Measurement of Neutron Activation Cross Sections for Energies up to 40 MeV for Natural Samples of Mg, Si, Ca, V, Cr, Cu and Zn

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The cross section measurements of 17 neutron induced reactions using an intense semi-monoenergetic neutron field were recently reported. Neutrons for energies up to 40 MeV were produced by a simple Be target system bombarded by protons of 9 different energies between 20 and 40 MeV.

Successively in this neutron field, we measured the cross sections of 15 reactions of $^{\mathrm{nat}}\mathrm{Mg}(n,Xnp)^{24}\mathrm{Na}$, $^{\mathrm{nat}}\mathrm{Si}(n,Xn2p)^{27}\mathrm{Mg}$, $^{\mathrm{nat}}\mathrm{Si}(n,Xnp)^{28}\mathrm{Al}$, $^{\mathrm{nat}}\mathrm{Si}(n,Xnp)^{29}\mathrm{Al}$, $^{\mathrm{nat}}\mathrm{Ca}(n,Xnp)^{42}\mathrm{K}$, $^{\mathrm{nat}}\mathrm{Ca}(n,Xn2p)^{46}\mathrm{Sc}$, $^{\mathrm{nat}}\mathrm{V}(n,Xn2p)^{47}\mathrm{Sc}$, $^{\mathrm{nat}}\mathrm{Cr}(n,Xnp)^{52}\mathrm{V}$, $^{\mathrm{nat}}\mathrm{Cr}(n,Xnp)^{52}\mathrm{V}$, $^{\mathrm{nat}}\mathrm{Cr}(n,Xnp)^{52}\mathrm{V}$, $^{\mathrm{nat}}\mathrm{Cr}(n,Xnp)^{52}\mathrm{V}$, $^{\mathrm{nat}}\mathrm{Cr}(n,Xnp)^{52}\mathrm{V}$, $^{\mathrm{nat}}\mathrm{Cn}(n,Xn2p)^{52}\mathrm{Mc}$, $^{\mathrm{nat}}\mathrm{Cn}(n,Xn2p)^{52}\mathrm{Mc}$, $^{\mathrm{nat}}\mathrm{Cn}(n,Xnp)^{54}\mathrm{Cu}$, $^{\mathrm{nat}}\mathrm{Cn}(n,Xnp)^{56}\mathrm{Cu}$, $^{\mathrm{nat}}\mathrm{Cn}(n,Xnp)^{56}\mathrm{Cu}$, $^{\mathrm{nat}}\mathrm{Cn}(n,Xnp)^{58}\mathrm{Mc}$,

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I. Introduction

Activation cross sections for high energy neutrons are needed as basic data for neutron dosimetry, radiation safety and material damage studies, and for nuclear physics. The interest in the cross-section data of neutron energy above 20 MeV is now increasing and the related conference reports⁽¹⁾⁽²⁾ and books ⁽⁸⁾⁽⁴⁾ have recently been published. Those cross section data are, however, still very poor and no evaluated data file exists, since preparing a monoenergetic neutron source for activation experiment in this energy region presents some difficulties⁽⁵⁾.

The authors recently published a paper (6) on the cross section measurements of 17 reactions of 23 Na $(n, 2n)^{22}$ Na, 27 Al $(n, \alpha)^{24}$ Na, 51 V $(n, \alpha)^{48}$ Sc, 51 V $(n, p)^{51}$ Ti, 50 Cr $(n, 3n)^{48}$ Cr, 50 Cr $(n, 2n)^{49}$ Cr, 55 Mn $(n, p\alpha)^{51}$ Ti, 55 Mn $(n, 4n)^{52}$ Mn, 58 Mn $(n, 2n)^{54}$ Mn, 63 Cu $(n, 3n)^{61}$ Cu, 63 Cu $(n, 2n)^{62}$ Cu, 65 Cu $(n, p)^{65}$ Ni, 64 Zn $(n, t)^{62}$ Cu, 64 Zn $(n, 3n)^{62}$ Zn, 64 Zn $(n, 2n)^{63}$ Zn, 197 Au $(n, 4n)^{194}$ Au and 197 Au $(n, 2n)^{196}$ Au using an intense semimonoenergetic neutron field (7) which has been established at the SF Cyclotron of the Institute for Nuclear Study, University of Tokyo. Neutrons for energies up to 40 MeV were pro-

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duced by a simple Be target system bombarded by protons of 9 different energies between 20 and 40 MeV. Since the spectrum of this neutron field is not purely monoenergetic, that is, a high energy peak coming from the $\mathrm{Be}(p,n)$ reaction is broad and a low energy tail coming from the Be target and also from the coolant water exists, an unfolding technique is necessary for obtaining an excitation function from measured activation rates of irradiated samples^{(8)~(10)}.

