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# Search for additional neutral MSSM Higgs bosons in the $\tau\tau$ final state in proton-proton collisions at $\sqrt{s} = 13$ TeV

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## Abstract

A search is presented for additional neutral Higgs bosons in the  $\tau\tau$  final state in proton-proton collisions at the LHC. The search is performed in the context of the minimal supersymmetric extension of the standard model (MSSM), using the data collected with the CMS detector in 2016 at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of  $35.9 \text{ fb}^{-1}$ . To enhance the sensitivity to neutral MSSM Higgs bosons, the search includes production of the Higgs boson in association with b quarks. No significant deviation above the expected background is observed. Model-independent limits at 95% confidence level (CL) are set on the product of the branching fraction for the decay into  $\tau$  leptons and the cross section for the production via gluon fusion or in association with b quarks. These limits range from 18 pb at 90 GeV to 3.5 fb at 3.2 TeV for gluon fusion and from 15 pb (at 90 GeV) to 2.5 fb (at 3.2 TeV) for production in association with b quarks. In the  $m_h^{\text{mod}+}$  scenario these limits translate into a 95% CL exclusion of  $\tan\beta > 6$  for neutral Higgs boson masses below 250 GeV, where  $\tan\beta$  is the ratio of the vacuum expectation values of the neutral components of the two Higgs doublets. The 95% CL exclusion contour reaches 1.6 TeV for  $\tan\beta = 60$ .

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# 1 Introduction

The discovery of a Higgs boson at the CERN LHC in 2012 [1–3] has provided evidence that spontaneous symmetry breaking, as proposed by the Brout–Englert–Higgs mechanism [4–9], may indeed be realized in nature. The determination of the properties of the new particle, based on the complete LHC Run-1 data set [10, 11], has revealed its consistency with the standard model (SM) Higgs boson, within the experimental accuracy. However several questions remain, concerning, for example, the underlying mechanism responsible for the symmetry breaking, or the exact form of the potential that breaks the symmetry. To address these questions one of the main tasks of the LHC is the further exploration of the Higgs sector. This includes the search for more complex structures, for example, in the form of more than one Higgs doublet. Supersymmetry (SUSY) [12, 13] is an example of a beyond the SM theory with a more complex Higgs sector. In the minimal supersymmetric standard model (MSSM) [14, 15] each particle of the SM is complemented by a SUSY partner, which has the same properties apart from its spin. The Higgs sector of the MSSM consists of two complex Higgs doublets,  $H_u$  and  $H_d$ , to provide masses for up- and down-type fermions. In the CP-conserving MSSM this leads to the prediction of five physical Higgs bosons: two charged Higgs bosons  $H^\pm$ , two neutral scalar Higgs bosons  $h$  and  $H$  (with masses  $m_h < m_H$ ) and one neutral pseudoscalar Higgs boson  $A$ . At tree-level in the MSSM, the masses of these five Higgs bosons and their mixing can be expressed in terms of the gauge boson masses and two additional parameters, which can be chosen as the mass of the  $A$ ,  $m_A$ , and the ratio of the vacuum expectation values of the neutral components of the two Higgs doublets

$$\tan \beta = \frac{\langle H_u^0 \rangle}{\langle H_d^0 \rangle} = \frac{v_u}{v_d}. \quad (1)$$

Dependencies on additional parameters of the SUSY breaking mechanism enter via higher-order corrections in perturbation theory. In the exploration of the MSSM Higgs sector these parameters are usually set to fixed values in the form of indicative benchmark scenarios [16] to illustrate certain properties of the theory. For values of  $m_A \gtrsim 300$  GeV, which seem to be favored by data [10, 11, 17, 18], the MSSM is close to the decoupling limit: the  $h$  usually takes the role of the observed SM-like Higgs boson at 125 GeV and the  $H$  and  $A$  are nearly degenerate in mass.

At leading-order (LO), the coupling of the  $H$  and the  $A$  to down-type fermions is enhanced by  $\tan \beta$  with respect to the expectation for an SM Higgs boson of the same mass, while the coupling to vector bosons and up-type fermions is suppressed. The enhanced coupling to down-type fermions makes searches for additional heavy neutral Higgs bosons that exploit final states containing such fermions particularly interesting. It also has consequences for the production: firstly, the production in association with  $b$  quarks dominates over the production via gluon fusion for large values of  $\tan \beta$ . Secondly, in gluon fusion production the kinematic properties of the Higgs boson change as a function of  $\tan \beta$  due to the increasing contribution of  $b$  quarks in the fermion loop. Diagrams for  $h$ ,  $H$ , and  $A$  production at LO are shown in Fig. 1.

Searches for additional heavy neutral Higgs bosons in the context of the MSSM were carried out in  $e^+e^-$  collisions at LEP [19] and in proton-antiproton collisions at the Tevatron [20–23]. At the LHC such searches have been carried out by the ATLAS and CMS Collaborations in the  $b$  quark [24, 25], dimuon [26, 27], and  $\tau\tau$  [17, 18, 26, 28–31] final states. The better experimental accessibility with respect to the  $b$  quark final state and the larger mass, and therefore larger coupling, with respect to the muon give the  $\tau\tau$  final state a leading role in these searches.

In this paper the results of a search for additional heavy neutral Higgs bosons in the con-

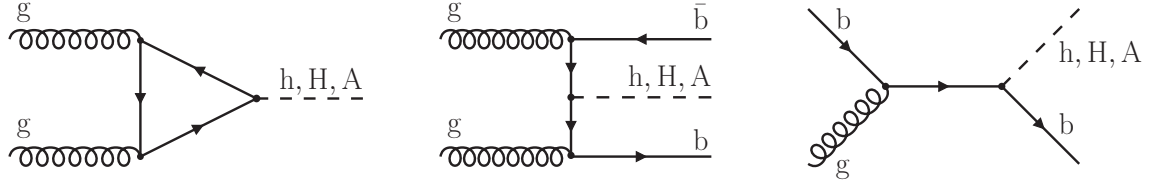


Figure 1: Diagrams for the production of neutral Higgs bosons (left) via gluon fusion and (middle and right) in association with  $b$  quarks. In supersymmetric extensions of the SM, the super-partners also contribute to the fermion loop, shown in the left panel. In the middle panel a pair of  $b$  quarks is produced from two gluons (the LO process in the four-flavor scheme). In the right panel the Higgs boson is radiated from a  $b$  quark in the proton (the LO process in the five-flavor scheme).

text of the MSSM are presented. They are based on the 2016 pp collision data set, taken at a center-of-mass energy of 13 TeV, by the CMS experiment, and correspond to an integrated luminosity of  $35.9 \text{ fb}^{-1}$ . The analysis is performed in four different  $\tau\tau$  final states:  $e\mu$ ,  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$ , where  $e$ ,  $\mu$  and  $\tau_h$  indicate  $\tau$  lepton decays into electrons, muons and hadrons respectively. For this analysis the most significant backgrounds are estimated from data, by using new techniques with respect to previous publications by CMS. Upper limits are presented on the product of the branching fraction for the decay into  $\tau$  leptons and the cross section for the production of a single narrow resonance via gluon fusion or in association with  $b$  quarks. In addition, exclusion contours in the  $m_A$ - $\tan\beta$  plane in selected MSSM benchmark scenarios are provided.

In Sections 2 and 3 the CMS detector and the event reconstruction are described. Section 4 summarizes the event selection and categorization. The event simulation and background estimation methods used for the analysis are described in Section 5. The signal extraction is discussed in Section 6, followed by a discussion of the systematic uncertainties in Section 7. Section 8 contains the results of the analysis. A summary is given in Section 9.

## 2 The CMS detector

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity coverage provided by the barrel and endcap detectors. Muons are detected in gas-ionization chambers embedded in the steel flux-return yoke outside the solenoid.

The silicon tracker measures charged particles within the pseudorapidity range  $|\eta| < 2.5$ . It consists of 1440 silicon pixel and 15 148 silicon strip detector modules. For nonisolated particles with a transverse momentum of  $1 < p_T < 10 \text{ GeV}$  and  $|\eta| < 1.4$ , the track resolutions are typically 1.5% in  $p_T$  and 25–90 (45–150)  $\mu\text{m}$  in the transverse (longitudinal) impact parameter [32]. The electron momentum is estimated by combining the energy measurement in the ECAL with the momentum measurement in the tracker. The momentum resolution for electrons with  $p_T \approx 45 \text{ GeV}$  from  $Z \rightarrow ee$  decays ranges from 1.7% for nonshowering electrons in the barrel region to 4.5% for showering electrons in the endcaps [33]. Muons are measured in the pseudorapidity range  $|\eta| < 2.4$ , with detection planes made using three technologies: drift tubes, cathode strip chambers, and resistive-plate chambers. Matching muons to tracks mea-

sured in the silicon tracker results in a relative  $p_T$  resolution for muons with  $20 < p_T < 100$  GeV of 1.3 to 2.0% in the barrel and better than 6% in the endcaps. The  $p_T$  resolution in the barrel is better than 10% for muons with  $p_T$  up to 1 TeV [34]. In the barrel section of the ECAL, an energy resolution of about 1% is achieved for unconverted or late-converting photons in the tens of GeV energy range. The remaining barrel photons have a resolution of better than 2.5% for  $|\eta| \leq 1.4$ . In the endcaps, the resolution of unconverted or late-converting photons is about 2.5%, while the remaining endcap photons have a resolution between 3 and 4% [35]. When combining information from the entire detector, the jet energy resolution amounts typically to 15% at 10 GeV, 8% at 100 GeV, and 4% at 1 TeV, to be compared to about 40, 12, and 5% obtained when the ECAL and HCAL calorimeters alone are used.

Events of interest are selected using a two-tiered trigger system [36]. The first level, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a time interval of less than 4  $\mu$ s. The second level, known as the high-level trigger, consists of a farm of processors running a version of the full event reconstruction software optimized for fast processing, and reduces the event rate to around 1.0 kHz before data storage.

A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [37].

### 3 Event reconstruction

The reconstruction of the pp collision products is based on the particle-flow (PF) algorithm as described in Ref. [38], combining the available information from all CMS subdetectors to reconstruct an unambiguous set of individual particle candidates. The particle candidates are categorized into electrons, photons, muons, and charged and neutral hadrons. During the 2016 data taking period the CMS experiment was operating with, on average, 23 pp collisions per bunch crossing. The fully recorded data of a bunch crossing defines an *event* for further processing. Collision vertices are obtained from reconstructed tracks using a deterministic annealing algorithm [39]. The reconstructed vertex with the largest value of summed physics-object  $p_T^2$  is taken to be the primary pp interaction vertex. The physics objects for this purpose are the jets, clustered using the jet finding algorithm [40, 41], as described below, with the tracks assigned to the vertex as inputs, and the associated missing transverse momentum, taken as the negative vector sum of the  $p_T$  of those jets. Any other collision vertices in the event are associated with additional soft inelastic pp collisions called *pileup*.

Electrons are reconstructed by combining clusters of energy deposits in the ECAL with hits in the tracker [33]. To increase their purity, reconstructed electrons are required to pass a multivariate electron identification discriminant, which combines information on track quality, shower shape, and kinematic quantities. For this analysis working points with an efficiency between 80 and 90% are used to identify electrons. Muons in the event are reconstructed by performing a simultaneous track fit to hits in the tracker and in the muon chambers [34]. The presence of hits in the muon chambers already leads to a strong suppression of particles misidentified as muons. Additional identification requirements on the track fit quality and the compatibility of individual track segments with the fitted track can reduce the misidentification rate further. For this analysis muon identification requirements with an efficiency of  $\approx 99\%$  are chosen. The contribution from backgrounds to the electron (muon) selection is further reduced by requiring the corresponding lepton to be isolated from any hadronic activity in the detector. This property is quantified by a relative isolation variable  $I_{\text{rel}}^{e(\mu)}$ , which starts from the sum of the

transverse momentum (energy) of all charged (neutral) particles,  $I_{\text{abs}}^{e(\mu)} = (\sum p_{T,i} + \sum E_{T,i})$  in a predefined cone of radius  $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$  around the lepton direction at the primary collision vertex, where  $\Delta\eta$  and  $\Delta\phi$  (measured in radians) correspond to the angular distance of the particle to the lepton in the  $\eta$  and  $\phi$  directions. The chosen cone size is  $\Delta R < 0.3$  (0.4) for electrons (muons). The lepton itself is not included in this calculation. To mitigate any distortions from pileup only those charged particles whose tracks are associated with the primary collision vertex are taken into account. The presence of neutral particles from pileup is estimated by summing the  $p_T$  of charged particles in the isolation cone whose tracks have been associated to pileup vertices, and multiplying this quantity by a factor of 0.5 to account for the approximate ratio of neutral to charged hadron production. The value obtained is subtracted from  $I_{\text{abs}}^{e(\mu)}$  and the result set to zero in case of negative values. Finally,  $I_{\text{abs}}^{e(\mu)}$  is divided by the  $p_T$  of the lepton to result in  $I_{\text{rel}}^{e(\mu)}$ .

