

# 1 Poisson equation in one dimension

In this section, we will solve the one-dimensional Poisson equation

$$u_{xx} = f(x) \quad (0 < x < 1)$$

subject to a source term  $f(x)$  and different boundary conditions at  $x = 0$  and  $x = 1$ . First, we will solve it with finite difference methods of first and second order on a uniform grid. Finally, we solve it on a non-uniform grid and investigate how adaptive mesh refinement (AMR) can be used to obtain accurate solutions by distributing fewer points more cleverly along the grid. a

First, consider Dirichlet and Neumann boundary conditions at opposite ends and the source given by

$$u_{xx} = x + \cos(2\pi x) \quad (0 < x < 1), \quad u(0) = a, \quad u_x(1) = b.$$

The analytical solution is

$$u(x) = C_1 + C_2x + \frac{1}{6}x^3 - \frac{1}{4\pi^2} \cos(2\pi x),$$

where the constants  $C_1$  and  $C_2$  are determined from the boundary conditions. To solve the equation numerically, we impose a uniform grid of  $M + 2$  points and step length  $h$  defined by

$$\begin{array}{ccccccc} x_0 = 0 & x_1 & x_2 & & x_m & & x_{M-1} & x_M & x_{M+1} = 1 \\ & \bullet & \bullet & & \bullet & & \bullet & \bullet & \bullet \\ & h & h & & & & h & h & \end{array}$$

To generate finite difference methods of both first and second order, we approximate the second derivative at interior points using the forward difference and central difference

$$\begin{aligned} u_{xx}(x_m) &= \frac{u_m - 2u_{m+1} + u_{m+2}}{h^2} + O(h^1) & (1 \leq m \leq M) \\ u_{xx}(x_m) &= \frac{u_{m-1} - 2u_m + u_{m+1}}{h^2} + O(h^2) & (1 \leq m \leq M-1). \end{aligned}$$

To handle the Dirichlet boundary condition  $u(0) = a$  at the left edge, we insert the trivial equation

$$1 \cdot u_0 = a.$$

To handle the Neumann boundary condition  $u_x(1) = b$  at the right edge to first or second order, we use

$$\begin{aligned} u_x(1) &= \frac{u_{M+1} - u_M}{h} + O(h^1) = b \\ u_x(1) &= \frac{\frac{1}{2}u_{M-1} - 2u_M + \frac{3}{2}u_{M+1}}{h} + O(h^2) = b. \end{aligned}$$

By writing all these equations in  $(M + 2) \times (M + 2)$ -matrix form  $AU = b$ , we obtain for example to second order

$$\begin{bmatrix} 1 & & & & \\ +1/h^2 & -2/h^2 & +1/h^2 & & \\ & \ddots & \ddots & \ddots & \\ & & +1/h^2 & -2/h^2 & +1/h^2 \\ & & +1/2h & -2/h & +3/2h \end{bmatrix} \begin{bmatrix} U_0 \\ U_1 \\ \vdots \\ U_M \\ U_{M+1} \end{bmatrix} = \begin{bmatrix} a \\ f(x_1) \\ \vdots \\ f(x_M) \\ b \end{bmatrix}$$

Some remarks:

- We could handle the Dirichlet boundary condition  $u(0) = a$  differently by treating  $U_0 = a$  as a known variable. The system of equations is equivalent if we remove the first row and column of  $A$  and the first entries in  $U$  and  $b$ , but simultaneously modify the entry  $f(x_1) \rightarrow f(x_1) - a/h^2$ . This approach is more consistent with treating  $U_0$  as a known variable, since its precise value is defined by the Dirichlet boundary condition. However, our approach of inserting a trivial equation  $1 \cdot U_0 = a$  keeps the matrix dimensions independent of boundary conditions and makes it easier to reason with how the discretized differential operator represented by  $A$  operates on the grid point  $U_0$  in the same way it operates on all other grid points.
- To handle different combinations of Dirichlet and Neumann boundary conditions at the ends, we simply replace the first or last rows of the matrix with the same type of equation. Note that if the Neumann boundary condition is imposed at the left boundary, the last row of the matrix above would have to be both reversed and negated.
- When Neumann boundary conditions are imposed at both ends, the solution is determined only up to a constant. To see this for a general Poisson boundary value problem, note that if  $u_{xx} = f(x)$ ,  $u_x(0) = a$  and  $u_x(1) = b$ , then also  $(u+C)_{xx} = f(x)$ ,  $(u+C)_x(0) = a$  and  $(u+C)_x(1) = b$  if  $C$  is only a constant. It can also be seen from the general solution for this particular source term that the constant  $C_1$  is undetermined when the solution is subject to boundary conditions that involve derivatives only. In this case, an additional constraint like  $u(0) = 0$  must be imposed to define a unique solution.

With this in mind, it is now straightforward to solve the Poisson equation subject to any combination of Dirichlet and Neumann boundary conditions at the ends to both first and second order.

