

Quantum Limits (Noise, Measurement, Metrology, and Information)

Compiled by Onri Jay Benally

December 2025

Conventions and scope. Unless otherwise stated, $S(\omega)$ denotes a (classical) one-sided power spectral density when used in RMS-in-bandwidth formulas (e.g., $\Delta I_{\text{rms}} = \sqrt{S_I \Delta f}$), while $\bar{S}(\omega)$ denotes a *symmetrized* (quantum) noise spectral density. When comparing to one-sided engineering conventions for stationary noise, a common conversion is $S^{(1)}(\omega > 0) = 2\bar{S}(\omega)$ (and similarly for cross-spectra). Noise temperature conventions can differ by constant factors if one uses the full quantum Nyquist expression for the source (this is definitional, not new physics). Several entries are *standard quantum limits* (SQLs): they are rigorous under stated assumptions (e.g., coherent probes, phase-preserving detection, linear response), and they can be surpassed by changing the assumptions (e.g., squeezed states, entanglement, quantum nondemolition measurement, or back-action evasion).

#	Limit/ bound name	Equation (LaTeX)	Category/ context	Key symbols, assumptions, and “how to beat”	Primary ref.
1	Poisson shot-noise current (RMS in bandwidth Δf)	$\Delta I_{\text{rms}} = \sqrt{2qI \Delta f}$	Electronic transport; photon/electron counting; Poisson statistics	I average current, q carrier charge, Δf measurement bandwidth. Equivalent one-sided current-noise PSD: $S_I^{(1)} = 2qI$. <i>Not an ultimate limit</i> : sub-Poisson sources (Fano factor < 1) reduce it.	(standard)
2	Coherent-beam number–phase noise constraint (optical shot noise/ SQL for phase readout)	$S_{\dot{N}\dot{N}}(\omega) S_{\phi\phi}(\omega) \geq \frac{1}{4}, S_{\phi\phi}^{\text{coh}}(\omega) = \frac{1}{4\dot{N}} = \frac{\hbar\omega}{4\bar{P}}$	Quantum optics; homodyne/heterodyne phase sensing; interferometry	\dot{N} photon flux, $\bar{P} = \hbar\omega\dot{N}$ optical power, ϕ phase. <i>Beatable</i> : phase squeezing (reduce $S_{\phi\phi}$) at cost of increased conjugate-quadrature noise; see squeezed-input interferometry.	[7, 6]
3	Phase-preserving amplifier quantum limit (Haus–Caves)	$n_{\text{add}} \geq \frac{1}{2} \ (G \gg 1), T_{\text{N}} \geq \frac{\hbar f}{2k_{\text{B}}}$	Linear amplification with equal treatment of both quadratures (“phase preserving”)	n_{add} added quanta referred to the input; T_{N} input-referred noise temperature; f signal frequency ($\omega = 2\pi f$). The second line matches Eq. (143) in [19]. <i>Definition caveat</i> : alternative noise-temperature conventions (e.g., using the full quantum equilibrium source noise) can introduce constants like $\ln 3$ while leaving the half-quantum added-noise physics unchanged. <i>Evasion</i> : phase-sensitive amplification can add ≈ 0 quanta to <i>one</i> quadrature.	[5, 7, 19]
4	Free-mass standard quantum limit (SQL) for displacement monitoring (time-domain form)	$(\Delta x)_{\text{SQL}} = \sqrt{\frac{2\hbar\tau}{m}}$	Continuous position readout of a free mass; gravitational-wave interferometers	m test-mass, τ averaging/integration time. <i>Interpretation</i> : optimal balance of measurement imprecision and radiation-pressure back-action in a standard (phase-preserving) readout. <i>Beatable</i> : squeezing, variational readout, speed-meters, QND strategies.	[6, 3]