For further characterization of the event all reconstructed PF objects are clustered into jets using the anti- $k_T$  jet clustering algorithm as implemented in FASTJET [40, 41] with a distance parameter of 0.4. To identify jets resulting from the hadronization of b quarks a re-optimized version of the *combined secondary vertex* b tagging algorithm, which exploits information from the decay vertices of long-lived hadrons, and the impact parameters of charged-particle tracks, in a combined discriminant, is used [42]. In the analysis a working point corresponding to a b jet identification efficiency of  $\approx 70\%$  and a misidentification rate for light quarks and gluons of 1% is chosen. Jets are also used as seeds for the reconstruction of hadronic  $\tau$  lepton decays. This is done by further exploiting the substructure of the jets, using the *hadrons-plus-strips* algorithm, as described in Refs. [43, 44]. For the analysis the decay into three charged hadrons and the decay into a single charged hadron accompanied by up to two neutral pions with  $p_T > 2.5 \text{ GeV}$  are used. The neutral pions are reconstructed as *strips* with dynamic size from reconstructed electrons and photons contained in the seeding jet, where the strip size varies as a function of the  $p_T$  of the electron or photon candidate. The hadronic  $\tau$  decay mode is then obtained by combining the charged hadrons with the strips. Since they do not carry color charge, high- $p_T$   $\tau$  leptons are expected to be isolated from any hadronic activity in the event as are high- $p_T$  electrons and muons. Furthermore, in accordance with its finite lifetime the charged decay products of the  $\tau$  lepton are expected to be slightly displaced from the primary collision vertex. To distinguish hadronic  $\tau$  lepton decays from jets originating from the hadronization of quarks or gluons a multivariate  $\tau_h$  identification discriminant is used [43]. It combines information on the hadronic activity in the detector in the vicinity of the  $\tau_h$  candidate with the reconstructed lifetime information from the tracks of the charged decay products. Of the predefined working points in Ref. [43] this analysis makes use of the Tight, Medium, and VeryLoose working points. These have efficiencies of 27% (Tight), 51% (Medium), and 71% (VeryLoose), for quark/gluon misidentification rates of less than  $4.4 \times 10^{-4}$  (Tight),  $3.3 \times 10^{-3}$  (Medium), and  $1.3 \times 10^{-2}$  (VeryLoose). Finally, requirements are imposed to reduce the misidentification of electrons and muons as hadronic  $\tau$  lepton decays. Also here predefined working points are used to discriminate against electrons, with efficiencies ranging from 65% (Tight) to 94% (VeryLoose) for electron misidentification rates between  $6.2 \times 10^{-4}$  (Tight) and  $2.4 \times 10^{-2}$  (VeryLoose). The misidentification rate of muons as hadronic  $\tau$  lepton decays is of  $\mathcal{O}(10^{-3})$ , for a  $\tau_h$  identification efficiency of 99%.

The missing transverse momentum vector  $\vec{p}_T^{\text{miss}}$ , defined as the negative vector sum of the  $p_T$  of all reconstructed PF objects, is also used to characterize the events. Its magnitude is referred to as  $p_T^{\text{miss}}$ . It is used for the discrimination of backgrounds that are expected to contain neutrinos with significantly more  $p_T$  than expected from the  $\tau\tau$  final state, such as W boson

production in association with jets (W+jets). It is furthermore used for the calculation of the final discriminating variable that is used for the statistical analysis, as detailed in Section 6.

## 4 Event selection and categorization

The four most sensitive final states of the  $\tau\tau$  pair are exploited:  $e\mu$ ,  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$ . The online selection for the  $e\tau_h$  ( $\mu\tau_h$ ) final state is based on the presence of at least one electron (muon) with  $p_T > 25$  (22) GeV and  $|\eta| < 2.1$  at trigger level. The online selection for the  $e\mu$  final state relies on a logical *or* of two lower threshold triggers that both require the presence of an electron and muon in the event with  $p_T > 23$  GeV for the higher- $p_T$  lepton and  $p_T > 12$  (8) GeV for the lower- $p_T$  electron (muon). In the  $\tau_h\tau_h$  final state, a trigger decision based on the presence of two hadronically decaying  $\tau$  leptons with  $p_T > 35$  GeV and  $|\eta| < 2.1$  is used.

Table 1: Kinematic selection of the  $\tau$  lepton decay products in the  $e\mu$ ,  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states. The expression “First (Second) object” refers to the final state label used in the first column.

Final state	First object	Second object
$e\mu^\dagger$	$p_T^e > 13 \text{ GeV},  \eta^e  < 2.5$	$p_T^\mu > 10 \text{ GeV},  \eta^\mu  < 2.4$
$e\tau_h$	$p_T^e > 26 \text{ GeV},  \eta^e  < 2.1$	$p_T^{\tau_h} > 30 \text{ GeV},  \eta^{\tau_h}  < 2.3$
$\mu\tau_h$	$p_T^\mu > 23 \text{ GeV},  \eta^\mu  < 2.1$	$p_T^{\tau_h} > 30 \text{ GeV},  \eta^{\tau_h}  < 2.3$
$\tau_h\tau_h$	$p_T^{\tau_h} > 40 \text{ GeV},  \eta^{\tau_h}  < 2.1$	

<sup>†</sup> For events passing only one trigger an additional requirement of  $p_T > 24$  GeV is applied on the higher- $p_T$  lepton candidate as explained in the text.

Requirements on the  $p_T$  and  $\eta$  of the reconstructed  $\tau$  lepton decay products are applied in the offline analysis as given in Table 1. In the  $e\mu$  final state an electron with  $p_T > 13$  GeV and  $|\eta| < 2.5$  and a muon with  $p_T > 10$  GeV and  $|\eta| < 2.4$  are required. If the event passed only one trigger the lepton identified with the higher- $p_T$  trigger object is required to have a  $p_T > 24$  GeV. This guarantees a trigger acceptance well above the turn-on of at least one of the triggers used. Both leptons are required to pass identification criteria as described in Section 3 and to be isolated according to  $I_{\text{rel}}^{e(\mu)} < 0.15$  (0.2). Events with additional electrons or muons fulfilling looser selection requirements than these are rejected.

In the  $e\tau_h$  ( $\mu\tau_h$ ) final state an electron (muon) with  $p_T > 26$  (23) GeV and  $|\eta| < 2.1$  and a  $\tau_h$  candidate with  $p_T > 30$  GeV and  $|\eta| < 2.3$  are required. The electron (muon) and the  $\tau_h$  candidate should fulfill the identification requirements as described in Section 3. The  $\tau_h$  candidate should pass the Tight working point of the  $\tau_h$  identification discriminant, the Tight (VeryLoose) working point of the discriminant to suppress electrons and the Loose (Tight) working point of the discriminant to suppress muons in the  $e\tau_h$  ( $\mu\tau_h$ ) case. In addition, the electron (muon) should be isolated according to  $I_{\text{rel}}^{e(\mu)} < 0.1$  (0.15). Events with additional electrons or muons fulfilling looser selection requirements are rejected.

In the  $\tau_h\tau_h$  final state two  $\tau_h$  candidates with  $p_T > 40$  GeV and  $|\eta| < 2.1$  are required. Both must pass the Medium working point of the  $\tau_h$  identification discriminant, the VeryLoose working point of the discriminant against electrons and the Loose working point of the discriminant against muons. Events with additional electrons or muons fulfilling looser requirements on identification, isolation and  $p_T$  than described for the  $e\tau_h$  or  $\mu\tau_h$  final state above are rejected.

In all cases the decay products of the two  $\tau$  leptons are required to be of opposite charge, separated by more than 0.5 in  $\Delta R$  and associated to the primary collision vertex within a distance of

0.045 cm in the transverse plane for electrons and muons and 0.2 cm along the beam axis for all final-state particles. The vetoing of additional electrons or muons helps with the suppression of backgrounds and ensures that no event will be categorized according to more than one  $\tau\tau$  final state. At most 0.8% of the selected events contain more  $\tau_h$  candidates than required for the corresponding final state. In this case, the  $\tau\tau$  pair with the most isolated final state products is chosen.

To increase the sensitivity of the analysis all selected events are further categorized: events with at least one jet with  $p_T > 20$  GeV and  $|\eta| < 2.4$  that passes the b tagging requirement as described in Section 3 are combined into a global b-tag category. This category is designed to target the production of the Higgs boson in association with b quarks. All other events are added to a global no b-tag category.

In the  $e\mu$  final state each event category is further split into three subcategories based on the quantity  $D_\zeta$ , introduced for the first time in Ref. [45], defined as

$$D_\zeta = p_\zeta^{\text{miss}} - 0.85 p_\zeta^{\text{vis}}; \quad p_\zeta^{\text{miss}} = \vec{p}_T^{\text{miss}} \hat{\zeta}; \quad p_\zeta^{\text{vis}} = (\vec{p}_T^e + \vec{p}_T^\mu) \hat{\zeta}, \quad (2)$$

where  $\vec{p}_T^{e(\mu)}$  corresponds to the transverse momentum vector of the electron (muon) and  $\hat{\zeta}$  to the bisectonal direction between the electron and the muon in the transverse plane. The linear combination of  $p_\zeta^{\text{miss}}$  and  $p_\zeta^{\text{vis}}$  has been chosen to optimize the sensitivity of the analysis in the  $e\mu$  final state. The variable  $D_\zeta$  is especially suited to suppress W+jets and  $t\bar{t}$  events, where the reconstructed lepton candidates and the direction of  $\vec{p}_T^{\text{miss}}$  are distributed more isotropically in the detector than for genuine  $\tau\tau$  signal events. The categories are defined as low- $D_\zeta$  ( $-50 < D_\zeta \leq -10$  GeV), medium- $D_\zeta$  ( $-10 < D_\zeta \leq 30$  GeV) and high- $D_\zeta$  ( $D_\zeta > 30$  GeV). In this way categories with different fractions of signal and  $t\bar{t}$  events can be exploited for the statistical analysis. The expected signal, for all masses tested, is mostly located in the medium- $D_\zeta$  subcategory.

In the  $e\tau_h$  ( $\mu\tau_h$ ) final state each global event category is further split into two subcategories based on the transverse mass of the electron or muon and  $p_T^{\text{miss}}$

$$m_T^{e(\mu)} = \sqrt{2 p_T^{e(\mu)} p_T^{\text{miss}} (1 - \cos \Delta\phi)}. \quad (3)$$

This transverse mass is used to discriminate between the signal and the backgrounds from W+jets and  $t\bar{t}$  events. In Eq. (3)  $p_T^{e(\mu)}$  refers to the  $p_T$  of the electron (muon) and  $\Delta\phi$  to the difference in the azimuthal angle between the electron (muon) and  $\vec{p}_T^{\text{miss}}$ . The categories are defined as tight- $m_T$  ( $m_T^{e(\mu)} < 40$  GeV) and loose- $m_T$  ( $40 < m_T^{e(\mu)} < 70$  GeV). The bulk of the signal events, particularly for the low-mass hypotheses, lie in the tight- $m_T$  subcategory. The loose- $m_T$  category has been added to increase the signal acceptance for mass hypotheses of  $m_{A,H} > 700$  GeV.

In combination this leads to 16 event categories entering the statistical analysis, complemented by three background control regions, as discussed in Section 5. In Fig. 2, the  $D_\zeta$  and  $m_T^\mu$  distributions are shown in the  $e\mu$  and  $\mu\tau_h$  final states respectively, before splitting the events into categories, indicating the corresponding subcategorization. A discussion of the composition of the expected background contributions is given in Section 6. A graphical representation of the complete event categorization is given in Fig. 3.



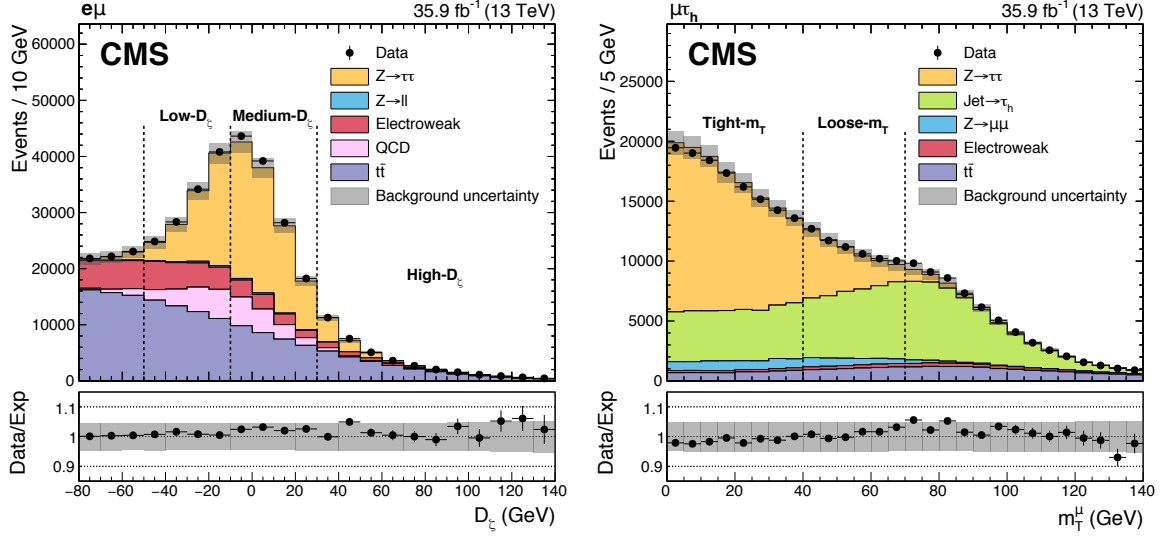


Figure 2: Observed and expected distributions of (left)  $D_\zeta$  in the  $e\mu$  final state and (right)  $m_T^\mu$  in the  $\mu\tau_h$  final state. The dashed vertical lines indicate the definition of the subcategories in each final state. The label “jet  $\rightarrow \tau_h$ ” indicates events with jets misidentified as hadronic  $\tau$  lepton decays, e.g.  $W$ +jets events, which are estimated from data as described in Section 5.2. A detailed description of the composition of the expected background is given in Section 6. The distributions are shown before any event categorization and prior to the fit used for the signal extraction. For these figures no uncertainties that affect the shape of the distributions have been included in the uncertainty model.

	No b-tag			b-tag		
$H \rightarrow \tau\tau \rightarrow e\mu$	Low- $D_\zeta$	Medium- $D_\zeta$	High- $D_\zeta$	Low- $D_\zeta$	Medium- $D_\zeta$	High- $D_\zeta$
$H \rightarrow \tau\tau \rightarrow e\tau_h$	Loose- $m_T$		Tight- $m_T$	Loose- $m_T$		Tight- $m_T$
$H \rightarrow \tau\tau \rightarrow \mu\tau_h$	Loose- $m_T$		Tight- $m_T$	Loose- $m_T$		Tight- $m_T$
$H \rightarrow \tau\tau \rightarrow \tau_h\tau_h$						
$Z \rightarrow \mu\mu$						
$t\bar{t}(e\mu)$						
	Signal region (SR)					
	Control region					

Figure 3: Overview of all event subcategories that enter the statistical analysis. Sixteen signal categories are complemented by three background control regions in the main analysis as described in Section 5.