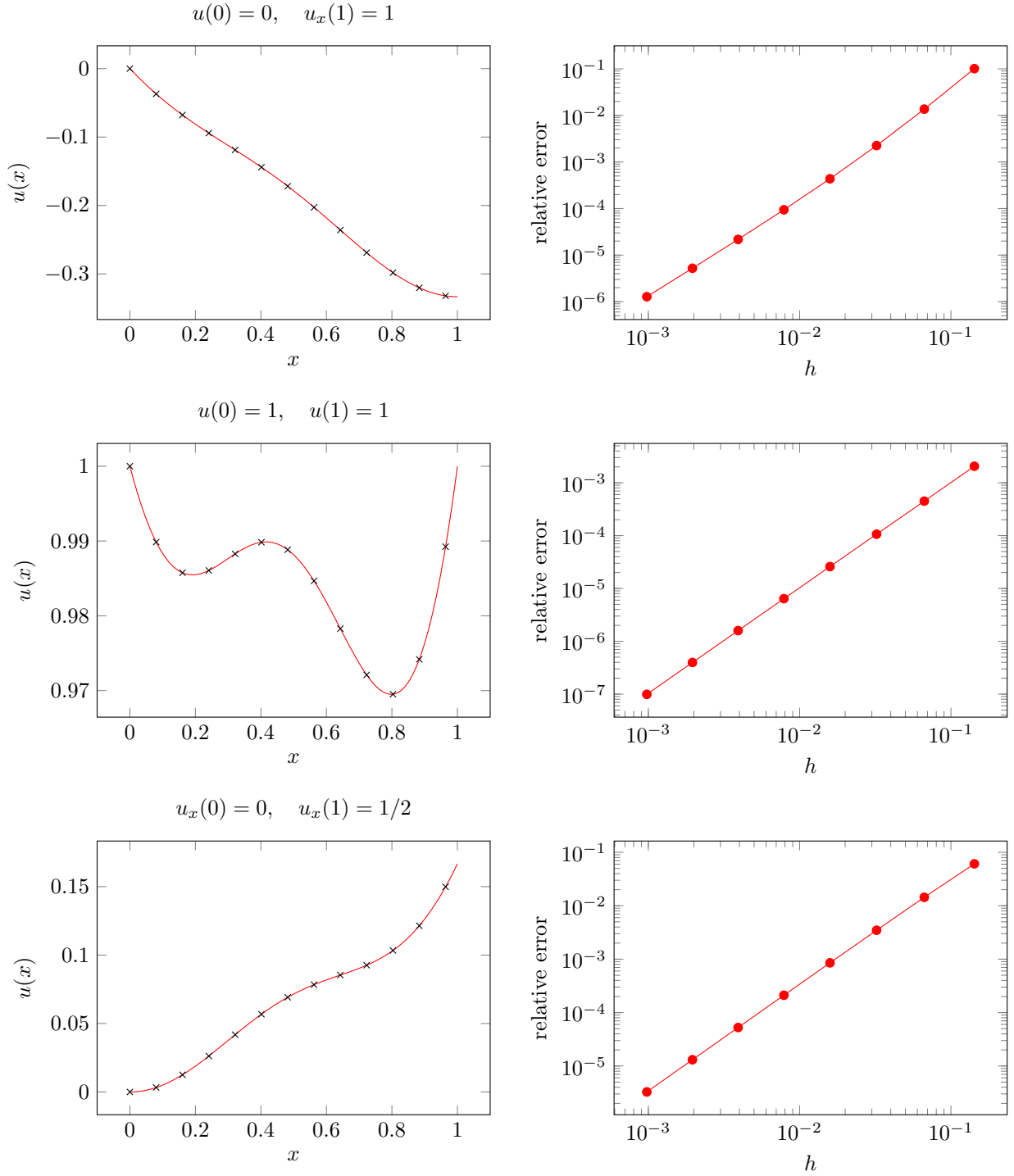


Figure 1: Analytical and numerical solutions (left) and convergence plots (right) for solutions to the Poisson equation subject to three different boundary conditions.



