

## Exercises PHY981 Spring 2015

The exercises are available at the beginning of the week and are to be handed in the week thereafter on **Wednesdays at noon**. This can be done electronically by sending your github link or by sending a scan of your notes, a pdf or postscript file, or ipython notebook or any format you prefer, by email to [hjensen@nscl.msu.edu](mailto:hjensen@nscl.msu.edu).

### Exercise 2: This is a numerical exercise and is optional

The program for finding the eigenvalues of the harmonic oscillator are in the github folder <https://github.com/NuclearStructure/PHY981/tree/master/doc/pub/spdata/programs>.

You can use this program to solve the exercises below, or write your own using your preferred programming language, be it python, fortran or c++ or other languages.

- Compute the eigenvalues of the five lowest states with a given orbital momentum and oscillator frequency  $\omega$ . Study these results as functions of the the maximum value of  $r$  and the number of integration points  $n$ , starting with  $r_{\max} = 10$ . Compare the computed ones with the exact values and comment your results.
- Plot thereafter the eigenfunctions as functions of  $r$  for the lowest-lying state with a given orbital momentum  $l$ .
- Replace thereafter the harmonic oscillator potential with a Woods-Saxon potential using the parameters discussed above. Compute the lowest five eigenvalues and plot the eigenfunction of the lowest-lying state. How does this compare with the harmonic oscillator? Comment your results and possible implications for nuclear physics studies.

### Exercise 3

Consider the Slater determinant

$$\Phi_{\lambda}^{AS}(x_1 x_2 \dots x_N; \alpha_1 \alpha_2 \dots \alpha_N) = \frac{1}{\sqrt{N!}} \sum_p (-)^p P \prod_{i=1}^N \psi_{\alpha_i}(x_i).$$

where  $P$  is an operator which permutes the coordinates of two particles. We have assumed here that the number of particles is the same as the number of available single-particle states, represented by the greek letters  $\alpha_1 \alpha_2 \dots \alpha_N$ .

- Write out  $\Phi^{AS}$  for  $N = 3$ .
- Show that

$$\int dx_1 dx_2 \dots dx_N \left| \Phi_{\lambda}^{AS}(x_1 x_2 \dots x_N; \alpha_1 \alpha_2 \dots \alpha_N) \right|^2 = 1.$$

- Define a general onebody operator  $\hat{F} = \sum_i^N \hat{f}(x_i)$  and a general twobody operator  $\hat{G} = \sum_{i>j}^N \hat{g}(x_i, x_j)$  with  $g$  being invariant under the interchange of the coordinates of particles  $i$  and  $j$ . Calculate the matrix elements for a two-particle Slater determinant

$$\langle \Phi_{\alpha_1 \alpha_2}^{AS} | \hat{F} | \Phi_{\alpha_1 \alpha_2}^{AS} \rangle,$$

and

$$\langle \Phi_{\alpha_1 \alpha_2}^{AS} | \hat{G} | \Phi_{\alpha_1 \alpha_2}^{AS} \rangle.$$

Explain the short-hand notation for the Slater determinant. Which properties do you expect these operators to have in addition to an eventual permutation symmetry?

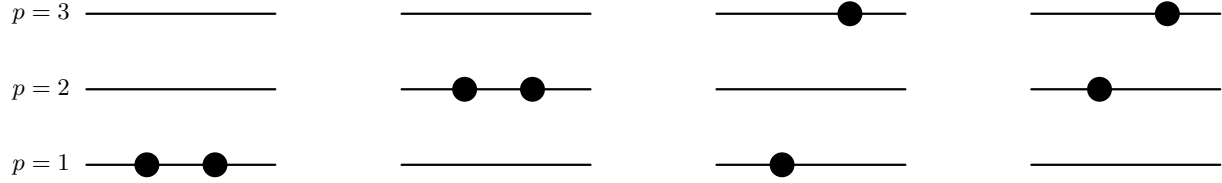


FIG. 1: Schematic plot of the possible single-particle levels with double degeneracy. The filled circles indicate occupied particle states. The spacing between each level  $p$  is constant in this picture. We show some possible two-particle states.

#### Exercise 4

We will now consider a simple three-level problem, depicted in the figure below. This is our first and very simple model of a possible many-nucleon (or just fermion) problem and the shell-model. The single-particle states are labelled by the quantum number  $p$  and can accommodate up to two single particles, viz., every single-particle state is doubly degenerate (you could think of this as one state having spin up and the other spin down). We let the spacing between the doubly degenerate single-particle states be constant, with value  $d$ . The first state has energy  $d$ . There are only three available single-particle states,  $p = 1$ ,  $p = 2$  and  $p = 3$ , as illustrated in the figure.

- How many two-particle Slater determinants can we construct in this space?
- We limit ourselves to a system with only the two lowest single-particle orbits and two particles,  $p = 1$  and  $p = 2$ . We assume that we can write the Hamiltonian as

$$\hat{H} = \hat{H}_0 + \hat{H}_I,$$

and that the onebody part of the Hamiltonian with single-particle operator  $\hat{h}_0$  has the property

$$\hat{h}_0 \psi_{p\sigma} = p \times d \psi_{p\sigma},$$

where we have added a spin quantum number  $\sigma$ . We assume also that the only two-particle states that can exist are those where two particles are in the same state  $p$ , as shown by the two possibilities to the left in the figure. The two-particle matrix elements of  $\hat{H}_I$  have all a constant value,  $-g$ . Show then that the Hamiltonian matrix can be written as

$$\begin{pmatrix} 2d - g & -g \\ -g & 4d - g \end{pmatrix},$$

and find the eigenvalues and eigenvectors. What is mixing of the state with two particles in  $p = 2$  to the wave function with two-particles in  $p = 1$ ? Discuss your results in terms of a linear combination of Slater determinants.

- Add the possibility that the two particles can be in the state with  $p = 3$  as well and find the Hamiltonian matrix, the eigenvalues and the eigenvectors. We still insist that we only have two-particle states composed of two particles being in the same level  $p$ . You can diagonalize numerically your  $3 \times 3$  matrix.

This simple model catches several birds with a stone. It demonstrates how we can build linear combinations of Slater determinants and interpret these as different admixtures to a given state. It represents also the way we are going to interpret these contributions. The two-particle states above  $p = 1$  will be interpreted as excitations from the ground state configuration,  $p = 1$  here. The reliability of this ansatz for the ground state, with two particles in  $p = 1$ , depends on the strength of the interaction  $g$  and the single-particle spacing  $d$ . Finally, this model is a simple schematic ansatz for studies of pairing correlations and thereby superfluidity/superconductivity in fermionic systems.