## 5 Event simulation and background estimation

A list of all SM backgrounds that contribute to the event selection described in Section 4 is given in Table 2. The most obvious background originates from Z boson production in the  $\tau\tau$  final state ( $Z \rightarrow \tau\tau$ ). Since the analysis is not sensitive to the CP-eigenvalue or spin of the Higgs boson, the signal can be distinguished from this background only by the difference in mass of the associated bosons. The same is true for  $Z \rightarrow \ell\ell$  events, where  $\ell$  refers to an electron or muon, if one of the leptons is misidentified as a hadronic  $\tau$  lepton decay. Similar arguments hold for  $t\bar{t}$  production (a dominant background especially in the  $e\mu$  final state), the production of single t quarks and vector boson pair production (WW, WZ, and ZZ). Common to all these backgrounds in the  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states is that they can be misinterpreted as signal events in two ways: firstly, if the final state contains one or more genuine  $\tau$  leptons or if an electron or muon in the final state is misinterpreted as a hadronic  $\tau$  lepton decay, and secondly, if one or more jets are misinterpreted as hadronic  $\tau$  lepton decays. In Table 2 the former is labeled as “ $\tau/\ell \rightarrow \tau_h$ ”, whilst the latter is labeled as “jet  $\rightarrow \tau_h$ ”. Typical misidentification probabilities are given in Section 3. Backgrounds due to W+jets or SM events comprised uniquely of jets produced through the strong interaction, referred to as quantum chromodynamics (QCD) multijet production, predominantly contribute to the event selection via the misidentification of jets as hadronic  $\tau$  lepton decays. The level to which each of these processes contributes to the event selection depends on the final state.

Table 2: Background processes contributing to the event selection, as given in Section 4. The first row corresponds to the SM Higgs boson in the  $\tau\tau$  final state, which is also taken into account in the statistical analysis. The further splitting of the processes in the second column refers only to final states that contain a  $\tau_h$  candidate. The label “MC” implies that the process is taken from simulation; the label “ $F_F$ ” implies that the process is determined from data using the fake factor method, as described in Section 5.2. The label “CR” implies that both the shape and normalization of QCD multijet events are estimated from control regions in data. The symbol  $\ell$  corresponds to an electron or muon.

Background process	Misidentification	$e\mu$	$e\tau_h$	$\mu\tau_h$	$\tau_h\tau_h$
$H \rightarrow \tau\tau$ (SM)		MC	MC	MC	MC
$Z \rightarrow \tau\tau$		MC <sup>†</sup>	MC <sup>†</sup>	MC <sup>†</sup>	MC <sup>†</sup>
$Z \rightarrow \ell\ell$	$\ell \rightarrow \tau_h$	MC	MC	MC	MC
	Jet $\rightarrow \tau_h$		$F_F$	$F_F$	$F_F$
Diboson+single t	$\tau/\ell \rightarrow \tau_h$	MC	MC	MC	MC
	Jet $\rightarrow \tau_h$		$F_F$	$F_F$	$F_F$
$t\bar{t}$	$\tau/\ell \rightarrow \tau_h$	MC <sup>†</sup>	MC <sup>†</sup>	MC <sup>†</sup>	MC <sup>†</sup>
	Jet $\rightarrow \tau_h$		$F_F$	$F_F$	$F_F$
W+jets	Jet $\rightarrow \tau_h$	MC	$F_F$	$F_F$	$F_F$
QCD multijet production	Jet $\rightarrow \tau_h$	CR	$F_F$	$F_F$	$F_F$

<sup>†</sup> Normalization from control region in data.

### 5.1 Event simulation

Drell–Yan events in the dielectron, dimuon, and  $\tau\tau$  final states, and W+jets events are generated at LO precision in the coupling strength  $\alpha_s$  [46], using the MADGRAPH5\_aMC@NLO 2.2.2

event generator [47]. To increase the number of simulated events in regions of high signal purity supplementary samples are generated with up to four outgoing partons in the hard interaction. For diboson production MADGRAPH5\_aMC@NLO is used at next-to-leading order (NLO) precision. For  $t\bar{t}$  and single  $t$  quark production samples are generated at NLO precision using POWHEG 2.0 [48–53]. For the interpretation of the results the expected contribution of the SM Higgs boson is taken into account; this process is simulated using POWHEG separately for the production via gluon fusion, vector boson fusion (VBF), or in association with a  $Z$  (ZH) or  $W$  (WH) boson. When compared to data and not modified by a control measurement in data, Drell–Yan,  $W$ +jets,  $t\bar{t}$ , and single  $t$  quark events in the  $tW$ -channel are normalized to their cross sections at next-to-next-to-leading order (NNLO) precision [54–56]. Single  $t$  quark production in the  $t$ -channel and diboson events are normalized to their cross sections at NLO precision [56, 57]. The gluon fusion signal process is simulated at LO precision using PYTHIA 8.212 [58]. For the statistical analysis the Higgs boson  $p_T$  distribution is weighted to NLO precision using POWHEG. To account for the multiscale nature of the process in the NLO plus parton shower POWHEG prediction, the  $p_T$  spectra corresponding to the contributions from the  $t$  quark alone, the  $b$  quark alone and the  $t\bar{b}$ -interference are each calculated separately, using a POWHEG damping factor set to the individual scales as discussed in Refs. [59–61]. For the model-independent limits the individual distributions are combined according to their contribution to the total cross section as expected for a CP-even Higgs boson with given mass in the SM. In the model-dependent interpretation in the MSSM, where the contributions of the individual distributions also depend on the model parameters, these contributions are obtained using POWHEG in the two Higgs doublet mode. Each distribution is scaled, depending on the model parameters, using the effective Yukawa couplings as predicted by the corresponding benchmark model, before all distributions are combined into one single prediction. In this context also the  $\tan\beta$  enhanced SUSY corrections to the  $b$  quark coupling are taken into account via the corresponding effective Yukawa coupling where appropriate. Other SUSY contributions have been checked to be less than a few percent and are neglected. The associated production with  $b$  quarks is simulated at NLO precision using MADGRAPH5\_aMC@NLO.

For the generation of all signal and background processes the NNPDF3.0 parton distribution functions (PDFs) are used, as described in Ref. [62]. The description of the underlying event is parametrized according to the CUETP8M1 tune [63]. Hadronic showering and hadronization, as well as the  $\tau$  lepton decays, are modeled using PYTHIA. For all simulated events the effect of the observed pileup is taken into account. For this purpose additional inclusive inelastic  $pp$  collisions are generated with PYTHIA and added to all simulated events according to the expected pileup profile. All events generated are passed through a GEANT4-based [64] simulation of the CMS detector and reconstructed using the same version of the CMS event reconstruction software as used for the data. The observed event yields in each event category and the composition of the expected background contributions to the selected events are given in Table 3.

## 5.2 Backgrounds estimated from data

A large fraction of the backgrounds outlined for the  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states in Table 2 can be attributed to jets misidentified as hadronic  $\tau$  lepton decays. For the signal extraction the shape and normalization of these backgrounds are estimated from control regions in data, using the “fake factor” method, as described in Ref [65]. In this approach the number of events for a certain background  $i$  due to  $\text{jet} \rightarrow \tau_h$  misidentification is estimated from a region that only differs from the signal region (SR) by modifying the  $\tau_h$  identification requirement. This region is referred to as the application region (AR). For this purpose the  $\tau_h$  identification is required

Table 3: Observed number of selected events ( $N_{\text{data}}$ ) and the relative contribution of the expected backgrounds in all event categories in the  $e\mu$ ,  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states. The relative contribution of the expected backgrounds is given in %, including the contribution of an SM Higgs boson with a mass of 125 GeV, and prior to the fit used for the signal extraction. In all but the  $e\mu$  final state, processes in which a jet is misidentified as a hadronic  $\tau$  lepton decay are subsumed into a common jet  $\rightarrow \tau_h$  background class which is estimated from data.

Category		$N_{\text{data}}$	$H \rightarrow \tau\tau$ (SM)	$Z \rightarrow \tau\tau$	$Z \rightarrow \ell\ell$	Diboson+single t	$t\bar{t}$	W+jets	QCD	
$e\mu$	No b-tag	Low- $D_\zeta$	85816	0.24	49.55	1.06	11.50	15.73	6.03	15.88
		Medium- $D_\zeta$	102143	0.37	69.90	0.75	4.97	8.22	1.94	13.85
		High- $D_\zeta$	18364	0.50	45.66	0.38	11.95	32.59	1.94	6.99
		Low- $D_\zeta$	42870	0.01	2.08	0.08	6.88	89.19	0.27	1.49
b-tag	Medium- $D_\zeta$	27803	0.03	4.37	0.12	6.57	86.63	0.14	2.14	
	High- $D_\zeta$	20208	0.03	1.72	0.01	7.00	91.09	0.16	> 0.01	
Category		$N_{\text{data}}$	$H \rightarrow \tau\tau$ (SM)	$Z \rightarrow \tau\tau$	$Z \rightarrow \ell\ell^t$	Diboson+single t	$t\bar{t}^t$	jet $\rightarrow \tau_h$		
$e\tau_h$	No b-tag	Tight- $m_\tau$	49120	0.66	55.82	12.72	0.95	1.24	28.61	
		Loose- $m_\tau$	29108	0.35	28.41	11.71	1.70	2.32	55.50	
	b-tag	Tight- $m_\tau$	3886	0.28	17.21	2.08	4.97	50.45	25.01	
		Loose- $m_\tau$	3657	0.10	5.58	1.65	5.65	57.51	29.52	
$\mu\tau_h$	No b-tag	Tight- $m_\tau$	124707	0.64	67.72	4.64	0.88	1.11	25.01	
		Loose- $m_\tau$	58534	0.38	34.62	5.77	1.96	2.57	54.70	
	b-tag	Tight- $m_\tau$	8778	0.28	19.33	0.98	4.80	50.11	24.50	
		Loose- $m_\tau$	7887	0.08	5.86	0.59	5.55	59.65	28.28	
$\tau_h\tau_h$	No b-tag		105545	0.45	14.94	2.08	0.28	0.23	82.02	
	b-tag		3416	0.52	20.80	1.03	2.65	18.75	56.25	

<sup>†</sup> Excluding events with jets misidentified as hadronic  $\tau$  lepton decays.

to fulfill the VeryLoose but not the Tight (Medium) working point of the discriminant in the  $e\tau_h/\mu\tau_h$  ( $\tau_h\tau_h$ ) final state. This region is primarily populated by events with jets misidentified as hadronic  $\tau$  lepton decays, with typical fractions of genuine  $\tau$  lepton decays at the level of a few percent or below. To arrive at an estimate for the number of events from background  $i$  due to  $\text{jet} \rightarrow \tau_h$  misidentification in the SR the number of events in the AR is then multiplied by the ratio

$$F_F^i = \frac{N_{\text{pass}}}{N_{\text{fail}}}, \quad (4)$$

where  $N_{\text{pass}}$  corresponds to the number of events that fulfill the Tight/Medium working point and  $N_{\text{fail}}$  to the number of events that fulfill the VeryLoose but not the Tight/Medium working point of the  $\tau_h$  identification discriminant. The number of events appearing in Eq. (4) is obtained from a dedicated determination region ( $\text{DR}_i$ ), which is orthogonal to the AR and SR, and dominated by background  $i$ . The contributions from backgrounds other than  $i$  are estimated from simulation and subtracted from the numerator and denominator of Eq. (4). For this purpose all corrections as described in Section 5.3 are applied to the simulation. The  $F_F^i$  can be different for different processes, for example, if the misidentified jet predominantly originates from a heavy flavor quark, a light flavor quark or gluon fragmentation.

The underlying assumption in this method is that the ratio of the number of events from background  $i$  in the SR to the number of events from the same background in the AR is equal to  $N_{\text{pass}}/N_{\text{fail}}$  in the  $\text{DR}_i$ . This can be ensured by determining  $F_F^i$  differentially as a function of several variables taking the most important potential kinematic dependencies into account. Residual biases can be removed by adequate corrections, which can be determined from control regions or from simulation. For the analysis the  $F_F^i$  are estimated from a fit to the measured values, as a function of the  $p_T$  of the  $\tau_h$  candidate in categories of the  $\tau_h$  decay mode, and the jet multiplicity, in bins of  $N_{\text{jet}} = 0$  or  $N_{\text{jet}} \geq 1$ . This is done independently in dedicated  $\text{DR}_i$  for the backgrounds due to QCD multijet,  $W$ +jets, and  $t\bar{t}$  events.

For the  $e\tau_h$  and  $\mu\tau_h$  final states  $\text{DR}_{\text{QCD}}$  is defined by the same selection as for the SR, but the electric charges of the  $\tau\tau$  pair are required to be of the same sign. To reduce the contamination from  $W$ +jets events the transverse mass is required to be  $m_T^{e(\mu)} < 40 \text{ GeV}$ , and the relative isolation requirement on the electron (muon) is changed to be  $0.05 < I_{\text{rel}}^{e(\mu)} < 0.15$  in both final states. The definition of  $\text{DR}_{W+\text{jets}}$  also uses the same selection as for the SR, but the requirement on the transverse mass is changed to  $m_T^{e(\mu)} > 70 \text{ GeV}$  to enrich this background and an additional requirement of the absence of  $b$  jets in the event is imposed to reduce the contamination from  $t\bar{t}$  events. In the  $e\tau_h$  and  $\mu\tau_h$  final states  $t\bar{t}$  production is a subdominant background with respect to  $W$ +jets and QCD multijet events. Since there is no sufficiently populated pure DR for  $t\bar{t}$  events covering a similar phase space as the SR, the  $F_F^{t\bar{t}}$  are estimated from simulation after the event selection and before the event categorization. Corrections to these factors are derived from a dedicated control region as described below.