#	Limit/ bound name	Equation (LaTeX)	Category/ context	Key symbols, assumptions, and “how to beat”	Primary ref.
5	Imprecision–back-action (Heisenberg) product for a linear continuous measurement (simplest form)	$\bar{S}_{xx}^{\text{imp}}(\omega) \bar{S}_{FF}^{\text{ba}}(\omega) \geq \frac{\hbar^2}{4}$	General linear measurement; quantum noise tradeoff	$\bar{S}_{xx}^{\text{imp}}$ imprecision noise, \bar{S}_{FF}^{ba} back-action force noise. <i>Assumption:</i> no useful imprecision–back-action correlations are exploited (cf. row 9). <i>Beatable in a narrow sense:</i> correlations and/or QND observables reshape total noise.	[3, 7, 16]
6	Heisenberg scaling for single-parameter phase estimation (entanglement-enabled)	$\Delta\phi \geq \frac{1}{\sqrt{\nu} N}$	Quantum metrology (unitary parameter encoding)	N quanta per probe state; ν independent repetitions. Often contrasted with shot-noise scaling $\Delta\phi \sim 1/\sqrt{\nu N}$. <i>Caveat:</i> decoherence/noise can destroy N -scaling, restoring shot-noise scaling in many realistic settings.	[9]
7	Shot-noise limit (SNL) for phase estimation (separable probes)	$\Delta\phi \geq \frac{1}{\sqrt{\nu N}}$	Classical-like metrology with independent probes (no entanglement)	N uncorrelated particles/photons per repetition; ν repetitions. <i>Beatable:</i> entanglement (Heisenberg scaling) or squeezing (improved constant factor).	[9]
8	Quantum Cramér–Rao bound (QCRB)	$\text{Var}(\hat{\theta}) \geq \frac{1}{\nu F_Q(\theta)}, \Delta\theta \geq \frac{1}{\sqrt{\nu F_Q(\theta)}}$	Ultimate single-parameter estimation bound; quantum Fisher information (QFI)	$F_Q(\theta)$ quantum Fisher information of ρ_θ ; ν repetitions. Shot-noise and Heisenberg scalings are corollaries of upper bounds on F_Q under different resource assumptions.	[4, 9]
9	General quantum noise inequality with imprecision–back-action correlations	$\bar{S}_{xx}^{\text{imp}}(\omega) \bar{S}_{FF}^{\text{ba}}(\omega) - \left \bar{S}_{xF}(\omega) \right ^2 \geq \frac{\hbar^2}{4}$	Linear-response detector theory; correlated quantum noise	\bar{S}_{xF} is the (symmetrized) cross-spectrum. <i>Meaning:</i> correlations can reduce total added noise at selected frequencies (variational measurement), though the inequality still constrains what is possible.	[7]
10	Quantum nondemolition (QND) measurement efficiency bound (measurement vs dephasing)	$\eta \equiv \frac{\Gamma_{\text{meas}}}{\Gamma_\varphi} \leq 1$	Continuous QND qubit readout; measurement chains	Γ_{meas} measurement rate; Γ_φ measurement-induced dephasing rate. $\eta = 1$ is the quantum-limited (ideal) case; technical loss/noise yields $\eta < 1$.	[7]
11	Landauer bound (bit erasure cost)	$Q_{\text{diss}} \geq k_B T \ln 2$	Information thermodynamics; quantum/classical computing hardware	Q_{diss} heat dissipated to a bath at temperature T by logically irreversible erasure. <i>Not “purely quantum”:</i> thermodynamic, but fundamental; relevant to quantum information engines and control costs.	[13]

