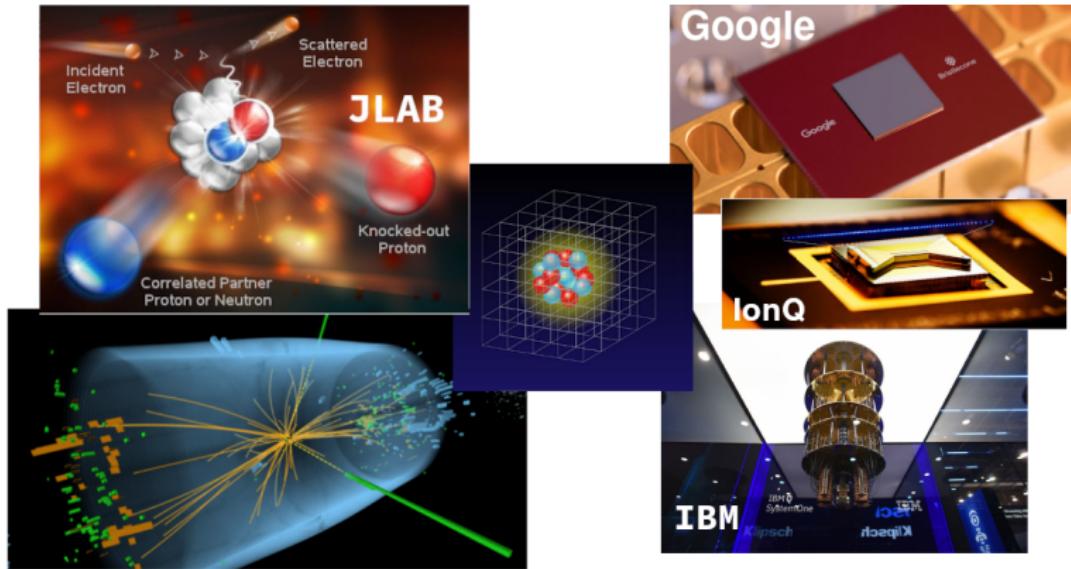


Nuclear Reactions on a Quantum Computer

Alessandro Roggero



DTP/TALENT School

ECT* – 3 July, 2025



Trento Institute for
Fundamental Physics
and Applications



The need for ab-initio many-body dynamics in NP

- ν scattering for supernovae explosion and NS cooling
- capture reactions for crust heating and nucleosynthesis
- cross sections for dark-matter discovery and neutrino physics
- transport properties of neutron star matter for X-ray emission

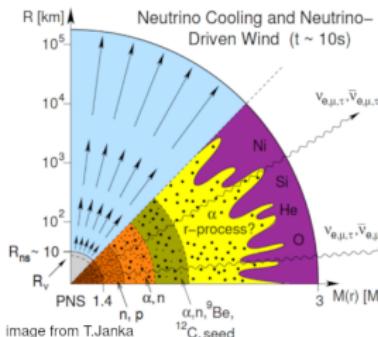


image from T.Janka

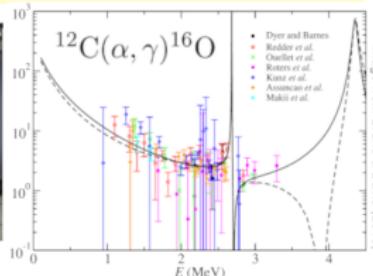
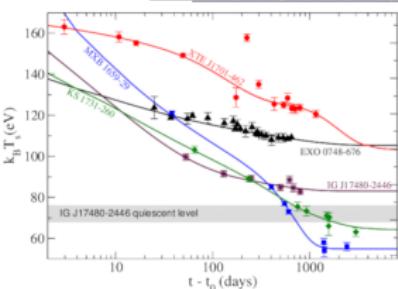
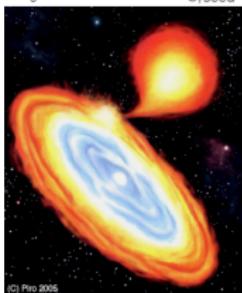
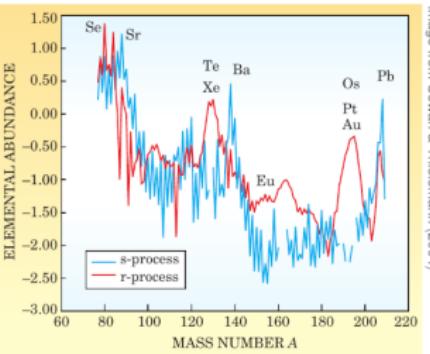
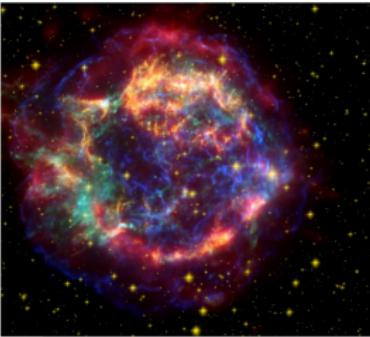
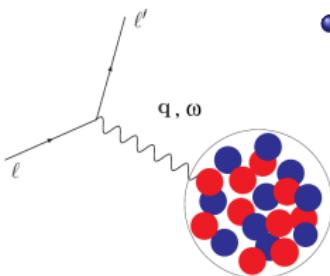


image from Cernan & Thielemann (2004)

image from Bigerne & Davids (2015)

Inclusive cross section and the response function

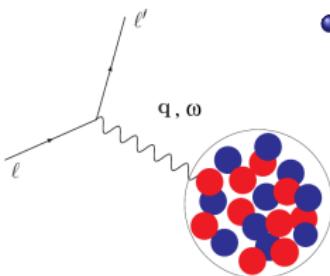


- xsection completely determined by response function

$$R_O(\omega) = \sum_f \left| \langle f | \hat{O} | \Psi_0 \rangle \right|^2 \delta(\omega - E_f + E_0)$$

- excitation operator \hat{O} specifies the vertex

Inclusive cross section and the response function



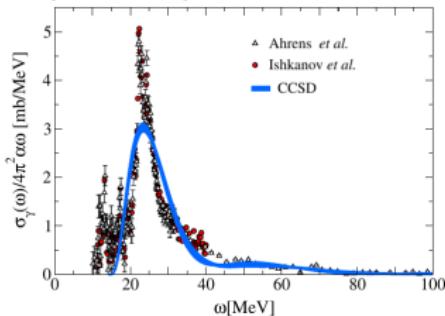
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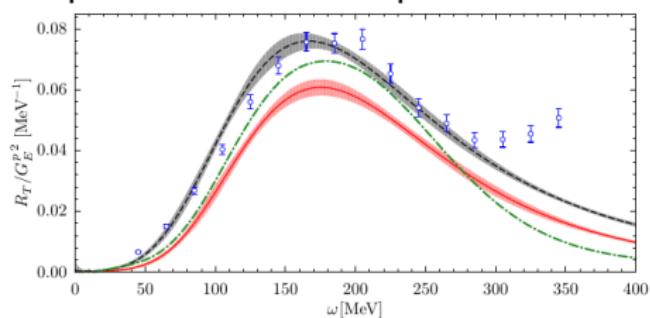
Extremely challenging classically for strongly correlated quantum systems

- dipole response of ^{16}O



Bacca et al. (2013) LIT+CC

- quasi-elastic EM response of ^{12}C



Lovato et al. (2016) GFMC

Exclusive cross sections in neutrino oscillation experiments



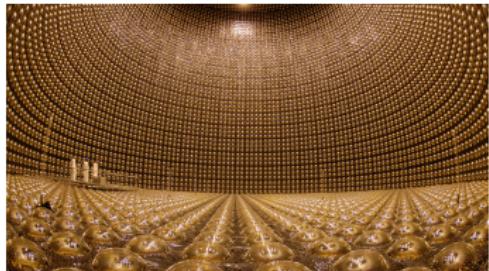
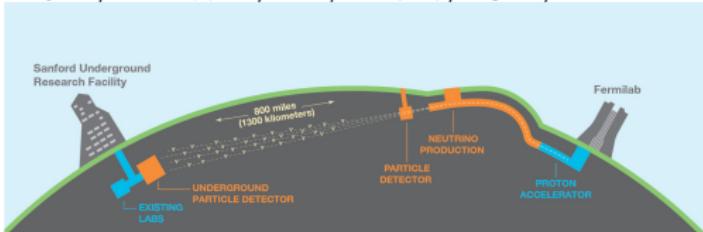
Goals for ν oscillation exp.

