

Project for the final oral examination of PHY981

Nucleon-nucleon scattering

The aim of this project is to solve the Lippman-Schwinger equation for two interacting nucleons and relate the obtained phase shifts with those extracted from the experimental cross sections.

We are going to solve the Schrödinger equation (SE) for the neutron-proton system in momentum space for positive energies E in order to obtain the phase shifts. We can rewrite the SE in momentum space as

$$\frac{k^2}{m}\psi_l(k) + \frac{2}{\pi} \int_0^\infty dq q^2 V_l(k, q) \psi_l(q) = E \psi_l(k). \quad (1)$$

Here we have used units $\hbar = c = 1$. This means that k has dimension energy. k is the relative momentum between the two particles. A partial wave expansion has been used in order to reduce the problem to an integral over the magnitude of momentum only. The subscript l refers therefore to a partial wave with a given orbital momentum l . To obtain the potential in momentum space we used the Fourier-Bessel transform (Hankel transform)

$$V_l(k, k') = \int j_l(kr) V(r) j_l(k'r) r^2 dr, \quad (2)$$

where j_l is the spherical Bessel function. We will just study the case $l = 0$, which means that $j_0(kr) = \sin(kr)/kr$.

For scattering states, $E > 0$, the corresponding equation to solve is the so-called Lippman-Schwinger equation. This is an integral equation where we have to deal with the amplitude $R(k, k')$ (reaction matrix) defined through the integral equation

$$R_l(k, k') = V_l(k, k') + \frac{2}{\pi} \mathcal{P} \int_0^\infty dq q^2 V_l(k, q) \frac{1}{E - q^2/m} R_l(q, k'), \quad (3)$$

where the total kinetic energy of the two incoming particles in the center-of-mass system is

$$E = \frac{k_0^2}{m}. \quad (4)$$

The symbol \mathcal{P} indicates that Cauchy's principal-value prescription is used in order to avoid the singularity arising from the zero of the denominator. We will discuss below how to solve this problem. Eq. (3) represents then the problem you will have to solve numerically.

The matrix $R_l(k, k')$ relates to the the phase shifts through its diagonal elements as

$$R_l(k_0, k_0) = -\frac{\tan \delta_l}{m k_0}. \quad (5)$$

From now on we will drop the subscript l in all equations.

In order to solve the Lippman-Schwinger equation in momentum space, we need first to write a function which sets up the mesh points. We need to do that since we are going to approximate an integral through

$$\int_a^b f(x) dx \approx \sum_{i=1}^N w_i f(x_i),$$

where we have fixed N lattice points through the corresponding weights w_i and points x_i . Start writing your main program by setting up the mesh points and the corresponding weights. Fix first the number of mesh points N . Use the included function *gauleg* in the file lib.cpp (see also lib.h) to set up the weights ω_i and the points k_i . Before you go on you need to keep in mind that *gauleg* uses the Legendre polynomials to fix the mesh points and weights. This means that the integral is for the interval $[-1, 1]$. Your integral is for the interval $[0, \infty]$. You will need to map the weights from *gauleg* to your interval. To do this, call first *gauleg(a, b, x[], w[], N)*, with $a = -1$, $b = 1$. It returns the mesh points x_i and weights w_i . You map these points over to the limits in your integral. You can then use the following mapping

$$k_i = \text{const} \times \tan \left\{ \frac{\pi}{4} (1 + x_i) \right\},$$

and

$$\omega_i = \text{const} \frac{\pi}{4} \frac{w_i}{\cos^2\left(\frac{\pi}{4}(1+x_i)\right)}.$$

If you choose dimension fm^{-1} for k , set $\text{const} = 1$. If you choose to work with MeV, set $\text{const} \sim 200$ ($\hbar c = 197$ MeVfm). You must choose the units to use.

The next step is to write a function which calculates the potential in momentum space. The potential we will use here is a parametrized potential between a proton and neutron for the partial wave 1S_0 , i.e., spin $S = 0$ and orbital momentum $l = 0$, a singlet S-state. This state does not have a bound state for the deuteron (only the triplet S-state has). The parametrized version of this potential fits the experimental phase-shifts. It is given by

$$V(r) = V_a \frac{e^{-ax}}{x} + V_b \frac{e^{-bx}}{x} + V_c \frac{e^{-cx}}{x} \quad (6)$$

with $x = \mu r$, $\mu = 0.7 \text{ fm}^{-1}$ (the inverse of the pion mass), $V_a = -10.463 \text{ MeV}$ and $a = 1$, $V_b = -1650.6 \text{ MeV}$ and $b = 4$ and $V_c = 6484.3 \text{ MeV}$ and $c = 7$. Find the potential in momentum space using Eq. (2) with $j_0(kr) = \sin(kr)/kr$. The transform of a potential on the form $V(r) = V_0 \exp(-\mu r)/r$ is

$$V(k', k) = \frac{V_0}{4k'k} \ln \left(\frac{(k' + k)^2 + \mu^2}{(k' - k)^2 + \mu^2} \right). \quad (7)$$

Write a function which calculates the expressions for the potential in momentum space.