Successively in this work, we measured the cross sections up to 40 MeV of 15 reactions of $^{\text{nat}}Mg(n, Xnp)^{24}Na, ^{\text{nat}}Si(n, Xn2p)^{27}Mg$, ^{nat}Si $(n, Xn p)^{28}$ Al, ^{nat}Si $(n, Xn p)^{29}$ Al, $(n, Xn p)^{42}K$, $^{nat}Ca(n, Xn p)^{43}K$, $^{nat}V(n, Xn2 p)$ ⁴⁶Sc, $^{\text{nat}}V(n, Xn2p)^{47}$ Sc, $^{\text{nat}}Cr(n, Xnp)^{52}V$, $^{\text{nat}}$ Cr $(n, Xnp)^{53}$ V, $^{\text{nat}}$ Cu $(n, Xn2p)^{62m}$ Co $(T_{1/2} =$ 14 min), $^{\text{nat}}$ Zn $(n, Xn2p)^{65}$ Ni, $^{\text{nat}}$ Zn $(n, Xnp)^{64}$ Cu, $^{\text{nat}}$ Zn $(n, Xn p)^{66}$ Cu and $^{\text{nat}}$ Zn $(n, Xn p)^{68}$ mCu $(T_{1/2}=3.8 \text{ min})$ by irradiating natural samples of Mg, Si, Ca, V, Cr, Cu and Zn in this neutron field. The cross-sections of these reactions were calculated by the ALICE/ LIVERMORE82 code(11), and they were adjusted to the measured reaction rates with the multi-parametric least-squares fitting (LSF). The LSF results were utilized as the initial guesses of two unfolding codes, SAND-II (12) and NEUPAC (13), which were used for the unfolding of the measured activation rates to excitation functions.

II. NEUTRON FIELD FOR ACTIVATION EXPERIMENT

Since a detailed description of the target system and of the neutron spectrum are given in a previous paper⁽⁷⁾, they are described briefly here. Proton beams of energies at every 2.5 MeV between 20 and 40 MeV were impinged to 1-mm-thick (E_p =20 to 37.5 MeV) and 2-mm-thick (E_p =40 MeV) Be targets, which were backed by a water coolant. Protons lose part of their energy at the Be target and totally stop in the water coolant.

The neutrons produced in the forward direction by this target system were measured with a 51-mm-dia.×51-mm-long NE-213 scintillation counter placed at about 1.1 m

from the target. Since the measured peak spectrum spread wider than the reality due to the energy resolution of the detection system, the measured spectrum of the peak energy region was replaced by the calculated spectrum of which peak area had been adjusted same as the measured one⁽⁷⁾. The neutron spectra are shown in Fig. 1 with the errors indicated in vertical bars. The digits in the figures indicate the incident proton energies (E_p) and average peak energies. Since the Be target is not thin, a high-energy peak due to the ${}^{9}\text{Be}(p, n)$ reaction is not sharp. A low-energy tail coming from the Be target and the water beam stopper is present.

III. MEASUREMENT OF INDUCED RADIOACTIVITY

The purity and the size of the irradiated samples are listed in **Table 1**. Since it is difficult to handle a metal Ca sample, a CaCO₃ sample was used instead. Oxygen and carbon nuclei in the sample produce radioactivities having much shorter half-lives than Ca, and this compound can be used without any problem. The other samples are all metals. The maximum size of the samples is 20 mm in diameter, and the flux depression of high-energy neutrons at the sample edge is estimated to be <3% in spite of their very sharp forward angular distribution⁽⁷⁾.

The activation samples except silicon were irradiated by turns for 1 h at 20 cm from the Be target. The Si sample was irradiated for 20 min. These irradiations were repeated for 9 proton energies.

