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Modelling Nonlinear optics with the Bloch-Messiah reduction

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Overview

Modelling Nonlinear optics with the Bloch-Messia

- What is nonlinear optics?
- Why do we care about it?
- Gaussian optics
- What I have been doing
- Outlook

Motivation quantum nonlinear optics

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The good

Spontaneous Parametric processes, SPDC, SFWM

- Heralded single photon sources
- Entangled photon pair generation (polarisation, spatial)

Kerr processes

- Self-Phase modulation (SPM), generating Bannana states (CV)
- Cross-Phase modulation (XPM) for sensing

The bad

Spontaneous parametric processes

- Generating more than two photons -> bad for quantum computing
- Understanding filtering
 All Kerr nonlinear processes
 - SPM -> Spectral broadening
 - XPM -> Unwanted phase shifts on single photons due to propagation of the pump

What do we mean by nonlinear optics?

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 Roughly processes that conserve energy but do not conserve photon number.

$$\vec{P} = \chi^{(1)}\vec{E}_1 + \chi^{(2)}\vec{E}_1\vec{E}_2 + \chi^{(3)}\vec{E}_1\vec{E}_2\vec{E}_3 + \dots$$
 (1)

Here we are going to talk about squeezing, i.e SPDC or SFWM, Hamiltonians are then of the form,

$$\hat{H} = A\hat{a}_{S}^{\dagger}\hat{a}_{I}^{\dagger}\hat{a}_{P} + h.c. \tag{2}$$

$$\hat{H} = A\hat{a}_S^{\dagger}\hat{a}_I^{\dagger}\hat{a}_P\hat{a}_P + h.c. \tag{3}$$

Gaussian Optics

Oliver Thomas, Dara McCutcheon, Will McCutcheon Using the undepleted pump approximation we can write the Hamiltonians as terms which are at most quadratic in creation and annihilation operators.

$$\hat{U} = \exp\left[-\frac{i}{\hbar} \left(P \int d\omega_{1} \int d\omega_{2} \ f(\omega_{1}, \omega_{2}) \ \hat{a}_{s}^{\dagger}(\omega_{1}) \hat{a}_{i}^{\dagger}(\omega_{2}) + h.c. \right) \right]$$
Power JSA Signal & Idler

 These are Gaussian transforms, they take Gaussian states to Gaussian states ¹

$$\begin{bmatrix} \vec{b} \\ \vec{b}^{\dagger} \end{bmatrix} = \mathbf{M} \begin{bmatrix} \vec{a} \\ \vec{a}^{\dagger} \end{bmatrix} \tag{5}$$

¹These are linear symplectic transforms which conviently can be written as a matrix

Types of Gaussian transformations

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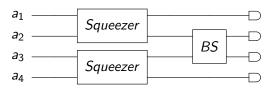


Figure: Two source HOM dip

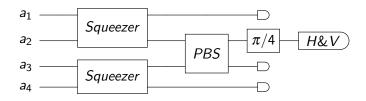


Figure: Type-1 Fusion gate

Schmidt decomposition

Oliver Thomas, Dara McCutcheon, Will We can re-write the Hamiltonian using a Schmidt-decomposition as,

$$P'F(\omega_1,\omega_2) = \sum_k r_k \psi_k(\omega_1) \phi_k(\omega_2)$$
 (6)

Where r_k is the Schmidt number, $\psi \& \phi$ are unitaries.

To solve this numerically we discretize the function and the Schmidt-decomposition is then the Singular value decomposition (SVD) of the JSA (F).

$$P'\mathbf{F}_{(\omega_1,\omega_2)} = \sum_{k} r_k \mathbf{U}_{(\omega_1,k)} \mathbf{V}_{(k,\omega_2)}^{\dagger}$$
 (7)

- ullet with $\psi_k(\omega_1)$ is the k-th row and ω_1 -th column of $oldsymbol{\mathsf{U}}_{(\omega_1,k)}$,
- with $\phi_k(\omega_2)$ is the ω_2 -th row and k-th column of $\mathbf{V}^{\dagger}_{(k,\underline{\omega}_2)}$

Joint Spectral Amplitudes (JSAs)

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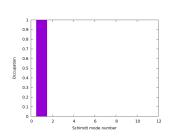


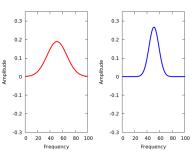
¹Moving to the rotating frame...

Seperable JSAs Schmidt modes

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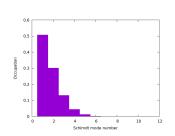
(a) Signal (red) and Idler (blue)

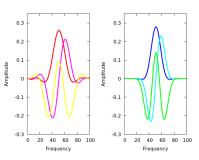
Non-separable JSAs

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Non-separable JSAs Schimdt modes

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(a) Signal (red) and Idler (blue)

Reducing the size of the state-space

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• The Schmidt decomposition lets us represent the system in a finite number of broadband modes $(\psi_k(\omega_1), \phi_k(\omega_2))$

0

Correlations in a HOM dip

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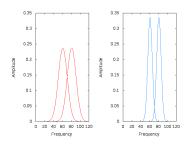


Figure: Two source HOM dip

Two squeezers JSA

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G(4) correlation function

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$$G^{(4)} = \frac{\left\langle \hat{a}_{1}^{\dagger} \hat{a}_{2}^{\dagger} \hat{a}_{3}^{\dagger} \hat{a}_{4}^{\dagger} \hat{a}_{1} \hat{a}_{2} \hat{a}_{3} \hat{a}_{4} \right\rangle}{\left\langle \hat{a}_{1}^{\dagger} \hat{a}_{1} \right\rangle \left\langle \hat{a}_{2}^{\dagger} \hat{a}_{2} \right\rangle \left\langle \hat{a}_{3}^{\dagger} \hat{a}_{3} \right\rangle \left\langle \hat{a}_{4}^{\dagger} \hat{a}_{4} \right\rangle} \tag{8}$$

Where,

$$a_i = \sum_j a_i(\omega_j) \tag{9}$$

Meaning we sum over all of the spectral modes of the spatial modes (1,2,3,4) separately. We end up with,

$$G^{(4)} = 1 - \left(\frac{2 \mid cosh(r) \mid^2}{\mid cosh(r) \mid^2 + \mid sinh(r) \mid^2} sin(\theta) cos(\theta)\right)^2$$
 (10)

G(4) correlation function

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Summary

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Nonlinear
optics with
the
Bloch-Messiah