- neutrino masses
- accurate mixing angles
- CP violating phase

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - \sin^2(2\theta)\sin^2\left(\frac{\Delta m^2 L}{4E_\nu}\right)$$

- need to use measured reaction products to constrain E_ν of the event

DUNE, MiniBooNE, T2K, Minerva, NO ν A,...



Idealized algorithm for exclusive processes at fixed q

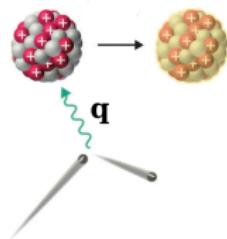
- prepare the target ground state



Roggero & Carlson (2019)

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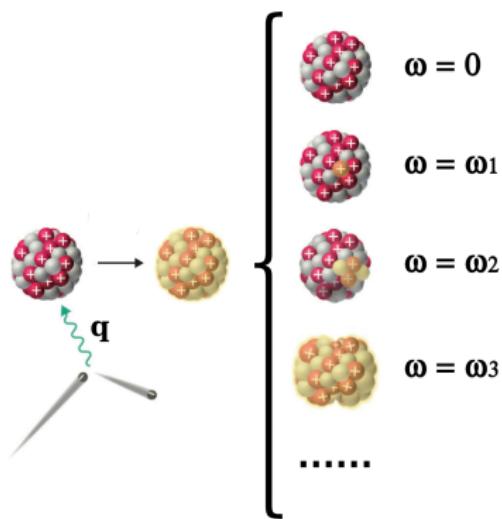
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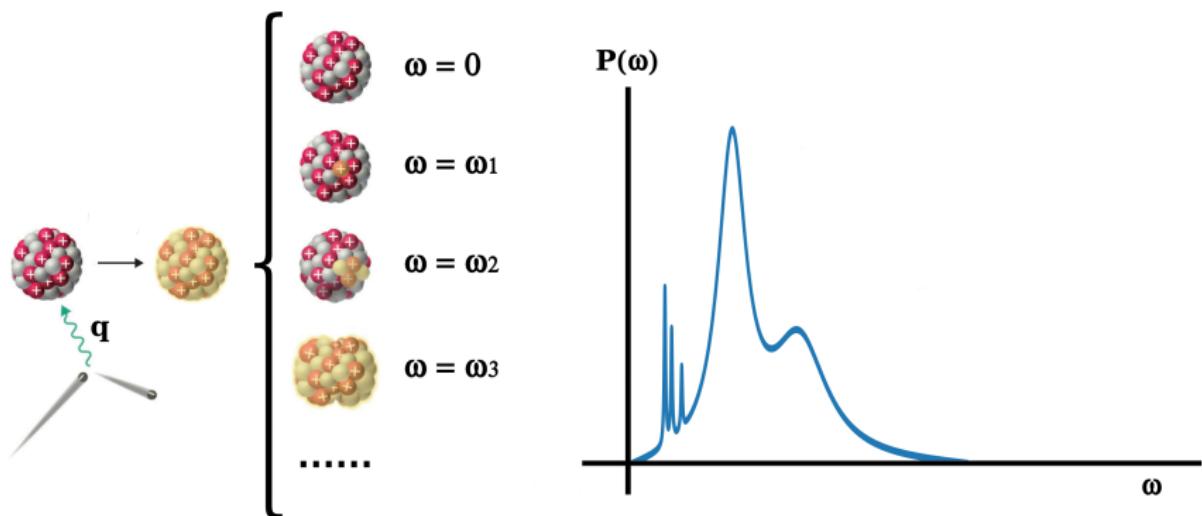