In the  $\tau_h\tau_h$  final state QCD multijet production is by far the dominant background. Therefore only  $\text{DR}_{\text{QCD}}$  is defined from the single requirement that the electric charges of the  $\tau\tau$  pair should be of the same sign. The  $F_F^{\text{QCD}}$  are then also used to estimate the background from  $W$ +jets and  $t\bar{t}$  events. Misidentified  $\tau_h$  candidates in the  $\tau_h\tau_h$  final state usually originate from  $\text{jet} \rightarrow \tau_h$  misidentification. In this final state a combinatorial effect arises from the fact that each event enters the AR twice, either if the leading  $\tau_h$  candidate fulfills the nominal  $\tau_h$  identification requirement and the subleading  $\tau_h$  candidate the inverted requirement or vice versa. Both cases are used with a weight of 0.5 to take this combinatorial effect into account. For the backgrounds from  $W$ +jets and  $t\bar{t}$  events typically one of the reconstructed  $\tau_h$  candidates originates from

a genuine  $\tau$  lepton decay, while the other one is due to a misidentified jet. The fraction of events with two misidentified jets is at most a few percent and thus well below the associated systematic uncertainties. Since there are no significant combinatorial effects involved, these events are considered with a weight of 1.

From the individually determined  $F_F^i$  a weighted factor  $F_F$  is obtained on an event-by-event basis from

$$F_F = \sum_i w_i F_F^i, \quad w_i = \frac{N_{AR}^i}{\sum_j N_{AR}^j}, \quad i, j \in \{\text{QCD}, W+\text{jets}, t\bar{t}\}, \quad (5)$$

where  $N_{AR}^i$  corresponds to the expected number of events for background  $i$  in the AR. The factor  $F_F$  is then applied to all events in the AR to obtain an estimate for the number and shape of the sum of QCD multijet,  $W$ +jets, and  $t\bar{t}$  events due to  $\text{jet} \rightarrow \tau_h$  misidentification. For this purpose the subdominant contributions from  $Z \rightarrow \ell\ell$ , diboson and single  $t$  quark events are subsumed into the  $W$ +jets estimate. The estimates for  $N_{AR}^{W+\text{jets}}$  and  $N_{AR}^{t\bar{t}}$  are taken from the simulation. The estimate for  $N_{AR}^{\text{QCD}}$  is obtained from the events in the AR after subtracting all other backgrounds. These estimates are cross-checked using a template fit to the data in the AR equivalent to the fit described in Section 6, but with the  $p_T^{\tau_h}$  distribution as the input shape. From the resulting distributions, the expected contribution from events with genuine hadronic  $\tau$  lepton decays or electrons or muons misidentified as hadronic  $\tau$  lepton decays are subtracted using the simulation. The principle of the method is outlined in Fig. 4.

For each of the backgrounds considered, corrections to the estimated  $F_F^i$  are determined in modified determination regions to account for residual biases of the method. These may originate from the finite number of events or the functional form of the fit with which the  $F_F^i$  are derived, and from subdominant dependencies that are neglected during the determination of the  $F_F^i$ . These effects are checked and corrected for in the  $\text{DR}_i$  themselves by comparing the actual number of events with the  $\tau_h$  candidate matching the Tight/Medium working point of the  $\tau_h$  identification discriminant to the number of events estimated from the method. Residual corrections are determined as a function of the invariant mass of the visible decay products of the  $\tau\tau$  pair,  $m_{\text{vis}}$ , and found to be compatible with unity within the statistical precision. This demonstrates that the main dependencies of the  $F_F^i$  are taken into account.

For  $F_F^{\text{QCD}}$  two additional corrections are applied: in the  $e\tau_h$  ( $\mu\tau_h$ ) final state a correction is obtained as a function of  $I_{\text{rel}}^{e(\mu)}$  by comparing the number of events matching the Tight/Medium working point of the  $\tau_h$  identification discriminant to the number of events estimated from the method in a control region equivalent to  $\text{DR}_{\text{QCD}}$ , with the only difference being that the initial requirement on  $I_{\text{rel}}^{e(\mu)}$  is dropped. This correction is found to be  $\mathcal{O}(10\%)$  and compatible with unity within one standard deviation of the statistical precision. In the  $\tau_h\tau_h$  final state a correction is derived as a function of the  $p_T$  of the other  $\tau_h$  candidate. This correction is found to range between a few percent and 20%. For all final states another correction is derived to account for the transition from  $\text{DR}_{\text{QCD}}$  with the same charge requirement on the  $\tau\tau$  pair to the SR with an opposite charge requirement. This correction is determined as a function of  $m_{\text{vis}}$  in a control region with  $0.1 < I_{\text{rel}}^e < 0.2$  ( $0.15 < I_{\text{rel}}^\mu < 0.25$ ) in the  $e\tau_h$  ( $\mu\tau_h$ ) final state and in a control region where the other  $\tau_h$  candidate matches the VeryLoose but fails the Medium working point of the  $\tau_h$  identification discriminant in the  $\tau_h\tau_h$  final state. In all final states the correction is found to be compatible with unity within one standard deviation of the statistical precision, which ranges from 10 to 20% in the  $e\tau_h$  and  $\mu\tau_h$  final states, respectively, and from a few percent to 10% in the  $\tau_h\tau_h$  final state.

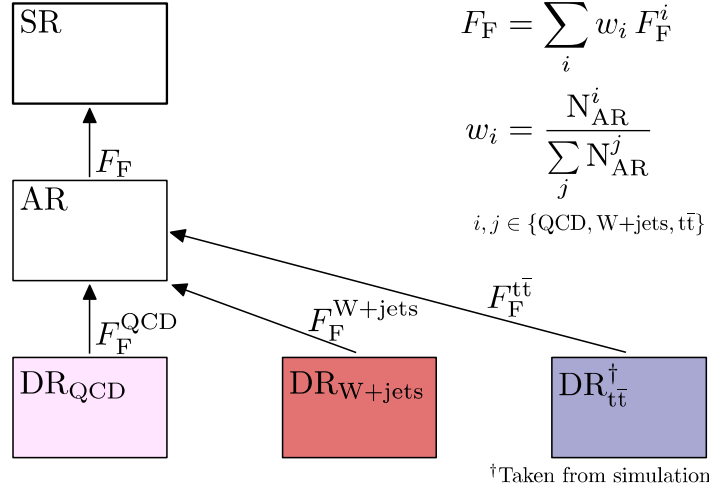


Figure 4: Schematic view of the determination and application of the  $F_F^i$  and  $F_F$  for the estimation of the background from QCD multijet, W+jets, and  $t\bar{t}$  events due to the misidentification of jets as hadronic  $\tau$  lepton decays. Note that  $DR_{t\bar{t}}^\dagger$  is taken from simulation.

In the  $e\tau_h$  ( $\mu\tau_h$ ) final state a correction is derived from simulation for  $F_F^{W+jets}$  as a function of  $m_T^{e(\mu)}$ . A residual dependency is expected from the selection requirements on  $p_T^{e(\mu)}$ : for low  $m_T^{e(\mu)}$  a value of  $p_T^{e(\mu)}$  above the thresholds of the offline selection will lead to a harder hadronic recoil and more jets in the event. This in turn may lead to less isolated  $\tau_h$  candidates especially at low  $p_T^{\tau_h}$ . The correction ranges from 10 to 30%, while usually compatible with unity within one standard deviation of the statistical precision. This correction is assumed to be the same for  $Z \rightarrow \ell\ell$ , diboson, and single  $t$  quark events.

The  $F_F^{t\bar{t}}$  are obtained from simulation and corrections are derived from a control region in data. This control region is characterized by the presence of at least one  $b$  jet and at least one lepton pair consisting of an isolated electron and an isolated muon in the event. Since this correction is found to be independent of  $p_T^{e(\mu)}$ ,  $m_T^{e(\mu)}$ , or  $m_{vis}$ , within the experimental precision, a common factor is used depending on the final state ( $e\tau_h$  or  $\mu\tau_h$ ) and the  $\tau_h$  decay mode. All corrections applied to the  $F_F^{QCD}$ ,  $F_F^{W+jets}$  and  $F_F^{t\bar{t}}$  are summarized in Table 4.

In the  $e\mu$  final state the background from QCD multijet events is estimated from an AR fulfilling the same selection requirements as the SR, however the charges of the leptons are required to be of the same sign. Extrapolation factors for the same charge to the opposite charge phase space are obtained in bins of the  $p_T$  of the two leptons and their separation in  $\Delta R$ . These extrapolation factors are derived in a DR region without event categorization, in which the isolation requirements on the leptons are chosen to be orthogonal to the SR. Finally, corrections are applied to account for the extrapolation into the exclusive event categories and for the extrapolation into the SR. The corrections for the extrapolation into the exclusive event categories are determined from the same DR, but inclusive in the  $p_T$  of, and separation between, the leptons. They are about 0.6 (1) for all  $b$ -tag (no  $b$ -tag) categories. The correction for the extrapolation into the SR is about 0.9 as determined from simulation.

Table 4: Corrections applied to the  $F_F^{\text{QCD}}$ ,  $F_F^{W+\text{jets}}$ , and  $F_F^{\text{t}\bar{\text{t}}}$  as described in the text. In the fourth column the source is indicated from which the correction is derived. The dependency  $p_T^{\tau_h}$  in the third line refers to the  $p_T$  of the  $\tau_h$  candidate that is assumed to originate from a genuine  $\tau$  lepton decay.

	Correction	Dependency	Source	$e\tau_h$	$\mu\tau_h$	$\tau_h\tau_h$
$F_F^{\text{QCD}}$	Nonclosure	$m_{\text{vis}}$	DR <sub>QCD</sub>	✓	✓	✓
	$I_{\text{rel}}^{e(\mu)}$ - , $p_T^{\tau_h}$ -dependent	$I_{\text{rel}}^{e(\mu)}$	DR <sub>QCD</sub> (w/o $I_{\text{rel}}^{e(\mu)}$ )	✓	✓	
		$p_T^{\tau_h\dagger}$	DR <sub>QCD</sub>			✓
	Opposite/Same charge	$m_{\text{vis}}$	Orthogonal iso./ID <sup>†</sup>	✓	✓	✓
$F_F^{W+\text{jets}}$	Nonclosure	$m_{\text{vis}}$	DR <sub>W+jets</sub>	✓	✓	
	$m_T^{e(\mu)}$ -dependent	$m_T^{e(\mu)}$	From simulation	✓	✓	
$F_F^{\text{t}\bar{\text{t}}}$	Nonclosure	$m_{\text{vis}}$	DR <sub>t\bar{t}</sub>	✓	✓	
	Data/Simulation	None	t\bar{t} enriched sideband	✓	✓	

<sup>†</sup> Refers to the  $\tau_h$  candidate that is assumed to originate from a genuine  $\tau$  lepton decay.

### 5.3 Backgrounds estimated from simulation

All other backgrounds, apart from the ones described in Section 5.2 are estimated from simulation. Corrections are derived to account for residual differences in the efficiency of the selected trigger paths, in the electron and muon tracking efficiency, and in the efficiency of the identification and isolation requirements, for electrons and muons. These corrections are obtained using the “tag-and-probe” method, as described in Ref. [66], with  $Z \rightarrow ee$  and  $Z \rightarrow \mu\mu$  events in bins of  $p_T$  and  $\eta$  of the probed electron or muon. They are usually not larger than a few percent. In a similar way, corrections are obtained for the efficiency of triggering on the hadronic  $\tau$  lepton decays in the  $\tau_h\tau_h$  final state and for the  $\tau_h$  identification efficiency. In this case the tag-and-probe method is applied to  $Z \rightarrow \tau\tau$  events in the  $\mu\tau_h$  final state.

The energy of a jet is corrected to the expected response of the jet at stable-hadron level, using corrections measured in bins of the jet  $p_T$  and  $\eta$ . These corrections are usually not larger than 10 to 15%. Residual data-to-simulation corrections are applied to the simulated samples. They usually range between sub-percent level at high jet  $p_T$  in the central part of the detector to a few percent in the forward region. A correction is applied to the direction and magnitude of the  $\vec{p}_T^{\text{miss}}$  vector based on the differences between the estimates of the hadronic recoil in  $Z \rightarrow \mu\mu$  events in data and simulation. This correction is applied to  $Z \rightarrow \tau\tau$ ,  $W+\text{jets}$ , and signal events, where a well-defined direction and magnitude of genuine  $\vec{p}_T^{\text{miss}}$  can be defined. The efficiency for genuine and misidentified b jets to pass the Medium working point of the b tagging discriminator is determined from data, using t\bar{t} events for genuine b jets and  $Z+\text{jets}$  events for jets predominantly originating from light-flavor quarks. Data-to-simulation corrections are obtained for these efficiencies and used to correct the number of b jets in the simulation, which translates into the number of events in the global b-tag and no b-tag event categories. In the  $e\mu$  final state data-to-simulation corrections are derived for the rate at which jets are misidentified as an electron or muon. These are determined as a function of the jet  $p_T$  from  $Z+\text{jets}$  events in the  $Z \rightarrow \ell\ell$  decay. They are applied to  $W+\text{jets}$  and diboson events, which form more than 90% of the expected background due to jet  $\rightarrow \ell$  misidentification in the  $e\mu$  final state, and where the flavor composition of jets is similar to that in the region in which the corrections are determined. Corrections are further applied to  $Z \rightarrow \mu\mu$  events in the  $\mu\tau_h$  and  $\tau_h\tau_h$  final states in which a muon is reconstructed as a  $\tau_h$  candidate and in  $Z \rightarrow ee$  events in the  $e\tau_h$  and  $\tau_h\tau_h$



final states in which an electron is reconstructed as a  $\tau_h$  candidate, to account for residual differences in the  $\ell \rightarrow \tau_h$  misidentification rate between data and simulation. Finally a correction, obtained from  $Z \rightarrow ee$  events, to the energy scale for electrons misidentified as hadronic  $\tau$  lepton decays is applied. Corresponding uncertainties in all these corrections are incorporated into the uncertainty model discussed in Section 7.