Figure 2: Nice fig

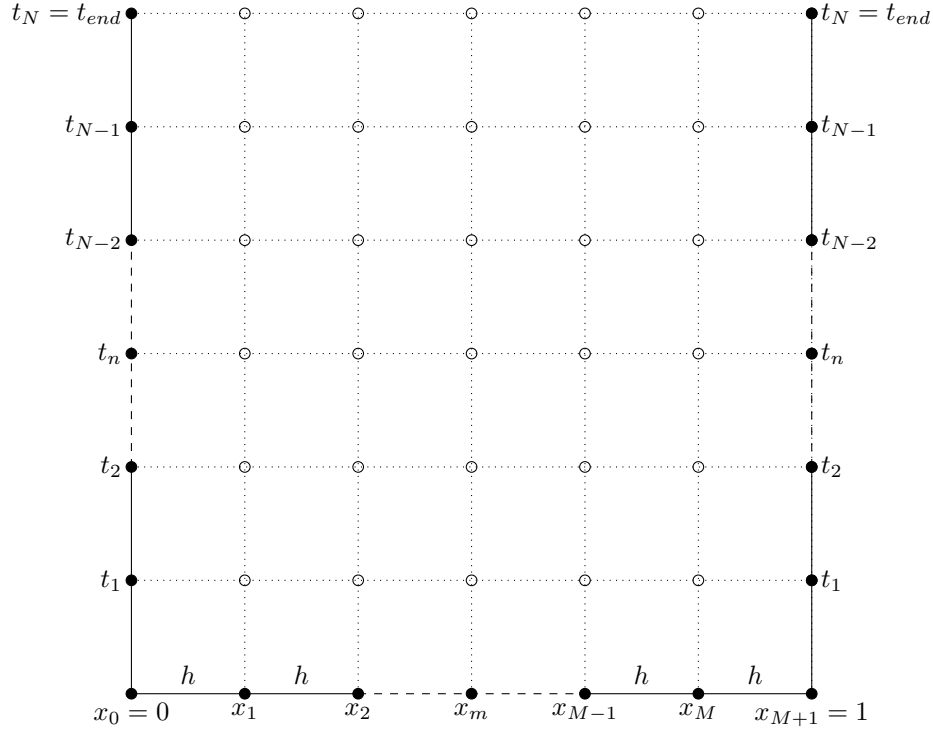
## 2 Heat equation in one dimension

In this section, we will solve one-dimensional differential equations with time evolution. Firstly, we will solve the one-dimensional heat equation with Neumann boundary conditions, using both the Backward Euler method and Crank-Nicolson's method. We will then analyse the error of our methods using a manufactured solution to which the analytical solution is known. For this, we will use both uniform and adaptive mesh refinement, and the different methods will be discussed. Finally, we will solve the inviscid Burgers' equation and evaluate the breaking which this equation exhibits.

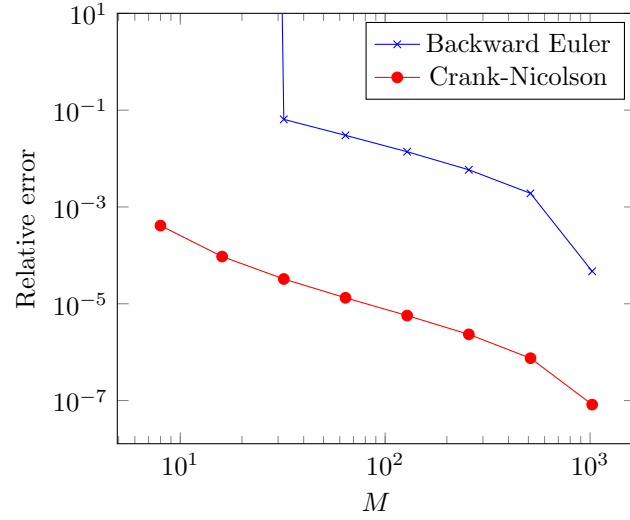
Firstly, we consider the one-dimensional heat equation with Neumann boundary conditions

$$u_t = u_{xx}, \quad u_x(0, t) = u_x(1, t) = 0, \quad u(x, 0) = 2\pi x - \sin(2\pi x),$$

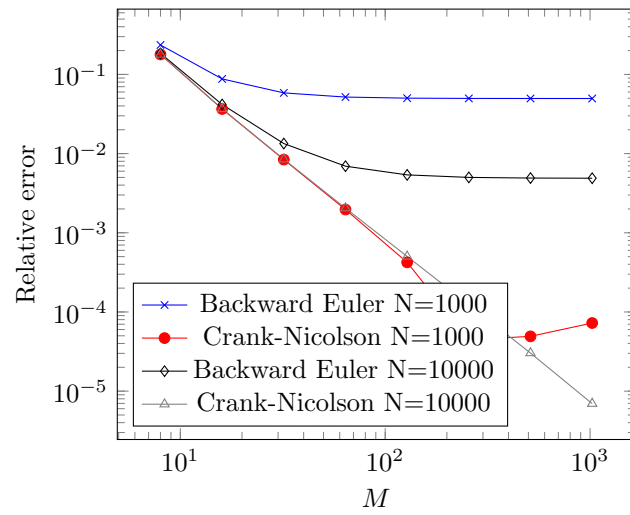
valid for  $x \in [0, 1] := \Omega$  and  $t > 0$ . The equation can be solved numerically by dividing the interval  $\Omega$  in  $M$  parts, as described in section 1. We also divide the time interval  $t \in [0, t_{end}]$  into parts, so that we approximate the solution at  $N$  finite times  $t_n = nk$ ,  $n = 1, \dots, N$ . For  $t = 0$ , the solution is the given initial condition.