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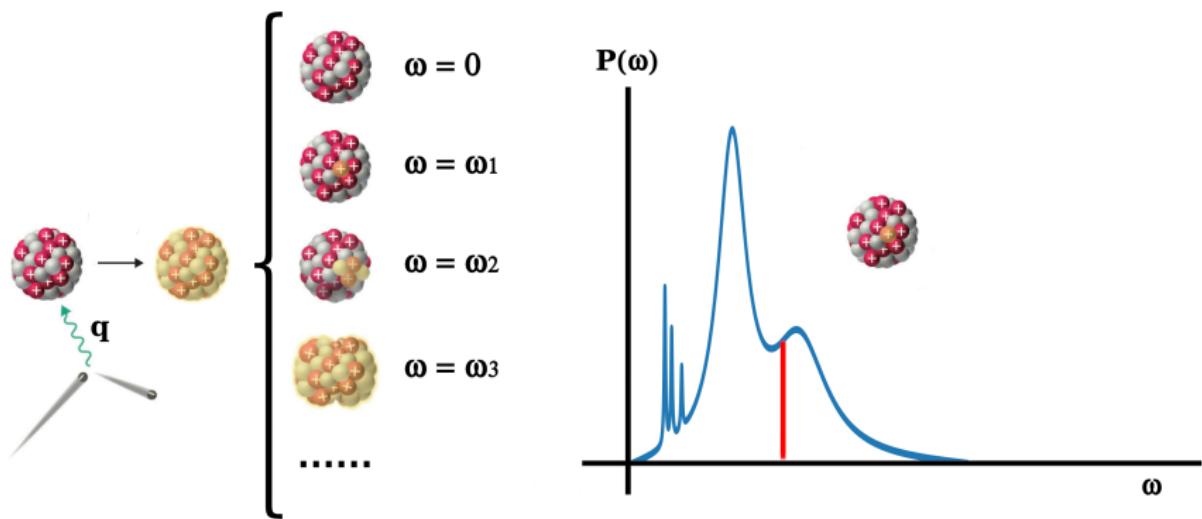
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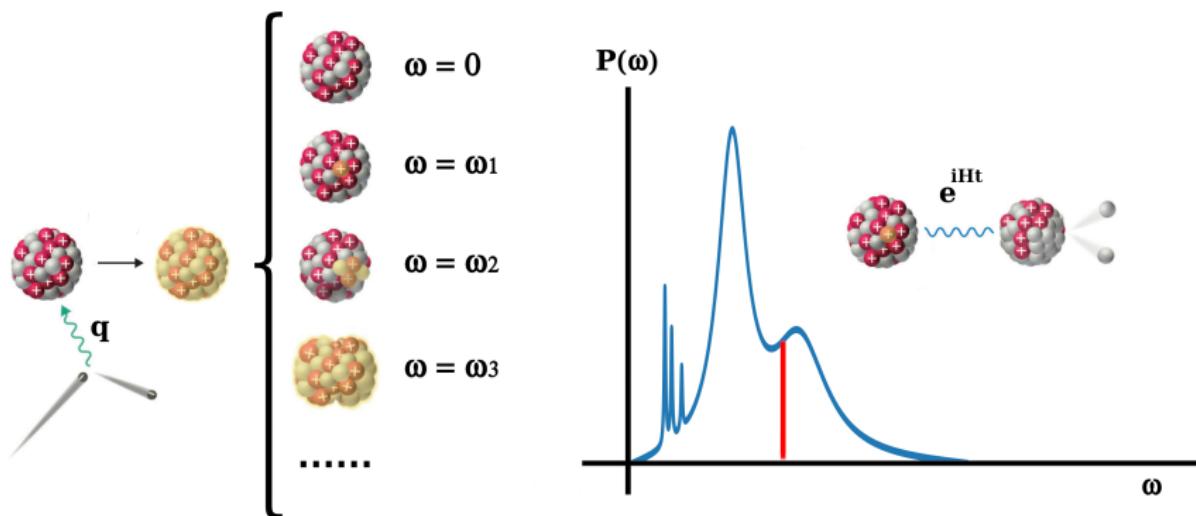
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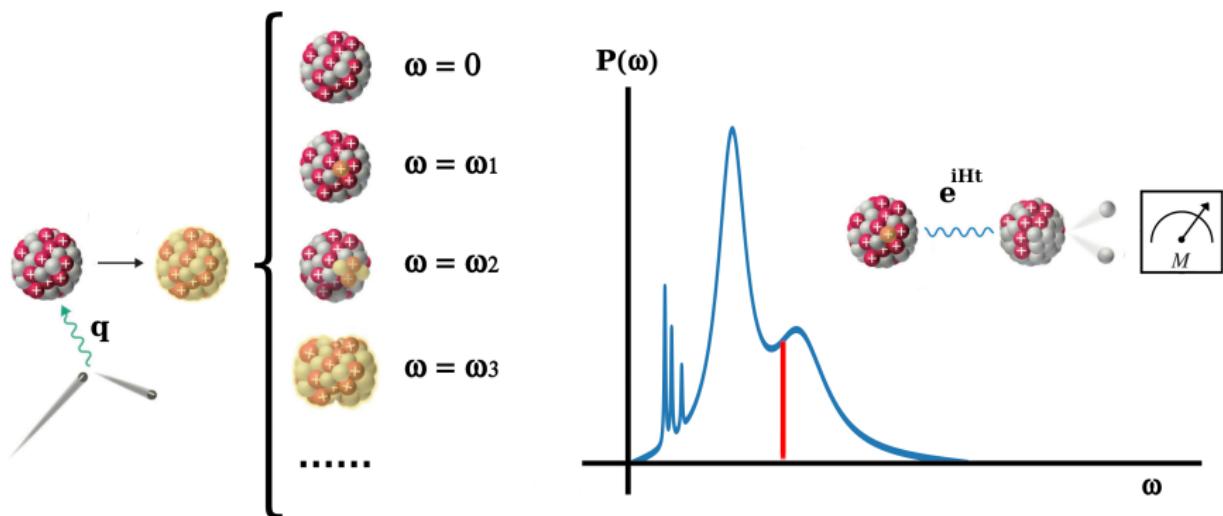
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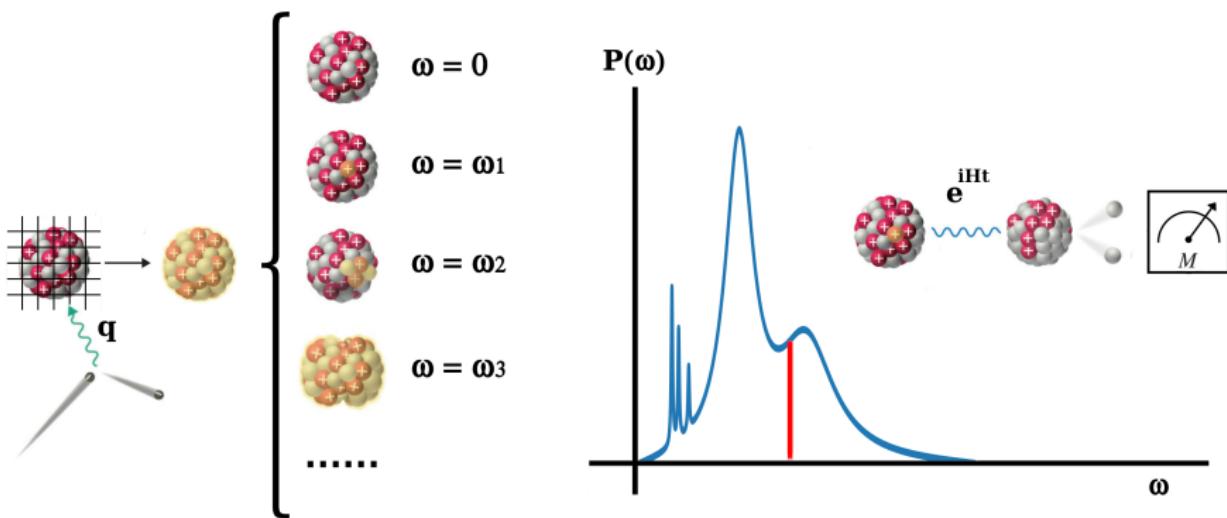
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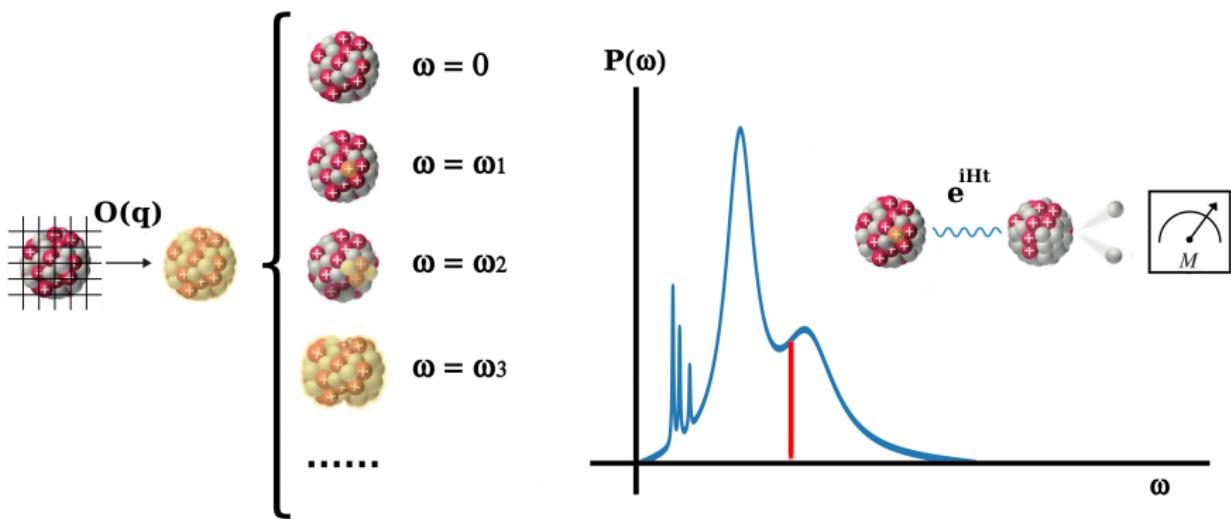
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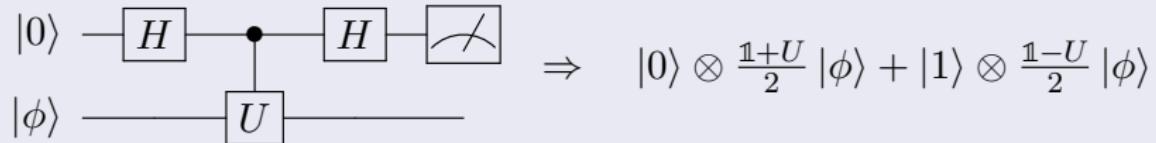


Roggero & Carlson (2019)

Can we apply a non-unitary operation?