#	Limit/ bound name	Equation (LaTeX)	Category/ context	Key symbols, assumptions, and “how to beat”	Primary ref.
12	Holevo bound (accessible classical information from a quantum ensemble)	$I_{\text{acc}}(X:Y) \leq \chi,$ $\chi \equiv S(\rho) - \sum_x p_x S(\rho_x), \quad \rho = \sum_x p_x \rho_x$	Quantum Shannon theory; communication; measurement limits	$S(\rho) = -\text{Tr}(\rho \log \rho)$ is von Neumann entropy. <i>Beatable?</i> Not in general without changing the task: χ is the task-dependent upper bound on accessible information for the ensemble.	[11, 20]
13	Helstrom bound (minimum error for binary quantum state discrimination)	$P_{e,\min} = \frac{1}{2} \left(1 - \ p_0 \rho_0 - p_1 \rho_1\ _1 \right)$	Quantum hypothesis testing; optimal measurements	$\ \cdot\ _1$ trace norm; priors p_0, p_1 ; states ρ_0, ρ_1 . <i>Meaning:</i> ultimate performance of any measurement for binary discrimination.	[10, 20]
14	Frequency-domain SQL for linear displacement sensing (uncorrelated imprecision and back-action)	$\bar{S}_{xx}^{\text{SQL}}(\omega) = \hbar \chi_{xx}(\omega) ,$ $S_{xx}^{(1),\text{SQL}}(\omega > 0) = 2\hbar \chi_{xx}(\omega) $	Continuous displacement measurement; optomechanics; interferometry; noise budgeting	$\chi_{xx}(\omega)$ is the mechanical susceptibility (response of x to force). <i>Assumptions:</i> linear response; phase-preserving measurement; no use of imprecision–back-action correlations. <i>Useful special case:</i> for a free mass, $ \chi_{xx}(\omega) = 1/(m\omega^2)$ so $S_{xx}^{(1),\text{SQL}}(\omega > 0) = 2\hbar/(m\omega^2)$. <i>Beatable:</i> variational readout (correlations), back-action evasion, and QND choices of measured observable.	[7, 3]
15	Quantum speed limit (QSL) for state evolution (Mandelstam–Tamm and Margolus–Levitin)	$\tau \geq \max \left(\frac{\pi \hbar}{2 \Delta E}, \frac{\pi \hbar}{2 \langle E \rangle - E_0} \right)$	Quantum control; gate-time lower bounds; bandwidth–time tradeoffs	ΔE is the energy standard deviation in the initial state; E_0 the ground-state energy. <i>Use:</i> lower-bounds how fast one can drive an isolated system to an orthogonal state (or, more generally, a target fidelity). <i>Caveat:</i> open-system dynamics and constrained controls modify practical limits; see modern QSL reviews.	[14, 15, 8]
16	Quantum Chernoff bound (QCB) for asymptotic binary state discrimination	$P_{e,\min}^{(n)} \leq \frac{1}{2} \exp \left[-n \xi_{\text{QCB}} \right],$ $\xi_{\text{QCB}} = -\ln \left(\min_{0 \leq s \leq 1} \text{Tr} [\rho_0^s \rho_1^{1-s}] \right)$	Quantum hypothesis testing; readout classification; sensing error exponents	n is the number of independent copies; ρ_0, ρ_1 are the hypotheses. <i>Meaning:</i> sets the optimal exponential rate at which error decays with repeated trials; complements the single-shot Helstrom bound (row 13).	[2, 20]