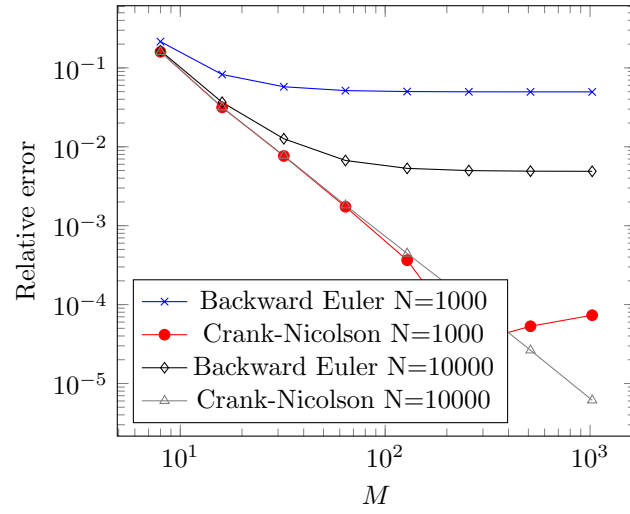
2a Convergence plot L2 discrete rel. error



2b Convergence plot L2 discrete rel err



2b Convergence plot l2 continous rel. err.



2b AMR Convergence plot L2 discrete rel err

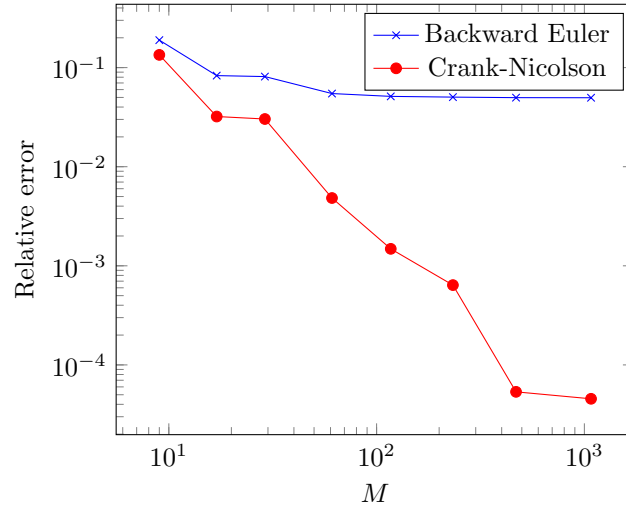


Figure 3: AMR

2b AMR Convergence plot l2 continous rel. err.

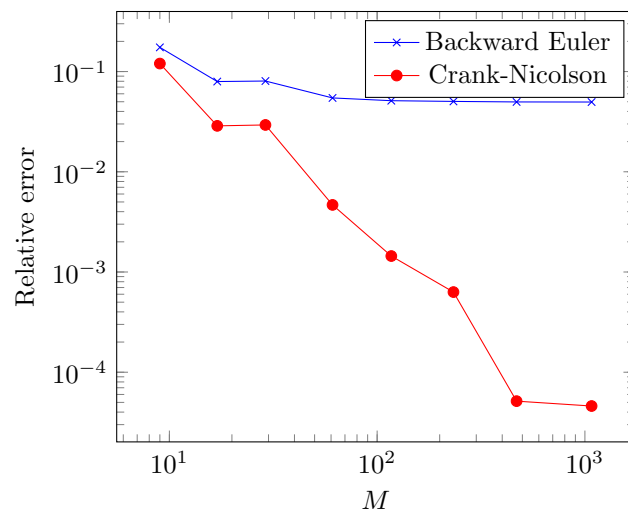


Figure 4: AMR



### 3 Laplace equation in two dimensions

## 4 Linearized Korteweg-De Vries equation in one dimension

In this section, we will study the one-dimensional linearized Korteweg-De Vries equation

$$\frac{\partial u}{\partial t} + \left(1 + \pi^2\right) \frac{\partial u}{\partial x} + \frac{\partial^3 u}{\partial x^3} = 0 \quad (t \geq 0) \quad (-L/2 \leq x \leq +L/2), \quad (1)$$

where the solution  $u = u(x, t)$  is subject to periodic boundary conditions

$$u(x + L, t) = u(x, t).$$

### 4.1 Analytical solution

As the solution is periodic in space at every time  $t$ , it can be expressed as a Fourier-series

$$u(x, t) = \sum_{n=-\infty}^{+\infty} c_n(t) \exp(ik_n x) \quad (2)$$

with wavenumbers  $k_n = 2\pi n/L$  and time-dependent coefficients  $c_n(t)$  that ensure spatial periodicity at all times. This can be derived formally by separation of variables. Inserting the Fourier series into equation (1) gives the condition

$$\sum_n \left( \dot{c}_n(t) + i \left( \left(1 + \pi^2\right) k_n - k_n^3 \right) c_n(t) \right) \exp(ik_n x) = 0$$

on the coefficients. Due to orthogonality of the Fourier basis functions  $\exp(ik_n x)$ , the sum can vanish only if all prefactors vanish separately. This gives a first-order differential equation for each coefficient with the solution

$$c_n(t) = c_n(0) \exp \left( -i \left( \left(1 + \pi^2\right) k_n - k_n^3 \right) t \right). \quad (3)$$