YES, but only with some probability

- this can be useful when excitation operator is not unitary

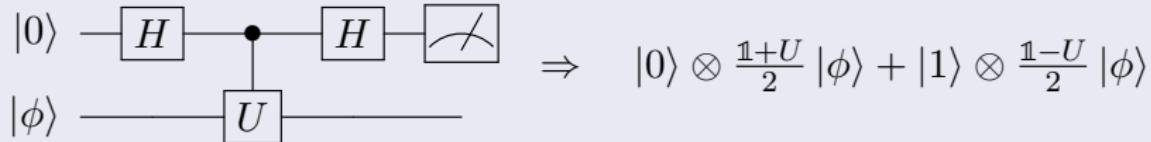


- we will measure $|0\rangle$ with $P_0 = \frac{1}{2} (1 + \mathcal{R}\langle\phi|U|\phi\rangle) \Rightarrow |\phi_0\rangle = \frac{1+U}{2\sqrt{P_0}} |\phi\rangle$

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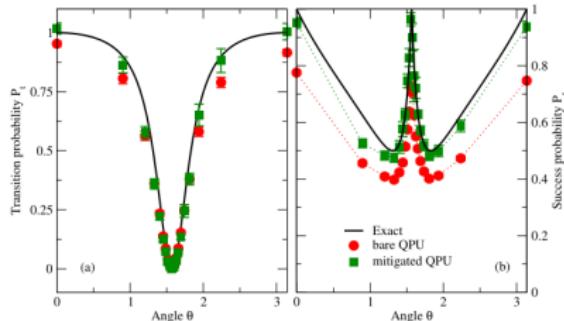
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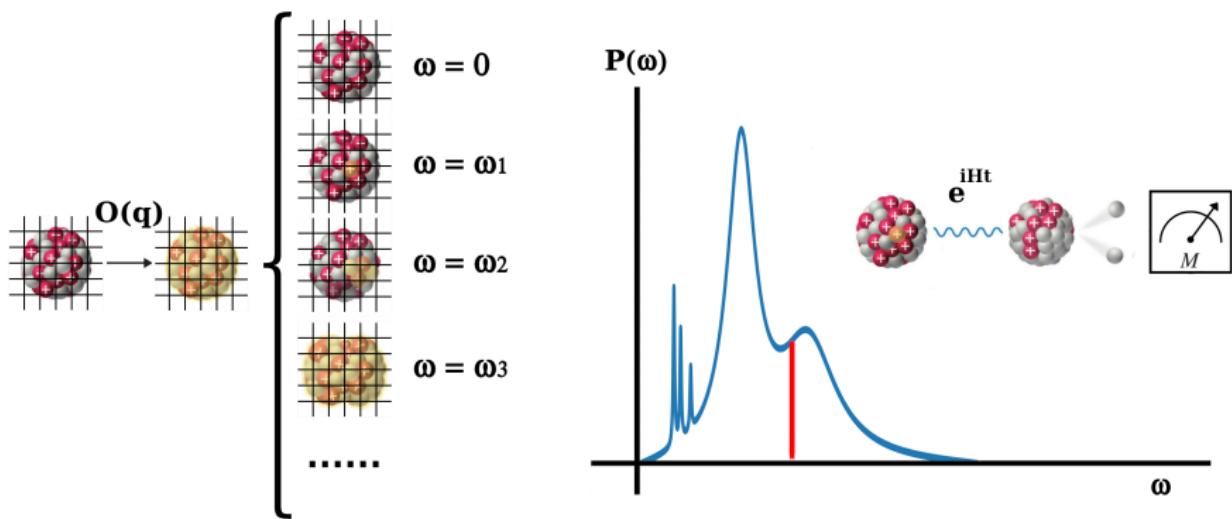
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- generalization to many U possible
- tested on IBM devices for $np \rightarrow d\gamma$
 - simple $M1$ transition $^1S_0 \rightarrow ^3S_1$
 - 2 qubits for $|\phi\rangle$, 3 for controls

AR, C.Gu, A.Baroni, T.Papenbrock (2020)



Quantum algorithm for exclusive processes at fixed q

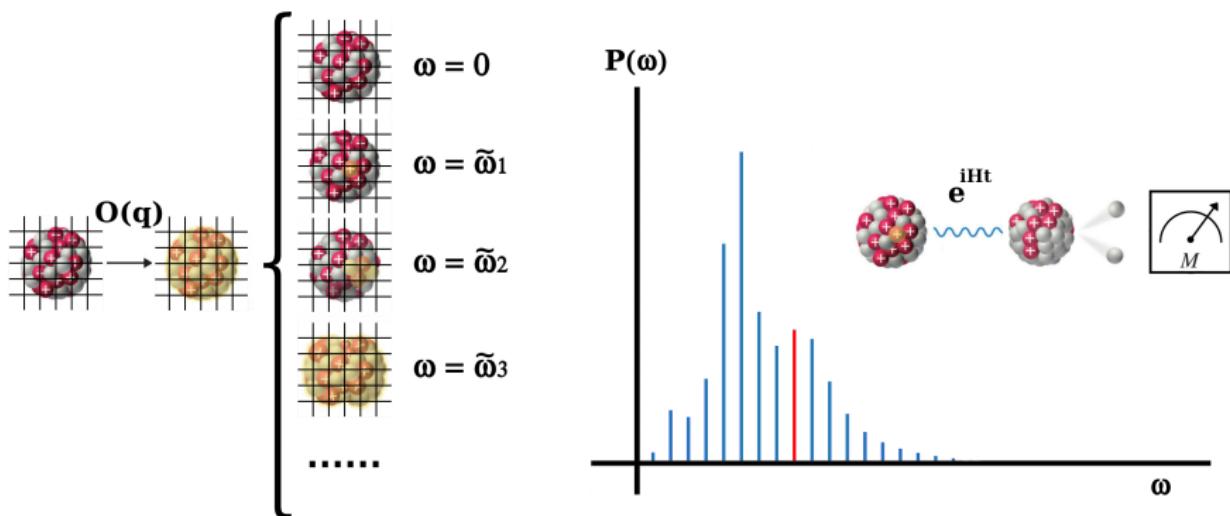
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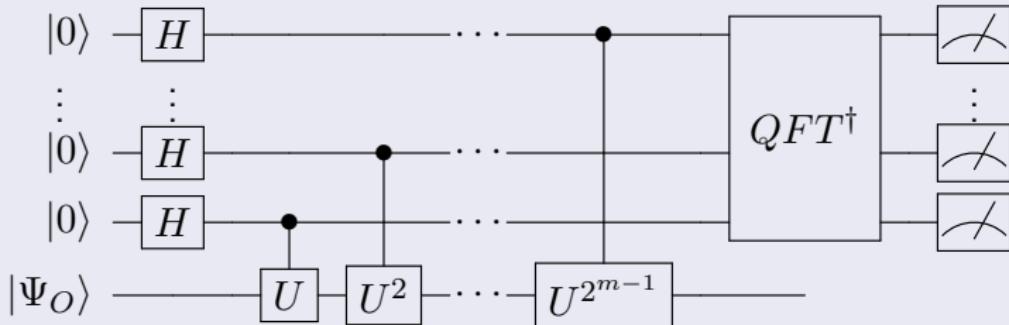
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Roggero & Carlson (2019)