The irradiated samples were placed on a Ge detector for γ -ray spectrometry. Each sample was measured several times with suitable intervals. The obtained pulse-height spectra were analyzed by the KEI-11EF program (14) to get photo-peak areas. The peak detection efficiencies were calculated by a Monte Carlo code, PEAK (15), which can evaluate the self-absorption of γ -rays in a source disk. The accuracy of the PEAK code had been experimentally checked for the Ge de-

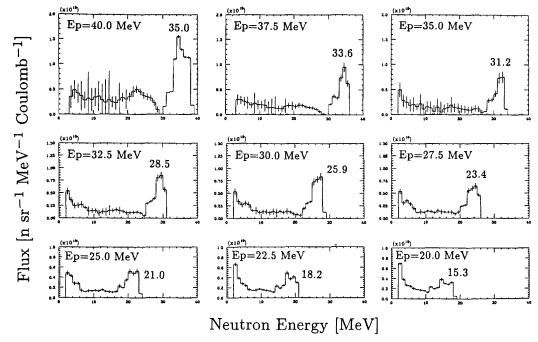


Fig. 1 Corrected neutron spectra of the irradiation fields

Table 1 Specification of the irradiated samples

Nuclide	Chemical form	Purity (%)	Size (mm)
Mg	Metal	99.99	$20.0 \mathrm{dia.} \times 2.0$
Si	Metal	99.999	17.0 dia.×2.0
Ca	CaCO ₃	99.999	$20.0\mathrm{dia.}\!\times\!2.0$
V	Metal	99.7	$5.0\times5.0\times5.0$
Cr	Metal	99.99	$5.0 \times 5.0 \times 5.0$
Cu	Metal	99.99	$20.0\mathrm{dia.}\!\times\!1.0$
Zn	Metal	99.9999	15.0 dia.×2.1

tector that was used and found to be of $<10\,\%$ error. The coincidence-summing corrections were performed by the SUMECC code (16). Total detection efficiencies, which are necessary for the SUMECC code, were calculated by the PEAK code. The data of decay schemes were derived from Ref. (17).

IV. CALCULATION OF CROSS SECTIONS

Since the neutron fields used in this activation experiment are not purely monoenergetic, excitation functions cannot be obtained directly, but can be obtained by using the unfolding method and the LSF method using the activation rates at 9 different semimono-

energetic neutron fields.

In LSF, considering the production of the relevant radionuclide through multi-reaction channels, initial values calculated by the ALICE/LIVERMORE82 code were adjusted to the measured reaction rates as shown in the following:

$$R = \sum_{i} \left[A_{i} - \sum_{f} \int_{0}^{E_{\text{max}}} \sigma_{f}(E) \Phi_{i}(E) dE \right]^{2}, \quad (1)$$

$$\sigma_{j}(E) = k_{j} g_{j}(E), \qquad (2)$$

where A_i : Activation rate at i'th neutron field

 $g_j(E)$: Calculated excitation function by ALICE for j'th

reaction channel $\Phi_i(E)$: Neutron flux shown in Fig. 1.

By determining the k_j value to make R minimum, the excitation function $\Sigma_j \sigma_j(E)$

In the use of the revised nuclear evaporation code, ALICE/LIVERMORE82⁽¹¹⁾, the geometry-dependent hybrid model was selected for precompound decay, since this model generated the most consistent results with the experiment in the test calculations. The

was obtained.

default values recommended by the authors of the code were chosen most often in the calculation.

The measured reaction rates were converted into excitation functions by the unfolding codes, SAND-II (12) and NEUPAC (13). The concept of the former code is the iterative perturbation method, and that of the latter is the variational principle, the so-called "J-1 method". These unfolding codes are usually used for neutron spectrum measurements by foil activations where the excitation functions of used foils are known. In contrast, in this work the neutron spectra were used as known input data, and unknown excitation functions were obtained.

These codes require an initial guess as input data, and it was obtained from the result of the LSF calculation and some empirical reference data as shown in **Table 2**. The initial guess of the $^{\rm nat}{\rm Mg}(n,\,Xn\,p)^{\rm 24}{\rm Na}$ reaction, for example, was obtained from the IAEA curve of $^{\rm 24}{\rm Mg}(n,\,p)$ cross section and the LSF result which were smoothly connected at 15 MeV.

