# No-core configuration interaction methods for ab initio nuclear theory

Patrick Fasano

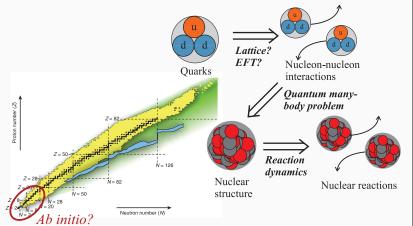
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ECT\* TALENT School 2017 University of Notre Dame

#### Goals of *ab initio* nuclear structure

First-principles understanding of nature Nuclei from QCD

What is the origin of simple patterns in complex nuclei?



#### **Ab Initio Nuclear Theory**

- Ab initio theory start with a realistic NN (and maybe 3N) interaction and solve the quantum many-body problem
- We want not only ground-state energies but many other low-energy observables:
  - energy spectra
  - spatial distributions
  - · momentum distributions
  - inputs for reaction theory (spectroscopic overlaps, spectroscopic factors, asymptotic normalization factors)
  - electromagnetic observables (transition probabilities, magnetic moments, form factors)
  - weak observables (Fermi and Gamow-Teller probabilities)
  - nuclear equation of state parameters

### **Ab Initio Nuclear Theory**

Many techniques for solving nuclear Hamiltonians have been devised<sup>1</sup>:

- Few-body:
  - Fadeev-Yakubowski Equation
  - Hyperspherical Harmonics
- Many-body:
  - Coupled-cluster
  - Quantum Monte Carlo
  - No-Core Configuration Interaction (NCCI) model

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Basic idea: configuration-interaction (CI) without a core!

All types of excitations are allowed.

Begin with an orthonormal single-particle basis:

$$\hat{h} | nljm \rangle = \epsilon_{nljm} | nljm \rangle$$

Of course, we must truncate to a finite number of single-particle states.

Construct a many-body basis of Slater determinants with good M:

$$|\Psi_{\alpha}\rangle = |\pi_{\alpha_1}\pi_{\alpha_2}\cdots\pi_{\alpha_Z}\nu_{\alpha_1}\nu_{\alpha_2}\cdots\nu_{\alpha_N}\rangle$$

Full Configuration Interaction: include all possible Slaters. Dimension blows up:

$N_{ m shell}$	2	3	4	5	6	7
M scheme						
<sup>4</sup> He	98	$3.06 \times 10^{3}$	$3.98 \times 10^{4}$	$3.14 \times 10^{5}$	$1.77 \times 10^{6}$	$7.84 \times 10^{6}$
<sup>6</sup> He	216	$6.51 \times 10^{4}$	$3.86 \times 10^{6}$	$9.80 \times 10^{7}$	$1.45 \times 10^{9}$	$1.47 \times 10^{10}$
<sup>6</sup> Li	293	$8.59 \times 10^{4}$	$5.08 \times 10^{6}$	$1.29 \times 10^{8}$	$1.91 \times 10^{9}$	$1.94 \times 10^{10}$
<sup>7</sup> Li	400	$3.60 \times 10^{5}$	$4.51 \times 10^{7}$	$2.05 \times 10^{9}$	$4.91 \times 10^{10}$	$7.50 \times 10^{11}$
<sup>8</sup> Be	518	$1.47 \times 10^{6}$	$3.96 \times 10^{7}$	$3.24 \times 10^{10}$	$1.26 \times 10^{12}$	$2.91 \times 10^{13}$
$^{10}B$	293	$1.34 \times 10^{7}$	$1.82 \times 10^{10}$	$5.02 \times 10^{11}$	$5.22 \times 10^{14}$	$2.78 \times 10^{16}$
<sup>12</sup> C	98	$8.22 \times 10^{7}$	$5.87 \times 10^{11}$	$5.50 \times 10^{14}$	$1.54 \times 10^{17}$	$1.90 \times 10^{19}$
<sup>16</sup> O	1	$8.12 \times 10^{8}$	$2.10 \times 10^{14}$	$2.51 \times 10^{18}$	$5.32 \times 10^{21}$	$3.59 \times 10^{24}$

T. Abe et al., "Benchmarks of the full configuration interaction, monte carlo shell model, and no-core full configuration methods", Phys. Rev. C  $\bf 86$ , 054301 (2012)

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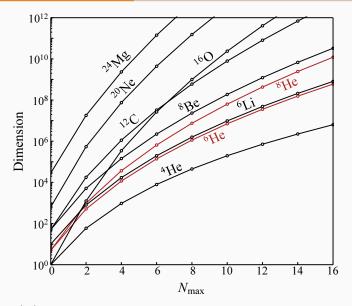
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Assign a weight to the Slater determinants by  $W_{\alpha} = \sum w_{\alpha_i}$ . Truncate based on weight of Slater determinant.

## **NCCI Basis Size**



#### $N_{max}$ Truncation

For the case where  $w = N = 2n + \ell$ , we define the  $N_{max}$  truncation.