QPE on general states



If we start with the excited state $|\Psi_O\rangle = \sum_j c_j^O |\phi_j\rangle$ we find

$$|\Phi_3\rangle = \sum_j c_j^O \sum_{q=0}^{2^m-1} \left(\frac{1}{2^m} \sum_{k=0}^{2^m-1} \exp\left(i \frac{2\pi k}{2^m} (2^m \phi_j - q)\right) \right) |q\rangle \otimes |\phi_j\rangle$$

The new probability becomes approximately S_O with resolution $\Delta\omega \approx 1/M$

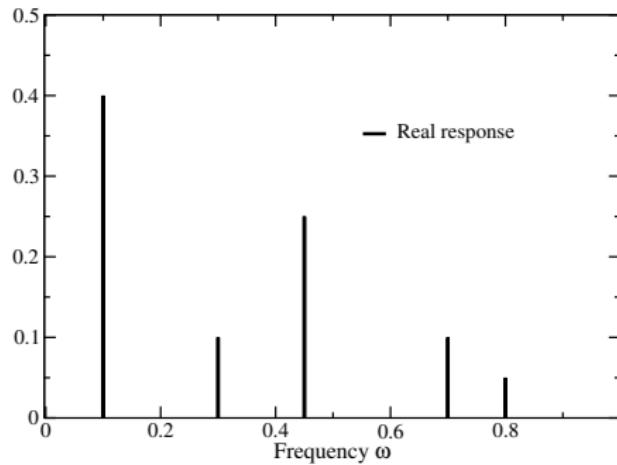
$$P(q) = \frac{1}{M^2} \sum_j |c_j^O|^2 \frac{\sin^2(M\pi(\phi_j - q/M))}{\sin^2(\pi(\phi_j - q/M))} \approx S_O \left(\omega = \frac{q}{M} \right)$$

Approximate response function with QPE

If we start with the excited state $|\Psi_O\rangle = \sum_j c_j^O |\phi_j\rangle$ we find, for $M = 2^m$

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- original response recovered for $M \rightarrow \infty$: $S_O(\omega) = \sum_j |c_j^O|^2 \delta(\phi_j - \omega)$

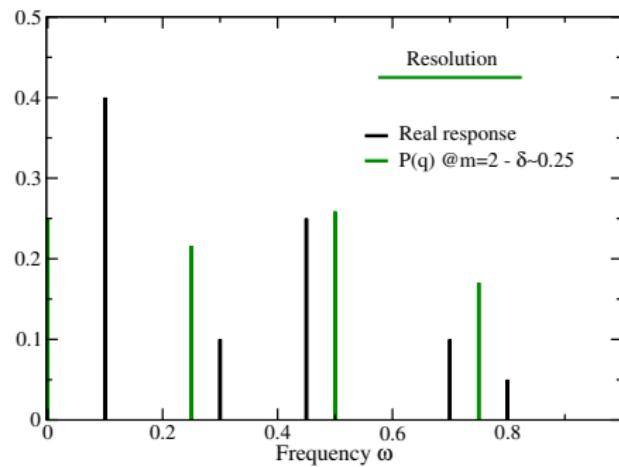


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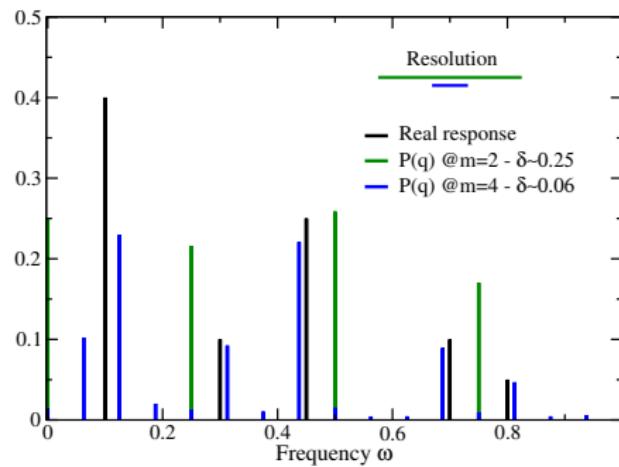


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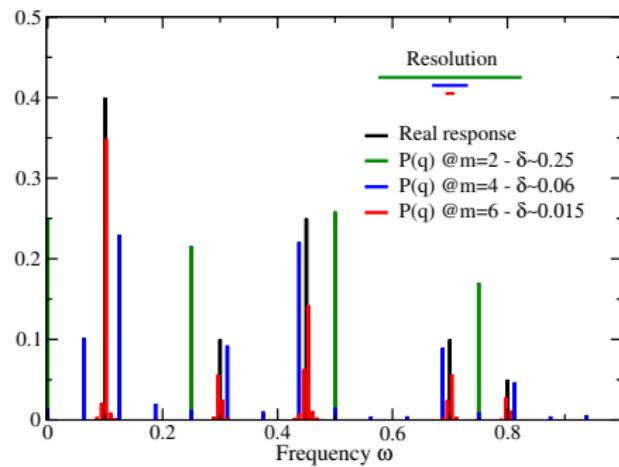


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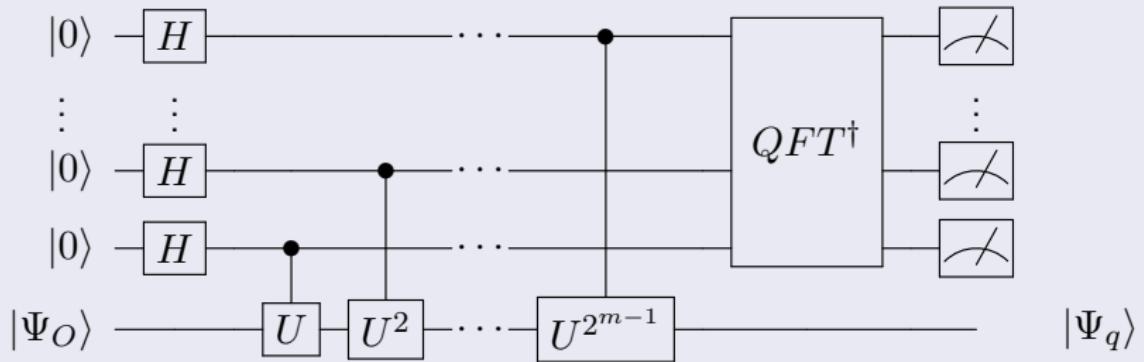
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QPE as state preparation