Deficiencies in the modeling of Drell–Yan events in the  $ee$ ,  $\mu\mu$  and  $\tau\tau$  final states are corrected for by a weighting of the simulated  $Z \rightarrow \mu\mu$  events to data in bins of  $p_T(\mu\mu)$  and  $m(\mu\mu)$ . The weights obtained are applied to the simulated events in all leptonic final states. For the statistical analysis the overall normalization of the background from  $Z \rightarrow \tau\tau$  events is furthermore constrained by dedicated control regions of  $Z \rightarrow \mu\mu$  events in each global event category, making use of the equal branching fractions for the  $Z$  boson decays into  $\tau$  leptons or muons, in the context of lepton universality. Theoretical uncertainties arising from residual kinematic differences between the selected dimuon and  $\tau\tau$  final states are incorporated into the uncertainty model. In addition all simulated  $t\bar{t}$  events are weighted to better match the top quark  $p_T$  distribution, as observed in data [67]. For the statistical analysis the overall normalization of this background is also constrained by a dedicated control region with an isolated electron, an isolated muon, and large  $p_T^{\text{miss}}$  in the final state, which is chosen to be orthogonal to the SR in the  $e\mu$  final state; this sample has a  $t\bar{t}$  purity of 85%. All control regions used for the statistical analysis are outlined in Fig. 3.

## 5.4 Cross-checks of background estimations

Two cross-checks are performed to give confidence in the background estimation. In a first cross-check all backgrounds apart from QCD multijet production and the normalization for  $W$ +jets events are taken from simulation. For this purpose all corrections as summarized in Section 5.3 are applied to all simulated events. This cross-check is performed in the  $e\tau_h$  and  $\mu\tau_h$  final states individually.

The  $W$ +jets prediction, prior to the statistical inference of the signal, is obtained by subtracting the small contribution of all other backgrounds except for QCD multijet and  $W$ +jets events from data in corresponding control regions requiring the charges of the  $\tau\tau$  pair to be of opposite (OS) or same sign (SS) and  $m_T^{e(\mu)} > 70$  GeV. An estimate for the normalization of the  $W$ +jets events can then be obtained from the following system of linear equations

$$\begin{aligned} N_{\text{data}}^{\text{SS}} &= N_{\text{QCD}}^{\text{SS}} + N_{W+\text{jets}}^{\text{SS}} \\ N_{\text{data}}^{\text{OS}} &= f_{\text{QCD}}^{\text{OS/SS}} N_{\text{QCD}}^{\text{SS}} + f_{W+\text{jets}}^{\text{OS/SS}} N_{W+\text{jets}}^{\text{SS}}, \end{aligned} \quad (6)$$

where  $N_{\text{data}}^{\text{SS(OS)}}$  corresponds to the number of events in the control regions, after subtracting the expected number of events for all other backgrounds, and  $f_{\text{QCD}(W+\text{jets})}^{\text{OS/SS}}$  is the expected OS to SS ratio for  $W$ +jets and QCD multijet events. For this estimate  $f_{W+\text{jets}}^{\text{OS/SS}}$  is obtained from the simulation and  $f_{\text{QCD}}^{\text{OS/SS}}$  from another control region with inverted isolation requirements on the electron or muon, as described below. An estimate for  $N_{W+\text{jets}}^{\text{SS}}$  can then be obtained from Eq. (6). From this the number of  $W$ +jets events in the SR can be inferred via  $f_{W+\text{jets}}^{\text{OS/SS}}$  and another extrapolation factor from the control region into the SR, which again is taken from simulation. To stay as close as possible to the kinematic regime in the signal regions an OS and an SS control region for the determination of  $N_{\text{data}}^{\text{OS}}$  and  $N_{\text{data}}^{\text{SS}}$  is defined, for each event subcategory in the  $e\tau_h$  and  $\mu\tau_h$  final states, as described in Section 4, amounting to eight control regions per final state. The shape of the final discriminating variable distribution used for the signal extraction is taken from simulation.

The shape and normalization of the QCD multijet background distributions prior to the signal extraction are obtained from control regions equivalent to the signal regions with the exception of a SS instead of an OS requirement on the charge of the selected  $\tau\tau$  pair. From the events in this control region all other expected backgrounds are subtracted using the normalization and shape information for the final discriminating variable distribution from simulation, with the exception of the normalization of W+jets events, which is obtained as described above. The extrapolation factors ( $f_{\text{QCD}}^{\text{OS/SS}}$ ) from the SS to OS selection are obtained from control regions, where in addition, to the corresponding charge requirement, the isolation requirement on the electron or muon is inverted. The extrapolation factors are then obtained from a fit to the data in the control regions similar to the one described in Section 6. To control the normalization of the W+jets and QCD multijet events the eight additional control regions per final state, as introduced above, are added to the fit for the signal extraction and the corresponding normalization uncertainties are incorporated into a modified uncertainty model.

In a second cross-check the background from  $Z \rightarrow \tau\tau$  events in the main analysis is replaced with the prediction obtained from the  $\mu \rightarrow \tau$  *embedding* method as used during the LHC Run-1 analyses and described, for example, in Refs. [68, 69]. In this process  $Z \rightarrow \mu\mu$  events are selected in data. The muons are then replaced by simulated  $\tau$  lepton decays with the same kinematics as the reconstructed muons. In this way the method relies only on the simulation of the well understood  $\tau$  lepton decay while all other parts of the event are obtained from data. As a consequence several data-to-simulation corrections as described in Section 5.3, which are of particular importance for the event categorization as well as for the shape of the final discriminating variable distribution, do not need to be applied for this process. This applies, for example, to corrections of the jet energy scale, b tagging efficiency, and  $\vec{p}_T^{\text{miss}}$ . This cross-check is applied in the  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states individually. Both the extrapolation factors from the inclusive event selection into the event subcategories, as well as the shapes of the final discriminating variable distribution for the signal extraction, as obtained from the simulation, are found to be in good agreement with the estimates as obtained from the embedding method, within the estimated uncertainties. In addition the uncertainties that are related to the experimental aspects of the  $\mu \rightarrow \tau$  embedding, which are orthogonal to the uncertainties in the estimate from simulation, are incorporated into a modified uncertainty model to replace several uncertainties for the estimate based on the simulation.

## 6 Statistical inference for the signal

The final discriminating variable used to search for a signal is the total transverse mass,  $m_T^{\text{tot}}$  [28], defined as

$$m_T^{\text{tot}} = \sqrt{m_T^2(p_T^{\tau_1}, p_T^{\tau_2}) + m_T^2(p_T^{\tau_1}, p_T^{\text{miss}}) + m_T^2(p_T^{\tau_2}, p_T^{\text{miss}})}, \quad (7)$$

where the pair  $(\tau_1, \tau_2)$  can be  $(e, \mu)$ ,  $(e, \tau_h)$ ,  $(\mu, \tau_h)$ , or  $(\tau_h, \tau_h')$ , and the transverse mass,  $m_T$ , between two objects with transverse momenta  $p_T$  and  $p'_T$ , and relative difference  $\Delta\phi$  in the azimuthal angle is given by:

$$m_T = \sqrt{2 p_T p'_T [1 - \cos(\Delta\phi)]}. \quad (8)$$

The input distributions to the statistical inference of the signal in a subset of the most sensitive event subcategories per final state are shown in Figs. 5 and 6. The expected  $m_T^{\text{tot}}$  distribution is represented by the stack of filled histograms in the upper panel of each subfigure, where each filled histogram corresponds to the estimated template distribution of the given SM process that has been taken into account for the analysis. For this purpose the fractions of QCD

multijet,  $W$ +jets, and  $t\bar{t}$  events contributing to the event selection by  $\text{jet} \rightarrow \tau_h$  misidentification are subsumed into one single contribution labeled as “ $\text{jet} \rightarrow \tau_h$ ”. The remaining fractions from  $W$ +jets, single  $t$  quark, and diboson events are subsumed into one single contribution labeled as “Electroweak”. The shaded band associated with the sum of filled histograms corresponds to the combination of all normalization and shape altering uncertainties in all background processes, taking into account all correlations as obtained from the fit used for the signal extraction. The ratio of the data points to the expectation from the sum of all filled histograms is shown in the lower panel of each subfigure; the statistical uncertainty in the data is represented by the error bars and the uncertainty in all background processes by the shaded band. The expected  $m_T^{\text{tot}}$  distribution for a signal of three neutral Higgs bosons from gluon fusion and the production in association with  $b$  quarks in the MSSM  $m_h^{\text{mod+}}$  scenario, discussed in Ref. [16], for  $m_A = 700$  GeV and  $\tan\beta = 20$  is also shown. The signal distribution reveals two peaking structures, related to the signal from the  $h$  at about 130 GeV, and the nearly mass degenerate  $H$  and  $A$  at 700 GeV.

To quantify the amount of signal a simultaneous binned maximum likelihood fit to the  $m_T^{\text{tot}}$  distributions in all event subcategories and all final states is performed. This is done under the background-only and several signal-plus-background hypotheses to search for potential excesses due to the presence of additional Higgs bosons over the known SM processes. For this purpose the SM Higgs boson is included in the background processes. The control regions, which have been designed to constrain the background from Drell–Yan and  $t\bar{t}$  events, are included in the likelihood model, resulting in a fit in sixteen event subcategories and three control regions, as outlined in Fig. 3.

The data are interpreted in two ways based on the ratio of the fitted likelihoods for the background-only and the tested signal-plus-background hypotheses. For each interpretation the model for the background processes is formed from the template distributions as shown, for example, in Figs. 5 and 6. In a first interpretation, which is meant to be as model-independent as possible, the signal model corresponds to a single resonance,  $\phi$ , with a width negligible compared to the experimental resolution. For this purpose, 28 simulated single narrow resonances with mass  $m_\phi$  between 90 GeV and 3.2 TeV in the gluon fusion and in association with  $b$  quarks are used. For both production modes the  $p_T$  spectrum of the  $\phi$  is simulated at NLO precision as described in Section 5.1. The signal is searched for in both production modes at the same time, using two freely varying parameters of interest for the fit to the data, one for each production mode. In a second interpretation, the simulated mass points are combined into the multi-resonance signal structure expected from each of the tested MSSM benchmark scenarios. This is done using the model predictions, as described in Sections 5.1 and 8, and a linear template morphing algorithm, as described in Ref. [70], to move the simulated mass points to their exact predicted values.

## 7 Systematic uncertainties

The uncertainty model comprises theoretical uncertainties, experimental uncertainties, and uncertainties due to the limited population of the template distributions used for the prediction of the background processes. The last group of uncertainties are most important for the high-mass Higgs boson searches. All systematic uncertainties are implemented in the form of nuisance parameters in the likelihood, which can be further constrained by the fit to the data. The following uncertainties are implemented as normalization uncertainties that leave the shape of the  $m_T^{\text{tot}}$  distributions unchanged:

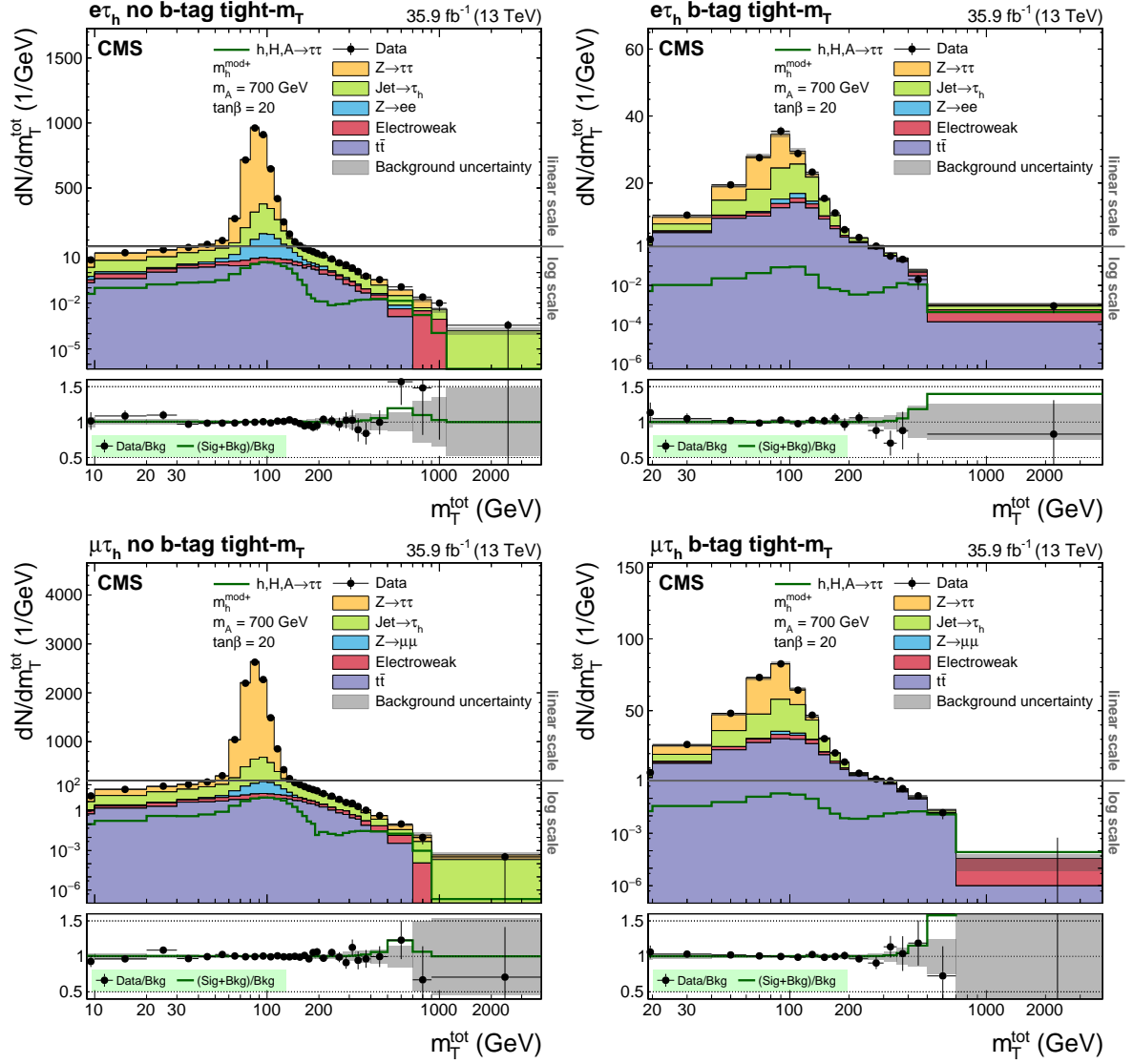


Figure 5: Distribution of  $m_T^{\text{tot}}$  in the global no b-tag (left) and b-tag (right) categories in the  $e\tau_h$  (upper row) and  $\mu\tau_h$  (lower row) final states. In all cases the most sensitive tight- $m_T$  event subcategory is shown. The gray horizontal line in the upper panel of each subfigure indicates the change from logarithmic to linear scale on the vertical axis.

