

# Upwind Derivative vs Downwind Derivative

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November 26, 2021

# Upwind and downwind method

Let us start from convective equation (square root of wave Eq):

$$\frac{\partial u(x, t)}{\partial t} + a \frac{\partial u(x, t)}{\partial x} = 0, \quad (0.1)$$

we assume  $a > 0$ , it has a solution:

$$u(x, t + \Delta) = u(x - a\Delta, t)$$

Upwind method utilizes upstream's information.

- Upwind ( $a > 0$ ):

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + a \frac{u_j^n - u_{j-1}^n}{\Delta x} = 0$$

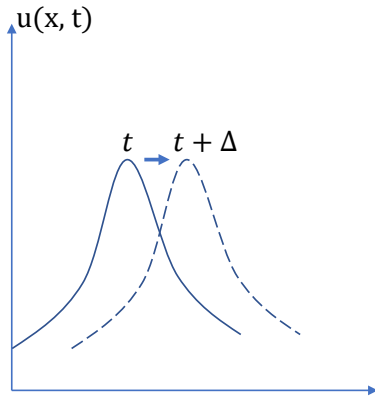
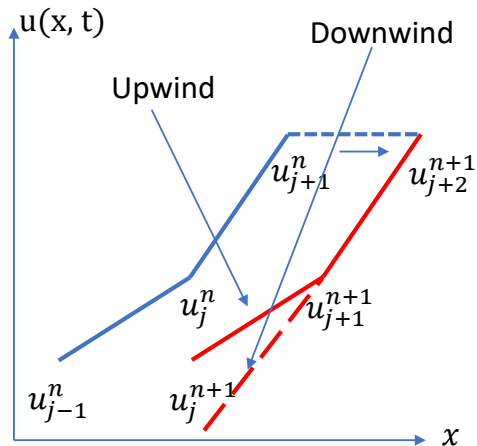
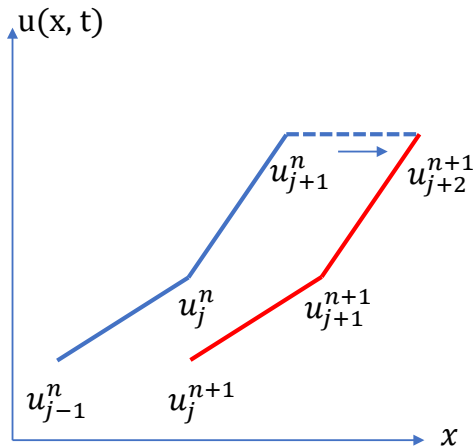


Figure: Convective Equation

# Upwind and downwind method Courant et al. (1952)<sup>1</sup>

Why upwind/downwind?



<sup>1</sup>Courant, Isaacson, and Rees proposed the CIR method

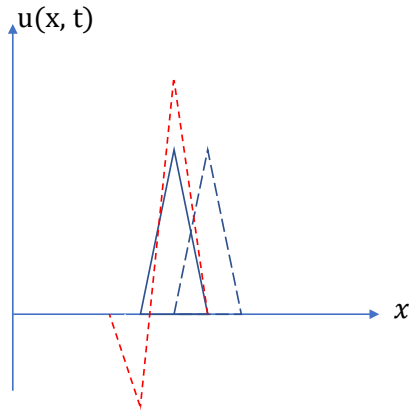
# Upwind and downwind method

## Why upwind/downwind?

We can also consider an extreme case that  $u_{n_1}^0 = 1, u_n^0 = 0, n \neq n_1$ . The wave behave morbidly.

- Intuitively, the information for the part uninfluenced cannot be utilized, they may be wrong.
- Quantitatively, it will lead to an error:  $|(u_{j+1}^n - u_j^n) - (u_j^n - u_{j-1}^n)|$ , for uninfluenced region, first part is generally zero while the second part is a finite one  $O(1)$ , or at least  $O(\Delta x)$ .
- **Remark:** there's another error comes from pseudo diffusion in convective Eq, it can be eliminated by choosing  $\Delta t = \frac{\Delta x}{a}$ .

$$\frac{\partial u(x, t)}{\partial t} + a \frac{\partial u(x, t)}{\partial x} = \frac{1}{2} a \Delta x \left( 1 - a \frac{\Delta t}{\Delta x} \right) \frac{\partial^2 u(x, t)}{\partial x^2}$$