The NEUPAC code evaluates the propagation of errors of initial guess values, neutron spectra, and reaction rates to the resultant excitation functions. Errors of the neutron spectra are about 10% in the peak region and about 30 to 50% in the low-energy tail as shown in Fig. 1, and errors of the activation rates are about 10%. The errors of the initial guess values are unknown. They were, however, set at 60%, since this value seemed to be adequate, that is, the results of the NEUPAC code adopting this value seemed to be reasonable with our experimental data and also with the empirical reference data. The errors of the initial guess values are most dominant on the propagated errors. Unfolded results are almost bound within the error band of the initial guess, and the fluctuations of the SAND-II results are usually larger than those of the NEUPAC results.

V. RESULTS AND DISCUSSION

Fifteen reactions of which cross sections were measured are listed in Table 2. These

(n, XnY p)-type reactions are the mixtures of several competing reactions which are also shown in Table 2, since natural samples used in this experiment contain more than one isotope.

The obtained cross section curves are shown in Figs. $2(a)\sim(0)$, where the results obtained by the unfolding processes performed by the SAND-II and the NEUPAC codes are shown in a thin zigzag solid line and a thin zigzag dotted line, respectively. The vertical bars attached on the NEUPAC results show errors estimated by this code. Results of the ALICE/LIVERMORE82 code were adjusted by the LSF and they are shown in a thin dashed line. If some reference data are available, they are also drawn in the same figure in a thick dashed line for ENDF-B/V (18), in a thick dash-dot line for JENDL-3(19), in a thick fine-dotted line for IAEA Tech. Rep. 273(3) and in circles for experimental values compiled by McLane, Dunford & Rose⁽⁴⁾, which is the fourth edition of the BNL-325 and is here referred to as BNL-325. The cross section data calculated with the GNASH code by Yamamuro (20) for the natZn $(n, Xn p)^{64}$ Cu and the ^{nat}Zn $(n, Xn p)^{66}$ Cu reactions are also shown in a thin dash-dot line. The results and discussions are given separately for each reaction in the following.

(1) $^{\text{nat}}Mg(n, Xnp)^{24}Na$ (**Fig. 2**(a))

This reaction can be considered to be a mixture of ${}^{24}\text{Mg}(n, p)$, ${}^{25}\text{Mg}(n, d)$ and ${}^{26}\text{Mg}$ (n, t) as shown in Table 2. In this article the (n, t) reaction, for example, means a mixture of the (n, t), (n, nd) and (n, 2np)reactions. The main channel of this reaction is ${}^{24}Mg(n, p)$ due to a large abundance of ²⁴Mg (78.99%), and a lower threshold energy (4.93 MeV) and a generally higher cross section value of the (n, p) reaction than those of other reactions (21). The SAND-II and NEUPAC results show good agreement within errors, except an unreasonable oscillation in the SAND-II result above 30 MeV. The LSF result also shows good agreement with these two unfolded results, but its rising up region below 10 MeV is shifted to higher energy than the others, reflecting the ALICE result.

Table 2 Observed productions, competing reactions, threshold energies, natural abundances, sources of initial guesses for unfolding process and recommendations for best result