#	Limit/ bound name	Equation (LaTeX)	Category/ context	Key symbols, assumptions, and “how to beat”	Primary ref.
17	Quantum Ziv–Zakai bound (QZZB) for mean-square error (Bayesian/global)	$\Sigma \geq \frac{1}{2} \int_0^\infty d\tau \tau \mathcal{V} \int_{-\infty}^\infty dx$ $2 \min[P_X(x), P_X(x + \tau)]$ $\times P_{e,\text{rel}}(x, x + \tau),$ $P_{e,\text{rel}}(x, x + \tau) \geq \frac{1}{2} \left[1 - \sqrt{1 - F(\rho_x, \rho_{x+\tau})} \right]$	Bayesian metrology; global parameter estimation; low-SNR/non-asymptotic regimes	Σ mean-square error (MSE); P_X prior density; \mathcal{V} is a “valley-filling” tightening operation used in ZZB/QZZB constructions. $P_{e,\text{rel}}$ is the minimum binary discrimination error probability between hypotheses x and $x + \tau$ with the <i>relative</i> priors induced by P_X . $F(\rho, \sigma) = (\text{Tr} \sqrt{\sqrt{\rho} \sigma \sqrt{\rho}})^2$ is the Uhlmann fidelity. <i>Why useful</i> : complements the (local) QCRB (row 8) when the likelihood is multimodal, the prior is non-negligible, or performance is far from the asymptotic regime. <i>“Beatable”</i> : only by changing the estimation task (prior), resources, or physical model; otherwise it is a lower bound on achievable mean-square error (MSE).	[18, 21]
18	Multi-parameter quantum Cramér–Rao bound (QCRB; QFIM/SLD form)	$\text{Cov}(\hat{\boldsymbol{\theta}}) \succeq \frac{1}{\nu} \mathbf{F}_Q^{-1},$ $[\mathbf{F}_Q]_{ij} = \text{Tr} \left[\rho_\theta \frac{L_i L_j + L_j L_i}{2} \right],$ $\partial_i \rho_\theta = \frac{L_i \rho_\theta + \rho_\theta L_i}{2}$	Multi-parameter metrology; imaging; simultaneous estimation of noncommuting generators	$\boldsymbol{\theta}$ parameter vector; \mathbf{F}_Q quantum Fisher information matrix (QFIM); L_i symmetric logarithmic derivatives (SLDs). <i>Attainability caveat</i> : unlike the single-parameter case, the SLD-QCRB can be <i>not simultaneously saturable</i> when optimal measurements for different parameters are incompatible. <i>What to use when incompatible</i> : the Holevo Cramér–Rao bound (row 19) is the standard tight benchmark for general multi-parameter problems.	[10, 17]
19	Holevo Cramér–Rao bound (HCRB; tight multi-parameter quantum precision bound)	$\text{Tr} [W \text{Cov}(\hat{\boldsymbol{\theta}})] \geq \frac{1}{\nu} C_H(W),$ $C_H(W) = \min_{\{X_i\}} \left\{ \text{Tr} [W \text{Re } Z] \right.$ $\left. + \left\ \sqrt{W} (\text{Im } Z) \sqrt{W} \right\ _1 \right\},$ $Z_{ij} = \text{Tr}(\rho_\theta X_i X_j)$	Multi-parameter metrology; simultaneous estimation; benchmark when measurements are incompatible	$W \succeq 0$ is a user-chosen weight matrix for the scalar figure of merit $\text{Tr}[W \text{Cov}]$. The minimization is over Hermitian operator sets $\{X_i\}$ satisfying unbiasedness constraints (e.g., $\text{Tr}[\rho_\theta X_i] = 0$ and $\text{Tr}[X_i \partial_j \rho_\theta] = \delta_{ij}$ in common formulations). $\ \cdot\ _1$ is the trace norm. <i>Use</i> : when the SLD-QCRB (row 18) is not jointly attainable, the HCRB is the standard “right” quantum benchmark for what is fundamentally possible.	[12, 1, 17]