All Slaters with a total number of oscillator quanta

$$N = \sum_{\alpha=1}^{A} N_{\alpha} \le N_0 + N_{max}$$

are included in the basis, where  $N_{\alpha}$  is the oscillator quantum number of the  $\alpha-th$  particle, and  $N_0$  is the number of oscillator quanta in the lowest configuration.

 $N_{max}$ -truncation has been preferred traditionally because it allows exact center-of-mass factorization, and can lead to faster convergence with respect to basis size than FCI-truncation.

#### **NCCI** Hamiltonian

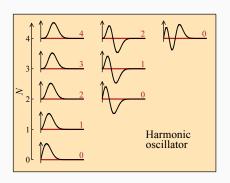
#### **No-Core Configuration Interaction Hamiltonian**

$$H = T_{intr} + V + aN_{c.m.}$$

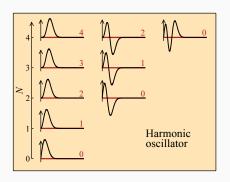
- T<sub>intr</sub> intrinsic kinetic energy ("upgraded" to a two-body operator)
- V realistic nucleon-nucleon (NN) interaction
- $N_{c.m.}$  center-of-mass number operator (Lawson term)

Take matrix elements of H from relative coordinates, Moshinksy transform into lab coordinates, and feed in as two-body matrix elements.

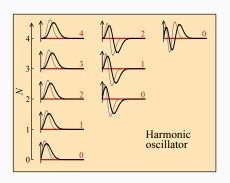
By completeness, a calculation in the infinite space  $\rightarrow$  independence from parameters in the single-particle basis (i.e.  $\hbar\omega$ ).



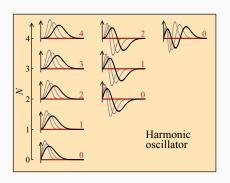
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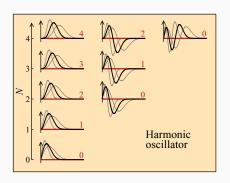
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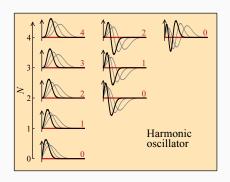
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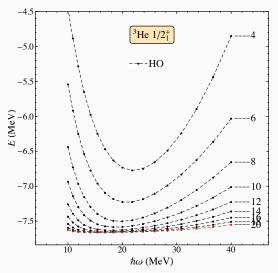


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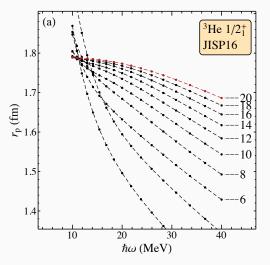


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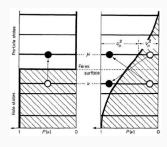


Ch. Constantinou et al., in preparation



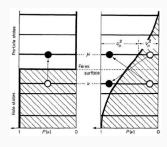
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- Attempt to formulate a "natural" basis for performing NCCI calculations.
- Define "natural" → minimum depletion of the Fermi sea.
- Run many-body calculation, diagonalize one-body static density matrix, use eigenvectors as similarity transformation on single-particle space.
- Minimizing depletion of Fermi sea, not minimizing energy!
- Built from many-body calculation, so "aware" of correlations.



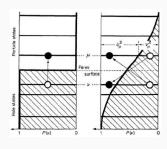
D. J. Rowe, *Nuclear collective motion: models and theory*, (World Scientific, Singapore, 2010)

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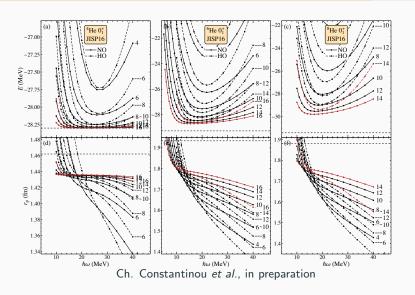


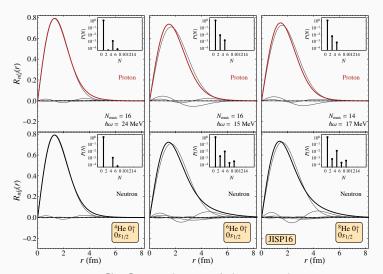
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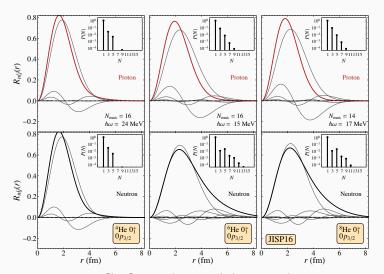


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### **Summary**

- Goal: Try to solve the many-body problem starting with a realistic NN (and 3N) interaction.
- Treat all nucleons on an equal footing; all nucleons in the "valence space". Truncate basis based on general weighting scheme.
- Convergence assessed based on independence from single-particle basis and many-body truncation.
- Picking better basis functions leads to better convergence!