# Upwind and downwind method

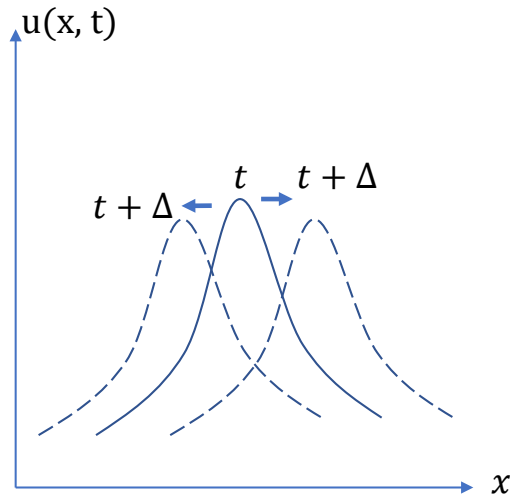
In practice, we consider a PDE:

$$\frac{\partial u}{\partial t} + \frac{f(u)}{\partial x} = 0$$

can be decomposed into:

$$\frac{\partial u}{\partial t} + \frac{f^+}{\partial x} + \frac{f^-}{\partial x} = 0$$

We introduce the upwind to avoid information from upstream getting stucked.



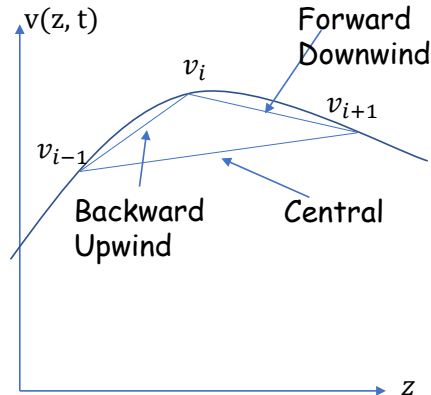
# Ben Moll's HJB and KF

Recall Ben's HJB:

$$\rho v(z) - v'(z)\mu(z) - \frac{1}{2}v''(z)\sigma^2(z) - \pi(z)$$

Standard Algorithm with upwind finite difference:

- **Forward approximation** (downwind):  $v'(z) \approx \frac{v_{j+1}^n - v_j^n}{\Delta z}$
- **Backward approximation** (upwind):  $v'(z) \approx \frac{v_j^n - v_{j-1}^n}{\Delta z}$
- Central difference diverges in convective Eq.



# The HJB in difference method

Explicit:

$$\frac{v_j^{n+1} - v_j^n}{\Delta} + \rho v_j^n = \pi(z) + \mu(z)1_{\mu(z)>0}\partial_F v_j^n + \mu(z)1_{\mu(z)<0}\partial_B v_j^n + \frac{1}{2}\sigma_j^2\partial^2 v_j^n$$

Or equivalently:

$$\frac{v_j^{n+1} - v_j^n}{\Delta} + \rho v_j^n = \pi(z) + \mu(z)1_{\mu(z)>0}\frac{v_{j+1}^n - v_j^n}{\Delta z} + \mu(z)1_{\mu(z)<0}\frac{v_j^n - v_{j-1}^n}{\Delta z} + \frac{1}{2}\sigma_j^2\frac{v_{j+1}^n - 2v_j^n + v_{j-1}^n}{\Delta z^2}$$

```
%Construct matrix
X = -min(mu,0)/dz + sig2/(2*dz2);           positive derivative/
Y = -max(mu,0)/dz + min(mu,0)/dz - sig2/dz2; negative derivative
Z = max(mu,0)/dz + sig2/(2*dz2);
A = spdiags(Y,0,I,I) + spdiags(X(2:I),-1,I,I) +
spdiags([0;Z(1:I-1)],1,I,I);
```

# Matrix Formulation

We reconstruct the matrix formulation of HJB in finite difference:

$$\frac{1}{\Delta}(v^{n+1} - v^n) + \rho v^n = \pi_n + \mathcal{A}^n v^n$$

where:

$$\mathcal{A}^n = \begin{bmatrix} y_1 & z_1 & 0 & \cdots & 0 \\ x_2 & y_2 & z_2 & \cdots & 0 \\ 0 & x_3 & y_3 & z_3 & 0 \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ 0 & \cdots & \cdots & x_N & y_N \end{bmatrix}$$

and

$$x_j = -\frac{\mu(z)1_{\mu < 0}}{\Delta a}, y_j = -\frac{\sigma_j^2}{\Delta z^2}, z_j = \frac{\mu(z)1_{\mu > 0}}{\Delta a}$$

If we have multi variables to deal with, i.e.,  $z, s$ , we should consider  $\mathcal{Z} \otimes \mathcal{S}$  and  $\mathcal{A}_z \otimes \mathcal{A}_s$ , it is not a triangular matrix any more.



# References

**Courant, Richard, Eugene Isaacson, and Mina Rees**, “On the solution of nonlinear hyperbolic differential equations by finite differences,” *Communications on pure and applied mathematics*, 1952, 5 (3), 243–255.