Production	Reaction	Threshold energy (MeV)	Natura abundan (%)		Source of unfolding initial guess	Recommendation for best result
Mg→ ²⁴ Na	24 Mg (n, p) * 25 Mg (n, d) 26 Mg (n, t)	4. 93 12. 54 15. 24	²⁴ Mg 78. ²⁵ Mg 10. ²⁶ Mg 11.	. 00	<15 MeV IAEA of ²⁴ Mg(n, p) >15 MeV LSF	<10 MeV NEUPAC >10 MeV LSF
Si→ ²⁷ Mg	$^{28}\text{Si}(n,2p)^{\dagger}$ $^{29}\text{Si}(n,^{3}\text{He})^{\dagger}$ $^{30}\text{Si}(n,\alpha)^{*}$	13.89 14.66 4.34	²⁸ Si 92. ²⁹ Si 4.		LSF	NEUPAC
Si→ ²⁸ Al	$^{28}Si(n, p)*$ $^{29}Si(n, d)$ $^{30}Si(n, t)$	4.00 10.46 14.94			<16 MeV IAEA of ²⁸ Si(n, p) >16 MeV LSF	NEUPAC
Si→ ²⁹ Al	$^{29}Si(n, p)*$ $^{30}Si(n, d)^{\dagger}$	3.00 11.66			LSF	SAND- II
Ca→ ⁴² K	42 Ca (n, p) * 43 Ca (n, d) † 44 Ca (n, t) † 46 Ca $(n, 2nt)$ 48 Ca $(n, 4nt)$	2.80 8.64 13.62 31.81 49.36	⁴³ Ca 0. ⁴⁴ Ca 2. ⁴⁶ Ca 0.	94 647 135 09 004 187	< 12 MeV IAEA of 42 Ca (n, p) > 12 MeV LSF	NEUPAC
Ca→ ⁴³ K	43 Ca (n, p) * 44 Ca (n, d) † 46 Ca (n, nt) 48 Ca $(n, 3nt)$	1.06 10.17 21.97 39.53			LSF	see text
V→ ⁴⁶ Sc	$^{50}V(n,n\alpha)^{\dagger}$ $^{51}V(n,2n\alpha)^{*}$	10. 09 21. 35	⁵⁰ V 0. ⁵¹ V 99.	. 25 75	LSF	LSF
V→⁴¹Sc	$^{50}\mathrm{V}(n,\alpha)^{\dagger}$ $^{51}\mathrm{V}(n,n\alpha)^{*}$	$\begin{array}{c} 0 \ (Q = 0.76) \\ 10.50 \end{array}$	`		LSF	LSF
Cr→ ⁵² V	$52 \text{Cr}(n, p)^*$ 53 Cr(n, d) 54 Cr(n, t)	3. 26 9. 08 12. 60	⁵² Cr 83. ⁵³ Cr 9.	. 35 . 79 . 50 . 36	<15 MeV IAEA of $^{52}\text{Cr}(n, p)$ >15 MeV LSF	<15 MeV SAND-II >15 MeV LSF
Cr→ ⁵³ V	53 Cr (n, p) * 54 Cr (n, d) †	2.69 10.32	0. 2.		<15 MeV BNL eye-guide of ⁵³ Cr(n, p) >15 MeV LSF	NEUPAC
$Cu \rightarrow 62mCo$ $T_{1/2} = 14 min$	63 Cu $(n, 2p)$ 65 Cu (n, α) *	10. 85 0. 21	⁶³ Cu 69. ⁶⁵ Cu 30.		<20 MeV IAEA of 66 Cu (n, α) >20 MeV LSF	NEUPAC
Zn→ ⁶⁵ Ni	66 Zn $(n, 2p)^{\dagger}$ 67 Zn $(n, ^{3}$ He $)$ 68 Zn $(n, \alpha)^{*}$ 70 Zn $(n, 2n\alpha)$	10. 44 9. 76 0 (Q=0. 77) 15. 15	⁶⁸ Zn 18.	9 10	LSF	<20 MeV SAND-II >20 MeV NEUPAC
Zn→ ⁶⁴ Cu	64 Zn (n, p) * 66 Zn (n, t) † 67 Zn (n, nt) 68 Zn $(n, 2nt)$ 70 Zn $(n, 4nt)$	0 (Q=0.20) 10.51 17.67 28.01 43.92	2.1 0.	-	$<17 \mathrm{MeV}$ IAEA of $^{64}\mathrm{Zn}(n,p)$ $>17 \mathrm{MeV}$ LSF	NEUPAC
Zn→ ⁶⁶ Cu	66 Zn $(n, p)*$ 67 Zn (n, d) 68 Zn $(n, t)^{\dagger}$ 70 Zn $(n, 2nt)$	1.89 6.79 10.78 26.70			LSF	NEUPAC
$Z_{n\rightarrow 68m}Cu$ $T_{1/2}=3.8 \text{ min}$	$^{68}Zn(n,p)^*$ $^{70}Zn(n,t)$	4. 63 11. 95			LSF	NEUPAC

^{*} Dominant reaction; † Supplemental observed reaction.

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In the peak region between 10 and 17 MeV, the LSF result gives the highest value and the SAND-II result gives the lowest. Our results are compared here with the 24 Mg(n, p) cross section values of IAEA, ENDF-B/V, JENDL-3 and BNL-325 which are multiplied by the natural abundance of 24 Mg. This comparison reveals that the NEUPAC result agrees well with these data in the rising up region below 10 MeV, and above 10 MeV the LSF result is the closest to these data. It is confirmed that the 24 Mg(n, p) reaction is dominant and contributions of other reactions are negligibly small.