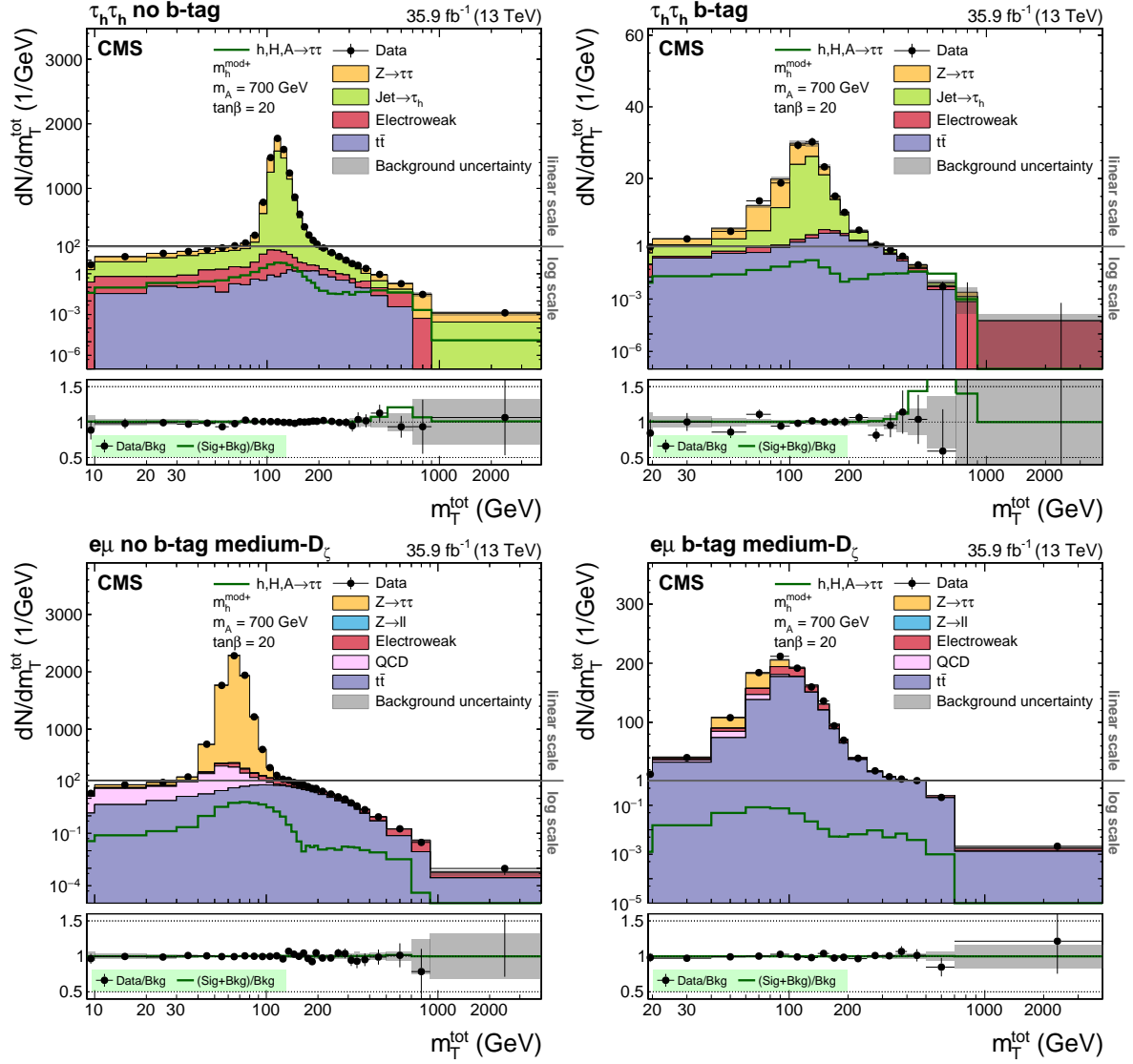


Figure 6: Distribution of  $m_T^{\text{tot}}$  in the global no b-tag (left) and b-tag (right) categories in the  $\tau_h \tau_h$  (upper row) and  $e\mu$  (lower row) final states. For the  $e\mu$  final state the most sensitive medium- $D_\zeta$  event subcategory is shown. The gray horizontal line in the upper panel of each subfigure indicates the change from logarithmic to linear scale on the vertical axis.

- The uncertainty in the integrated luminosity measurement is 2.5% [71]. It is applied to all processes that have been estimated from simulation.
- The uncertainties in the measurement of the identification, isolation, and trigger efficiencies are found to amount to 2% both for electrons and muons, adding all individual contributions in quadrature. These uncertainties are applied to all processes that are estimated from simulation.
- Uncertainties in the measurement of the probability of electrons ( $e \rightarrow \tau_h$ ) and muons ( $\mu \rightarrow \tau_h$ ) to be misidentified as hadronic  $\tau$  lepton decays are applied to the fraction of simulated Drell–Yan events with light leptons being misidentified as hadronic  $\tau$  lepton decays in the  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states. The uncertainty in the  $e \rightarrow \tau_h$  misidentification probability amounts to 11 (3)% in the  $e\tau_h$  ( $\tau_h\tau_h$ ) final state. The uncertainty in the  $\mu \rightarrow \tau_h$  misidentification probability is 12 (5)% in the  $\mu\tau_h$  ( $\tau_h\tau_h$ ) final state.
- The uncertainty in the  $\tau_h$  identification efficiency is found to be 5% per  $\tau_h$  candidate. It is factorized into a 4 (8)% part that is correlated and a 3 (6)% part that is uncorrelated across all final states containing hadronic  $\tau$  lepton decays in the  $e\tau_h$  and  $\mu\tau_h$  ( $\tau_h\tau_h$ ) final states. A 7% uncertainty in the  $\tau_h$  trigger efficiency measurement is added to the uncorrelated part in the  $\tau_h\tau_h$  final state. The uncertainties related to the  $\tau_h$  reconstruction and identification are applied to all processes that have been estimated from simulation and that contain genuine hadronic  $\tau$  lepton decays.
- The uncertainty in the jet energy scale affects the number of events entering each category. It is applied to all processes estimated from simulation and ranges from 1 to 6%, depending on the final state and subcategory. Similarly, uncertainties in the rate with which both light-flavor jets and genuine  $b$  jets pass the  $b$  tagging discriminator selection are applied to all processes estimated from simulation. These uncertainties range from 1 to 5%.
- Uncertainties in the resolution and response of the  $p_T^{\text{miss}}$  are derived as part of the determination of the recoil corrections. This leads to uncertainties ranging from 1 to 5% that are incorporated for all processes estimated from simulation and to which recoil corrections are applied. These are all signal processes, Drell–Yan production and  $W$ +jets events. For the single  $t$  quark, diboson and  $t\bar{t}$  backgrounds, which do not have recoil corrections applied, jet energy scale and unclustered energy scale variations are propagated to the  $p_T^{\text{miss}}$ , also leading to uncertainties ranging from 1 to 5%.
- The uncertainty in the background yield from single  $t$  quark and diboson production amounts to 5%, based on CMS measurements [72, 73]. In the  $e\mu$  final state, where the  $W$ +jets contribution is taken from simulation, the theoretical uncertainty in the cross section calculation is 4%. Due to the inclusion of the  $Z \rightarrow \mu\mu$  and the  $t\bar{t}$  control regions in the model for the statistical inference of the signal, which control the Drell–Yan and  $t\bar{t}$  normalization respectively, no theoretical cross section uncertainties are applied for these processes. However, uncertainties are applied to the  $Z \rightarrow \tau\tau$ ,  $Z \rightarrow \ell\ell$ , and  $t\bar{t}$  processes in all signal categories to account for the extrapolation from the control region to the signal regions. The  $Z \rightarrow \tau\tau$  extrapolation uncertainties range from 1 to 7%. The extrapolation uncertainties for  $Z \rightarrow \ell\ell$  events are 4%. The extrapolation uncertainties from the  $t\bar{t}$  control region to the signal regions are found to be below 1%. An additional uncertainty of 1% is, however, applied in the  $t\bar{t}$  control region to account for fluctuations in the variables used to select the events in this control region. The uncertainty in the estimation of the backgrounds

in the  $DR_i$ , which are taken from the simulation and subtracted from the data, for the determination of the  $F_F^i$  amounts to 3 (4)% in the  $e\tau_h$  and  $\mu\tau_h$  ( $\tau_h\tau_h$ ) final states.

- Since the background from QCD multijet events in the  $e\mu$  final state is determined from a control region, uncertainties that account for the statistical uncertainty in the data and the subtracted backgrounds in this control region are applied. In addition, this background is subject to uncertainties related to the extrapolation from the control region to the signal regions. An overall 30% extrapolation uncertainty is applied, in addition to category-dependent uncertainties ranging from 4 to 29%, in the measurement of the OS to SS transfer factor.
- Theoretical uncertainties in the acceptance of signal events in the associated production with b quarks are obtained from variations of the renormalization ( $\mu_r$ ) and factorization ( $\mu_f$ ) scales and of the internal generator matching scale  $Q_{sh}$  related to parton showering. The scales  $\mu_r$  and  $\mu_f$  are varied by factors of 0.5 and 2. The scale uncertainty is obtained from the envelope of the six variations of  $\mu_r$  and  $\mu_f$ , as recommended in Ref. [74]. Depending on the tested mass it ranges between  $-4\%$  (for 90 GeV),  $-0.4\%$  (for 500 GeV), and  $-2.5\%$  (for 3.2 TeV) in the b-tag categories, and  $0.8\%$  (for 90 GeV),  $0.3\%$  (for 500 GeV), and  $2.0\%$  (for 3.2 TeV) in the no b-tag categories. The scale  $Q_{sh}$  is varied by factors of  $1/\sqrt{2}$  and  $\sqrt{2}$ . The resulting uncertainty ranges between  $-13.2\%$  (for 90 GeV),  $-4.6\%$  (for 500 GeV), and  $-1.8\%$  (for 3.2 TeV) in the b-tag categories, and  $2.6\%$  (for 90 GeV),  $2.9\%$  (for 500 GeV), and  $1.4\%$  (for 3.2 TeV) in the no b-tag categories. The uncertainty from the variation of  $\mu_r$  and  $\mu_f$ , and the uncertainty from the variation of  $Q_{sh}$  are added linearly, following the recommendation in Ref. [74].
- For the parameter scan in the model interpretations, theoretical uncertainties due to the different choices of the factorization and renormalization scales in the signal predictions are included. The MSTW2008 [75] PDFs are used for the calculation of the production cross sections. The uncertainties in the choice for the PDFs are calculated following the recommended prescription given in Refs. [75, 76]. The uncertainties are evaluated separately for each  $m_A$ - $\tan\beta$  point. They vary between 15 and 25%.
- For all results shown in the following the SM Higgs boson production is taken into account in the likelihood ratio. Uncertainties due to different choices of the renormalization and factorization scales for the calculation of the production cross section of the SM Higgs boson amount to 3.9% for gluon fusion, 0.4% for VBF, 2.8% for ZH, and 0.5% for WH production. Uncertainties due to different choices for the PDFs and  $\alpha_s$  amount to 3.2% for gluon fusion, 2.1% for VBF, 1.6% for ZH, and 1.9% for WH production. The procedure for deriving these uncertainties is further described in Ref. [74].

The following systematic uncertainties allow correlated changes across bins that alter the shape of the  $m_T^{\text{tot}}$  input distributions, and are referred to as shape uncertainties hereafter:

- In the  $e\mu$  final state, shape uncertainties are applied to all processes with jets misidentified as electrons or muons to account for the uncertainties in the  $\text{jet} \rightarrow e$  and  $\text{jet} \rightarrow \mu$  misidentification probability. The size of these uncertainties depends on the jet  $p_T$ , with a minimum uncertainty of 13 (10)% for electrons (muons).
- Three independent uncertainties are applied on the energy scale for genuine  $\tau$  leptons decaying hadronically; for the decay into a single charged hadron with and without neutral pions and the decay into three charged hadrons. Each uncertainty is 1.2%. They affect both the normalization and the shape of the  $m_T^{\text{tot}}$  distribution for

the signal, the  $Z \rightarrow \tau\tau$ ,  $t\bar{t}$  and diboson backgrounds containing genuine  $\tau$  leptons in the  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states.