- before the ancilla measurement we have

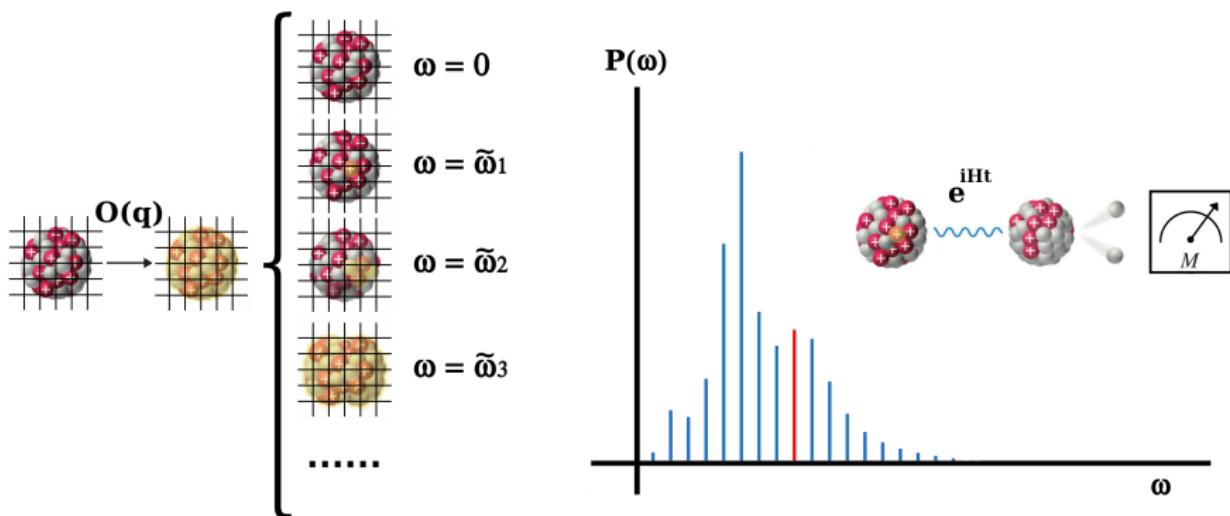
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- after measuring the integer value q the system qubits are left in

$$|\Psi_q\rangle = \frac{1}{M\sqrt{P(q)}} \sum_j c_j^O \frac{\sin(M\pi(\phi_j - \frac{q}{M}))}{\sin(\pi(\phi_j - \frac{q}{M}))} |\phi_j\rangle \approx \sum_{|\phi_j - \frac{q}{M}| \lesssim \frac{1}{M}} c_j^O |\phi_j\rangle$$

Quantum algorithm for exclusive processes at fixed q

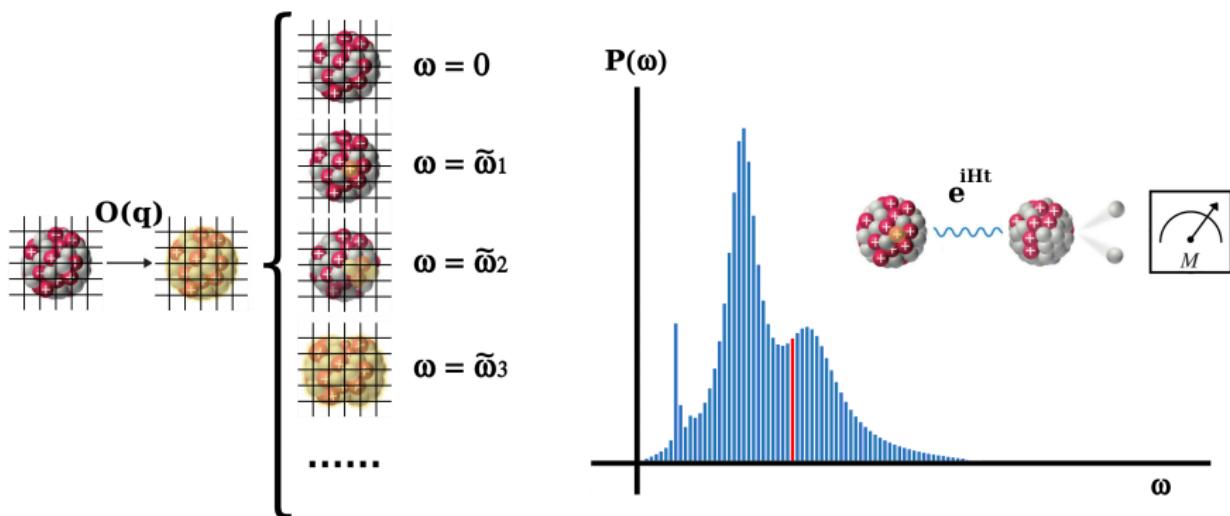
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Roggero & Carlson (2019)

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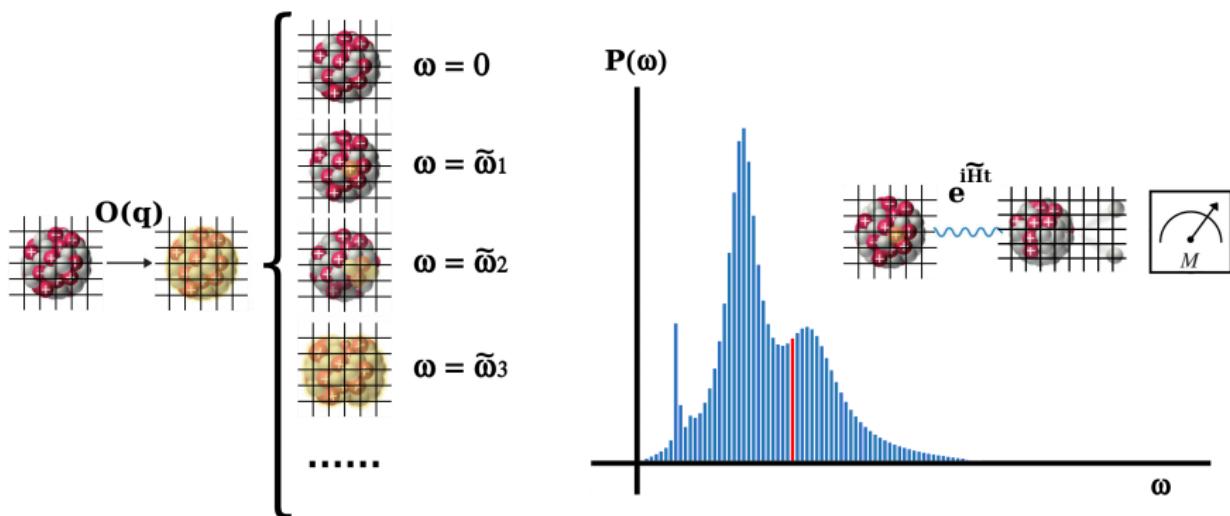
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- further approximate time evolution to let system decay
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Roggero & Carlson (2019)

Progress on time evolution for NP

As a starting point we can use simple nuclear interactions without pions

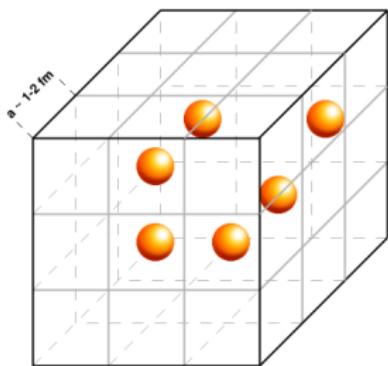


Progress on time evolution for NP

As a starting point we can use simple nuclear interactions without pions



We then place nucleons on spatial lattice to regularize them



Minimal setup

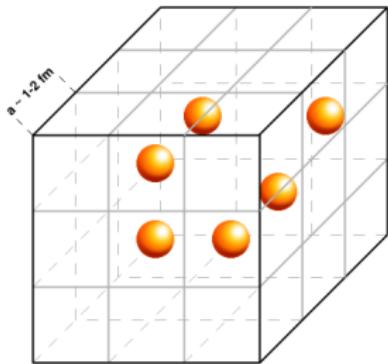
- 10^3 lattice with spacing $a \approx 1 - 2 \text{ fm}$
- 4 spin-isospin states for each particle
 - we need at least 4000 orbitals
- for energy resolution $\Delta\omega$ we need total evolution time $T \approx 1/\Delta\omega$