(2) $^{\text{nat}}\text{Si}(n, Xn2p)^{27}\text{Mg}$ (Fig. 2(b))

The LSF was performed with the results of the ALICE calculations for the 28 Si(n, 2p), 29 Si $(n, ^3$ He) and 30 Si (n, α) reactions. The SAND-II and NEUPAC results agree well each other within unfolded errors, except a strong decrease of the SAND-II result above 35 MeV. The LSF result is slightly larger below 20 MeV and smaller above 20 MeV than the two unfolded results. The excitation

functions have a peak at about 15 MeV which corresponds to the $^{30}Si(n,\alpha)$ reaction and a swelling above 20 MeV. Our SAND-II and NEUPAC results are close to the IAEA data and lie in the middle of the scattered data of BNL-325. The swelling may be due to the $^{28}Si(n,2p)$ reaction.

(3) $^{\text{nat}}\text{Si}(n, Xn p)^{28}\text{Al}$ (**Fig. 2**(c))

This reaction was considered to consist of ${}^{28}\mathrm{Si}(n,\,p)$, ${}^{29}\mathrm{Si}(n,\,d)$ and ${}^{30}\mathrm{Si}(n,\,t)$. The ${}^{28}\mathrm{Si}(n,\,p)$ reaction is dominant because of the large abundance of ${}^{28}\mathrm{Si}$ (92.23%), the lower threshold energy (4.0 MeV), and generally higher cross section value of $(n,\,p)$ reaction than the others. The SAND-II and NEUPAC results agreed within the unfolded errors, and they are also close to the LSF result, except its rising up region. The IAEA curve gives about 60% higher value at the peak than our unfolded results which agree with the lowest data of BNL-325.

(4) $^{\text{nat}}\text{Si}(n, Xnp)^{29}\text{Al}$ (Fig. 2(d))

This reaction was considered to be a mixture of the $^{29}Si(n, p)$ and the $^{30}Si(n, d)$ reaction

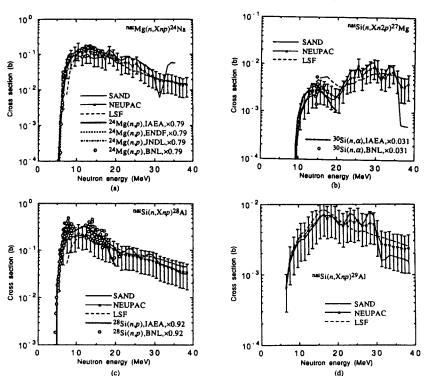


Fig. $2(a) \sim (d)$ Obtained excitation function curves

tions. The LSF, the SAND-II and the NEUPAC results are all in good agreement within unfolded errors. The excitation function has two peaks; a big peak around 16 MeV corresponds to the $^{29}\mathrm{Si}(n,\,p)$ reaction and a small peak around 25 MeV to the $^{30}\mathrm{Si}(n,\,d)$ reaction. No available reference data exist for this reaction.

(5) $^{\text{nat}}\text{Ca}(n, Xn p)^{42}\text{K}$ (**Fig.** 2(e))

The LSF was performed with the ALICE results of the $^{42}\text{Ca}(n, p)$, $^{43}\text{Ca}(n, d)$ and $^{44}\text{Ca}(n, t)$ reactions. Although the two unfolded results show oscillation, especially the SAND-II result, which may be caused by a statistical error of the measured activation rates, these three results are close each other, except a shift of LSF to higher energy above the threshold energy. The excitation functions have three peaks; the isolated lowest-energy peak around 15 MeV corresponds to the $^{42}\text{Ca}(n, p)$ reaction having a low threshold energy of 2.80 MeV. Our results are in good agreement below 15 MeV with the IAEA and the BNL-325 data of the $^{42}\text{Ca}(n, p)$ cross

section.

(6) $^{\text{nat}}\text{Ca}(n, Xn p)^{43}\text{K}$ (**Fig. 2**(f))

The LSF result is larger below 15 MeV and smaller between 15 and 35 MeV than the SAND-II and the NEUPAC results which are close each other. Only at 15 MeV, our result is compared with the experimental data of the 43 Ca(n, p) and the 44 Ca(n, d) reactions by multiplying the natural abundances of 43 Ca and 44 Ca. The main component of the obtained excitation function is the 43 Ca(n, p) reaction cross section. The BNL-325 data of 43 Ca(n, p) are largely deviated, and it is therefore difficult to decide which cross section result is the best among our calculations.