- An asymmetric uncertainty of  $+5\% \times p_T[\text{TeV}]$  and  $-35\% \times p_T[\text{TeV}]$  is applied to account for the extrapolation in the  $\tau_h$  identification efficiency estimate, which is mostly determined by low- $p_T$  hadronic  $\tau$  lepton decays close to the  $Z$  boson peak, to higher- $p_T$  regimes of the  $\tau$  leptons that are particularly relevant for the high-mass signal hypotheses. The  $p_T$  of the  $\tau_h$  candidate is scaled by the corresponding factor. This uncertainty is applied to the signal, the  $Z \rightarrow \tau\tau$ ,  $t\bar{t}$ , and diboson backgrounds containing genuine  $\tau$  leptons in the  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states.
- In the  $e\tau_h$  final state, an uncertainty in the energy scale of electrons misidentified as hadronic  $\tau$  lepton decays is applied, split into a 1 (0.5)% uncertainty in the correction for the decay mode with one charged hadron with (without) neutral pions. This uncertainty is only applied to the  $Z \rightarrow ee$  process where one of the electrons is misidentified as a hadronic  $\tau$  lepton decay.
- In the  $e\mu$  final state, an uncertainty in the electron energy scale is applied that amounts to 1% in the barrel and 2.5% in the endcaps. In the  $e\tau_h$  final state this uncertainty is covered by the uncertainty in the energy scale of the  $\tau_h$  candidate.
- An uncertainty in the correction of the  $p_T$  of the top quarks in simulated  $t\bar{t}$  events is applied that corresponds to 100% of the correction as discussed in Section 5.3. It affects this background in all signal regions and in the  $t\bar{t}$  control region. It is further constrained by the  $t\bar{t}$  control region described in Section 4.
- Five uncertainties are included to cover the uncertainty in the reweighting method used to improve the simulation of Drell–Yan events as described in Section 5.3. These uncertainties include the propagation of the 0.2% muon energy scale uncertainty to the derived weights and the propagation of a 6%  $t\bar{t}$  cross section uncertainty, which affects the simulated  $t\bar{t}$  background that needs to be subtracted in the  $Z \rightarrow \mu\mu$  selection. Since the reweighting is obtained prior to the statistical inference for the signal this is not coupled to the  $t\bar{t}$  control region. In addition, the statistical uncertainties in the measured weights are found to be nonnegligible in three of the bins used to derive the correction, which leads to three additional shape uncertainties related to the reweighting procedure.

In the  $\mu\tau_h$ ,  $e\tau_h$ , and  $\tau_h\tau_h$  final states, the following shape uncertainties related to the fake factor method are applied to those background components that are estimated by this method:

- Statistical uncertainties in the estimate of the  $F_F^i$  in the  $DR_i$  are obtained from the uncertainties of the fit used to parametrize the  $F_F^i$ . They amount to 4% in the  $\mu\tau_h$  final state and range between 4 and 7% (2 and 3%) in the  $e\tau_h$  ( $\tau_h\tau_h$ ) final states.
- In the  $e\tau_h$  and  $\mu\tau_h$  final states, uncertainties are taken into account in the corrections due to the finite number of events or omitted dependencies during the determination of the  $F_F^i$ . This is done for all backgrounds considered. Additional uncertainties are taken into account in all process specific corrections that are applied to the  $F_F^i$ . For  $F_F^{\text{QCD}}$  these are the correction of the extrapolation from the SS to the OS region and the correction as a function of the lepton isolation. For  $F_F^{W+\text{jets}}$  this is the correction as a function of  $m_T^{e(\mu)}$ . For  $F_F^{t\bar{t}}$  this is the data-to-simulation correction in the dedicated control region. All these uncertainties are added in quadrature for each corresponding background and vary between 7 and 10% and between 5 and 7% in the  $e\tau_h$  and  $\mu\tau_h$  final states respectively.



- In the  $\tau_h \tau_h$  final state, uncertainties are taken into account in the corrections due to the finite number of events or omitted dependencies during the determination of the  $F_F^i$ . Additional uncertainties in the correction of the SS to OS extrapolation as a function of the  $p_T$  of the other  $\tau_h$  candidate, in the estimate of the fractions of  $W$ +jets, Drell–Yan, and  $t\bar{t}$  events with one jet misidentified as a hadronic  $\tau$  lepton decay, and in the use of  $F_F^{QCD}$  for the estimation of the  $W$ +jets and  $t\bar{t}$  contributions to the total  $\text{jet} \rightarrow \tau_h$  background are taken into account. When added in quadrature, these additional uncertainties are of the order of 10%.

The shape uncertainties related to the fake factor method are factorized into a pure shape and pure normalization part. The normalization terms of the statistical uncertainties are added in quadrature for each individual category in each final state and applied as normalization uncertainties.

In addition, uncertainties due to the limited population of the template distributions used for the prediction of the background processes are taken into account by allowing each bin of each background template to vary within its statistical uncertainty. These uncertainties are uncorrelated across the bins of the input distributions. An overview of all uncertainties that have been taken into account in the likelihood model used for the statistical analysis is given in Table 5.

## 8 Results

The complete model, to extract the signal, results in a likelihood function of the form

$$\mathcal{L}(\{k_i\} | \mu s(\theta) + b(\theta)) = \prod_i \mathcal{P}(k_i | \mu s_i(\theta) + b_i(\theta)) \prod_j \mathcal{C}(\hat{\theta}_j | \theta_j), \quad (9)$$

where  $i$  labels all bins of the input distributions with event numbers  $k_i$  in all event subcategories and control regions and  $j$  all nuisance parameters. The term  $\theta_j$  corresponds to a given nuisance parameter,  $\mu$  to a scaling parameter for a given signal  $s_i$ , and  $b_i$  to the prediction of all backgrounds in bin  $i$ . The function  $\mathcal{P}(k_i | \mu s_i(\theta) + b_i(\theta))$  corresponds to a Poisson distribution,  $\mathcal{C}(\hat{\theta}_j | \theta_j)$  to the probability density function used to implement the uncertainty related to the nuisance parameter  $\theta_j$ , and  $\hat{\theta}_j$  to the estimate for  $\theta_j$  from the fit to the data. All distributions shown in Figs. 5 and 6 are after an MSSM  $m_h^{\text{mod+}}$  signal-plus-background hypothesis, corresponding to  $m_A = 700 \text{ GeV}$  and  $\tan \beta = 20$ , has been fitted to the data. No signal is observed in the investigated mass range between 90 GeV and 3.2 TeV and upper limits on the presence of a signal are set in the two interpretations of the data as discussed in Section 6. This is done following the modified frequentist approach as described in Refs. [77, 78], using the same definition of the test statistic as in the search for the SM Higgs boson [79, 80]:

$$q_\mu = -2 \ln \left( \frac{\mathcal{L}(\{k_i\} | \mu s(\hat{\theta}_\mu) + b(\hat{\theta}_\mu))}{\mathcal{L}(\{k_i\} | \hat{\mu} s(\hat{\theta}_{\hat{\mu}}) + b(\hat{\theta}_{\hat{\mu}}))} \right), \quad 0 \leq \hat{\mu} \leq \mu, \quad (10)$$

where the hat again indicates the estimate from the fit to the data and the index of  $q_\mu$  indicates that the fit to the data has been performed for a fixed value of  $\mu$ . In the large number limit the distribution of  $q_\mu$  can be approximated by analytic functions, from which the median and the uncertainty contours can be obtained as described in Ref. [81].

In the first interpretation of the data 95% confidence level (CL) upper limits are set on the product of the branching fraction for the decay into  $\tau$  leptons and the cross section for the

Table 5: Overview of the systematic uncertainties used in the likelihood model for the statistical inference of the signal. The label “MC” refers to all processes that are obtained from simulation, the label “ $F_F$ ” refers to all backgrounds that are obtained from the fake factor method. Values in parentheses correspond to additional uncertainties correlated across final states or event categories. Detailed descriptions are given in Section 7.

Uncertainty	$e\mu$	$e\tau_h$	$\mu\tau_h$	$\tau_h\tau_h$	Process	Shape	Variation
Integrated luminosity	✓	✓	✓	✓	MC	—	2.5%
Jet $\rightarrow$ e mis-ID	✓	—	—	—	MC	✓	13%
Jet $\rightarrow$ $\mu$ mis-ID	✓	—	—	—	MC	✓	10%
$e/\mu$ -trigger, ID, isolation	✓	✓	—	—	MC	—	2%
	✓	—	✓	—	MC	—	2%
$e \rightarrow \tau_h$ mis-ID	—	✓	—	—	$Z \rightarrow ee$	—	11%
	—	—	—	✓	$Z \rightarrow ee$	—	3%
$\mu \rightarrow \tau_h$ mis-ID	—	—	✓	—	$Z \rightarrow \mu\mu$	—	12%
	—	—	—	✓	$Z \rightarrow \mu\mu$	—	5%
$\tau_h$ -trigger	—	—	—	✓	MC	—	7%
$\tau_h$ -ID	—	✓	✓	—	MC	—	3 (4)%
	—	—	—	✓	MC	—	6 (8)%
$\tau_h$ -ID (high $p_T$ )	—	✓	✓	✓	MC	✓	$p_T$ dep.
$\tau_h$ energy scale	—	✓	✓	✓	MC	✓	1.2%
$e \rightarrow \tau_h$ energy scale	—	✓	—	—	$Z \rightarrow ee$	✓	0.5–1.0%
e energy scale	✓	—	—	—	MC	✓	1.0–2.5%
Jet energy scale	✓	✓	✓	✓	MC	—	1–6%
b tagging	✓	✓	✓	✓	MC	—	1–5%
$p_T^{\text{miss}}$ resp./res.	✓	✓	✓	✓	MC	—	1–5%
	✓	✓	✓	✓	Diboson	—	5%
Bkgr. in signal categories	✓	✓	✓	✓	Single t	—	5%
	✓	—	—	—	W+jets	—	4%
	✓	✓	✓	✓	$Z \rightarrow \tau\tau$	—	1–7%
	✓	✓	✓	✓	$Z \rightarrow \ell\ell$	—	4%
Sideband extrapolation	✓	✓	✓	✓	$t\bar{t}$	—	1%
	✓	—	—	—	QCD	—	4–29 (30)%
	✓	✓	✓	✓	$t\bar{t}$	✓	100%
Top quark $p_T$ reweighting	✓	✓	✓	✓	$Z \rightarrow \tau\tau, \ell\ell$	✓	See text
Z reweighting of LO MC	—	✓	✓	—	MC	—	3%
	—	—	—	✓	MC	—	4%
Bkgr. in $DR_{\text{QCD/W+jets}}$	—	✓	—	—	$F_F$	✓	4–7%
	—	—	✓	—	$F_F$	✓	4%
	—	—	—	✓	$F_F$	✓	2–3%
	—	✓	—	—	$F_F$	✓	7–10%
$F_F^i$ corrections	—	—	✓	—	$F_F$	✓	5–7%
	—	—	—	✓	$F_F$	✓	10%
b-associated signal acceptance	✓	✓	✓	✓	Signal	—	3.2–16.5%
PDF/scale	✓	✓	✓	✓	Signal	—	15–25%
	✓	✓	✓	✓	SM Higgs	—	0.5–3.2%

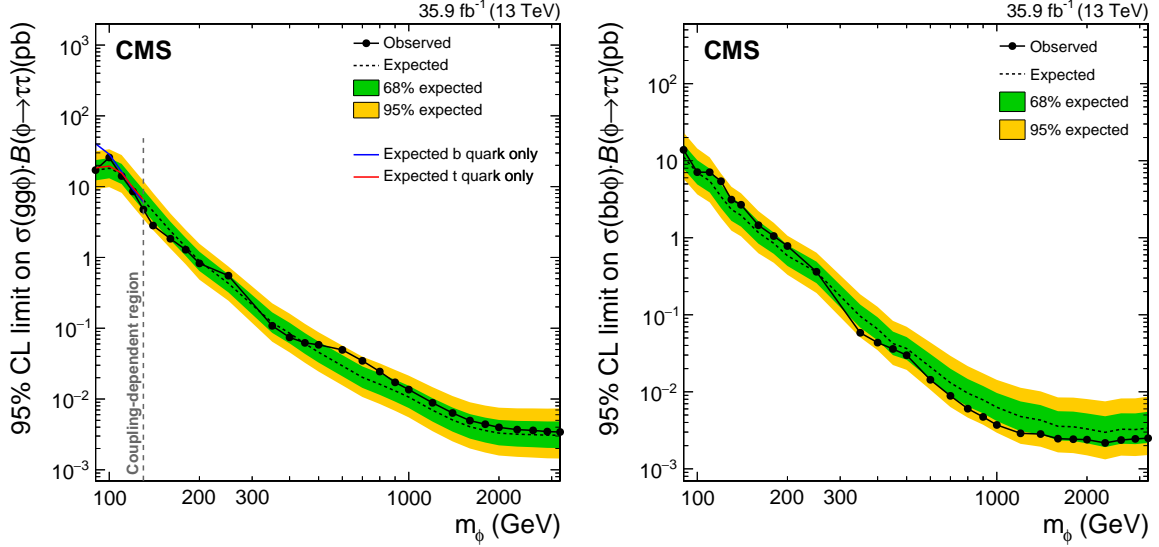


Figure 7: Expected and observed 95% CL upper limits for the production of a single narrow resonance,  $\phi$ , with a mass between 90 GeV and 3.2 TeV in the  $\tau\tau$  final state (left) for the production via gluon fusion ( $gg\phi$ ) and (right) in association with b quarks ( $bb\phi$ ). The expected median of the exclusion limit is shown by the dashed line. The dark green and bright yellow bands indicate the 68 and 95% confidence intervals for the variation of the expected exclusion limit. The black dots correspond to the observed limits. In the left panel the expected exclusion limits for the cases where (blue continuous line) only the b quark and (red continuous line) only the t quark are taken into account in the fermion loop are also shown. Left of the dashed vertical line the two different assumptions lead to visible differences in the expected exclusion limit.

production of a single narrow width resonance,  $\phi$ , via gluon fusion or in association with b quarks. In Fig. 7 these limits are shown as a function of  $m_\phi$ . For the determination of the limit on one process, e.g., gluon fusion, the normalization for the corresponding other process, e.g., associated production with b quarks, is treated as a freely varying parameter in the signal-plus-background fit that is performed prior to the limit calculation. The expectation for an SM Higgs boson at 125 GeV is taken into account in the SM backgrounds. For both production modes the  $p_T$  spectrum of the  $\phi$  is estimated at NLO precision in  $\alpha_s$ , as described in Section 5.1. Differences in the sensitivity of the analysis only occur at low masses, where the  $p_T$  of the  $\phi$  significantly contributes to the  $p_T$  of its decay products. In the figure this is emphasized by adding the median for the expected limit using either only the b quark or only the t quark for the modeling of the  $\phi$   $p_T$  spectrum. For the production via gluon fusion the expected limits range between 18 pb at  $m_\phi = 90$  GeV and 3.5 fb at  $m_\phi = 3.2$  TeV. For the production in association with b quarks they range between 15 pb (at  $m_\phi = 90$  GeV) and 2.5 fb (at  $m_\phi = 3.2$  TeV). In both cases, the excluded cross section falls with increasing mass, before becoming constant at around 1 TeV. No significant deviation from the expectation is observed. When restricted to the  $e\tau_h$ ,  $\mu\tau_h$ , or  $\tau_h\tau_h$  final state, the results obtained from the cross-checks summarized in Section 5.4 are compatible with the results obtained from the main analysis described in this paper. A scan of the likelihood for this signal model is also performed, as a function of the gluon fusion cross section and the cross section for the associated production with b quarks, for the tested mass points. A representative subset of this likelihood scan at six mass points is shown in Fig. 8.