### 4.2 Numerical solution method

To find a numerical solution  $U_m^n = U(x_m, t_n) \approx u(x_m, t_n) = u_m^n$  of the Korteweg-De Vries equation, we will discretize it with central differences in space and integrate over time with the Euler method and the Crank-Nicholson method. For the first order spatial derivative, we use the central difference

$$\frac{\partial u_m^n}{\partial x} \approx \frac{\delta u_m^n}{2\Delta x} = \frac{u_{m+1}^n - u_{m-1}^n}{2\Delta x}.$$

We repeat the same finite difference three times to approximate the third order spatial derivative as

$$\frac{\partial^3 u_m^n}{\partial x^3} \approx \frac{\delta^3 u_m^n}{(2\Delta x)^3} = \frac{u_{m+3}^n - 3u_{m+1}^n + 3u_{m-1}^n - u_{m-3}^n}{8\Delta x^3}.$$

Inserting these approximations into equation (1), we get the intermediate result

$$\frac{\partial u_m^n}{\partial t} \approx F(u^n) = - \left(1 + \pi^2\right) \frac{u_{m+1}^n - u_{m-1}^n}{2\Delta x} - \frac{u_{m+3}^n - 3u_{m+1}^n + 3u_{m-1}^n - u_{m-3}^n}{8\Delta x^3}.$$

For later convenience, we write the Euler method and Crank-Nicholson method collectively with the theta method. This gives the final system of difference equations for the numerical solution

$$\frac{U_m^{n+1} - U_m^n}{\Delta t} = (1 - \theta)F(U^n) + \theta F(U^{n+1}), \quad (4)$$

where the Euler method or the Crank-Nicholson method is obtained by inserting  $\theta = 0$  or  $\theta = 1/2$ , respectively. In matrix form, the system can be written

$$(I - \theta \Delta t A) U^{n+1} = (I - (1 - \theta) \Delta t A) U^n, \quad (5)$$

where  $U^n = [U_0^n \ \dots \ U_{M-1}^n]^T$  and  $A =$

$$\frac{-1}{2\Delta x} \begin{bmatrix} 0 & +1 & & & & & -1 \\ -1 & 0 & +1 & & & & \\ & -1 & 0 & +1 & & & \\ & & -1 & 0 & +1 & & \\ & & & \ddots & \ddots & \ddots & \\ & & & & -1 & 0 & +1 \\ & & & & & -1 & 0 & +1 \\ & & & & & & -1 & 0 \\ +1 & & & & & & & -1 \end{bmatrix} - \frac{1 + \pi^2}{\Delta x^3} \begin{bmatrix} 0 & -3 & 0 & +1 & & & -1 & 0 & +3 \\ +3 & 0 & -3 & 0 & +1 & & & -1 & 0 \\ 0 & +3 & 0 & -3 & 0 & +1 & & & -1 \\ -1 & 0 & +3 & 0 & -3 & 0 & +1 & & \\ & \ddots & \ddots & \ddots & \ddots & \ddots & \ddots & \ddots & \\ & & -1 & 0 & +3 & 0 & -3 & 0 & +1 \\ +1 & & & -1 & 0 & +3 & 0 & -3 & 0 \\ 0 & +1 & & & -1 & 0 & +3 & 0 & -3 \\ -3 & 0 & +1 & & & -1 & 0 & +3 & 0 \end{bmatrix}.$$

where we have imposed periodic boundary conditions  $U_m^n = U_{m+M}^n$  by simply wrapping the spatial derivative stencils around the matrix.

We then solve the system by preparing  $U^0$  from the initial condition  $u(x, 0)$  and solve equation (5) repeatedly to step forward in time. Note that with the constant time step  $\Delta t$ , all matrices in equation (5) are constant in time, and the process of solving the system many times can be accelerated by for example LU-factorizing the matrix on the left side.

Next, we test our numerical solution on the problem defined by the initial condition  $u(x, 0) = \sin(\pi x)$  on  $x \in [-1, +1]$  with  $L = 2$ . The Fourier series of the analytical solution then has nonzero coefficients  $c_{\pm 1}(0) = \pm 1/2i$  and wavenumbers  $k_{\pm 1} = \pm \pi$ , which give rise to the analytical solution  $u(x, t) = \sin(\pi(x - t))$  when inserted into equation (2). As shown in figure 5, the solution represents a sine wave traveling with velocity 1 to the right.

In figure 6, we compare snapshots of the numerical solution at  $t = 1$  from the Euler method and the Crank-Nicholson method. Note that the Crank-Nicholson method approaches the exact solution with as little as  $N = 10$  time steps and a few hundred spatial grid points  $M$ , while the Euler method produces garbage with