(7) $^{\text{nat}}V(n, Xn2p)^{46}Sc$ (Fig. 2(g))

This reaction can be considered to be a sum of ${}^{50}\mathrm{V}(n,\,n\alpha)$ and ${}^{51}\mathrm{V}(n,\,2n\alpha)$. The abundance of ${}^{50}\mathrm{V}$ is very small (0.25%), and the latter reaction having threshold energy of 21.35 MeV may be dominant at higher energies. The three results of LSF, SAND-II and NEUPAC are very close altogether. Since

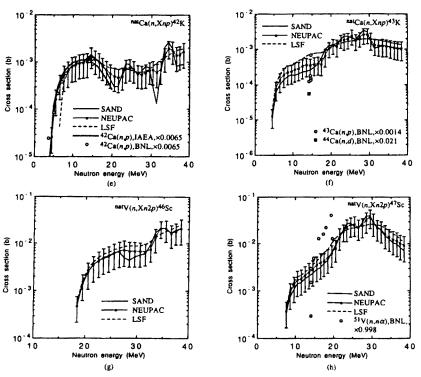


Fig. 2(e)~(h) Obtained excitation function curves

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no experimental data exist, our cross section data are unique.

(8) $^{\text{nat}}V(n, Xn2p)^{47}Sc$ (Fig. 2(h))

This reaction is dominated at higher energies by the 51 V(n, $n\alpha$) reaction of which threshold energy is 10.50 MeV due to the small natural abundance of 50 V. The three results of LSF, SAND-II and NEUPAC are close altogether, but the LSF result gives the smoothest curve without small twin peaks between 20 and 30 MeV which are seen in the SAND-II and the NEUPAC results. Our results are compared with the BNL-325 data of the 51 V(n, $n\alpha$) reaction multiplied by the natural abundance of 51 V. The four cross section values at the neutron energies between 16.0 and 19.6 MeV given by Bormann et al. $^{(22)}$ are 3 to 5 times larger than our measurement.

(9) $^{\text{nat}}Cr(n, Xnp)^{52}V$ (Fig. 2(i))

Since 52 Cr has a large abundance of 83.79% and its (n, p) reaction has a low threshold energy of 3.26 MeV, this production is dominated by the 52 Cr(n, p) reaction. Excluding a rising up region of the LSF result, three

results of LSF, SAND-II and NEUPAC show good agreement within the errors, although some oscillation can be seen in the SAND-II and NEUPAC results. Our results agree with the IAEA data of 52 Cr(n, p). In the energy region between 15 and 20 MeV, the LSF result is close to the higher values of BNL-325 data, while on the other hand, the SAND-II and NEUPAC results are close to the lower values of BNL-325 data. Contribution of the 53 Cr(n, d) reaction is negligibly small as seen in Fig. 2(i).

(10) $^{\text{nat}}Cr(n, Xnp)^{53}V$ (**Fig. 2**(j))

This reaction consists of two reactions of 53 Cr(n, p) and 54 Cr(n, d). Excluding a rising up region of the LSF result, our three results show good agreement within the errors. The unfolded results can be divided into a peak and a bump at a boundary of about 20 MeV. The peak around 15 MeV indicates good agreement with the BNL-325 data of 53 Cr(n, p) cross section multiplied by the natural abundance of 53 Cr. The small bump above 20 MeV might be a contribution of the 54 Cr

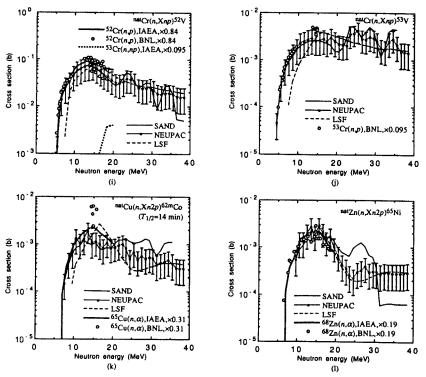


Fig. $2(i) \sim (1)$ Obtained excitation function curves

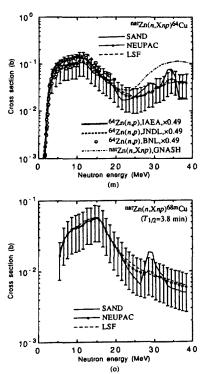
(n, d) reaction.