The principal value in Eq. (3) is rather tricky to evaluate numerically, mainly since computers have limited precision. We will here use a subtraction trick often used when dealing with singular integrals in numerical calculations. We introduce first the calculus relation

$$\int_{-\infty}^{\infty} \frac{dk}{k - k_0} = 0. \quad (8)$$

It means that the curve $1/(k - k_0)$ has equal and opposite areas on both sides of the singular point k_0 . If we break the integral into one over positive k and one over negative k , a change of variable $k \rightarrow -k$ allows us to rewrite the last equation as

$$\int_0^{\infty} \frac{dk}{k^2 - k_0^2} = 0. \quad (9)$$

We can use this to express a principal values integral as

$$\mathcal{P} \int_0^{\infty} \frac{f(k)dk}{k^2 - k_0^2} = \int_0^{\infty} \frac{(f(k) - f(k_0))dk}{k^2 - k_0^2}, \quad (10)$$

where the right-hand side is no longer singular at $k = k_0$, it is proportional to the derivative df/dk , and can be evaluated numerically as any other integral.

We can then use the trick in Eq. (10) to rewrite Eq. (3) as

$$R(k, k') = V(k, k') + \frac{2}{\pi} \int_0^{\infty} dq \frac{q^2 V(k, q) R(q, k') - k_0^2 V(k, k_0) R(k_0, k')}{(k_0^2 - q^2)/m}. \quad (11)$$

This is the equation you are going to solve numerically in order to calculate the phase shifts of Eq. (5). We are interested in obtaining $R(k_0, k_0)$.

How do we proceed in order to solve Eq. (11)?

1. Using the mesh points k_j and the weights ω_j , we can rewrite Eq. (11) as

$$R(k, k') = V(k, k') + \frac{2}{\pi} \sum_{j=1}^N \frac{\omega_j k_j^2 V(k, k_j) R(k_j, k')}{(k_0^2 - k_j^2)/m} - \frac{2}{\pi} k_0^2 V(k, k_0) R(k_0, k') \sum_{n=1}^N \frac{\omega_n}{(k_0^2 - k_n^2)/m}. \quad (12)$$

This equation contains now the unknowns $R(k_i, k_j)$ (with dimension $N \times N$) and $R(k_0, k_0)$.

2. We can turn Eq. (12) into an equation with dimension $(N + 1) \times (N + 1)$ with a mesh which contains the original mesh points k_j for $j = 1, N$ and the point which corresponds to the energy k_0 . Consider the latter as the 'observable' point. The mesh points become then k_j for $j = 1, n$ and $k_{N+1} = k_0$.
3. With these new mesh points we define the matrix

$$A_{i,j} = \delta_{i,j} - V(k_i, k_j)u_j, \quad (13)$$

where δ is the Kronecker δ and

$$u_j = \frac{2}{\pi} \frac{\omega_j k_j^2}{(k_0^2 - k_j^2)/m} \quad j = 1, N \quad (14)$$

and

$$u_{N+1} = -\frac{2}{\pi} \sum_{j=1}^N \frac{k_0^2 \omega_j}{(k_0^2 - k_j^2)/m}. \quad (15)$$

The first task is then to set up the matrix A for a given k_0 . This is an $(N + 1) \times (N + 1)$ matrix. It can be convenient to have an outer loop which runs over the chosen observable values for the energy k_0^2/m . *Note that all mesh points k_j for $j = 1, N$ must be different from k_0 . Note also that $V(k_i, k_j)$ is an $(N + 1) \times (N + 1)$ matrix.* Write a small function which sets up A .

4. With the matrix A we can rewrite Eq. (12) as a matrix problem of dimension $(N + 1) \times (N + 1)$. All matrices R , A and V have this dimension and we get

$$A_{i,l} R_{l,j} = V_{i,j}, \quad (16)$$

or just

$$AR = V. \quad (17)$$

5. Since you already have defined A and V (these are stored as $(N + 1) \times (N + 1)$ matrices) Eq. (17) involves only the unknown R . We obtain it by matrix inversion, i.e.,

$$R = A^{-1}V. \quad (18)$$

Thus, to obtain R , you will need to set up the matrices A and V and invert the matrix A . To do that you must call the function `matinv` in the program library `lib.cpp` (see also `lib.h`). With the inverse A^{-1} , performing a matrix multiplication with V results in R .

With R you can then evaluate the phase shifts by noting that

$$R(k_{N+1}, k_{N+1}) = R(k_0, k_0), \quad (19)$$

and you are done.

You can choose to read k_0 from file or screen, or set up a loop over chosen values of k_0 and for each k_0 solve Eq. (18).

1. When you have $R(k, k')$ for the given potential, evaluate now the phase-shifts using

$$R(k_0, k_0) = -\frac{\tan \delta}{mk_0}.$$

2. Compare the phase shifts for the potential of Eq. (6) with the experimental phase shifts that can be found in the article of the Nijmegen group in Physical Review C **48**, 792 (1993). Alternatively look up their webpage at <http://nn-online.org/>

The code can easily be extended to include realistic nucleon-nucleon interactions. Let me know if this is of interest, I will then provide a function (fortran) for the nucleon-nucleon force.