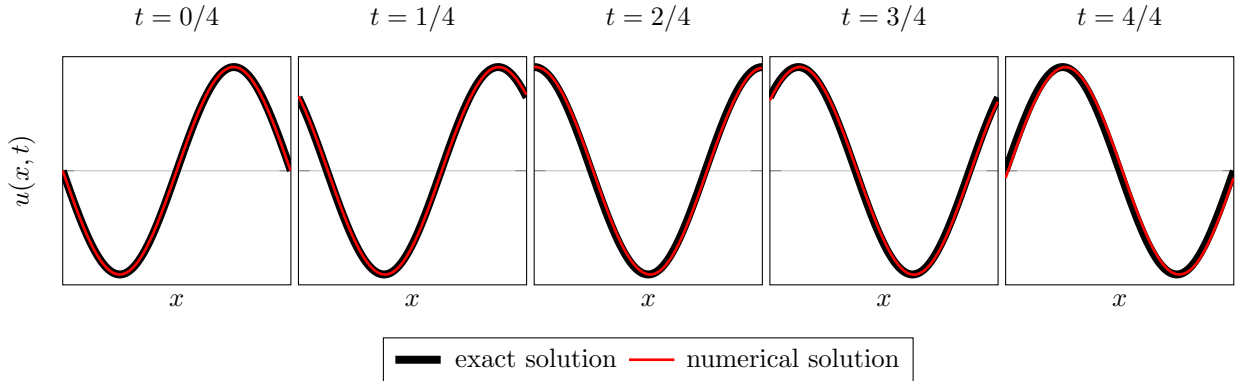


Figure 5: Comparison between the time evolution of the exact solution  $u(x, t) = \sin(\pi(x - t))$  and the numerical solution from the Crank-Nicholson method with  $\Delta x = 1/799$  and  $\Delta t = 1/99$ .

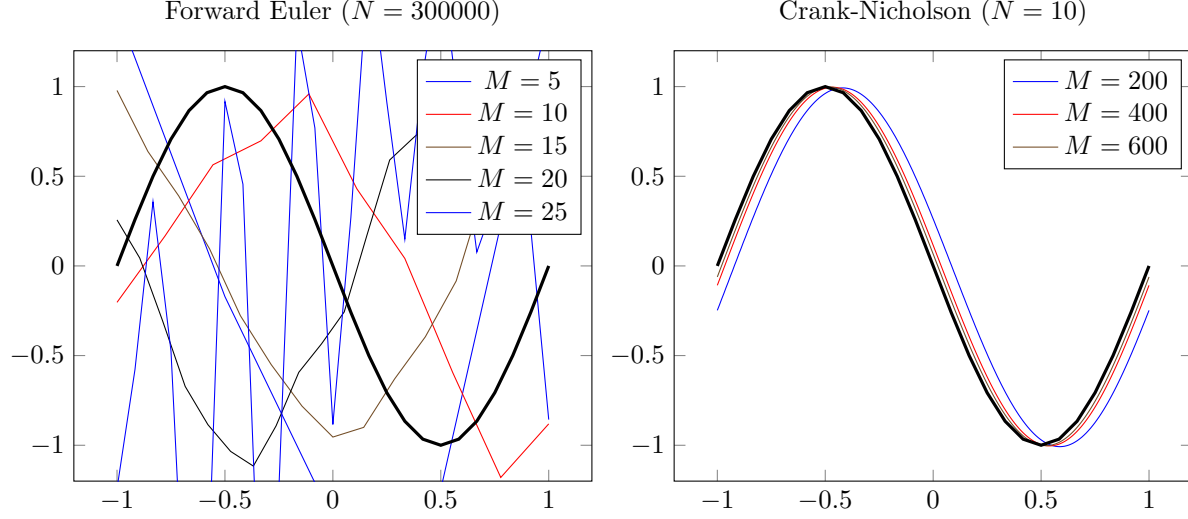


Figure 6

as many as  $N = 300000$  time steps and becomes less stable as the number of spatial grid points increases. The convergence plot of the error at  $t = 1$  in figure 7 supports our suspicions, showing the stable second order (first order???) nature of the Crank-Nicholson method and the instability of the Euler method.

### 4.3 Stability analysis

Motivated by the examples of the Euler method and the Crank-Nicholson method, we perform a Von Neumann analysis of their stability. Just like the exact solution, the numerical solution is subject to periodic boundary conditions in space and can therefore be expanded in a Fourier series

$$U_m^n = U(x_m, t_n) = \sum_l C_l^n \exp(ik_l x_m). \quad (6)$$

Consider now a single Fourier mode  $C_l^n \exp(ik_l x_m)$  in this series. Inserting it into equation (4), dividing by  $\exp(ik_l x_m)$  and expanding exponentials using Euler's identity gives

$$\frac{C_l^{n+1} - C_l^n}{\Delta t} = i \left( (1 - \theta) C_l^n + \theta C_l^{n+1} \right) f(k_l), \quad \text{where } f(k_l) = \left( - \left( 1 + \pi^2 \right) \frac{\sin(k_l h)}{h} - \frac{\sin^3(k_l h)}{h^3} \right).$$

Now look at the amplification factor  $G_l = C_l^{n+1}/C_l^n$  of Fourier mode  $l$  over one time step. With  $\theta = 1/2$ , the Crank-Nicholson method gives

$$G_l = \frac{1 + i\Delta t f(x)/2}{1 - i\Delta t f(x)/2} \implies |G_l| = 1. \quad (7)$$

The amplitude of all Fourier modes is thus preserved over time independently of  $\Delta t$  and  $\Delta x$ , and we say the Crank-Nicholson method is **unconditionally stable**.