(11) $^{\text{nat}}\text{Cu}(n, Xn2p)^{\text{62m}}\text{Co} (T_{1/2} = 14 \text{ min})$ (**Fig.** 2(k))

This reaction consists of two reactions of 63 Cu(n, 2p) and 65 Cu (n, α) . The results given by LSF, SAND-II and NEUPAC are slightly different each other. In the energy region below 15 MeV, the NEUPAC result shows good agreement with the IAEA data of 65 Cu (n, α) and we therefore consider that the NEUPAC result may be most appropriate.

(12) $^{\text{nat}}$ Zn $(n, Xn2 p)^{65}$ Ni (Fig. 2(1))

This reaction consists of three reactions of 66 Zn(n, 2p), 67 Zn(n, 3 He) and 68 Zn(n, α). The three results of LSF, SAND-II and NEUPAC agree within the error below 20 MeV, but they show large discrepancy above 20 MeV. At a large peak below 20 MeV, the SAND-II result agrees well with the IAEA curve and the BNL-325 data of the 68 Zn(n, α) cross section. The second small bump at about 25 MeV which are seen in SAND-II and NEUPAC might be due to the 66 Zn(n, 2p) reaction.



(13) $^{\text{nat}}$ Zn(n, Xnp) 64 Cu (**Fig.** 2(m))

This reaction almost consists of two reactions of 64 Zn(n, p) and 66 Zn(n, t). The LSF result is higher than the unfolded results below 15 MeV and is very close to them above 15 MeV. The BNL-325 data of 64 Zn(n, p) are classified into two groups in the peak region between 5 and 15 MeV, the higher and the lower groups. The lower group agrees with our SAND-II and NEUPAC results, and the evaluated values of JENDL-3 and IAEA. The first peak below 20 MeV is certainly due to the 64 Zn(n, p) reaction from this comparison and the second peak above 20 MeV may be due to the 66 Zn(n, t) reaction. The natZn $(n, Xnp)^{64}$ Cu cross sections was calculated with the GNASH code by Yamamuro. His calculated result in Fig. 2(m) agrees with our results below 20 MeV, but gives about a factor of 3 higher values above 30 MeV.

(14) $^{\text{nat}}$ Zn(n, Xnp) 66 Cu (Fig. 2(n))

This reaction consists of three reactions of ${}^{66}\text{Zn}(n, p)$, ${}^{67}\text{Zn}(n, d)$ and ${}^{68}\text{Zn}(n, t)$. Our three results show good agreement within

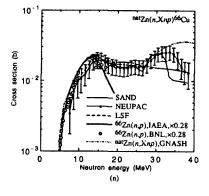


Fig. $2(m)\sim(0)$ Obtained excitation function curves

the error below $28\,\text{MeV}$. The peak below $20\,\text{MeV}$ corresponds to $^{66}\text{Zn}(n,\,p)$ and our results indicate there a good agreement with the BNL-325 data. The IAEA curve is slightly higher below $12\,\text{MeV}$. The peak in the unfolded results around $33\,\text{MeV}$ may be due to the $^{67}\text{Zn}(n,\,d)$ and the $^{68}\text{Zn}(n,\,t)$ reactions. The GNASH result by Yamamuro gives good agreement with our results below $33\,\text{MeV}$.

(15)
$$^{\text{nat}}Zn(n, Xnp)^{68m}Cu \ (T_{1/2}=3.8 \text{ min})$$

(**Fig.** 2(o))

This reaction consists of two reactions of 68 Zn(n, p) and 70 Zn(n, t), but the former is the dominant reaction due to larger abundance of 68 Zn (18.8%) and low threshold energy (4.63 MeV) compared to the latter. Our three results are close each other, except a big bump of SAND-II around 30 MeV which may come from the statistical error of the measured reaction rate. No reference data exist for this reaction and our results are the first experimental data.

Three excitation functions were obtained for every reaction by using the LSF method, and the SAND-II and the NEUPAC unfolding codes. It depends on reactions what method seems to give the most reasonable result, and recommended methods giving the best results are listed in Table 2.

The numerical data will be sent on request.

VI. SUMMARY

The excitation functions of 15 activation reactions were obtained for neutron energies up to 40 MeV by unfolding the activation rates measured at 9 different semimonoenergetic neutron fields, succeeding to our previous paper which gave 17 activation reaction cross sections. Although the estimated error of our results is not small for many reactions, we gave useful experimental cross-section data for these activation reactions above 20 MeV. Our results are the first data for three activation reactions.

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