wick's theory for three scalar product will be

$$\begin{aligned} T(\phi(x_1) \phi(x_2) \phi(x_3)) = & : \phi(x_1) \phi(x_2) \phi(x_3) : + \overline{\phi(x_1) \phi(x_2)} : \phi(x_3) : + \overline{\phi(x_2) \phi(x_3)} : \phi(x_1) : \\ & + \overline{\phi(x_3) \phi(x_1)} : \phi(x_2) : \end{aligned}$$

And we know

$$\overline{\phi(x_1) \phi(x_2)} = \Delta_F(x_1 - x_2) = \overline{\phi(x_2) \phi(x_1)} \quad (18)$$

Therefore

$$\begin{aligned} T(\phi(x_1) \phi(x_2) \phi(x_3)) = & : \phi(x_1) \phi(x_2) \phi(x_3) : + \Delta_F(x_1 - x_2) : \phi(x_3) : + \Delta_F(x_2 - x_3) : \phi(x_1) : \\ & + \Delta_F(x_3 - x_1) : \phi(x_2) : \end{aligned} \quad (19)$$

Wick's theorem for four scalar product

$$\begin{aligned}
T(\phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4)) = & \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) : + \overline{\phi(x_1)\phi(x_2)} : \phi(x_3)\phi(x_4) : \\
& + \overline{\phi(x_2)\phi(x_3)} : \phi(x_1)\phi(x_4) : + \overline{\phi(x_3)\phi(x_4)} : \phi(x_2)\phi(x_1) : \\
& + \overline{\phi(x_4)\phi(x_1)} : \phi(x_2)\phi(x_3) : + \overline{\phi(x_1)\phi(x_3)} : \phi(x_2)\phi(x_4) : \\
& + \overline{\phi(x_2)\phi(x_4)} : \phi(x_1)\phi(x_3) : + \overline{\phi(x_1)\phi(x_2)} \overline{\phi(x_3)\phi(x_4)} \\
& + \overline{\phi(x_1)\phi(x_3)} \overline{\phi(x_2)\phi(x_4)} + \overline{\phi(x_1)\phi(x_4)} \overline{\phi(x_2)\phi(x_3)}
\end{aligned}$$

$$\begin{aligned}
\therefore T(\phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4)) = & \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) : + \Delta_F(x_1 - x_2) : \phi(x_3)\phi(x_4) : \\
& + \Delta_F(x_2 - x_3) : \phi(x_1)\phi(x_4) : + \Delta_F(x_3 - x_4) : \phi(x_1)\phi(x_2) : \\
& + \Delta_F(x_4 - x_1) : \phi(x_2)\phi(x_3) : + \Delta_F(x_1 - x_3) : \phi(x_2)\phi(x_4) : \\
& + \Delta_F(x_2 - x_4) : \phi(x_1)\phi(x_3) : + \Delta_F(x_1 - x_2)\Delta_F(x_3 - x_4) \\
& + \Delta_F(x_1 - x_3)\Delta_F(x_2 - x_4) + \Delta_F(x_1 - x_4)\Delta_F(x_2 - x_3) \\
& (20)
\end{aligned}$$