The Euler method has  $\theta = 0$  and gives

$$G_l = 1 + i\Delta t f(x) \implies |G_l| = \sqrt{1 + \Delta t^2 f(k_l)^2}. \quad (8)$$

Since  $|\sin(k_l h)| \leq 1$  for all  $k_l$ , we can bound  $f(k_l)$  by

$$|f(k_l)| \leq \frac{(1 + \pi^2)}{\Delta x} + \frac{1}{\Delta x^3} = \frac{1}{\Delta x^3} \left( (1 + \pi^2) \Delta x^2 + 1 \right) \leq \frac{1}{\Delta x^3} \left( (1 + \pi^2) L^2 + 1 \right).$$

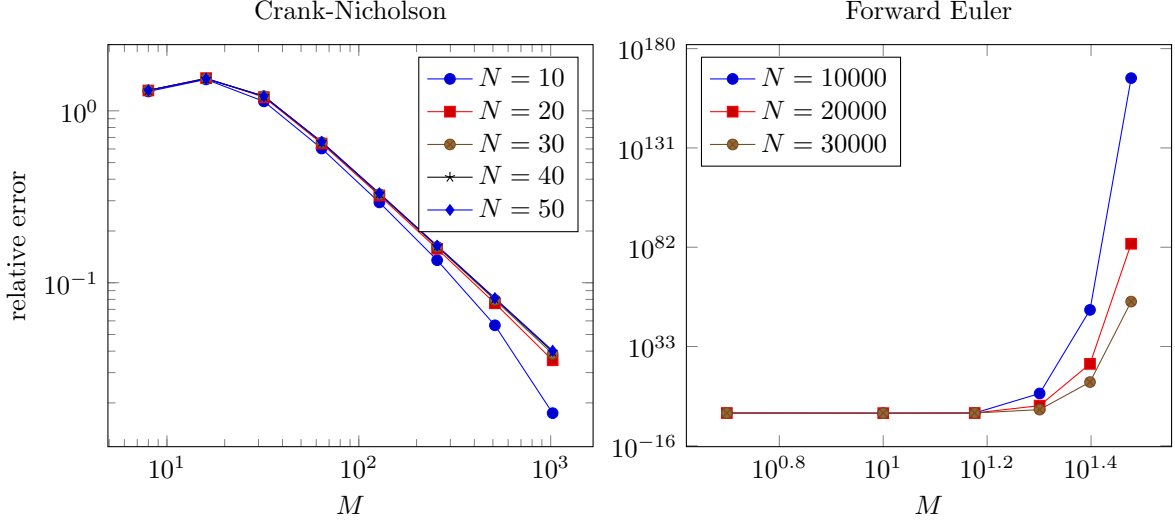


Figure 7: Convergence plot at  $t = 1$ , comparison between Crank-Nicholson and Forward Euler

Then  $|G_l| = \sqrt{1 + O(\Delta t^2 / \Delta x^6)}$ . The Von Neumann stability criterion  $|G_l| \leq 1 + O(\Delta t)$  [1] is attained only with  $\Delta t \leq O(\Delta x^6)$ , corresponding to *extremely small* time steps. Thus, while the Euler method in theory is **conditionally stable**, it is unstable for all practical time steps. The Crank-Nicholson method is far superior, as it remains stable while allowing much greater spatial resolution and larger time steps.

#### 4.4 Time evolution of norm

The stability of the finite difference methods can be even better illustrated by investigating the time evolution of the  $L_2$ -norm of the solution. To this end, we will first show that the  $L_2$ -norm of the analytical solution is preserved over time. Then we will show the time evolution of the  $L_2$ -norm of numerical solutions.

The  $L_2$ -norm of the analytical solution is defined as

$$\|u(x, t)\|_2 = \left( \frac{1}{2} \int_{-L/2}^{+L/2} |u(x, t)|^2 dx \right)^{1/2}.$$

To understand why it is constant in time, insert the Fourier expansion equation (2) for  $u(x, t)$  to get

$$\int_{-1}^{+1} dx |u(x, t)|^2 = \sum_{m,n} c_m(t) c_n^*(t) \underbrace{\int_{-L/2}^{+L/2} \exp(i(k_m - k_n)x) dx}_{L\delta_{mn}} = L \sum_n |c_n(t)|^2.$$

Since the Fourier coefficients (3) have constant magnitude  $|c_n(t)| = |c_n(0)|$ , the  $L_2$ -norm is the same at any time  $t$ . Note that it is the *odd* number of spatial derivatives in equation (1) that gives the exponential in equation (3) a purely imaginary argument and preserves the amplitude of each Fourier mode.