Acronym Glossary

Acronym	Expansion
RMS	Root-mean-square
MSE	Mean-square error
PSD	Power spectral density
SQL	Standard quantum limit
SNL	Shot-noise limit
QCRB	Quantum Cramér–Rao bound
QFI	Quantum Fisher information
QFIM	Quantum Fisher information matrix
QND	Quantum nondemolition
TLS	Two-level system
ZPF	Zero-point fluctuations
QSL	Quantum speed limit
QCB	Quantum Chernoff bound
MT	Mandelstam–Tamm
ML	Margolus–Levitin
ZZB	Ziv–Zakai bound
QZZB	Quantum Ziv–Zakai bound
SLD	Symmetric logarithmic derivative
HCRB	Holevo Cramér–Rao bound

References

- [1] Francesco Albarelli et al. “Evaluating the Holevo Cramér–Rao bound for multiparameter quantum metrology”. In: *Physical Review Letters* 123.20 (2019), p. 200503. DOI: 10.1103/PhysRevLett.123.200503. URL: <https://doi.org/10.1103/PhysRevLett.123.200503>.
- [2] Koenraad M. R. Audenaert et al. “Discriminating States: The Quantum Chernoff Bound”. In: *Physical Review Letters* 98.16 (2007), p. 160501. DOI: 10.1103/PhysRevLett.98.160501. URL: <https://doi.org/10.1103/PhysRevLett.98.160501>.
- [3] Vladimir B. Braginsky and Farid Ya. Khalili. *Quantum Measurement*. Cambridge University Press, 1992. ISBN: 9780521419284. URL: <https://doi.org/10.1017/CB09780511622748>.
- [4] Samuel L. Braunstein and Carlton M. Caves. “Statistical distance and the geometry of quantum states”. In: *Physical Review Letters* 72.22 (1994), pp. 3439–3443. DOI: 10.1103/PhysRevLett.72.3439. URL: <https://doi.org/10.1103/PhysRevLett.72.3439>.
- [5] Carlton M. Caves. “Quantum limits on noise in linear amplifiers”. In: *Physical Review D* 26.8 (1982), pp. 1817–1839. DOI: 10.1103/PhysRevD.26.1817. URL: <https://doi.org/10.1103/PhysRevD.26.1817>.
- [6] Carlton M. Caves. “Quantum-mechanical noise in an interferometer”. In: *Physical Review D* 23.8 (1981), pp. 1693–1708. DOI: 10.1103/PhysRevD.23.1693. URL: <https://doi.org/10.1103/PhysRevD.23.1693>.
- [7] A. A. Clerk et al. “Introduction to quantum noise, measurement, and amplification”. In: *Reviews of Modern Physics* 82.2 (2010), pp. 1155–1208. DOI: 10.1103/RevModPhys.82.1155. URL: <https://doi.org/10.1103/RevModPhys.82.1155>.
- [8] Sebastian Deffner and Steve Campbell. “Quantum speed limits: from Heisenberg’s uncertainty principle to optimal quantum control”. In: *Journal of Physics A: Mathematical and Theoretical* 50.45 (2017), p. 453001. DOI: 10.1088/1751-8121/aa86c6. URL: <https://doi.org/10.1088/1751-8121/aa86c6>.
- [9] Vittorio Giovannetti, Seth Lloyd, and Lorenzo Maccone. “Quantum Metrology”. In: *Physical Review Letters* 96.1 (2006), p. 010401. DOI: 10.1103/PhysRevLett.96.010401. URL: <https://doi.org/10.1103/PhysRevLett.96.010401>.
- [10] Carl W. Helstrom. *Quantum Detection and Estimation Theory*. Academic Press, 1976. ISBN: 9780123400505.
- [11] A. S. Holevo. “Bounds for the quantity of information transmitted by a quantum communication channel”. In: *Problems of Information Transmission* 9.3 (1973), pp. 177–183.
- [12] Alexander S. Holevo. *Probabilistic and Statistical Aspects of Quantum Theory*. Edizioni della Normale, 2011.
- [13] R. Landauer. “Irreversibility and Heat Generation in the Computing Process”. In: *IBM Journal of Research and Development* 5.3 (1961), pp. 183–191. DOI: 10.1147/rd.53.0183. URL: <https://doi.org/10.1147/rd.53.0183>.
- [14] L. Mandelstam and I. Tamm. “The uncertainty relation between energy and time in non-relativistic quantum mechanics”. In: *Journal of Physics (USSR)* 9 (1945), pp. 249–254.
- [15] Norman Margolus and Lev B. Levitin. “The maximum speed of dynamical evolution”. In: *Physica D: Nonlinear Phenomena* 120.1–2 (1998), pp. 188–195. DOI: 10.1016/S0167-2789(98)00054-2. URL: [https://doi.org/10.1016/S0167-2789\(98\)00054-2](https://doi.org/10.1016/S0167-2789(98)00054-2).
- [16] Vivishek Sudhir. *Quantum Limits on Measurement and Control of a Mechanical Oscillator*. Springer International Publishing, 2018. DOI: 10.1007/978-3-319-69431-3. URL: <https://doi.org/10.1007/978-3-319-69431-3>.
- [17] Magdalena Szczykulska, Tillmann Baumgratz, and Animesh Datta. “Multi-parameter quantum metrology”. In: *Advances in Physics: X* 1.4 (2016), pp. 621–639. DOI: 10.1080/23746149.2016.1230476. URL: <https://doi.org/10.1080/23746149.2016.1230476>.

- [18] Mankei Tsang. “Ziv–Zakai Error Bounds for Quantum Parameter Estimation”. In: *Physical Review Letters* 108.23 (2012), p. 230401. DOI: 10.1103/PhysRevLett.108.230401. arXiv: 1111.3568 [quant-ph]. URL: <https://doi.org/10.1103/PhysRevLett.108.230401>.
- [19] Florian Vigneau et al. “Probing quantum devices with radio-frequency reflectometry”. In: *Applied Physics Reviews* 10.2 (2023), p. 021305. DOI: 10.1063/5.0088229. URL: <https://doi.org/10.1063/5.0088229>.
- [20] Mark M. Wilde. *Quantum Information Theory*. Cambridge University Press, 2013. ISBN: 9781107034259.
- [21] Jacob Ziv and Moshe Zakai. “Some lower bounds on signal parameter estimation”. In: *IEEE Transactions on Information Theory* 15.3 (1969), pp. 386–391. DOI: 10.1109/TIT.1969.1054301. URL: <https://doi.org/10.1109/TIT.1969.1054301>.

License

Copyright © 2025 Onri Jay Benally

This document is licensed under the **Creative Commons Attribution 4.0 International License (CC BY 4.0)**.

You are free to:

- **Share** — copy and redistribute the material in any medium or format.
- **Adapt** — remix, transform, and build upon the material for any purpose, even commercially.

Under the following terms:

- **Attribution** — You must give appropriate credit to the author (*Onri Jay Benally*), provide a link to the license, and indicate if changes were made. You may do so in any reasonable manner, but not in any way that suggests the licensor endorses you or your use.

To view a copy of this license, visit:

<http://creativecommons.org/licenses/by/4.0/>