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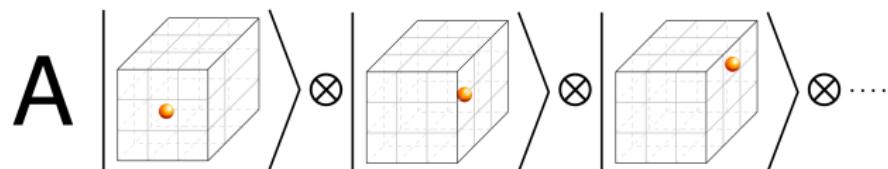
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- $10^{11} - 10^{12}$ operations and ≈ 4000 qubits [Roggero et al. PRD (2020)]
- $10^9 - 10^{11}$ operations and ≈ 6000 qubits [J.Watson et al. arXiv:2312.05344]

Progress on time evolution for NP II

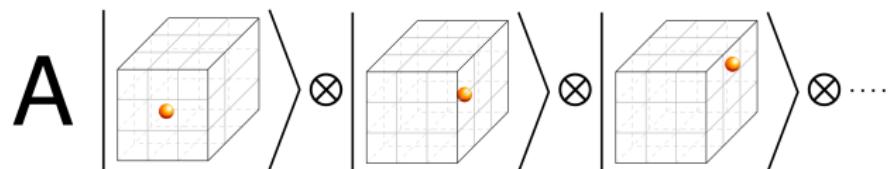
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→ IDEA: why not use first quantization instead?



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Turns out 1st quant. can beat 2nd quant. in both memory and gates

$$C_{2nd}^T = \mathcal{O} \left(\frac{t^{3/2}}{\sqrt{\epsilon}} \sqrt{N} V \right) \quad C_{1st}^T = \mathcal{O} \left(\frac{t^{3/2}}{\sqrt{\epsilon}} N^{7/2} \log^2(V) \right)$$

Roggero, Spagnoli, Lissoni (in prep.)

Progress on time evolution for NP III

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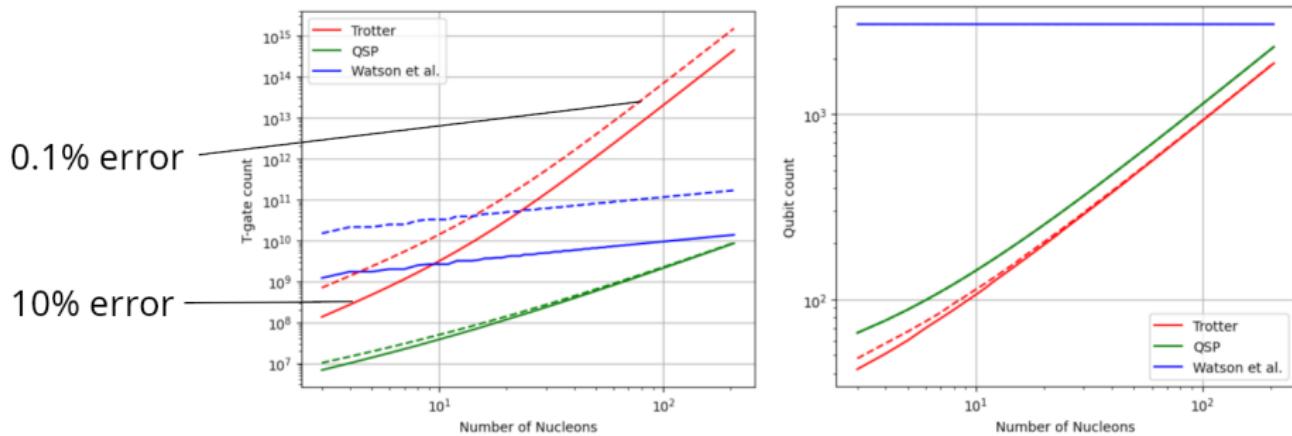
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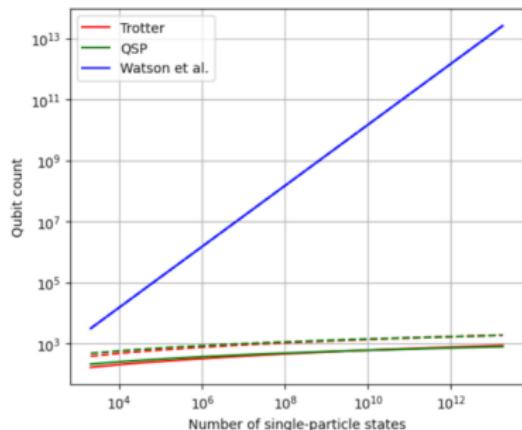
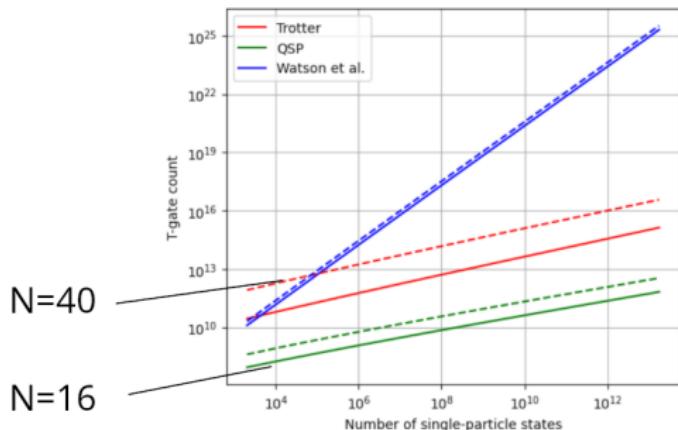
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Recap on nuclear reaction on a quantum computer

- ① need to prepare the target ground state (VQE, QPE, Rodeo, QSP, ...)
- ② apply the excitation operator (LCU seems reasonably good)
- ③ apply an energy projection with resolution $\Delta\omega$ (QPE, Rodeo, GIT, ...)
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Classical computers are really good at doing steps 1 and 2

Quantum computers are really good at doing steps 3 and 4 (maybe 5?)