We now investigate the norm of the numerical solution with the initial gaussian  $u(x, 0) = \exp(x^2/0.1)$ . The time evolution illustrated in figure 8 shows how multiple Fourier modes are activated. In figure 9, we show how the norm of the numerical solution evolves over time. The Euler method diverges even with tiny time steps, reflecting the amplification factor  $G_l > 1$  found in equation (8). In contrast, the Crank-Nicholson method is always stable and preserves the norm of the solution, reflecting the amplification factor  $G_l = 1$  found in equation (7).

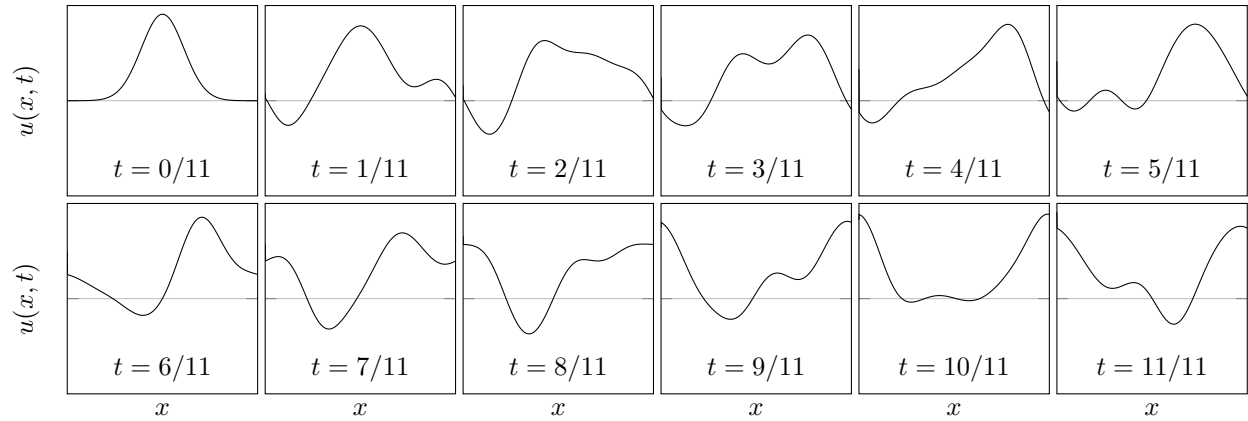


Figure 8: Time evolution of a initial gaussian  $u(x, 0) = \exp(x^2/0.1)$  computed from the Crank-Nicholson method on a grid with  $M = 800$  points in space and  $N = 100$  points in time.

The stability and norm preservation of the Crank-Nicholson method makes it the method of choice for problems like this, where the analytical solution is known to have the same property.

## References

- [1] Brynjulf Owren: *TMA4212 Numerical solution of partial differential equations with finite difference methods* (2017) [<http://www.math.ntnu.no/emner/TMA4212/2020v/notes/master.pdf>]

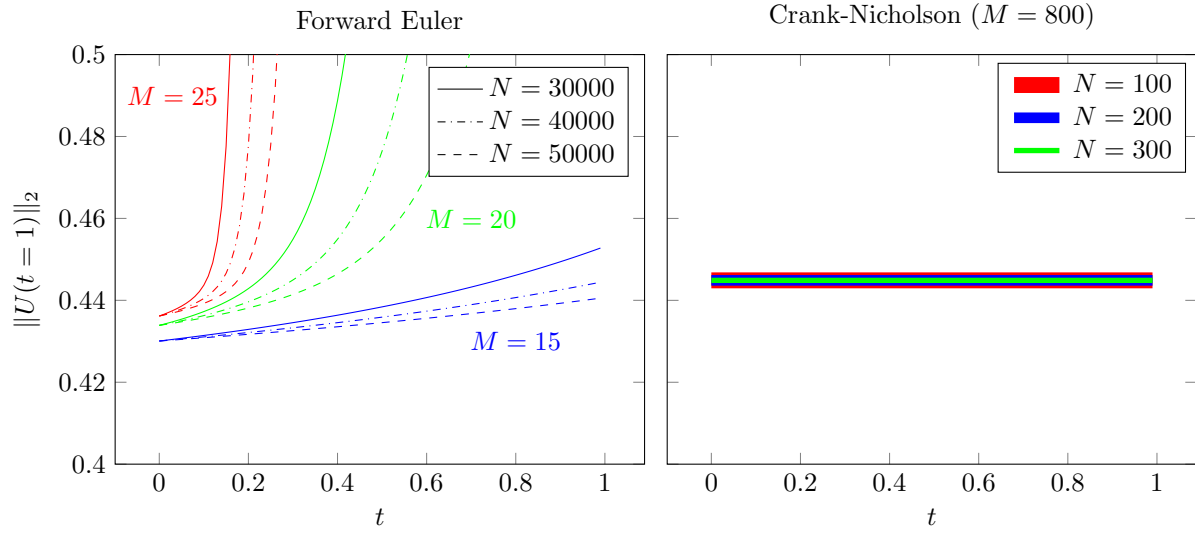


Figure 9: Time evolution of the discrete  $L_2$ -norm of the initial gaussian  $u(x,0) = \exp(x^2/0.1)$  computed from the Crank-Nicholson method and the Euler-method on different grids.