In the second interpretation of the data, exclusion contours in the  $m_A$ - $\tan\beta$  plane are determined for two representative benchmark scenarios of the MSSM, the  $m_h^{\text{mod}+}$  and the hMSSM [82–

84]. The  $m_h^{\text{mod}+}$  scenario is compatible with the observation of the Higgs boson at 125 GeV, which is interpreted as the  $h$  within the theoretical uncertainties in  $m_h$  of  $\pm 3$  GeV [85, 86]. The phenomenological hMSSM also incorporates the observed Higgs boson at 125 GeV, interpreting it as the  $h$ . The uncertainties in the mass measurement are then used in turn to estimate the main radiative corrections to predict the masses and couplings of the remaining MSSM Higgs bosons. For the determination of the exclusion contours the model predictions as provided by the LHC Higgs Cross Section Working Group [74, 87] are used. Inclusive cross sections for the production via gluon fusion are calculated using the program SUSHI (v1.4.1) [88], including NLO QCD corrections in the context of the MSSM [89–94], as well as NNLO QCD corrections for the top quark contribution to the fermion loop in the heavy top quark limit [95–99], and electroweak effects from light quarks [100, 101]. For associated production with  $b$  quarks four-flavor scheme NLO QCD calculations [102, 103] and five-flavor scheme NNLO QCD calculations, as implemented in SUSHI based on BBH@NNLO [104], are combined using the Santander matching scheme [105]. The Higgs boson masses and mixing, and the effective Yukawa couplings for the  $m_h^{\text{mod}+}$  scenario, are calculated using the FEYNHIGGS 2.10.2 [85, 106–109] code. The branching fraction of the MSSM Higgs bosons to  $\tau$  leptons is calculated with FEYNHIGGS for the  $m_h^{\text{mod}+}$  scenario and using the program HDECAY 6.40 [110] for the hMSSM scenario.

The simulated single neutral Higgs boson signals are combined into a multiresonance signal model for the given values of  $m_A$  and  $\tan \beta$ , taking into account the predictions for the mass, production cross sections, and branching fraction into  $\tau$  leptons for each of the neutral Higgs bosons. For each value of  $m_A$  and  $\tan \beta$ , using a fine-grain scan, a maximum likelihood fit to the data is performed under the background-only and the signal-plus-background hypotheses using the likelihood of Eq. (9) with a test statistic that is slightly different from Eq. (10). The numerator remains the same, with a fixed value of  $\mu = 1$ , and corresponds to the signal prediction for the given value of  $m_A$  and  $\tan \beta$ . However no signal strength parameter is included in the denominator; the model is thus fixed to the background-only prediction. Note that the SM Higgs boson is added to the non Higgs boson background processes. This turns the likelihood ratio into a comparison between the MSSM and the SM Higgs sector hypotheses, and ensures a well defined problem even when the analysis becomes sensitive to the observed Higgs boson at 125 GeV. In such a situation a test of the MSSM hypothesis against a background hypothesis ignoring the SM Higgs boson would be based on a wrong null-hypothesis. The median and confidence intervals for the expected exclusion contour are determined from pseudo-experiments. In Fig. 9 the observed and expected 95% CL exclusion contours for the MSSM  $m_h^{\text{mod}+}$  and the hMSSM scenarios are shown. The exclusion contours reach up to 1.6 TeV, extending the excluded mass range by almost a factor of two in  $m_A$  compared to the previous CMS publication using the same final state [18]. In both scenarios the exclusion contours extend down to values of  $\tan \beta \approx 6$  for values of  $m_A \lesssim 250$  GeV. For the  $m_h^{\text{mod}+}$  scenario, those parts of the parameter space in which  $m_h$  deviates by more than  $\pm 3$  GeV from the mass of the observed Higgs boson at 125 GeV are indicated by a red hatched area.

In the low mass region the exclusion contour is similar to the previous CMS publication, while a higher sensitivity might be expected. This can be attributed to three main factors: the choice of single lepton triggers in the  $e\tau_h$  and  $\mu\tau_h$  final states together with the higher instantaneous luminosity leads to the need for higher  $p_T$  thresholds at the trigger level and therefore reduced signal acceptance; the change of the discriminating variable from the estimate of the fully reconstructed  $\tau\tau$  mass to  $m_T^{\text{tot}}$  provides more sensitivity for high masses, but slightly less sensitivity for lower masses; and finally the prediction of the kinematic distributions of the signal at NLO precision reveals a generally softer  $p_T$  spectrum for the gluon fusion production mode, which dominates for low values of  $\tan \beta$ . Over the whole mass range the observed exclusion contours

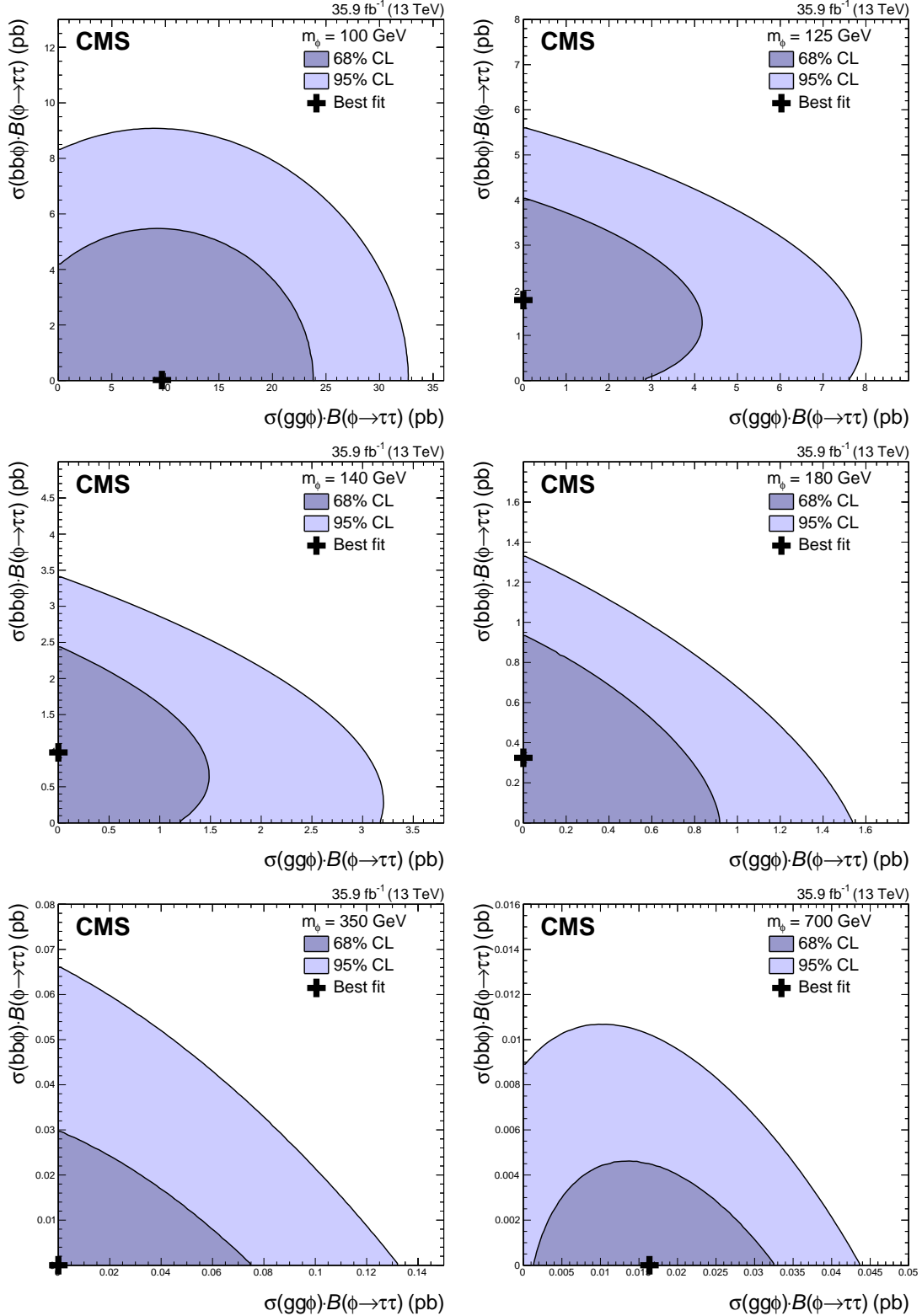


Figure 8: Scan of the likelihood function for the search in the  $\tau\tau$  final state for a single narrow resonance,  $\phi$ , produced via gluon fusion ( $gg\phi$ ) or in association with  $b$  quarks ( $bb\phi$ ). A representative subset of the mass points tested at (upper left) 100 GeV, (upper right) 125 GeV, (middle left) 140 GeV, (middle right) 180 GeV, (lower left) 350 GeV, and (lower right) 700 GeV is shown. Note that in the fits the signal strengths are not allowed to become negative.

follow the expectation with the largest deviations still contained in the 95% confidence interval for the variation of the expected exclusion.

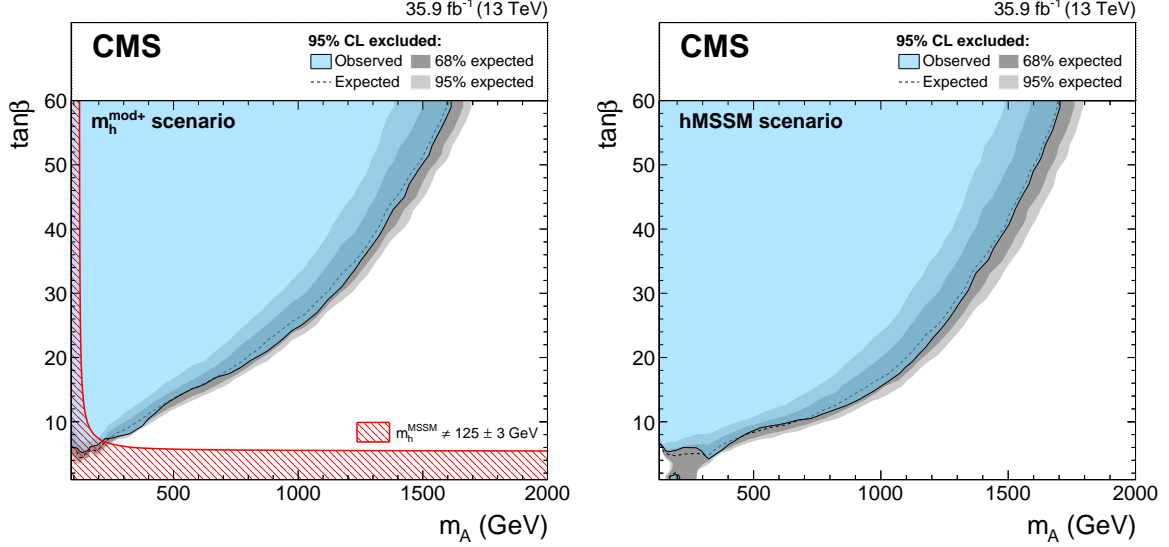


Figure 9: Expected and observed 95% CL exclusion contour (left) in the MSSM  $m_h^{\text{mod}+}$  and (right) in the hMSSM scenarios. The expected median is shown as a dashed black line. The dark and bright gray bands indicate the 68 and 95% confidence intervals for the variation of the expected exclusion. The observed exclusion contour is indicated by the colored blue area. For the  $m_h^{\text{mod}+}$  scenario, those parts of the parameter space, where  $m_h$  deviates by more than  $\pm 3 \text{ GeV}$  from the mass of the observed Higgs boson at  $125 \text{ GeV}$  are indicated by a red hatched area.

## 9 Summary

A search for additional heavy neutral Higgs bosons in the decay into two  $\tau$  leptons in the context of the minimal supersymmetric standard model (MSSM) has been presented. This search has been performed in the most sensitive  $e\mu$ ,  $e\tau_h$ ,  $\mu\tau_h$ , and  $\tau_h\tau_h$  final states of the  $\tau\tau$  pair, where  $\tau_h$  indicates a hadronic  $\tau$  lepton decay. No signal has been found. Model-independent limits at 95% confidence level have been set for the production of a single narrow resonance decaying into a pair of  $\tau$  leptons. These range from  $18 \text{ pb}$  at  $90 \text{ GeV}$  to  $3.5 \text{ fb}$  at  $3.2 \text{ TeV}$  for production via gluon fusion and from  $15 \text{ pb}$  (at  $90 \text{ GeV}$ ) to  $2.5 \text{ fb}$  (at  $3.2 \text{ TeV}$ ) for production in association with  $b$  quarks. Finally 95% confidence level exclusion contours have been provided for two representative benchmark scenarios, namely the  $m_h^{\text{mod}+}$  and the hMSSM scenarios. In these two scenarios the presence of a neutral heavy MSSM Higgs boson up to  $m_A \lesssim 250 \text{ GeV}$  is excluded for  $\tan\beta$  values above 6. The exclusion contour reaches  $1.6 \text{ TeV}$  for  $\tan\beta = 60$ .

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