If we only need inclusive cross section, and thus the response function $S(\omega)$, we can stop at step 3. Most (all?) energy filters do something like

$$\langle \Psi_{ex} | \delta_{\Delta\omega}(H - \omega) | \Psi_{ex} \rangle = \sum_{k=0}^M c_k(\omega, \Delta\omega) \langle \Psi_{ex} | P_k(H) | \Psi_{ex} \rangle$$

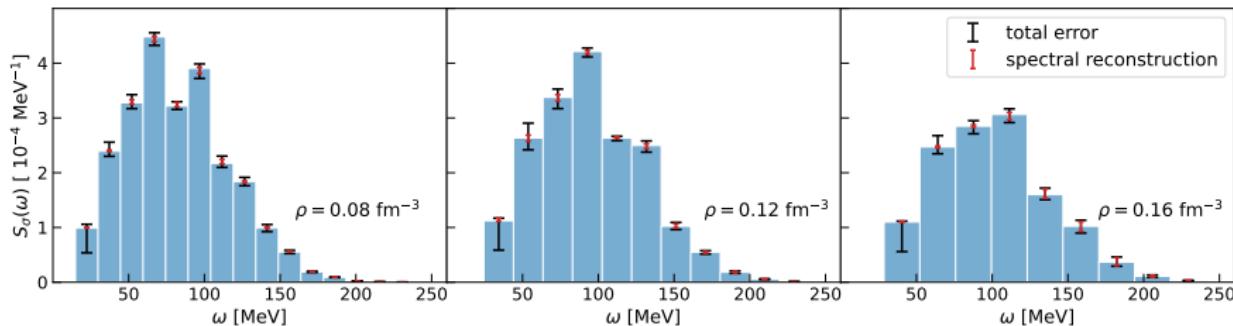
QPE and Rodeo use Fourier polynomials, GIT can also use Chebyshev

Classical simulation of reactions

Since we want Chebyshev moments, why not get them classically instead?

With EOM-CC we can get a reasonably good approximation of them in a very efficient way: 5k moments for $N = 114$ particles in $\approx 7\text{k}$ CPU hours

Sobczyk, Jiang, Roggero (2025)

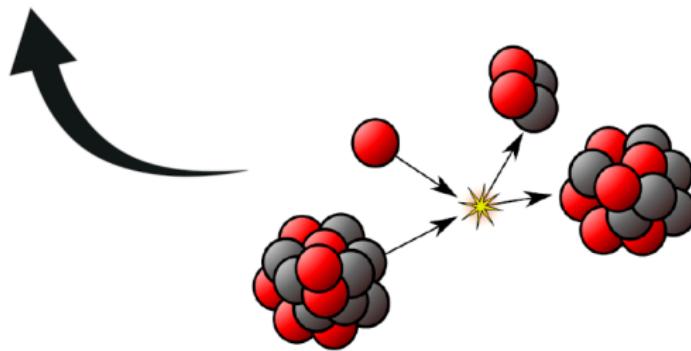
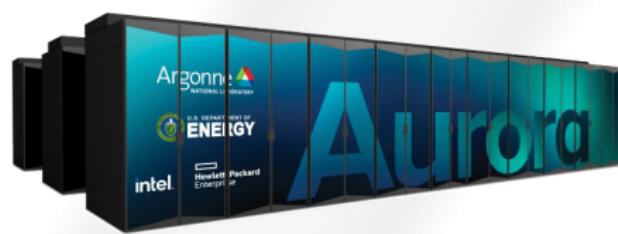


First ab-initio calculation of the frequency dependent spin response of neutron matter with controllable errors, important for ν processes

- exactly same post-processing we would use on a quantum computer

Simulations of nuclear dynamics

Classical simulations

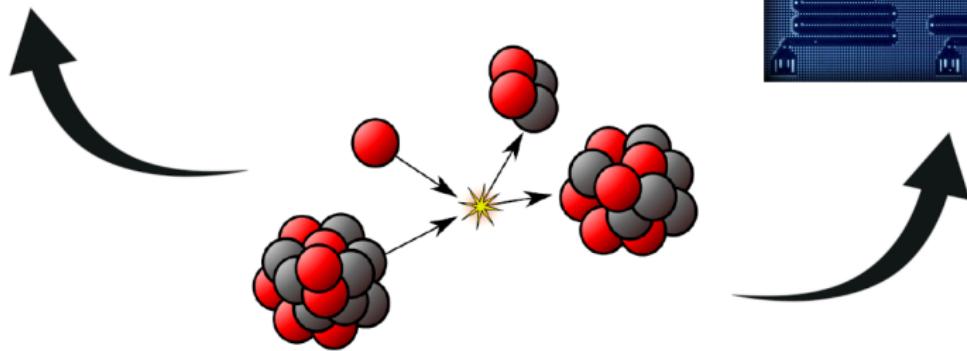
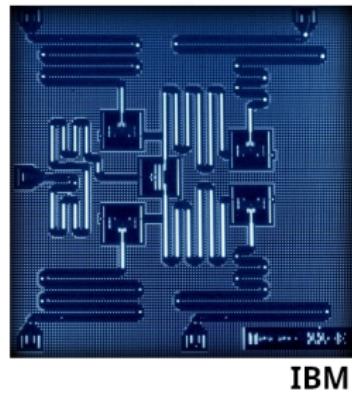


Simulations of nuclear dynamics

Classical simulations

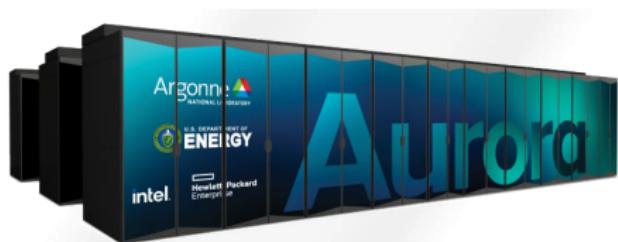


Quantum simulations

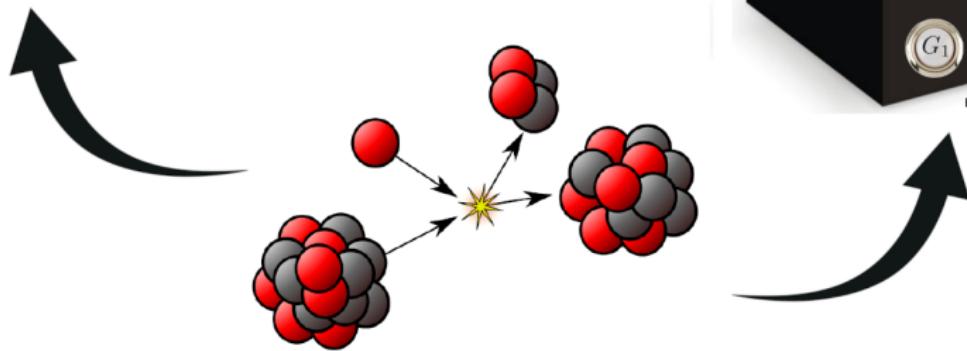


Simulations of nuclear dynamics

Classical simulations



Quantum simulations

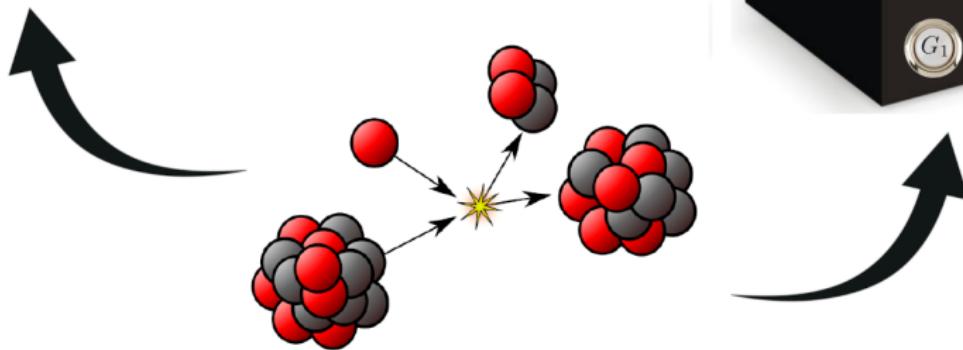


Simulations of nuclear dynamics

Classical simulations



Quantum simulations



Simulations of nuclear dynamics

Classical simulations



Quantum simulations

